

Solutions

Midterm Exam: E&M II 504

Spring 2018

Problem 1 [20 points]

Consider a particle with charge q and mass m moving in a uniform magnetic field $\vec{B} = B\hat{z}$ which is allowed to change adiabatically with time. Let us assume for simplicity that the initial velocity of the particle, $\vec{v}_\perp(0)$, is perpendicular to \vec{B} . The adiabatic condition means that the magnetic field changes negligibly during any one orbit, so that the particle orbit remains nearly circular.

1. Find the orbital angular speed $\vec{\omega}_B$ of the particle's motion. Write the result in vector form (i.e. find both amplitude and direction of $\vec{\omega}_B$).
2. Show that the orbital magnetic moment of the charged particle is given by

$$\vec{m} = -\frac{mv_\perp^2}{2} \frac{\vec{B}}{B^2}, \quad (1)$$

where $\vec{v}_\perp(t)$ is the component of the particle velocity perpendicular to \vec{B} .

3. Show that the power delivered to the charge by the time-dependent field is

$$P = \frac{mv_\perp^2}{2} \frac{d}{dt}(\ln B). \quad (2)$$

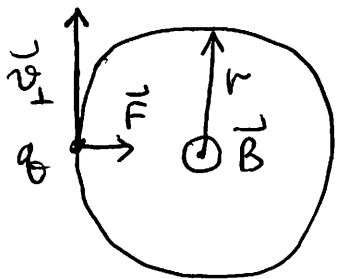
4. Use the result above to show that \vec{m} is invariant.

1. The Lorentz force is acting on the charge: $\frac{mv_\perp^2}{r} = qv_\perp B$.
radius of the orbit

Since $\omega_B = \frac{v_{\perp}}{r}$,

$$m \omega_B = q_B B, \text{ or}$$

$$\omega_B = \frac{q_B B}{m}$$



$$\vec{F} = q \vec{v}_{\perp} \times \vec{B}$$

Positive q revolves clockwise around \vec{B} :

$$\vec{\omega}_B = - \frac{q \vec{B}}{m}$$

$$2. \quad \vec{m} = \int \vec{r} \times \vec{A} = - \left(\frac{q \omega_B}{2\pi} \right) \pi r^2 \hat{z} = - \left(\frac{q \omega_B}{2\pi} \right) \pi \left(\frac{v_{\perp}}{\omega_B} \right)^2 \hat{z} =$$

$$\frac{q}{T} = \frac{q \omega_B}{2\pi} = - \frac{m v_{\perp}^2}{2} \frac{\vec{B}}{B^2}$$

$$3. \quad \vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \Rightarrow \oint \vec{E} \cdot d\vec{l} = - \frac{d\Phi_B}{dt}, \text{ or}$$

$$E(2\pi r) = - \pi r^2 \frac{dB}{dt} \Rightarrow E = - \frac{r}{2} \frac{dB}{dt}$$

$\vec{E} \uparrow \uparrow \vec{v}_{\perp}$

Then the power $P = q \vec{E} \cdot \vec{v}_{\perp} = q |E| v_{\perp} =$

$$= q \left(\frac{r}{2} \frac{dB}{dt} \right) v_{\perp} = \frac{m v_{\perp}^2}{2} \frac{d}{dt} (\log B)$$

$r = \frac{v_{\perp} m}{q B}$

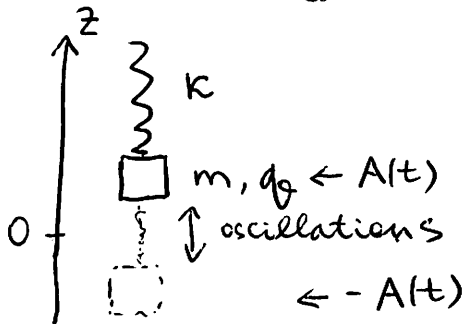
4. Note that here the particle's KE is not conserved: $\frac{d}{dt} \left(\frac{m v_{\perp}^2}{2} \right) = \frac{m v_{\perp}^2}{2} \frac{d}{dt} (\log B),$

$$\frac{d}{dt} \left(\log \left(\frac{m v_{\perp}^2}{2} \right) \right) = \frac{d}{dt} (\log B) \Rightarrow \frac{m v_{\perp}^2}{2} = B \times \text{constant}$$

But $\vec{m} = - \frac{m v_{\perp}^2}{2} \frac{\vec{z}}{B} = \underline{\underline{\text{constant} \cdot \hat{z}}}$ invariant!

Problem 2 [10 points]

Consider a simple harmonic oscillator which consists of a particle with charge q and mass m connected to an ideal spring with a spring constant κ (there is no gravity). The particle is set oscillating with the initial amplitude $A(0)$. Assuming non-relativistic motion, find the radiated power per unit solid angle and the total radiated power. Find $E(t)$, where E is the total energy of the harmonic oscillator and t is time.



$$m \ddot{z} = -\kappa z \quad \text{gives}$$

$$\omega^2 = \frac{\kappa}{m}$$

$$z(t) = A \cos(\omega t)$$

assume
[initial phase = 0]

The dipole moment $p(t) = qz(t) =$
 $= q_0 A \cos(\omega t) = q_0 A \operatorname{Re} \{ e^{-i\omega t} \}$.

In Jackson, we used the convention

$$p(\vec{x}, t) = p(\vec{x}) e^{-i\omega t}, \quad (\operatorname{Re} \text{ part implied})$$

which leads to $p = q_0 A$.

(9.23)

$$\kappa = \frac{\omega^2}{c} = \frac{1}{c} \sqrt{\frac{\kappa}{m}}$$

$$\text{Then } \frac{dP}{d\Omega} = \frac{c^2 z_0}{32\pi^2} \kappa^4 p^2 \sin^2 \theta =$$

power
per unit
solid
angle

$$= \frac{z_0}{32\pi^2 c^2} \frac{\kappa^2}{m^2} q^2 A^2 \sin^2 \theta$$

$$\left\{ \begin{aligned} z_0 &= \sqrt{\frac{\mu_0}{\epsilon_0}}, \\ c &= \frac{1}{\sqrt{\mu_0 \epsilon_0}} \end{aligned} \right.$$

$$\begin{aligned}
 \overset{\text{total power}}{\uparrow} P &\stackrel{(9.24)}{=} \frac{c^2 z_0 k^4}{12\pi} q_0^2 A^2 = \frac{q_0^2 A^2 z_0 k^2}{12\pi c^2 m^2} = \\
 &\frac{z_0}{c^2} = \frac{1}{\epsilon_0 c^3} = \frac{q_0^2 A^2 k^2}{12\pi \epsilon_0 c^3 m^2} . \\
 &=
 \end{aligned}$$

As the oscillator is losing energy through radiation, A decreases:

$$\frac{dE}{dt} = - \frac{q_0^2 A^2 k^2}{12\pi \epsilon_0 c^3 m^2}, \text{ where}$$

$$E(t) = \frac{1}{2} k A^2(t).$$

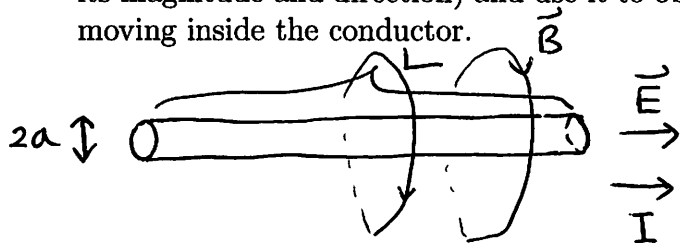
$$\text{So, } \frac{dE}{dt} = - \frac{q_0^2 k}{6\pi \epsilon_0 c^3 m^2} E, \text{ or}$$

$$E(t) = E(0) e^{-t/\tau}, \text{ where}$$

$$\tau = \frac{6\pi \epsilon_0 c^3 m^2}{q_0^2 k} .$$

Problem 3 [10 points]

A long cylindrical ohmic conductor of radius a and length L carries a current I due to a constant potential difference V between its ends. Calculate the Poynting vector \vec{S} (both its magnitude and direction) and use it to obtain work per unit time done on the charges moving inside the conductor.



$$|\vec{E}| = \frac{V}{L},$$

$$|\vec{B}| = \frac{\mu_0 I}{2\pi a}.$$

$$\vec{S} = \vec{E} \times \vec{H} \Rightarrow |\vec{S}| = \frac{VI}{2\pi aL}.$$

\vec{S} is pointing \perp to the surface of the wire & away from it.

Recall that in a steady-state situation,

$$\underbrace{\int \vec{\nabla} \cdot \vec{S} d^3x}_{=} = - \underbrace{\int \vec{J} \cdot \vec{E} d^3x}_{\substack{\text{work per unit} \\ \text{time}}},$$

$\int_S \vec{S} \cdot \vec{n} ds = VI$, which is the total power dissipated into heat in the system

\uparrow
 surface S