

HW 10 solution

Prob. 5.5

$$(a) \quad H = \begin{cases} \frac{p_1^2 + p_2^2}{2m} & \text{for } 0 < x_1, x_2 < a \\ \infty & \text{else} \end{cases}$$

$$\psi(x_1, x_2) = \frac{\sqrt{2}}{a} \left[\sin(\pi x_1/a) \sin(2\pi x_2/a) - \sin(2\pi x_1/a) \sin(\pi x_2/a) \right] \\ \text{for } 0 < x_1, x_2 < a.$$

$$\begin{aligned} H \psi(x_1, x_2) &= -\frac{\hbar^2}{2m} \left[\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} \right] \psi(x_1, x_2) \\ &= -\frac{\hbar^2}{2m} \cdot \frac{\sqrt{2}}{a} \left[-\left(\frac{\pi}{a}\right)^2 \sin(\pi x_1/a) \sin(2\pi x_2/a) \right. \\ &\quad \left. + \left(\frac{2\pi}{a}\right)^2 \sin(2\pi x_1/a) \sin(\pi x_2/a) \right. \\ &\quad \left. - \left(\frac{2\pi}{a}\right)^2 \sin(\pi x_1/a) \sin(2\pi x_2/a) \right. \\ &\quad \left. + \left(\frac{\pi}{a}\right)^2 \sin(2\pi x_1/a) \sin(\pi x_2/a) \right] \\ &= \frac{\hbar^2}{2m} \frac{\sqrt{2}}{a} \left[\left(\frac{\pi}{a}\right)^2 + \left(\frac{2\pi}{a}\right)^2 \right] \left[\sin(\pi x_1/a) \sin(2\pi x_2/a) \right. \\ &\quad \left. - \sin(2\pi x_1/a) \sin(\pi x_2/a) \right] \\ &= \frac{\hbar^2}{2m} \left(\frac{\pi}{a}\right)^2 \cdot 5 \cdot \psi(x_1, x_2) \\ &= 5K \psi(x_1, x_2) \end{aligned}$$

Thus $\psi(x_1, x_2)$ is the eigenfunction of the Hamiltonian with the eigenvalue $5K$.

(b) distinguishable

2nd excited state

$$\psi_{22} = \frac{2}{a} \sin\left(\frac{2\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right), \quad E_{22} = 8K$$

3rd excited state

$$\psi_{13} = \frac{2}{a} \sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{3\pi x_2}{a}\right), \quad E_{13} = 10K$$

$$\psi_{31} = \frac{2}{a} \sin\left(\frac{3\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right), \quad E_{31} = 10K$$

Identical boson

2nd excited state

$$\psi_{22} = \frac{2}{a} \sin\left(\frac{2\pi x_1}{a}\right) \sin\left(\frac{2\pi x_2}{a}\right), \quad E_{22} = 8K$$

3rd excited state

$$\psi_{13} = \frac{\sqrt{2}}{a} \left[\sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{3\pi x_2}{a}\right) + \sin\left(\frac{3\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right]$$

$$, \quad E_{13} = 10K$$

Identical fermion

1st excited state

$$\psi_{13} = \frac{\sqrt{2}}{a} \left[\sin\left(\frac{\pi x_1}{a}\right) \sin\left(\frac{3\pi x_2}{a}\right) - \sin\left(\frac{3\pi x_1}{a}\right) \sin\left(\frac{\pi x_2}{a}\right) \right]$$

$$, \quad E_{13} = 10K$$

2nd excited state

$$\psi_{23} = \frac{\sqrt{2}}{\alpha} \left[\sin\left(\frac{2\pi x_1}{\alpha}\right) \sin\left(\frac{3\pi x_2}{\alpha}\right) - \sin\left(\frac{3\pi x_1}{\alpha}\right) \sin\left(\frac{2\pi x_2}{\alpha}\right) \right]$$

$$E_{23} = 13K$$

Prob. 5.7.

(a) distinguishable

$$\psi(x_1, x_2, x_3) = \psi_a(x_1) \psi_b(x_2) \psi_c(x_3)$$

(b) identical bosons

There are six ways to order the states as shown

	x_1	x_2	x_3
a	$\swarrow$	b	$-c$
	$\searrow$	c	$-b$
b	$\swarrow$	a	$-c$
	$\searrow$	c	$-a$
c	$\swarrow$	a	$-b$
	$\searrow$	b	$-a$

let's use a simplified notation such that

$$\psi_a(x_1) \psi_b(x_2) \psi_c(x_3) \equiv \psi_{abc}$$

Then

$$\psi(x_1, x_2, x_3) = \frac{1}{\sqrt{6}} (\psi_{abc} + \psi_{acb} + \psi_{bac} + \psi_{bca} + \psi_{cab} + \psi_{cba})$$

where the normalization constant is simply determined from $\sqrt{\# \text{ of terms}}$

c) Identical fermions

Similar as in (b), but this time sign should be changed for every interchange of the states. If we start with "+" sign for ψ_{abc} , then any odd # of label changes from abc should have a "-" sign and even # of change should have a "+" sign, as shown:

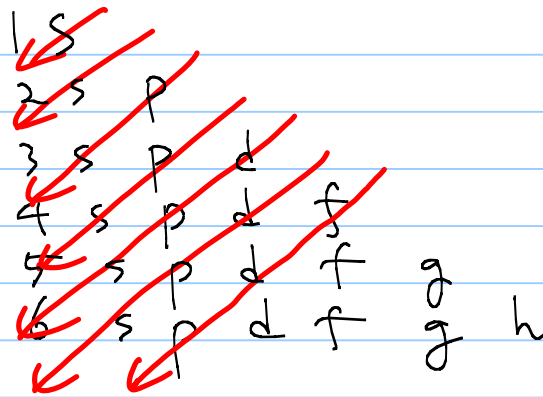
x_1	x_2	x_3	sign
a	b	c	+
	c	b	-
b	a	c	-
	c	a	+
c	a	b	+
	b	a	-

Thus

$$\psi(x_1, x_2, x_3) = \frac{1}{\sqrt{6}} \left(\psi_{abc} - \psi_{acb} - \psi_{bac} + \psi_{bca} + \psi_{cab} - \psi_{cba} \right)$$

3.

Tb has 65 electrons



$$1s^2 2s^2 2p^6 3s^2 3p^6 \quad 4s^2 3d^{10} 4p^6 5s^2$$

18 18

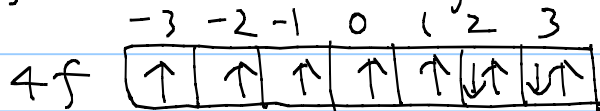
$$4d^{10} 5p^6 6s^2 4f^9$$

18 11

$$\Rightarrow 18 + 18 + 18 + 11 = 65$$

Outer most partially-filled subshell is

4f⁹. Thus using the Hund's rules



$$S = \frac{1}{2} \times 5 = \frac{5}{2}$$

$$L = (-3) + (-2) + (-1) + (0) + 1 + 2 \times 2 + 3 \times 2 = 5$$

$$J = L + S = 5 + \frac{5}{2} = \frac{15}{2}$$

↑
more than half-filled

$$\therefore {}^{25+1}L_J = {}^6H_{\frac{15}{2}}$$

Prob. 5.16 (a) $E_F = \frac{\hbar^2}{2m} (3\rho\pi^2)^{\frac{2}{3}}$

$$\rho = \frac{N}{V}$$

$$(\text{density}) d = \frac{\text{Mass}}{V} = \frac{N \cdot \frac{M_A}{N_A}}{V}$$

$$M_A \equiv \text{atomic mass of copper} = 63.5 \text{ g}$$

$$N_A = \text{Avogadro's number} = 6.02 \times 10^{23}$$

$$\Rightarrow d = \rho \cdot \frac{M_A}{N_A}$$

$$\Rightarrow \rho = \frac{d \cdot N_A}{M_A} = \frac{8.96 \text{ g/cm}^3 \cdot 6.02 \times 10^{23}}{63.5 \text{ g}}$$

$$= 8.49 \times 10^{22} / \text{cm}^3 = 8.49 \times 10^{28} / \text{m}^3$$

$$m = 9.11 \times 10^{-31} \text{ kg}$$

$$\hbar = 1.05 \times 10^{-34} \text{ J s}$$

$$E_F = \frac{(\hbar \cdot 8.49 \times 10^{28} / \text{m}^3)^2}{2 \cdot 9.11 \times 10^{-31} \text{ kg}} \left(3 \cdot 8.49 \times 10^{28} \text{ m}^{-3} \cdot \pi^2 \right)^{\frac{2}{3}}$$

$$\times \frac{\text{eV}}{1.60 \times 10^{-19} \text{ J}} = 7.0 \text{ eV}$$

(b)

$$E_F = \frac{1}{2} m v^2$$

$$\Rightarrow v = \sqrt{\frac{2 E_F}{m}} = \sqrt{\frac{2 \cdot 7.0 \text{ eV}}{0.51 \times 10^6 \text{ eV}/c^2}} = 5.2 \times 10^3 \cdot c$$

$$= 1.6 \times 10^6 \text{ m/s}$$

Because $\frac{v}{c} = 5.2 \times 10^{-3}$, the electrons are non-relativistic

(c)

$$T = \frac{E_F}{k_B} = \frac{7.0 \text{ eV}}{8.62 \times 10^{-5} \text{ eV/K}} = 8.2 \times 10^4 \text{ K}$$

$$(d) \quad P = \frac{(3\pi^2)^{2/3} \hbar^2}{5m} \rho^{5/3} = \frac{(3\pi^2)^{2/3} (1.05 \times 10^{-31})}{5 \cdot 9.11 \times 10^{-31}} \\ \times (8.49 \times 10^{28}) \text{ N/m}^2 = 3.8 \times 10^{10} \text{ N/m}^2$$

P. 6.2 (a) $E_n^0 = (n + \frac{1}{2}) \hbar \omega_0$, the unperturbed,

$$E_n = (n + \frac{1}{2}) \hbar \omega = (n + \frac{1}{2}) \hbar \sqrt{\frac{(1+\epsilon)k}{m}} \\ = (n + \frac{1}{2}) \hbar \omega_0 \sqrt{1+\epsilon} \\ = (n + \frac{1}{2}) \hbar \omega_0 \left(1 + \frac{\epsilon}{2} - \frac{1}{8} \epsilon^2 \right)$$

↑
expansion up to 2nd order

$$(b) \quad H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \\ = \frac{p^2}{2m} + \frac{1}{2} m \left(\frac{(1+\epsilon)k}{m} \right) x^2 \\ = \frac{p^2}{2m} + \frac{1}{2} m \omega_0^2 x^2 + \frac{1}{2} \epsilon m \omega_0^2 x^2 \\ = H^0 + \underbrace{\epsilon V^0}_{H'}$$

, where $V^0 \equiv \frac{1}{2} m \omega_0^2 x^2$ is the unperturbed potential energy.

$$E_n^1 = \langle \psi_n^0 | H' | \psi_n^0 \rangle \\ = \epsilon \langle \psi_n^0 | V^0 | \psi_n^0 \rangle \\ \rightarrow = \frac{\epsilon}{2} \cdot (n + \frac{1}{2}) \hbar \omega_0$$

from Example 2.5 of Griffiths, thus up to the 1st order

$$E_n = E_n^0 + E_n^1 = (n + \frac{1}{2}) \hbar \omega_0 \left(1 + \frac{\epsilon}{2} \right)$$

$\therefore$ consistent with (a) up to the 1st order