

Intermediate Quantum Mechanics

Lecture 13 Notes (3/4/15)

Simple Harmonic Oscillator II

Matrix elements of \hat{a} and \hat{a}^\dagger

- We'll now calculate the constants that result when the raising and lowering operators act on states

$$\begin{aligned} \hat{a} |E_n\rangle &= c_n |E_n\rangle & \Rightarrow & \langle E_n | \hat{a}^\dagger = c_n^* \\ \langle E_n | \hat{a}^\dagger \hat{a} |E_n\rangle &= n & \Rightarrow & |c_n|^2 = n \end{aligned}$$

We can absorb an overall phase into the eigenstates and make c_n real. Then $c_n = \sqrt{n}$.

- We have then

$$\begin{aligned} \hat{a} |E_n\rangle &= \sqrt{n} |E_{n-1}\rangle & \hat{a}^\dagger |E_{n-1}\rangle &= \sqrt{n} |E_n\rangle & \text{or} & \hat{a}^\dagger |E_n\rangle &= \sqrt{n+1} |E_{n+1}\rangle \\ \langle E_{n-1} | \hat{a} |E_n\rangle &= \langle E_n | \hat{a}^\dagger |E_{n-1}\rangle = \sqrt{n} & \langle E_{n+1} | \hat{a}^\dagger |E_n\rangle &= \sqrt{n+1} \end{aligned}$$

Ground eigenstate wave function

- We will now find the ground state wave function.

$$\hat{a} |E_0\rangle = |\emptyset\rangle \Rightarrow (\hat{x}' + i\hat{p}') |E_0\rangle = |\emptyset\rangle \Rightarrow \langle x | (\hat{x}' + i\hat{p}') |E_0\rangle = \langle x | \emptyset \rangle$$

This gives a simple first order differential equation:

$$\left(x' + \frac{d}{dx'} \right) u_0(x') = 0$$

- It's easy to see that the solution to this equation is:

$$u_0(x') = A_0' e^{-x'^2/2}$$

Using $x' = \sqrt{\frac{m\omega}{\hbar}} x$, we have: $u_0(x) = A_0 e^{-m\omega x^2/2\hbar}$

First excited state

- We now solve for the first excited state $u_1(x)$.

$$\hat{a}^\dagger |E_0\rangle = |E_1\rangle \Rightarrow (\hat{x}' - i\hat{p}') |E_0\rangle = |E_1\rangle \Rightarrow \langle x | (\hat{x}' - i\hat{p}') |E_0\rangle = \langle x | E_1 \rangle$$

This gives the first order differential equation:

$$\left(x' - \frac{d}{dx'} \right) u_0(x') = u_1(x')$$

- From the differential equation for $u_0(x')$ above we have that:

$$\frac{d}{dx'} u_0(x') = -x' u_0(x')$$

which gives: $u_1(x') = 2x' u_0(x') = A_1' x' e^{-x'^2/2}$

Using $x' = \sqrt{\frac{m\omega}{\hbar}} x$, we have: $u_1(x) = A_1 x e^{-m\omega x^2/2\hbar}$

- We can now get $u_2(x')$ by solving

$$\left(x' - \frac{d}{dx'}\right) u_1(x') = u_2(x')$$

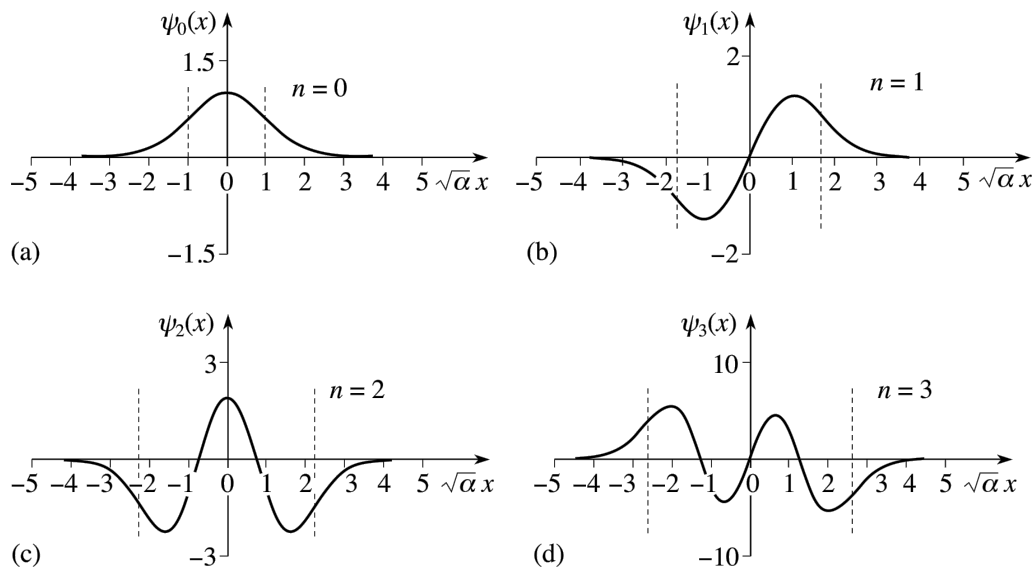
and continue on getting all of the $u_n(x')$'s. We would find:

$$u_n(x') = H_n(x') e^{-x'^2/2}$$

where $H_n(x')$ is the Hermite polynomial of order n . The even n Hermite polynomials have only even power terms in x' and the odd n Hermite polynomials have only odd power terms in x' . If you are interested, you can find a listing of the Hermite polynomials on wikipedia or in almost any quantum mechanics text book.

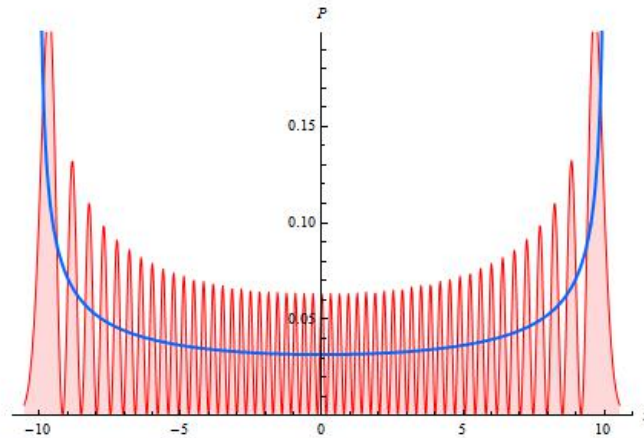
Plots of wave functions and probability distributions

- The plots below show the $n = 0$, $n = 1$, $n = 2$ and $n = 3$ wave functions.



- The plot below shows the probability distribution $u_n^*(x)u_n(x)$ for $n = 50$. The blue curve is the classical expectation. It is the position probability distribution that you would get if you randomly took photographs of a mass on the end of a

spring. For a very high energy states as $n \rightarrow \infty$, the position distribution of the quantum harmonic oscillator approaches the classical case.



Calculating SHO matrix elements

- Using:

$$\hat{x}' = \frac{1}{\sqrt{2}} (\hat{a}^\dagger + \hat{a}) \quad \text{and} \quad \hat{p}' = \frac{i}{\sqrt{2}} (\hat{a}^\dagger - \hat{a})$$

and the orthonormality of the energy eigenstates, we can evaluate the matrix elements of any powers of \hat{x}' or \hat{p}' between energy eigenstates.

- For example,

$$\begin{aligned} \langle E_m | \hat{x}' | E_n \rangle &= \frac{1}{\sqrt{2}} \langle E_m | (\hat{a}^\dagger + \hat{a}) | E_n \rangle \\ &= \frac{1}{\sqrt{2}} \left(\sqrt{n+1} \langle E_m | E_{n+1} \rangle + \sqrt{n} \langle E_m | E_{n-1} \rangle \right) \\ &= \sqrt{\frac{n+1}{2}} \delta_{m,n+1} + \sqrt{\frac{n}{2}} \delta_{m,n-1} \end{aligned}$$