



# An analytic model for the spatial clustering of dark matter haloes

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# Motivations

- Galaxies form around haloes
  - But haloes are much easier to model
- Reinforce and illuminate simulations
- A model that agrees with observations may
  - Predict dark mass distribution from visible mass
  - Constrain cosmological parameters

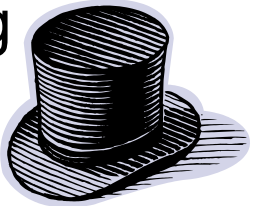


# Assumptions

- Fluctuations in the initial mass distribution are Gaussian
- Einstein-de Sitter universe
  - $\Omega=1, \Lambda=0$
  - “may be readily extended”
- Derived from Press Schechter (PS) model
- Assumes spherically symmetric halo formation

# Initial conditions and considerations

overdensity  $\delta(x) \equiv \frac{\rho(x) - \bar{\rho}}{\bar{\rho}}$

Windowing function 

↓

windowed overdensity  $\delta(x; R) = \int W(|x - y|; R) \delta(y) d^3 y$

RMS mass fluctuations  $\Delta^2(x; R) = \int P(k) \hat{W}^2(k; R) d^3 y$

probability distribution for  $\delta$   $p(\delta; R) d\delta = (2\pi)^{-1/2} \exp\left(-\frac{\delta^2(R)}{2\Delta^2(R)}\right) \frac{d\delta}{\Delta(R)}$

■ Note: critical dependence on  $v = \delta/\Delta$

# Am I in a halo?

- From PS theory: The probability that a random mass element is part of a halo  $> M_1$  is twice the probability that a surrounding sphere of mass  $M_1$  in the initial conditions has lineary extrapolated overdensity  $> \delta_c$

$$F(M_1, z_1) = 2 \int_{(1+z_1)\delta_c}^{\infty} p(\delta; R_1) d\delta = \text{erf} \left( \frac{v_1}{\sqrt{2}} \right)$$

leading to

$$n(M_1, z_1) dM_1 = -\sqrt{\frac{2}{\pi}} \frac{\bar{\rho}}{M_1} v_1 \frac{d \ln \Delta_1}{d \ln M_1} \exp\left(-\frac{v_1^2}{2}\right) \frac{dM_1}{M_1}$$



# Halo number density

- Number of  $M_1$  haloes at  $z_1$  with comoving radius  $R_0$  and overdensity  $\delta_0$

$$N(1|0)dM_1 = \frac{M_0}{M_1} f(1|0) \frac{d\Delta_1^2}{dM_1} dM_1$$

for  $f(1|0)$  the fraction of the mass in a region of initial radius  $R_0$ ,  $\delta_0$ ,  $z_1$  contained in dark haloes of mass  $M_1$

This formalism due to Bond, J. R., Coles S., Efstathiou G., Kaiser N., 1991, ApJ, 379, 440

- Notation:  $N(1|0) \equiv N(\Delta_1, \delta_1 | \Delta_0, \delta_0)$



# Spherical Collapse Model

- Concentric shells move as unit
  - Cross only right before collapse through zero
  - Mass interior to a given shell is constant
    - $(1 + \delta)R^3 = R_0^3$ , with  $\delta$  the present overdensity
- Gives an expression for  $R(R_0, \delta_0, z)$

# Bias Function

- Halo overdensity may be derived from previous expressions: 
$$\delta_h(1|0) = \frac{N(1|0)}{n(M_1, z_1)V} - 1$$

- $V = 4\pi R^3/3$ ,  $R_0 = R(1+\delta)^{1/3}$

- Given the above dynamics we may express this in the limit where  $R_0 \gg R_1$  and  $|\delta_0| \ll \delta_1$  as

$$b(M_1, z_1)\delta \equiv \delta_h(1|0) = \left(1 + \frac{v_1^2 - 1}{\delta_1}\right)\delta$$

with  $b(M_1, z_1)$  the “bias relation”



# And Finally...

- Average (Eulerian) cross-correlation between dark halo number density and mass within sphere

$$\overline{\xi}_{hm}(R, M_1, z_1) = \left\langle \delta_h(1|0)\delta \right\rangle_R = \frac{1}{n(M_1, z_1)V} \times \int_{-1}^{\infty} N(1|0)p(\delta; R)\delta d\delta$$

Average over spheres of radius R at z=0

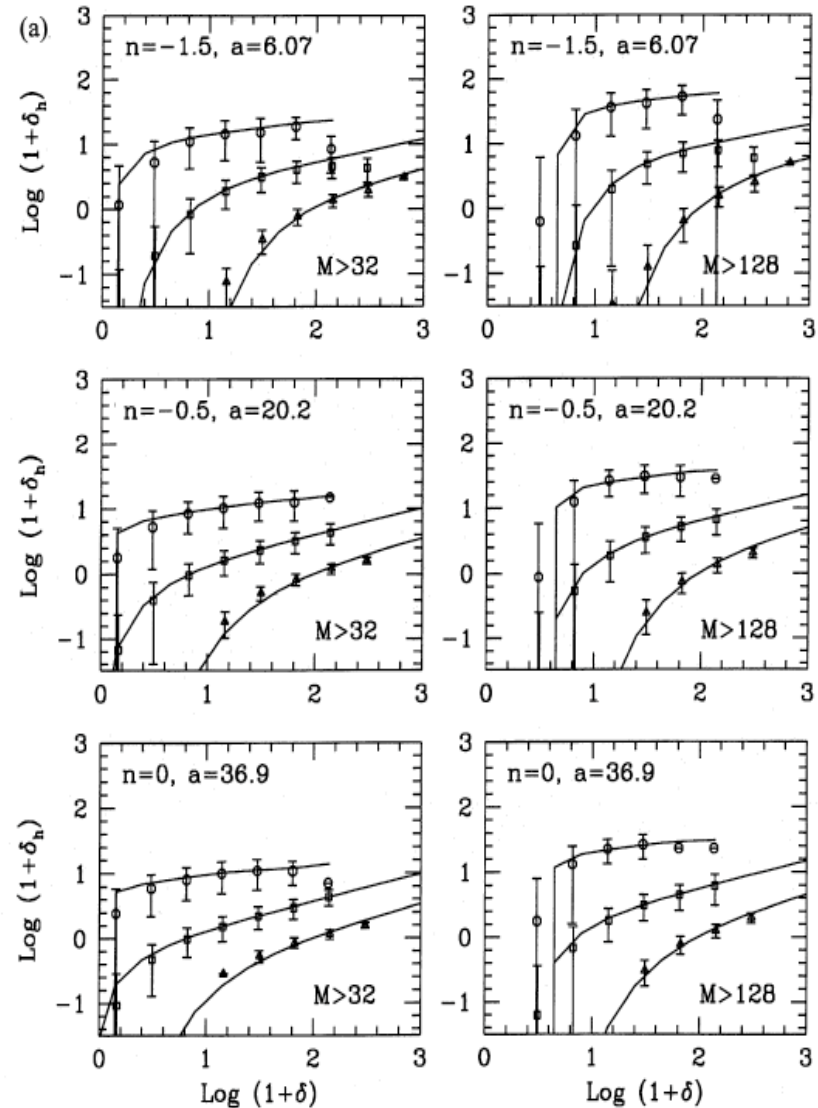


# N-Body Simluations

- P<sup>3</sup>M method (See Efstathiou G., et al. 1988, MNRAS, 235, 715)
- N = one million bodies
- Periodic, cubic lattice
  - Spatial resolution  $L/2500$  for lattice edge length  $L$
- Time is measured by a expansion factor  $a$ 
  - $a=1$  at  $t=0$
- Normalized power spectrum  $P(k) \propto k^n$
- Haloes identified via “friends-of-friends” algorithm
  - Other algorithms also tested
- Sample random spheres to gather statistics

# Comparison to Simulation

- Bias relation prediction
  - Overdensity of haloes  $\delta_h$  versus mass overdensity  $\delta$
  - Three curves for spheres of radius  $R/L = 0.02, 0.05, 0.13$ 
    - Data shifted rightward two decades for presentation
- Very good agreement



# Comparison to Simulation

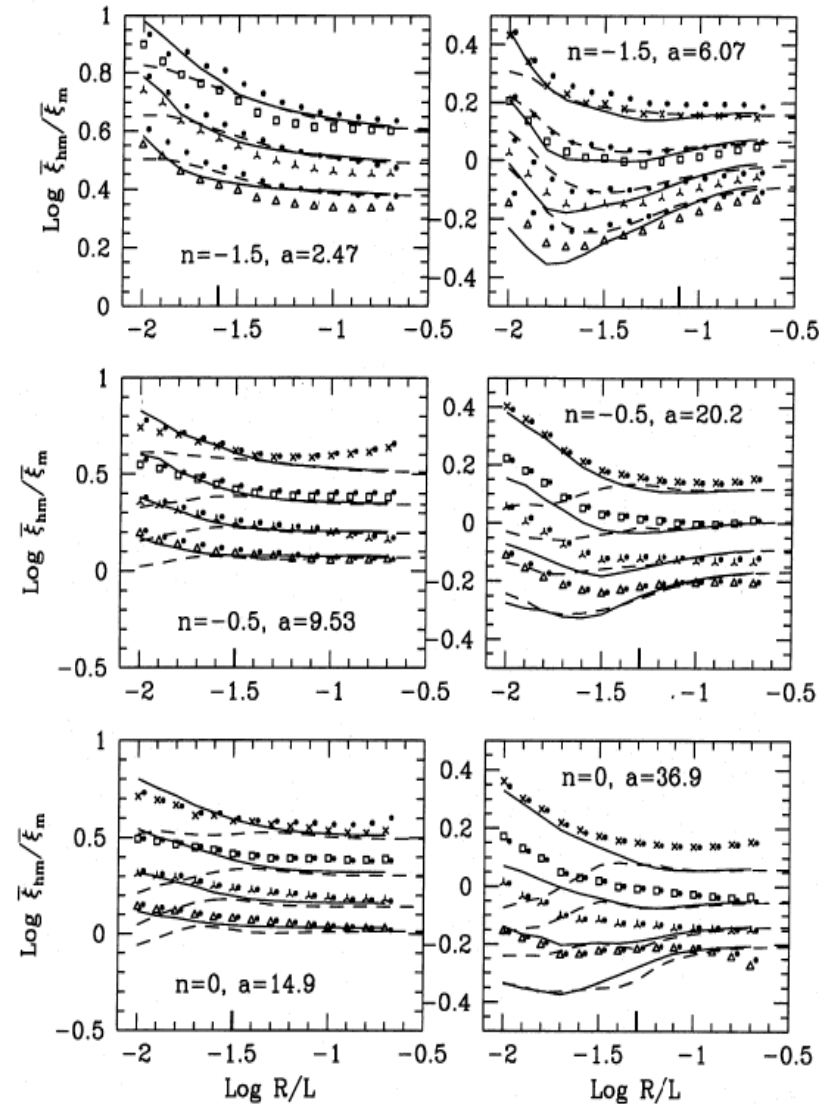
- Mass-Halo correlation

- Top to bottom,  $M \geq (256, 128, 64, 32)$
- Must assume a distribution for  $\delta$

- Solid curves use the observed distribution
- Dashed curves use lognormal approximation

- Worse agreement for  $\xi_m > 1$

- Such cases bias haloes toward initially underdense regions
- Suggests failure in spherical collapse model





# Results

- Halo bias increases with  $z$

- Older low-mass haloes can become more strongly clustered than current massive haloes
  - Old clusters may be less massive but still tight
- Clustering does not in itself predict mass

- Analytic expression for bias factor

- Given some knowledge of  $P(k)$  for initial conditions we can use observed clustering to infer underlying dark mass

$$b(M_1, z_1)\delta = \delta_h(1|0) = \left(1 + \frac{v_1^2 - 1}{\delta_1}\right)\delta$$