

HW 7

$$\begin{aligned} \text{G 4.1 (a)} \quad [x, p_x] f(\vec{r}) &= \frac{\hbar}{i} \left(x \frac{\partial}{\partial x} - \frac{\partial}{\partial x} x \right) f(\vec{r}) \\ &= \frac{\hbar}{i} \left(x \frac{\partial}{\partial x} f(\vec{r}) - f(\vec{r}) - x \frac{\partial}{\partial x} f(\vec{r}) \right) \\ &= i\hbar f(\vec{r}) \quad \therefore [x, p_x] = i\hbar \\ \text{similarly, } [y, p_y] &= [z, p_z] = i\hbar \end{aligned}$$

Obviously, $[p_x, x] = -[x, p_x]$

So $[r_i, p_i] = -[p_i, r_i] = i\hbar$

$$\begin{aligned} \text{Now } [x, p_y] f(\vec{r}) &= \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - \frac{\partial}{\partial y} x \right) f(\vec{r}) \\ &= \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} f(\vec{r}) - x \frac{\partial}{\partial y} f(\vec{r}) \right) = 0 \end{aligned}$$

So $[x, p_y] = 0$, and similarly
 $[r_i, p_j] = 0$ if $i \neq j$

Combining all these, $[r_i, p_j] = -[p_i, r_j] = i\hbar \delta_{ij}$

Since $xy = yx$, $xz = zx$ and $yz = zy$
trivially $[r_i, r_j] = 0$.

Also, since $\frac{\partial}{\partial x} \frac{\partial}{\partial y} = \frac{\partial}{\partial y} \frac{\partial}{\partial x}$, etc,

$[p_i, p_j] = 0$

(b) Eq. 3.71 should be valid for any dimension s , because we did not rely on any particular dimension in the derivation.

$$\text{Then } \frac{d\langle x \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{H}, x] \rangle$$

$$[\hat{H}, x] = \left[\frac{p_x^2 + p_y^2 + p_z^2}{2m} + V(\vec{r}), x \right]$$

$$\text{Since } [p_y^2, x] = [p_z^2, x] = 0$$

$$\text{and } [V(x, y, z), x] = 0,$$

$$[\hat{H}, x] = \frac{1}{2m} [p_x^2, x] = \frac{1}{2m} (p_x [p_x, x] + [p_x, x] p_x) = \frac{1}{2m} (-2i\hbar p_x) = -i\hbar \frac{p_x}{m}$$

$$\therefore \frac{d\langle x \rangle}{dt} = \frac{i}{\hbar} (-i\hbar \frac{\langle p_x \rangle}{m}) = \frac{\langle p_x \rangle}{m}$$

$$\text{Similarly } \frac{d\langle y \rangle}{dt} = \frac{\langle p_y \rangle}{m} \text{ and } \frac{d\langle z \rangle}{dt} = \frac{\langle p_z \rangle}{m}$$

$$\therefore \underline{\frac{d\langle \vec{r} \rangle}{dt} = \frac{\langle \vec{p} \rangle}{m}}$$

$$\text{Now } \frac{d\langle p_x \rangle}{dt} = \frac{i}{\hbar} \langle [\hat{H}, p_x] \rangle$$

$$[\hat{H}, p_x] = \left[\frac{p_x^2 + p_y^2 + p_z^2}{2m} + V(\vec{r}), p_x \right] \\ = [V(\vec{r}), p_x] = i\hbar \frac{\partial}{\partial x} V(\vec{r})$$

$$\text{Thus } \frac{d\langle P_x \rangle}{dt} = \left\langle -\frac{\partial V}{\partial x} \right\rangle$$

$$\text{Similarly } \frac{d\langle P_y \rangle}{dt} = \left\langle -\frac{\partial V}{\partial y} \right\rangle \text{ and } \frac{d\langle P_z \rangle}{dt} = \left\langle -\frac{\partial V}{\partial z} \right\rangle$$

$$\text{Thus } \frac{d\langle \vec{P} \rangle}{dt} = \langle -\nabla V \rangle$$

(c) From Eq 3.62,

$$\sigma_A^2 \sigma_B^2 \geq \left(\frac{1}{2i} \langle [\hat{A}, \hat{B}] \rangle \right)^2$$

$$\sigma_x^2 \sigma_{p_x}^2 \geq \left(\frac{1}{2i} \langle [X, P_x] \rangle \right)^2$$

$$= \left(\frac{\hbar}{2} \right)^2$$

$$\Rightarrow \sigma_x \sigma_{p_x} \geq \frac{\hbar}{2}$$

$$\text{Similarly } \sigma_y \sigma_{p_y} \geq \frac{\hbar}{2} \text{ and } \sigma_z \sigma_{p_z} \geq \frac{\hbar}{2}$$

Since $[X, P_y] = 0$, etc,

there is no restriction on $\sigma_{x_i} \sigma_{p_j}$ if $i \neq j$.

4.2 T.I.S.E. within the box is

$$\frac{-\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi(x, y, z)$$

$$= E \psi(x, y, z)$$

With $\psi(x, y, z) = X(x) Y(y) Z(z)$

$$\left(\frac{d^2}{dx^2}X\right)YZ + \left(\frac{d^2}{dy^2}Y\right)XZ + \left(\frac{d^2}{dz^2}Z\right)XY = -\frac{2mE}{\hbar^2}XYZ$$

With $\sqrt{\frac{2mE}{\hbar^2}} \equiv k$ and dividing both sides by XYZ , we have

$$\frac{1}{X} \frac{d^2X}{dx^2} + \frac{1}{Y} \frac{d^2Y}{dy^2} + \frac{1}{Z} \frac{d^2Z}{dz^2} = -k^2$$

For this equation to be valid for any (x,y,z) each term should be a constant

$$\frac{1}{X} \frac{d^2X}{dx^2} = -k_x^2$$

$$\frac{1}{Y} \frac{d^2Y}{dy^2} = -k_y^2, \text{ and } \frac{1}{Z} \frac{d^2Z}{dz^2} = -k_z^2$$

$$\text{Then } k_x^2 + k_y^2 + k_z^2 = k^2.$$

Now the boundary condition requires

$$X(0)=X(a)=Y(0)=Y(a)=Z(0)=Z(a)=0$$

The general solution for each single variable differential equation is

$$X(x) = A \sin(k_x x) + B \cos(k_x x)$$

$$\Rightarrow X(0)=X(a)=0 \Rightarrow B=0, k_x a = n_x \pi$$

$$\Rightarrow X(x) = A \sin\left(\frac{n_x \pi}{a} x\right)$$

$$\int_0^a |X(x)|^2 dx = 1 \quad \text{with } n_x = 1, 2, 3, \dots \Rightarrow A = \sqrt{\frac{2}{a}}$$

Using similar methods for $Y(y)$ and $Z(z)$,

The full stationary states are

$$\psi_{(n_x, n_y, n_z)}(x, y, z) = \left(\frac{2}{a}\right)^{\frac{3}{2}} \sin\left(\frac{n_x \pi}{a} x\right) \sin\left(\frac{n_y \pi}{a} y\right) \sin\left(\frac{n_z \pi}{a} z\right)$$

$$\text{with } n_x = 1, 2, \dots$$

$$n_y = 1, 2, \dots, \quad n_z = 1, 2, \dots$$

The corresponding energies are

$$\begin{aligned} E &= \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2}{2m} \left(\left(\frac{n_x \pi}{a}\right)^2 + \left(\frac{n_y \pi}{a}\right)^2 + \left(\frac{n_z \pi}{a}\right)^2 \right) \\ &= \frac{\hbar^2 \pi^2}{2ma^2} (n_x^2 + n_y^2 + n_z^2) \end{aligned}$$

(b) (n_x, n_y, n_z) combination in the order of increasing energies are

$$(1, 1, 1) : d=1 \quad E_1 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 3$$

, here "d" stands for degeneracy.

$$\left. \begin{array}{l} (2, 1, 1) \\ (1, 2, 1) \\ (1, 1, 2) \end{array} \right\} d=3, \quad E_2 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 6$$

$$\left. \begin{array}{l} (2, 2, 1) \\ (2, 1, 2) \\ (1, 2, 2) \end{array} \right\} d=3, \quad E_3 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 9$$

$$\left. \begin{array}{l} (3, 1, 1) \\ (1, 3, 1) \\ (1, 1, 3) \end{array} \right\} d=3, \quad E_4 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 11$$

$$(2, 2, 2), \quad d=1, \quad E_5 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 12$$

$$\left. \begin{array}{l} (3, 2, 1) \\ (3, 1, 2) \\ (2, 3, 1) \\ (2, 1, 3) \\ (1, 3, 2) \\ (1, 2, 3) \end{array} \right\} d=6, \quad E_6 = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 14$$

$$E_7 \Rightarrow (3, 2, 2) \quad (9+4+4=17)$$

$$E_8 \Rightarrow (4, 1, 1) \quad (16+1+1=18)$$

$$E_9 \Rightarrow (3, 3, 1) \quad (9+9+1=19)$$

$$E_{10} \Rightarrow (4, 2, 1) \quad (16 + 4 + 1 = 21)$$

$$E_{11} \Rightarrow (3, 3, 2) \quad (9 + 9 + 4 = 22)$$

$$E_{12} \Rightarrow (4, 2, 2) \quad (16 + 4 + 4 = 24)$$

$$E_{13} \Rightarrow (4, 3, 1) \quad (16 + 9 + 1 = 26)$$

$$E_{14} \Rightarrow \begin{cases} (3, 3, 3) & (9 \times 3 = 27) \\ (5, 1, 1) & (25 + 1 + 1 = 27) \end{cases}$$

Degeneracy of $E_{14} = 1 + 3 = 4$
: unlike the other energy values,
 E_{14} is composed of two number
combinations $(3, 3, 3)$ and $(5, 1, 1)$.

Prob. 4.13

$$(a) \quad \langle r \rangle_{\psi_{100}} = \langle R_{10} | r | R_{10} \rangle \underbrace{\langle Y_0^0 | Y_0^0 \rangle}_{=1}$$

$\uparrow$
 $R_{10} Y_0^0$

$$= \int_{r=0}^{\infty} r^3 R_{10}^2 dr = 4 \bar{a}^3 \int_0^{\infty} r^3 e^{-2r/a} dr$$

$$= 4 \bar{a}^3 3! \left(\frac{a}{2}\right)^4 = \frac{3}{2} a$$

$$\langle r^2 \rangle_{\psi_{100}} = \langle R_{10} | r^2 | R_{10} \rangle = \int_0^{\infty} r^4 R_{10}^2 dr$$

$$= 4 \bar{a}^3 \int_0^{\infty} r^4 e^{-2r/a} dr = 4 \bar{a}^3 4! \left(\frac{a}{2}\right)^5$$

$$= \underline{3a^2}$$

(b) $\langle x \rangle_{\psi_{100}} = 0$, because ψ_{100} is spherically symmetric and " x " is "odd" w.r.t. $x \rightarrow -x$.

$$\langle x^2 \rangle_{\psi_{100}} = \frac{1}{3} \langle r^2 \rangle_{\psi_{100}} = \underline{a^2}$$

$$(c) \langle x^2 \rangle_{\psi_{211}} = \langle r^2 \sin^2 \theta \cos^2 \phi \rangle_{\psi_{211}}$$

$$= \int_0^{\infty} r^4 |R_{21}|^2 dr \iint |Y_1^1|^2 \sin^3 \theta \cos^2 \phi d\theta d\phi$$

$$\text{Radial part} = \frac{\bar{a}^3}{24} \int_0^{\infty} r^4 \frac{r^2}{a^2} \exp(-r/a) dr$$

$$= \frac{a^{-5}}{24} \int_0^{\infty} r^6 \exp\left(-\frac{r}{a}\right) dr$$

$$= \frac{a^{-5}}{24} 6! a^7 = \underline{30 a^2}$$

$$\text{Angular part} = \frac{3}{8\pi} \int_0^{\pi} \sin^5 \theta d\theta \int_0^{2\pi} \cos^2 \phi d\phi$$

$$= \frac{3}{8\pi} \left(2 \cdot \frac{2 \cdot 4}{1 \cdot 3 \cdot 5} \right) \left(\frac{1}{2} \cdot 2\pi \right)$$

$$= \frac{2}{5}$$

$$\therefore \langle r^2 \rangle \psi_{211} = \underline{12 a^2}$$

$$4.14 \quad \psi(r) = \psi_{100}(r) = R_{10} Y_0^0$$

$P(r) dr$ = probability of finding the electron between r and $r+dr$

$$= \int_0^{\pi} \int_0^{2\pi} |\psi_{100}(r)|^2 d^3 \vec{r} = |R_{10}|^2 r^2 dr \cdot \langle Y_0^0 | Y_0^0 \rangle$$

$$= |R_{10}|^2 r^2 dr$$

$$\Rightarrow P(r) = |R_{10}|^2 r^2 = 4 a^3 \exp\left(-\frac{2r}{a}\right) \cdot r^2$$

Maximum of $P(r)$ occurs when

$$\frac{dP(r)}{dr} = 0 \Rightarrow -\frac{2}{a} \exp\left(-\frac{2r}{a}\right) r^2 + 2 \exp\left(-\frac{2r}{a}\right) r = 0$$

$$\Rightarrow r \left(-\frac{r}{a} + 1 \right) = 0$$

$$\Rightarrow r=0 \text{ or } r=a$$

$$P(r=0) = 0, \quad P(r=a) = 4\tilde{a}^3 e^{-2} a^2 = \frac{4}{ae^2}$$

$\Rightarrow P(r=0)=0$ must be the minimum and $P(r=a) = \frac{4}{ae^2}$ must be the maximum

Thus the most probable value of r is $r=a$

4.15

$$\begin{aligned} (a) \quad \Psi(\vec{r}, t) &= \frac{1}{\sqrt{2}} (\psi_{2,1,1} e^{-i\frac{E_2}{\hbar}t} + \psi_{2,1,-1} e^{-i\frac{E_2}{\hbar}t}) \\ &= \frac{1}{\sqrt{2}} (\psi_{2,1,1} + \psi_{2,1,-1}) e^{-i\frac{E_2}{\hbar}t} \\ &= \frac{1}{\sqrt{2}} R_{2,1} (Y_{1,1} + Y_{1,-1}) e^{-i\frac{E_2}{\hbar}t}, \text{ where } E_2 = -\frac{13.6 \text{ eV}}{4} \end{aligned}$$

$$(b) \quad \langle V \rangle = -\frac{e^2}{4\pi\epsilon_0} \left\langle \frac{1}{r} \right\rangle$$

$$\begin{aligned} \left\langle \frac{1}{r} \right\rangle &= \langle \Psi(\vec{r}, t) | \frac{1}{r} | \Psi(\vec{r}, t) \rangle \\ &= \frac{1}{2} \langle \psi_{2,1,1} + \psi_{2,1,-1} | e^{i\frac{E_2}{\hbar}t} \cdot \frac{1}{r} e^{-i\frac{E_2}{\hbar}t} | \psi_{2,1,1} + \psi_{2,1,-1} \rangle \\ &= \frac{1}{2} \langle \psi_{2,1,1} + \psi_{2,1,-1} | \frac{1}{r} | \psi_{2,1,1} + \psi_{2,1,-1} \rangle \end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \langle R_{21} | \frac{1}{r} | R_{21} \rangle \langle Y_{11} + Y_{1-1} | Y_{11} + Y_{1-1} \rangle \\
&= \frac{1}{2} \int_0^\infty \frac{1}{24} a^3 \frac{r^2}{a^2} e^{-r/a} \cdot \frac{1}{r} \cdot r^2 dr \\
&\quad \times \left[\underbrace{\langle Y_{11} | Y_{11} \rangle}_{=1} + \cancel{\langle Y_{11} | Y_{1-1} \rangle} + \cancel{\langle Y_{1-1} | Y_{11} \rangle} + \underbrace{\langle Y_{1-1} | Y_{1-1} \rangle}_{=1} \right] \\
&= \frac{1}{24} \cdot a^5 \int_0^\infty r^3 e^{-r/a} dr \\
&= \frac{1}{24} \cdot a^5 \cdot 3! \cdot a^4 = \frac{1}{4a}
\end{aligned}$$

Thus

$$\langle V \rangle = -\frac{e^2}{4\pi\epsilon_0} \cdot \frac{1}{4a} = \frac{1}{2} \left(-\frac{1}{2a} \frac{e^2}{4\pi\epsilon_0} \right)$$

$$= \frac{E_1}{2}$$

From Eq. 4.70
and Eq. 4.72

= -6.8 eV, independent of time.