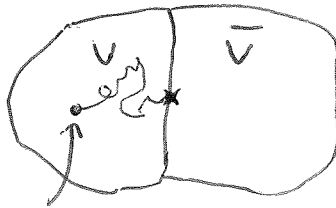


Lecture 3B

Consider random var. $X(t)$ [scalar or vector] ; define $p(x, t | y, 0)$ as before.

~~Consider~~ ~~after~~

Consider



starting point y

(FPT)

1st passage time: time needed to reach $\bar{V}$ for the 1st time. For ex.: a chemical reaction which need ^{activation} energy E_c to occur:

$$V = \{E < E_c\}, \quad \bar{V} = \{E \geq E_c\}$$

Prob. that $X(t)$ is still in V at time t is given by $p_V(y, t) = \int_V dx p(x, t | y, 0)$

Define $\eta(y, t)$ to be prob. density for FPT:

if $\tilde{t}$ is the 1st time $X(t)$ reaches $\bar{V}$, given $X(0) = y \in V$, $\eta(y, t)dt$ is the prob. that $t < \tilde{t} < t + dt$.

Note that if $X(t)$ is in V @ time t , @ $t + dt$ it either stays in V or makes 1st passage.

So, $p_v(y, t) - p_v(y, t+dt) = \eta(y, t)dt$, or

$$\boxed{\eta(y, t) = - \frac{\partial p_v(y, t)}{\partial t}}$$

FPT moments:

$$\langle t^n(y) \rangle = \int_0^\infty dt t^n \eta(y, t) \stackrel{\text{by parts}}{=} \downarrow$$

$$= n \int_0^\infty dt t^{n-1} p_v(y, t). \quad n=1, 2, \dots$$

Now, 1D case:  $X(t)$ takes these values; endpoints are absorbing

Theoretically, could solve the FP eq'n w/ absorbing boundary conditions; however, it's easier to ~~define~~ define $\varphi(x, t)$ - the prob. that FPT is $< t$, given $X(0) = x$.

Define $p(y, dt | x, 0)dy$ - the prob. that

$y \leq X(dt) \leq y+dy$, given that

$X(0) = x$. [homogeneous process, $p(y, t+dt | x, t) = p(y, dt | x, 0)$]

Then $\varphi(x, t+dt) = \int_0^A dy \underbrace{p(y, dt | x, 0)}_{\substack{x \rightarrow y \text{ transition} \\ \text{for } \forall y \in (0, A); \\ \text{this transition} \\ \text{does not lead to} \\ \text{1st passage}}} \underbrace{\varphi(y, t)}_{\substack{\text{prob. that} \\ \text{FPT is } < t}}$

Expand $\Psi(y, t)$ around x :

$$\Psi(y, t) = \Psi(x, t) + \frac{\partial \Psi(x, t)}{\partial x} (y-x) + \frac{1}{2} \frac{\partial^2 \Psi(x, t)}{\partial x^2} (y-x)^2 + \dots$$

Then

$$\begin{aligned} \Psi(x, t+dt) &= \Psi(x, t) \underbrace{\int_0^A dy p(y, dt|x, 0)}_{=1} + \\ &+ \frac{\partial \Psi(x, t)}{\partial x} \int_0^A dy (y-x) p(y, dt|x, 0) + \\ &+ \frac{1}{2} \frac{\partial^2 \Psi(x, t)}{\partial x^2} \int_0^A dy (y-x)^2 p(y, dt|x, 0) + \dots, \end{aligned}$$

$$\lim_{dt \rightarrow 0} \frac{\Psi(x, t+dt) - \Psi(x, t)}{dt} = \frac{\partial \Psi(x, t)}{\partial t} =$$

$$= A(x) \frac{\partial \Psi(x, t)}{\partial x} + \frac{B(x)}{2} \frac{\partial^2 \Psi(x, t)}{\partial x^2} \quad (*)$$

$$A(x) = \lim_{dt \rightarrow 0} \frac{1}{dt} \int_0^A dy (y-x) p(y, dt|x, 0),$$

$$B(x) = \lim_{dt \rightarrow 0} \frac{1}{dt} \int_0^A dy (y-x)^2 p(y, dt|x, 0).$$

"another" backward eq'n in 1D

No t -dependence in $A(x)$ & $B(x)$ b/c the process is homogeneous.

BCs: $\Psi(0, t) = 1$, $\Psi(A, t) = 1$.
absorbing ends

Compute moments of $\Psi(x, t)$ [easier than solving $(*)$ directly]

$$\text{Define } \mu_j(x) = \int_0^\infty dt t^j \frac{\partial \Psi(x, t)}{\partial t}$$

$$\Psi(x, t+dt) - \Psi(x, t) = \tilde{\eta}(x, t) dt, \text{ or}$$

$$\tilde{\eta}(x, t) = \frac{\partial \Psi(x, t)}{\partial t}$$

Now, apply $\int_0^\infty dt t^j \frac{\partial}{\partial t}$ to (*):

$$\begin{aligned} \underline{j=1}: \quad \int_0^\infty dt t \underbrace{\frac{\partial}{\partial t} \left(\frac{\partial \psi}{\partial t} \right)}_{\frac{d\psi}{dt}} & \stackrel{\text{by parts}}{=} \frac{\partial \psi}{\partial t} t \Big|_0^\infty - \int_0^\infty dt \frac{\partial \psi}{\partial t} = \\ & = -\underbrace{\psi(x, +\infty)}_1 + \underbrace{\psi(x, 0)}_0 = -1 \\ \mu_1(0) = \mu_1(A) = 0 \quad [\text{absorbing}] \end{aligned}$$

$$\text{So, } A(x) \frac{d\mu_1(x)}{dx} + \frac{B(x)}{2} \frac{d^2\mu_1(x)}{dx^2} = -1,$$

Likewise, $\underline{j=2}$:

$$\begin{aligned} \int_0^\infty dt t^2 \frac{\partial}{\partial t} \left(\frac{\partial \psi}{\partial t} \right) & = \frac{\partial \psi}{\partial t} (t^2) \Big|_0^\infty - 2 \int_0^\infty dt t \frac{\partial \psi}{\partial t} = \\ & = -2\mu_1(x). \end{aligned}$$

$$\text{So, } A(x) \frac{d\mu_2}{dx} + \frac{B(x)}{2} \frac{d^2\mu_2}{dx^2} = -2\mu_1.$$

In general,

$$A(x) \frac{d\mu_j}{dx} + \frac{B(x)}{2} \frac{d^2\mu_j}{dx^2} = -j\mu_{j-1}. \quad \underline{j=2,3,\dots}$$

$$\text{BCs: } \mu_j(0) = \mu_j(A) = 0.$$

These are ordinary differential eq's, can be solved (!)

For example, for $\mu_1(x)$:

$$\mu_1(x) = 2 \frac{\int_0^A e^{-u(y)} dy \int_0^y dz [e^{u(z)} / B(z)]}{\int_0^A e^{-u(y)} dy} \int_0^x e^{-u(y)} dy -$$

$$- 2 \int_0^x e^{-u(y)} dy \int_0^y dz [e^{u(z)} / B(z)]$$

$[B] = \frac{L^2}{T}$
 $C_1 (= \text{const})$
 $[C_1] = L/D = \frac{T}{L}$
 $[C_1] = T \quad "L^2/T$

$$\mu_1(0) = 0 \quad (\text{obviously}) \quad \boxed{u(x) = 2 \int_0^x dy \frac{A(y)}{B(y)}} \quad \left[\frac{A}{B} \right] = \frac{1}{L}, [u] \text{ dimless}$$

$$\mu_1(A) = 2 \int_0^A e^{-u(y)} dy \int_0^y z dz \dots$$

Finally,

$$\frac{d^2 \mu_1}{dx^2} = -C_1 e^{-u(x)} \left[\underbrace{2 \frac{A(x)}{B(x)}}_{u'(x)} \right] + 2 e^{-u(x)} \left[\underbrace{2 \frac{A(x)}{B(x)}}_{u'(x)} \right] x$$

$$= 2 \frac{A(x)}{B(x)} e^{-u(x)} \left[\cancel{2 \int_0^x} 2 \int_0^x dz \frac{e^{u(z)}}{B(z)} - C_1 \right] -$$

Now, $\frac{B(x)}{2} \frac{d^2 \mu_1}{dx^2} + A(x) \frac{d\mu_1}{dx} = A(x) e^{-u(x)} \left[2 \int_0^x dz \frac{e^{u(z)}}{B(z)} - C_1 \right] - 1 + A(x) C_1 e^{-u(x)} - 2A(x) \times$
 $\times e^{-u(x)} \int_0^x dz \frac{e^{u(z)}}{B(z)} = -1$ -5-

So, the sol'n is valid.

Finally, consider $A(x)=0$ & $B(x)=2D$,
 D - diff'n const

$$\text{Then } \frac{d^2 \mu_1}{dx^2} = -\frac{1}{D}$$

With absorbing BCs, $\leftarrow$ at both 0 & A

$$\mu_1(x) = \frac{1}{2D} x(A-x) \text{ is the sol'n.}$$

$$\uparrow \text{MFPT}(0)$$

Now consider diffusion with drift:

$$A(x) = -v, B = 2D$$

Note: only $x=0$ is absorbing now

$$\text{From above, } \mu_1(x) = C_1 \int_0^x dy e^{\frac{vy}{D}} - 2 \int_0^x dy e^{\frac{vy}{D}} \int_0^y dz \frac{e^{-\frac{vz}{D}}}{2D} \quad \ominus$$

$$\mu(x) = -\frac{vx}{D} \Rightarrow$$

$$\ominus C_1 \frac{D}{v} [e^{\frac{vx}{D}} - 1] - \frac{1}{D} \int_0^x dy e^{\frac{vy}{D}} \left(-\frac{D}{v}\right) [e^{-\frac{vy}{D}} - 1]$$

$$= C_1 \frac{D}{v} [e^{\frac{vx}{D}} - 1] + \frac{1}{v} \int_0^x dy [1 - e^{\frac{vy}{D}}] =$$

$$= C_1 \frac{D}{v} [e^{\frac{vx}{D}} - 1] + \frac{1}{v} \left\{ x - \frac{D}{v} [e^{\frac{vx}{D}} - 1] \right\}$$

$\mu_1(0) = 0$ is satisfied for $\forall C_1$.
 Now, consider $\frac{vA}{D} \gg 1$, then $\leftarrow$ for some sufficiently large A

$$\mu_1(A) \approx C_1 \frac{D}{v} e^{\frac{vA}{D}} + \frac{A}{v} - \frac{D}{v^2} e^{\frac{vA}{D}} \quad \text{or}$$

$$C_1 = \frac{\int_0^A \int_0^A \frac{e^{-\frac{vz}{D}}}{2D} dz dy}{\int_0^A \int_0^A \frac{e^{-\frac{vy}{D}}}{2D} dy dx} = \frac{1}{v} \quad (\text{see below})$$

$$\text{Now, } \mu_1(x) = \frac{D}{v^2} [e^{\frac{vx}{D}} - 1] + \frac{x}{v} - \frac{D}{v^2} [e^{\frac{vx}{D}} - 1] = \frac{x}{v} \quad \text{indep. of } D (!)$$

Indeed,

$$\begin{aligned} C_1 &= 2 \frac{\int_0^A dy e^{\frac{vy}{D}} \int_0^y dz \frac{e^{-\frac{vz}{D}}}{2D}}{\int_0^A dy e^{\frac{vy}{D}}} = \\ &= \frac{1}{D} \frac{\int_0^A dy e^{\frac{vy}{D}} \left(-\frac{D}{v}\right) [e^{-\frac{vy}{D}} - 1]}{\left(\frac{D}{v}\right) [e^{\frac{vA}{D}} - 1]} = \\ &= -\frac{1}{D} \frac{\int_0^A dy [1 - e^{\frac{vy}{D}}]}{e^{\frac{vA}{D}} - 1} = \frac{1}{D} \frac{\left(\frac{D}{v}\right) [e^{\frac{vA}{D}} - 1] - A}{e^{\frac{vA}{D}} - 1} \approx \\ &\approx \frac{1}{v} \end{aligned}$$