

Maxwell's fridge

[coupled to T_h] $\left\{ \begin{array}{l} T_h > T_c \\ -u \\ -d \end{array} \right.$ $\Delta E = E_u - E_d > 0$

\rightarrow $\boxed{101101101} \dots \rightarrow$ $\frac{R_{d \rightarrow u}}{R_{u \rightarrow d}} = e^{-\beta_h \Delta E} (k_B T_h)^{-1}$ } intrinsic to demon
 T_c

As discussed previously,

$$\left\{ \begin{array}{l} R_{d \rightarrow u} = \gamma(1-\epsilon), \\ R_{u \rightarrow d} = \gamma(1+\epsilon), \end{array} \right. \quad \epsilon = \tanh \frac{\beta_h \Delta E}{2}, \quad 0 < \epsilon < 1$$

Bit states can change only through bit-demon interactions, of duration

τ : $0d \leftrightarrow 1u$. These transitions are coupled to T_c :

$$\frac{R_{\phi d \rightarrow 1u}}{R_{1u \rightarrow \phi d}} = e^{-\beta_c \Delta E} (k_B T_c)^{-1}$$

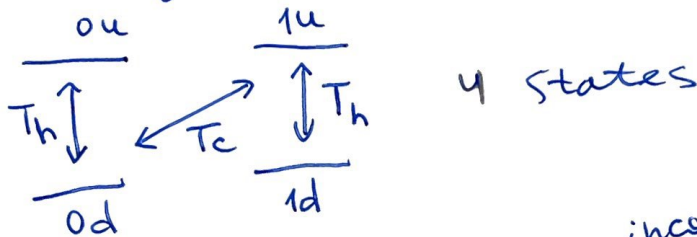
$$\left\{ \begin{array}{l} R_{\phi d \rightarrow 1u} = 1 - \omega, \\ R_{1u \rightarrow \phi d} = 1 + \omega, \end{array} \right. \quad \omega = \tanh \frac{\beta_c \Delta E}{2}, \quad 0 < \omega < 1$$

Note that $\epsilon \equiv \frac{\omega - \epsilon}{1 - \omega \epsilon} = \frac{e^{\beta_c \Delta E} - 1}{e^{\beta_c \Delta E} + 1} - \frac{e^{\beta_h \Delta E} - 1}{e^{\beta_h \Delta E} + 1} = \frac{(e^{\beta_c \Delta E} - 1)(e^{\beta_h \Delta E} + 1) - (e^{\beta_h \Delta E} - 1)(e^{\beta_c \Delta E} + 1)}{(e^{\beta_c \Delta E} + 1)(e^{\beta_h \Delta E} + 1) - (e^{\beta_c \Delta E} - 1)(e^{\beta_h \Delta E} - 1)} =$

$$\begin{aligned} \square & \frac{e^{\beta_c \Delta E} - e^{\beta_h \Delta E} - e^{\beta_h \Delta E} + e^{\beta_c \Delta E}}{e^{\beta_c \Delta E} + e^{\beta_h \Delta E} + e^{\beta_c \Delta E} + e^{\beta_h \Delta E}} = \\ & = \frac{e^{\beta_c \Delta E} - e^{\beta_h \Delta E}}{e^{\beta_c \Delta E} + e^{\beta_h \Delta E}} = \tanh \frac{(\beta_c - \beta_h) \Delta E}{2} \end{aligned}$$

Note that $0 < \epsilon < 1$

Finally, as before ^{we have} $\delta = p_0 - p_1$:



Assume that each ^{incoming} bit is in state ϕ : $\delta = 1$.

Then the system is in $(0u, 0d)$ states to begin with. After τ , if we end up in $(0u, 0d)$ again: all $0d \rightarrow 1u$ transitions were balanced by $1u \rightarrow 0d$ transitions \Rightarrow \Rightarrow no energy exchanged with the T_c reservoir. If, however, we end up in $(1u, 1d) \Rightarrow \Delta E$ withdrawn from the cold reservoir.

Thermal rectification: energy withdrawn from T_c but never delivered to it.

In the end, flux of energy $T_c \rightarrow T_h$.

The energy exchange events imprinted on the outgoing bit stream.

More generally: $\begin{cases} \delta > 0 \text{ (excess of 0's)} : T_c \rightarrow T_h \text{ flow} \\ \delta < 0 \text{ (excess of 1's)} : T_h \rightarrow T_c \text{ flow} \end{cases}$



at ss, $\phi = p_1' - p_1 = \frac{\delta - \delta'}{2}$

Average energy transfer per interaction interval: $Q_{c \rightarrow h} = \phi \Delta E$.

If $Q_{c \rightarrow h} > 0 \Rightarrow$ acts like a demon

Now, define

$$S(\delta) = - \sum_{i=0,1} p_i \log p_i = - \frac{1-\delta}{2} \log \frac{1-\delta}{2} - \frac{1+\delta}{2} \log \frac{1+\delta}{2}$$

↑ information content per bit

$$\Delta S_B \equiv S(\delta') - S(\delta) = \underbrace{S(\delta - 2\phi) - S(\delta)}$$

One can show that

$$\phi(\delta, \epsilon, \gamma, \omega, \tau) = \frac{\delta - \epsilon}{2} \underbrace{\eta(\Lambda)}_{> 0}$$

As before,

$$T_{2 \times 2} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} e^{RT} \begin{pmatrix} p_0 & 0 \\ 0 & p_0 \\ p_1 & 0 \\ 0 & p_1 \end{pmatrix} \left\{ \begin{array}{l} \begin{pmatrix} p_u \\ p_d \end{pmatrix} \\ \begin{pmatrix} p_0 \\ p_1 \end{pmatrix} \end{array} \right. \begin{array}{l} \text{demon} \\ \text{bit} \end{array}$$

$$C \rightarrow \phi: T_{2 \times 2} \begin{pmatrix} p_u \\ p_d \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} p_0 & 0 \\ 0 & p_0 \\ p_1 & 0 \\ 0 & p_1 \end{pmatrix} \begin{pmatrix} p_u \\ p_d \end{pmatrix} = \begin{pmatrix} p_u \\ p_d \end{pmatrix}, \text{ as expected}$$

$\begin{pmatrix} p_u p_0 \\ p_d p_0 \\ p_u p_1 \\ p_d p_1 \end{pmatrix}$

$$R = \begin{matrix} & \text{ou} & \text{od} & \text{1u} & \text{1d} \\ \begin{pmatrix} \times & \gamma(1-\epsilon) & 0 & 0 \\ \gamma(1+\epsilon) & \times & 1+\omega & 0 \\ 0 & 1-\omega & \times & \gamma(1-\epsilon) \\ 0 & 0 & \gamma(1+\epsilon) & \times \end{pmatrix} & \Rightarrow & e^{R\tau} \end{matrix}$$

$$\mathcal{J}_{2 \times 2} \begin{pmatrix} p_u \\ p_d \end{pmatrix}_{ss} = \begin{pmatrix} p_u \\ p_d \end{pmatrix}_{ss}$$

$$\begin{pmatrix} p_o' \\ p_i' \end{pmatrix} = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 1 \end{pmatrix} e^{R\tau} \begin{pmatrix} p_u^{ss} & p_o \\ p_d^{ss} & p_o \\ p_u^{ss} & p_i \\ p_d^{ss} & p_i \end{pmatrix} \Rightarrow \text{get } \phi$$

$$\phi = \frac{\delta - \epsilon}{2} \underbrace{\eta(\Lambda)}_{\text{can be shown } > 0} \text{ in general,}$$

$$\phi \rightarrow \frac{\delta - \epsilon}{2} [1 - e^{-(1-\epsilon\omega)\tau}] \text{ as } \gamma \rightarrow \infty.$$

$$\text{So, } \text{sign}(\phi) = \text{sign}(\delta - \epsilon).$$

$$\delta > \epsilon \Rightarrow \phi > 0 \Rightarrow Q_{c \rightarrow h} = \phi \Delta E > 0.$$

A surplus of ϕ 's prevails over the T difference: flow of energy $T_c \rightarrow T_h$.

Now consider

$$\frac{d\vec{p}}{dt} = R\vec{p}.$$

As $\tau \uparrow$, the system relaxes to SS:

$$R\vec{p}_{ss} = 0, \text{ where}$$

$$\begin{pmatrix} -\gamma(1+\epsilon) & \gamma(1-\epsilon) & 0 & 0 \\ \gamma(1+\epsilon) & (\omega-1) - \gamma(1-\epsilon) & 1+\omega & 0 \\ 0 & 1-\omega & -(1+\omega) - \gamma(1+\epsilon) & \gamma(1-\epsilon) \\ 0 & 0 & \gamma(1+\epsilon) & -\gamma(1-\epsilon) \end{pmatrix} \begin{pmatrix} 1 \\ \mu \\ \mu\nu \\ \mu^2\nu \end{pmatrix} = 0$$

$$\mu = \frac{1+\epsilon}{1-\epsilon}, \quad \nu = \frac{1-\omega}{1+\omega}$$

$$\begin{pmatrix} -\gamma(1+\epsilon) + \mu\gamma(1-\epsilon) \\ \gamma(1+\epsilon) + \mu(\omega-1) - \mu\gamma(1-\epsilon) + \mu\nu(1+\omega) \\ \mu(1-\omega) - \mu\nu(1+\omega) - \mu\nu\gamma(1+\epsilon) + \mu^2\nu\gamma(1-\epsilon) \\ \mu\nu\gamma(1+\epsilon) - \mu^2\nu\gamma(1-\epsilon) \end{pmatrix} =$$

$$= \begin{pmatrix} -\gamma(1+\epsilon) + \gamma(1+\epsilon) \\ \gamma(1+\epsilon) + \frac{1+\epsilon}{1-\epsilon}(\omega-1) - \gamma(1+\epsilon) + \frac{1+\epsilon}{1-\epsilon}(1-\omega) \\ \frac{1+\epsilon}{1-\epsilon}(1-\omega) - \frac{1+\epsilon}{1-\epsilon}(1-\omega) - \cancel{\mu\nu\gamma(1+\epsilon)} + \mu\nu\gamma(1+\epsilon) \\ \mu\nu\gamma(1+\epsilon) - \mu\nu\gamma(1+\epsilon) \end{pmatrix} =$$

$$= \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \text{ as desired.}$$

With norm'n:

$$\vec{p}_{ss} = \begin{pmatrix} 1 \\ \mu \\ \mu\nu \\ \mu^2\nu \end{pmatrix} \frac{1}{(1+\mu)(1+\mu\nu)}$$

Note that $p_{ss,ij} = p_{i,ss}^D p_{j,ss}^B$

$$\begin{cases} \vec{p}_{ss}^D = \frac{1}{1+\mu} \begin{pmatrix} 1 \\ \mu \end{pmatrix} \\ \vec{p}_{ss}^B = \frac{1}{1+\mu\nu} \begin{pmatrix} 1 \\ \mu\nu \end{pmatrix} \end{cases}$$

Define KL distance:

$$D(\vec{p} | \vec{p}_{ss}) = \sum_{k \in \{0u, 0d, 1u, 1d\}} p_k \log \frac{p_k}{p_{k,ss}} \geq 0$$

as \vec{p} approaches \vec{p}_{ss} , $D \downarrow$:

$$\frac{d}{dt} D(\vec{p} | \vec{p}_{ss}) \leq 0.$$

\downarrow

Define $\begin{cases} \vec{p}_0 \leftarrow \text{demon-bit distr's} \\ \vec{p}_\tau \leftarrow \text{@ } t=0 \text{ \& } \tau \end{cases}$

$$D(\vec{p}_0 | \vec{p}_{ss}) - D(\vec{p}_\tau | \vec{p}_{ss}) \geq 0 \quad (*)$$

Then

$$\underbrace{\sum_k p_{0,k} \log p_{0,k} - \sum_k p_{\tau,k} \log p_{\tau,k}}_{S_{\tau} - S_0} - \left[\sum_k p_{0,k} \log p_{ss,k} - \sum_k p_{\tau,k} \log p_{ss,k} \right] =$$

$$= S_{\tau} - S_0 + \sum_k (p_{\tau,k} - p_{0,k}) \log p_{k,ss} =$$

$$= S_{\tau} - S_0 + \sum_{\substack{i \in \{u,d\} \\ \text{demon states}}} (p_{\tau,i}^D - p_{0,i}^D) \log p_{i,ss}^D +$$

$$+ \sum_{j \in \{0,1\}} (p_{\tau,j}^B - p_{0,j}^B) \log p_{j,ss}^B$$

The joint entropy of the demon-bit system: $S' = S^D + S^B - \underbrace{I(D;B)}_{\geq 0}$

mutual info, decreases the entropy for correlated systems

$$S' = - \sum_i p_i(D,B) \log p_i(D,B) \quad \text{joint prob.}$$

$$= - \sum_i p_i(D|B) p_i(B) [\log p_i(B) + \log p_i(D|B)]$$

$$\Rightarrow S(X;Y) = S(X) + S(Y) - I(X;Y)$$

$$I_0(D;B) = 0$$

$$I_{\tau}(D;B) \text{ may be } > 0 \text{ (demon-bit correlations)}$$

$$I(X;Y) = \sum_{x,y} p(x,y) \log \frac{p(x,y)}{p(x)p(y)}$$

$$S_{\tau}^D = S_0^D \leftarrow \text{demon}_{in}^{ss}$$

$$S_0, S_{\tau} - S_0 = \Delta S_B - I_{\tau}(D; B)$$

also, $\vec{p}_0^D = \vec{p}_{\tau}^D$ in SS:

(*) gives

$$\Delta S_B - I_{\tau}(D; B) + \sum_{j \in \{0,1,3\}} (p_{\tau,j}^B - p_{0,j}^B) \log p_{j,SS}^B \geq 0$$

$$\text{At SS, } p_{\tau,0}^B - p_{0,0}^B = - (p_{\tau,1}^B - p_{0,1}^B) = -\Phi$$

$$S_0, \sum_{j \in \{0,1,3\}} (p_{\tau,j}^B - p_{0,j}^B) \log p_{j,SS}^B =$$

$$= -\Phi \log \frac{1}{1+\mu\nu} + \Phi \log \frac{\mu\nu}{1+\mu\nu} =$$

$$= \Phi \log \mu\nu = \Phi \log \frac{(1+\epsilon)(1-\omega)}{(1-\epsilon)(1+\omega)} =$$

$$= \Phi \log \frac{R_{u \rightarrow d}}{R_{d \rightarrow u}} + \Phi \log \frac{R_{d \rightarrow u}}{R_{u \rightarrow d}} =$$

$$= \Phi \beta_h \Delta E - \Phi \beta_c \Delta E = (\beta_h - \beta_c) \underbrace{\Phi \Delta E}_{Q_{c \rightarrow h}}$$

So, $\Delta S_B + Q_{c \rightarrow h} (\beta_h - \beta_c) \geq I_{\tau}(D; B) > 0$
 strong "Clausius statement"

Thus if $\Phi > 0 \Rightarrow Q_{c \rightarrow h} > 0$, ($\Rightarrow Q_{c \rightarrow h} (\beta_h - \beta_c) > 0$)
 we obtain: $\Delta S_B > 0$

"information is written
 to the bit stream"

Now if $\delta > \delta' > 0 \Rightarrow \Phi < 0 \Rightarrow Q_{c \rightarrow h} = \underbrace{\Phi}_{< 0} \underbrace{\Delta E}_{> 0} < 0$,

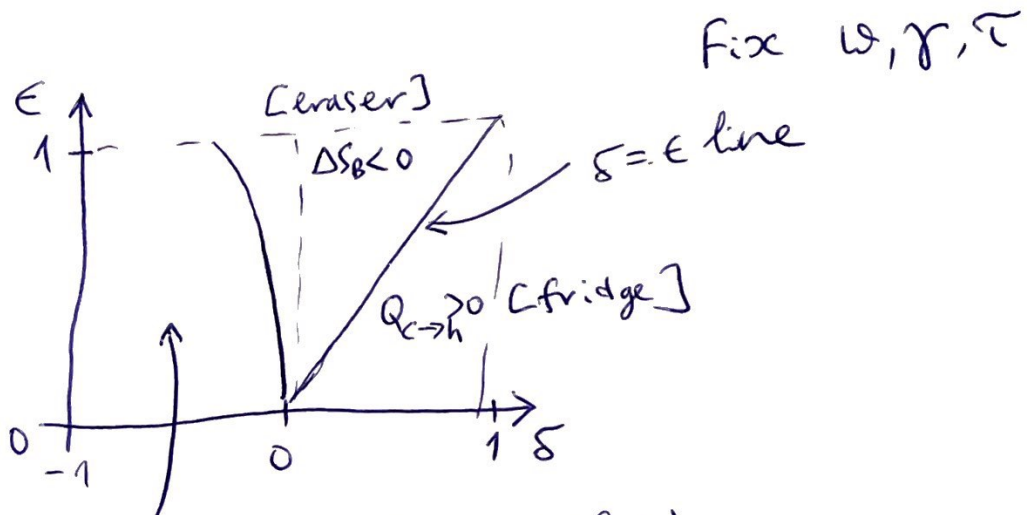
energy flows from $T_h \rightarrow T_c$.

$$Q_{c \rightarrow h} (\beta_h - \beta_c) > 0 \Rightarrow \Delta S_B < 0$$

$$\text{s.t. } Q_{c \rightarrow h} (\beta_h - \beta_c) + \Delta S_B \geq 0$$

$\Delta S_B < 0$ here b/c $\delta > \delta' > 0$
 $\underbrace{\beta_c}_{\Phi < 0} \underbrace{\frac{\delta - \delta'}{2}}$

So, information is erased from the bit stream.



$$\begin{cases} Q_{c \rightarrow h} < 0 \text{ (hot} \rightarrow \text{cold)} \\ \Delta S_B > 0 \end{cases}$$

"both thermodynamic & information resources are consumed"