

Measuring The Elliptic Genus

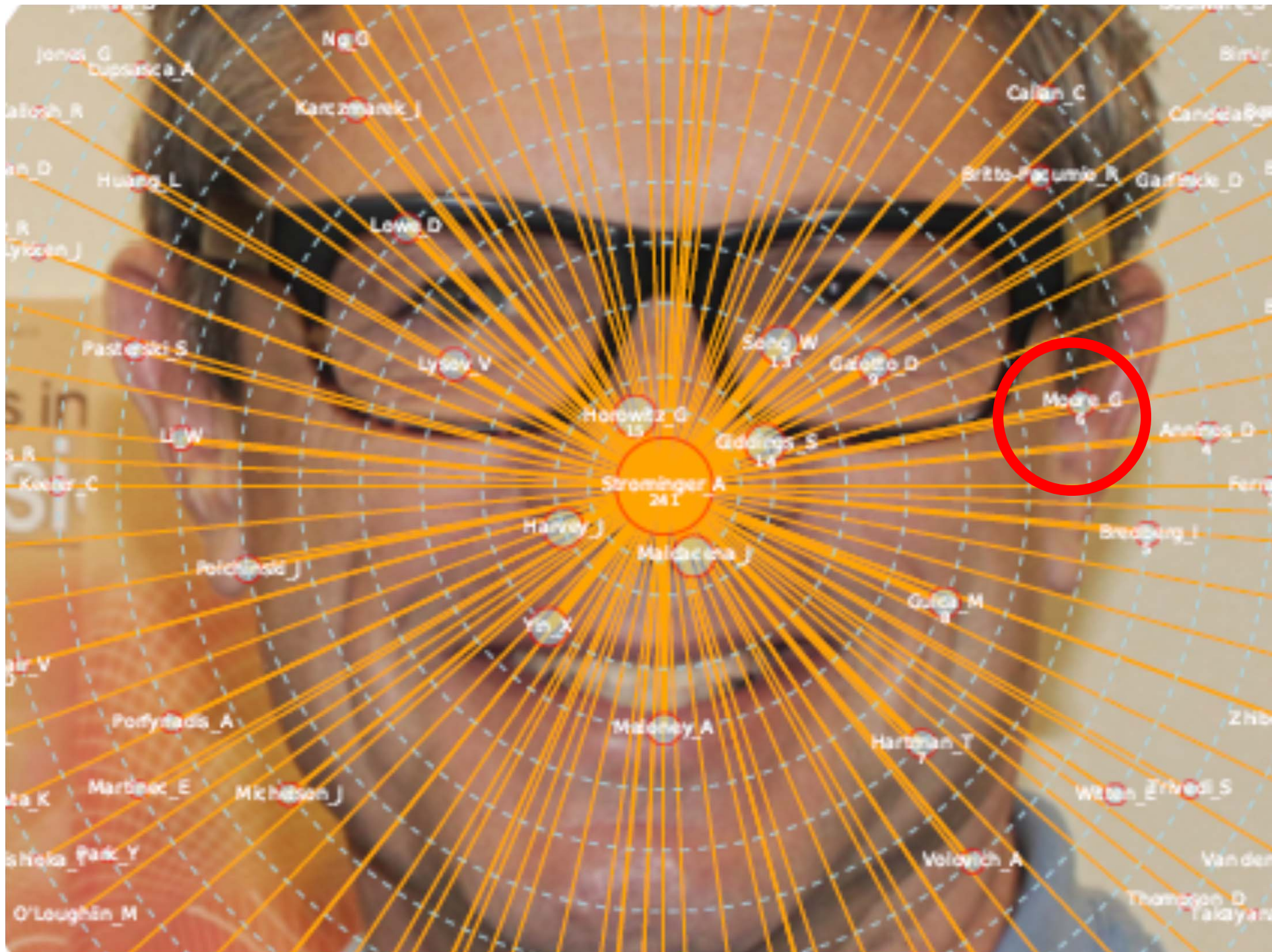
Gregory Moore

Rutgers

AndyFest, Harvard, July 31, 2015



Nicolaus Copernicus of Ithaka



Infinite Genus Swallows The Dilaton!

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This talk has its origins in the D1D5-system, where Andy has done such great work.

It was an important milestone on the road to AdS/CFT

$$\mathcal{C}_M = \text{Sym}^M(X)$$

$$X = K3, T4$$

Holographically dual to IIB strings on

$$AdS_3 \times S^3 \times X$$

(we won't say anything about $X = S^3 \times S^1 \dots$)

Some recent activity has centered on question:

“Do more general sequences $\{\mathcal{C}_M\}$ have holographic duals with weakly coupled gravity?”

This talk: AdS3/CFT2. CFT's are unitary and (for simplicity) have $c = 6M$ (4,4) supersymmetry:

Put necessary conditions on partition functions $Z(\mathcal{C}_M)$ for a holographic dual of an appropriate type to exist.

Keller; Hartman, Keller, Stoica; Haehl, Rangamani; Belin, Keller, Maloney;



Miranda
Cheng



Nathan
Benjamin



Shamit
Kachru



Natalie
Paquette

Our paper: Apply criterion of existence of a
Hawking-Page phase transition to the elliptic genus.

Criterion goes back to Witten and
a paper of Maldacena-Strominger.

AdS₃ Black Holes and a Stringy Exclusion Principle

Juan Maldacena and Andrew Strominger

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Abstract

The duality relating near-horizon microstates of black holes obtained as orbifolds of a subset of AdS_3 to the states of a conformal field theory is analyzed in detail. The

In each modular region of the τ plane there is a unique lowest (negative) action instanton. Defining

$$S_{min}(\tau) \equiv \min \left[\frac{i\pi k}{2} \left(\frac{a\tau + b}{c\tau + d} - \frac{a\bar{\tau} + b}{c\bar{\tau} + d} \right) \right], \quad (3.23)$$

the leading semiclassical approximation to the partition function is

$$Z(\tau) = e^{-S_{min}(\tau)}. \quad (3.24)$$

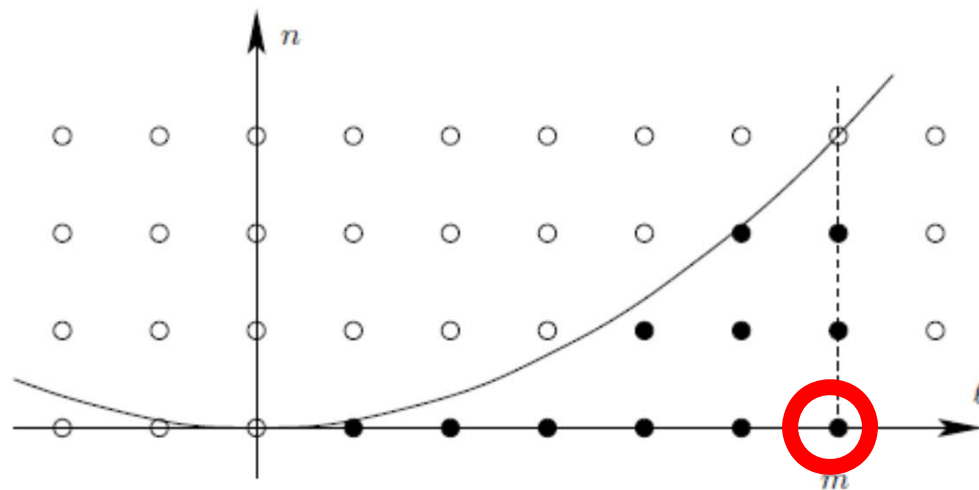
At low temperatures (large β) the partition function is dominated by (3.20) corresponding to a thermal gas in AdS_3 . At higher temperatures there is a transition to the black hole phase in which (3.16) dominates. This is a sharp first order phase transition in the limit $k \rightarrow \infty$ [20].

Reminder On Elliptic Genera

$$\mathcal{E}(\tau, z; \mathcal{C}) = \text{Tr}_{\mathcal{H}_{RR}} q^{L_0 - c/24} e^{2\pi i z J_0} \bar{q}^{\bar{L}_0 - c/24} e^{i\pi(J_0 - \bar{J}_0)}$$

$$\mathcal{E}(\tau, z; \mathcal{C}) = \sum_{n, \ell \in \mathbb{Z}} c(n, \ell; \mathcal{C}) q^n y^\ell$$

Modular object: Weak Jacobi form of weight zero and index m .



Extreme Polar Coefficient

$$\mathcal{E}(\tau, z; \mathcal{C}) = e(\mathcal{C})y^m + \dots$$

Benjamin et. al. put constraints on coefficients of elliptic genera of a sequence $\{e_M\}$ so that it exhibits HP transition. A corollary:

A necessary condition for $\{e_M\}$ to exhibit a HP transition is that

$e(\mathcal{C}_M)$ has at most polynomial growth in M for $M \rightarrow \infty$

Just a necessary condition.

Shamit's Question

“How likely is it for a sequence of CFT's $\{ \mathcal{C}_M \}$ to have a holographic dual with weakly coupled gravity? ”

We'll now make that more precise, and give an answer.

Zamolodchikov Metric

Space of CFT's is thought to have a topology. So we can speak of continuous families and connected components.

At smooth points the space is thought to be a manifold and there is a canonical isomorphism:

$$\Psi : V^{1,1}(\mathcal{C}) \rightarrow T_{\mathcal{C}}\mathcal{M}$$

$$\frac{\partial}{\partial t} |_0 \mathcal{S}[t] = \int \mathcal{O} \quad \Psi(\mathcal{O}) = \frac{\partial}{\partial t} |_0 = v \in T_{\mathcal{C}}\mathcal{M}$$

$$\langle \mathcal{O}(z_1) \mathcal{O}(z_2) \rangle := g_{\mathcal{Z}}(v, v) \frac{d^2 z_1 d^2 z_2}{|z_1 - z_2|^4}$$

Strategy

Suppose we have an ensemble \mathcal{E} of (4,4) CFTs:

$$\mathcal{E} = \coprod_M \mathcal{E}_M \quad \mathcal{E}_M = \coprod_\alpha \mathcal{E}_{M,\alpha}$$

$$\sum_\alpha \text{vol}_Z(\mathcal{E}_{M,\alpha}) < \infty$$

Then use the Z-measure to define a probability density on \mathcal{E}_M for fixed M.

Strategy – 2/2

Now suppose $\{ \mathcal{E}_M \}$ is a sequence drawn from \mathcal{E} .

$$p_M(\kappa, \ell) := \sum_{\mathbf{e} \leq \kappa M^\ell} \frac{\text{vol}(\mathbf{e}; \mathcal{E}_M)}{\text{vol}(\mathcal{E}_M)}$$

$$\wp(\ell) := \lim_{M \rightarrow \infty} p_M(\kappa, \ell)$$

$\wp(\ell)$ probability that a sequence drawn from \mathcal{E} has extremal polar coefficient growing at most like a power M^ℓ

Multiplicative Ensembles

$e(\mathcal{C})$ is constant on each component: $\mathcal{C} \in \mathcal{E}_{M,\alpha}$

$$e(\mathcal{C}_1 \times \mathcal{C}_2) = e(\mathcal{C}_1)e(\mathcal{C}_2)$$

Definition: A ***multiplicative ensemble*** satisfies:

$$\text{vol}(\mathcal{C}_1 \times \mathcal{C}_2) = \text{vol}(\mathcal{C}_1)\text{vol}(\mathcal{C}_2)$$

Definition: A CFT \mathcal{e} in a multiplicative ensemble is ***prime*** if it is not a product of CFT's (even up to deformation) each of which has $m > 0$.

A Generating Function

$\mathcal{C}_{m,\alpha}$ prime CFT's with $c=6m$, $\alpha = 1, \dots, f_m$

$$e(m, \alpha) = e(\mathcal{C}_{m,\alpha})$$

$$v(m, \alpha) = \text{vol}_Z(\mathcal{C}(m, \alpha))$$

$$\prod_{m=1}^{\infty} \prod_{\alpha=1}^{f_m} \frac{1}{1 - v(m, \alpha) e(m, \alpha)^{-s} q^m} = 1 + \sum_{M=1}^{\infty} \xi(s; M) q^M$$

$$\xi(s; M) = \sum_{\mathbf{e}=1}^{\infty} \frac{\text{vol}(\mathbf{e}; M)}{e^s}$$

Some Representative (?) Ensembles

We do not know what the space of (4,4) CFT's is

We do not even know how to classify compact hyperkähler manifolds !

$$S^m K3 := \text{Hilb}^m(K3)$$

$$S^m T4 := \text{Hilb}^{m+1}(T4)/T4$$

$$\mathcal{E} = \{ (S^1 X)^{n_1} \times \cdots \times (S^r X)^{n_r} \}$$

$$X = K3 \quad X = T4$$

$$X \in \{K3, T4\}$$

Moduli Spaces Of The Prime CFTs -1/2

These ensembles are multiplicative.

Primes: $S^m X$

$$Q_{r,s} = II^{r+8s,r}$$

$$\mathcal{N}_{r+8s,r} = O_{\mathbb{Z}}(Q_{r,s}) \backslash O_{\mathbb{R}}(Q_{r,s}) / O(r+8s) \times O(r)$$

$$\mathcal{M}(X) = \begin{cases} \mathcal{N}_{4,4} & X = T4 \\ \mathcal{N}_{20,4} & X = K3 \end{cases}$$

Moduli Spaces Of The Prime CFTs -2/2

One can derive $\mathcal{M}(S^m(X))$ using the ATTRACTOR MECHANISM:

Dijkgraaf;
Seiberg & Witten

Begin with $O(5,21)$ (or $O(5,5)$) moduli space of supergravity

Consider the subgroup fixing a primitive vector $u \in \mathbb{R}^{r+8s,r}$

The conjugacy class only depends on $u^2 = 2m$

$$O_{\mathbb{Z}}(Q_{r,s}, m) \backslash O_{\mathbb{R}}(Q_{r,s}, m) / O(r + 8s) \times O(r - 1)$$

$S^m(X)$ for $m > 1$ has "extra" hypermultiplet of blowup modes

$$(r, s) = \begin{cases} (5, 0) & X = T4 \\ (5, 2) & X = K3 \end{cases}$$

A Generating Function

$\mathcal{C}_{m,\alpha}$ prime CFT's with $c=6m$, $\alpha = 1, \dots, f_m$



$$e(m, \alpha) = e(\mathcal{C}_{m,\alpha})$$

$$\prod_{m=1}^{\infty} \prod_{\alpha=1}^{f_m} \frac{1}{1 - v(m, \alpha) e(m, \alpha)^{-s} q^m} = 1 + \sum_{M=1}^{\infty} \xi(s; M) q^M$$

$$\xi(s; M) = \sum_{\mathbf{e}=1}^{\infty} \frac{\text{vol}(\mathbf{e}; M)}{e^s}$$

Extreme Polar Coefficient

From the formula for the partition functions of symmetric product orbifolds we easily find

$$e(S^m K3) = m + 1$$

For the torus we have to work harder, since the elliptic genus vanishes.

Fortunately, that work has been done, thanks to

Counting BPS Blackholes in Toroidal Type II String Theory

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(Never let Andy submit papers!)

$$e(S^m T4) = m + 1$$

A Generating Function

$\mathcal{C}_{m,\alpha}$ prime CFT's with $c=6m$, $\alpha = 1, \dots, f_m$



$$v(m, \alpha) = \text{vol}_Z(\mathcal{C}(m, \alpha))$$

$$\prod_{m=1}^{\infty} \prod_{\alpha=1}^{f_m} \frac{1}{1 - v(m, \alpha) e(m, \alpha)^{-s} q^m} = 1 + \sum_{M=1}^{\infty} \xi(s; M) q^M$$

$$\xi(s; M) = \sum_{\mathbf{e}=1}^{\infty} \frac{\text{vol}(\mathbf{e}; M)}{e^s}$$

Volumes

Using results from number theory, especially the “mass formulae” of Carl Ludwig Siegel, one can -- with some nontrivial work -- compute the Z-volumes of these spaces. For example:

$$\text{vol}_Z(K3) =$$

$$\pi^{-40} \frac{(131)(283)(593)(617)(691)^2(3617)(43867)}{2^{40} \cdot 3^{34} \cdot 5^{15} \cdot 7^9 \cdot 11^5 \cdot 13^4 \cdot 17^3 \cdot 19^3 \cdot 23}$$

$$\cong 1.66 \times 10^{-61}$$

$$\text{vol}_Z(S^M K3) = \rho M^{42} f_{13}(M)$$

$$f_{13}(M) = \prod_{p|M} \frac{1 - p^{-12-12e_p(M)}}{1 - p^{-12}}$$

$$\rho = \pi^{-42} \frac{(103)(131)(283)(593)(617)(691)(3617)(43867)(2294797)}{2^{51} \cdot 3^{35} \cdot 5^{15} \cdot 7^{10} \cdot 11^5 \cdot 13^4 \cdot 17^3 \cdot 19^3 \cdot 23^2}$$

$$\cong 5.815 \times 10^{-63}$$

Result For Probabilities

$$\begin{aligned} H_\ell(s) &:= \lim_{M \rightarrow \infty} (M+1)^{\ell s} \frac{\xi(s; M)}{\xi(0; M)} \\ &= \lim_{M \rightarrow \infty} \sum_{\mathbf{e}=M+1}^{2^M} \frac{\text{vol}(\mathbf{e}; M)}{\text{vol}(M)} \left(\frac{(M+1)^\ell}{\mathbf{e}} \right)^s \\ &\geq \lim_{M \rightarrow \infty} \kappa^{-s} p_M(\kappa, \ell) \geq 0 \end{aligned}$$

Result For Probabilities

$$H_\ell(s) := \lim_{M \rightarrow \infty} (M + 1)^{\ell s} \frac{\xi(s; M)}{\xi(0; M)}$$

Claim: The limit exists for all nonnegative ℓ and

$$H_\ell(s) = \begin{cases} 1 & s = 0 \\ 0 & s > 0 \end{cases}$$
$$\Rightarrow \wp(\ell) = 0$$

Conclusion: Almost every sequence $\{e_M\}$ does not have a holographic dual

Proof – 1/3

$$\prod_{m=1}^{\infty} \prod_{\alpha=1}^{f_m} \frac{1}{1 - v(m, \alpha) e(m, \alpha)^{-s} q^m} = 1 + \sum_{M=1}^{\infty} \xi(s; M) q^M$$

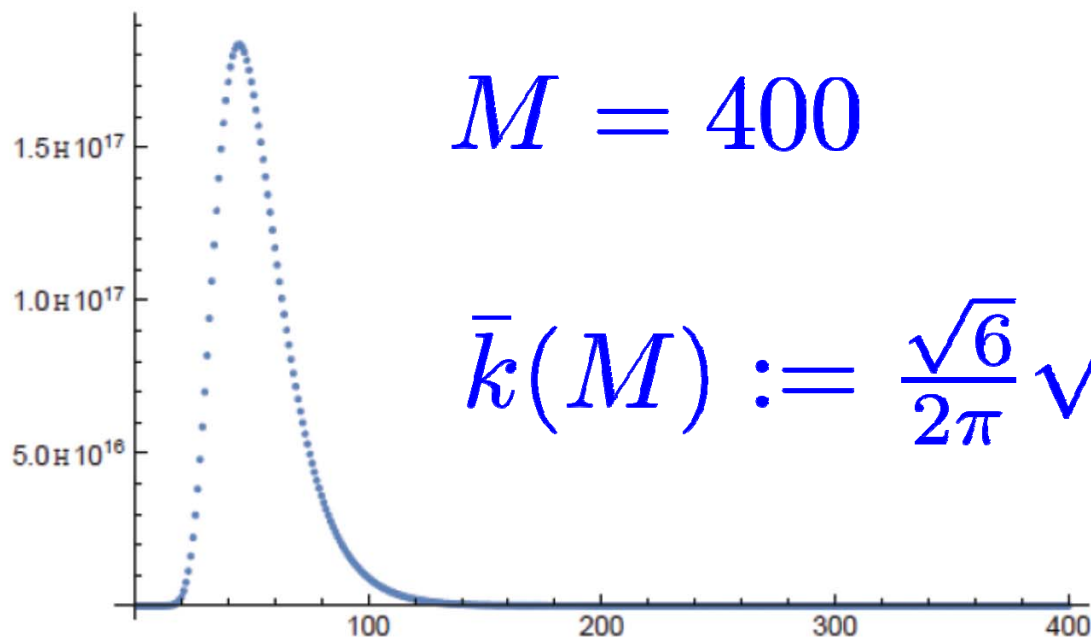
So $\xi(s; M)$ is a sum over partitions:

$$M = \lambda_1 + \cdots + \lambda_k$$

$$\lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_k$$

Statistics Of Partitions

For large M the distribution of partitions into k parts is sharply peaked:



$$\bar{k}(M) := \frac{\sqrt{6}}{2\pi} \sqrt{M} \log M$$

Moreover the “typical” partition has “most” parts of order:

$$\lambda_j \cong \sqrt{M}$$

Proof – 3/3

$\xi(s; M)$ is dominated by:

$$\xi(s; M) \cong V(\sqrt{M})^{\sqrt{M}} (\sqrt{M})^{-s\sqrt{M}} e^{2\pi\sqrt{\frac{M}{6}}}$$

$$H_\ell(s) := \lim_{M \rightarrow \infty} (M+1)^{\ell s} \frac{\xi(s; M)}{\xi(0; M)}$$

$$= \lim_{M \rightarrow \infty} (M+1)^{\ell s} (\sqrt{M})^{-s\sqrt{M}}$$

$$= \begin{cases} 1 & s = 0 \\ 0 & s > 0 \end{cases}$$

Some Wild Speculation:

(Discussions with Shamit Kachru and Alex Maloney)

Siegel Mass Formula:

Two lattices Γ_1 and Γ_2 are in the same genus if

$$\Gamma_1 \oplus S \cong \Gamma_2 \oplus S$$

Even unimodular lattices of rank $8n$ form a single genus, and:

$$\sum_{\alpha} \frac{1}{|\text{Aut}\Gamma_{\alpha}|} = \prod_p \alpha_p$$

A Natural Ensemble & Measure

Consider the ensemble of holomorphic CFT's.

(What would they be dual to? Presumably some version of chiral gravity in 3d!)

Holomorphic CFTs have $c = 24n$

They are completely rigid

$$Z = \sum_{\alpha} \frac{1}{|\text{Aut}(\mathcal{C}_{\alpha})|} < \infty$$

$$\mu(\mathcal{C}_{\alpha}) = \frac{1}{Z} \frac{1}{|\text{Aut}(\mathcal{C}_{\alpha})|}$$

Some Wild Speculation – 3/4

$$\mathcal{E}_N = \{ \mathcal{C}(\Gamma, G) \mid G \subset \text{Aut}(\Gamma) \quad \& \quad \Gamma \in II^{24N} \}$$

(Speculation: This set exhausts the set of $c=24N$ holomorphic CFTs.)

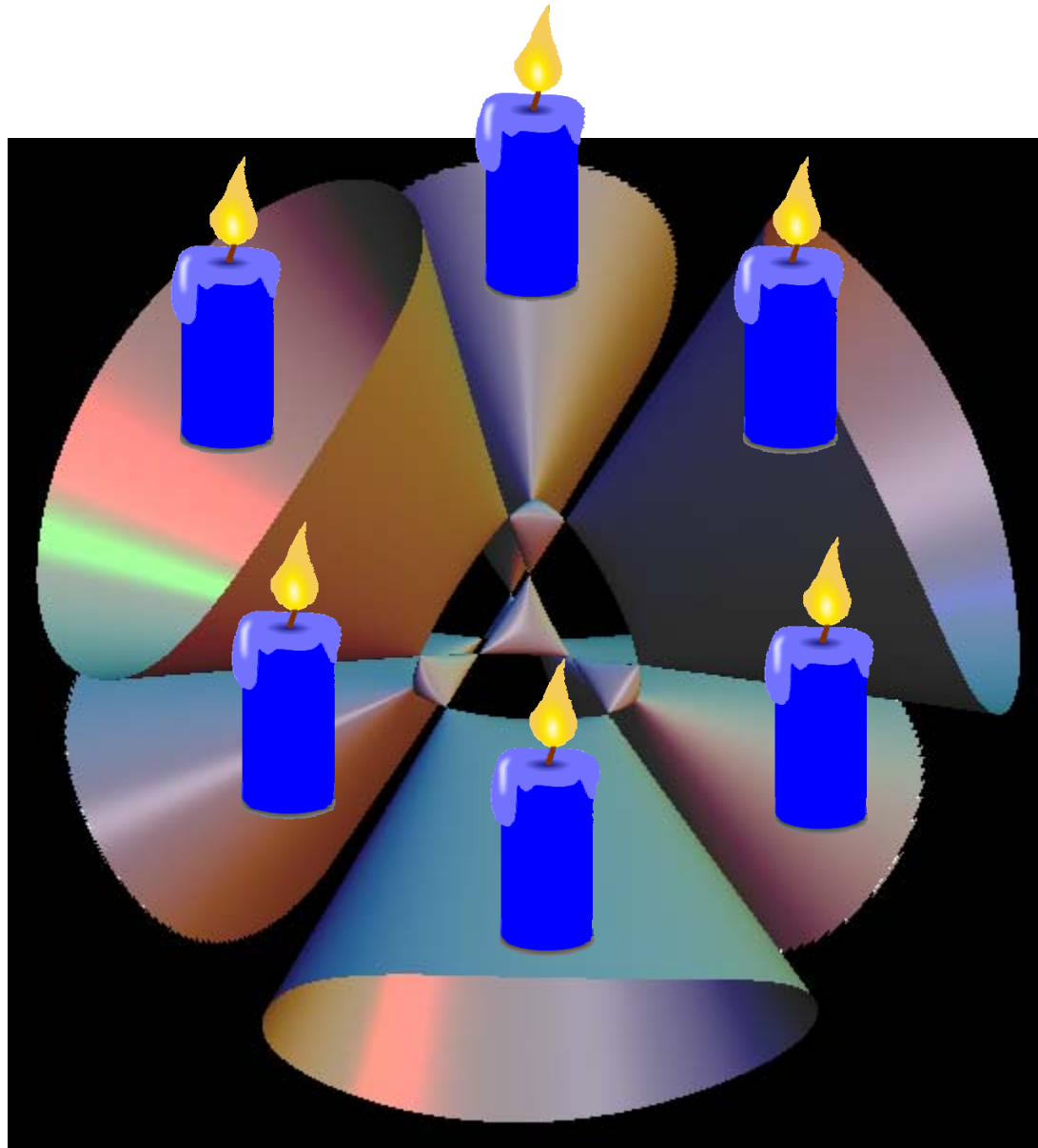
Speculation: Using results on the mass formula for lattices with nontrivial automorphism we can again prove that sequences $\{ \mathcal{E}_M \}$ with a holographic dual are measure zero.

Even Wilder Speculation – 4/4

Define a “genus” to be an equivalence class under tensoring with a lattice theory of chiral scalar fields.

$$\sum_{\alpha} \frac{1}{|\text{Aut}\mathcal{C}(\Gamma, G)_{\alpha}|} = \prod_p \alpha_p^{CFT}$$

Where the local densities are computed by counting automorphisms of the vertex operator algebra localized at a prime p .



HAPPY BIRTHDAY ANDY!!