

Supplementary material for “Predicting polarization and nonlinear dielectric response of arbitrary perovskite superlattice sequences”

Xifan Wu¹, Massimiliano Stengel², Karin M. Rabe³, and David Vanderbilt³

¹*Chemistry Department, Princeton University, Princeton, NJ 08544-0001, USA*

²*Materials Department, University of California,*

Santa Barbara, Santa Barbara, CA 93106-5050, USA and

³*Department of Physics and Astronomy, Rutgers University, Piscataway, NJ 08854-8019, USA*

(Dated: May 21, 2008)

In this supplementary material, we present the form of the cluster expansions for the layer polarizations of AO and TiO₂ layers, and give the fitted values of the effective cluster interaction coefficients for the layer polarizations and the lattice constants.

As explained in the main text, we separate the D -dependence of the layer polarization of an AO layer into an odd part $p_{\text{AO}}^{(-)}(D)$ (reflecting the inversion-symmetry-conserving part of its environment) and an even part $p_{\text{AO}}^{(+)}(D)$ (reflecting the inversion-symmetry-breaking character), and similarly for TiO₂ layers. Then our model for AO layers is

$$\begin{aligned} p_{\text{AO}}^{(-)} &= \mathcal{J} + \mathcal{J}_0 s_0 + \mathcal{J}'_0 t_0 + \mathcal{J}_1 (s_{\bar{1}} + s_1) + \mathcal{J}'_1 (t_{\bar{1}} + t_1) \\ &+ \mathcal{J}_2 (s_{\bar{2}} + s_2) + \mathcal{J}_{01} (s_{\bar{1}} s_0 + s_0 s_1) \\ &+ \mathcal{J}_{02} (s_{\bar{2}} s_0 + s_0 s_2) + \mathcal{J}_{\bar{1}1} s_{\bar{1}} s_1 \\ &+ \mathcal{J}_{12} (s_{\bar{2}} s_{\bar{1}} + s_1 s_2) \end{aligned} \quad (1)$$

and

$$\begin{aligned} p_{\text{AO}}^{(+)} &= \tilde{\mathcal{J}}_1 (s_{\bar{1}} - s_1) + \tilde{\mathcal{J}}'_1 (t_{\bar{1}} - t_1) \\ &+ \tilde{\mathcal{J}}_{01} (s_{\bar{1}} s_0 - s_0 s_1) + \tilde{\mathcal{J}}'_{01} (t_{\bar{1}} s_0 - s_0 t_1) \\ &+ \tilde{\mathcal{J}}'_{10} (s_{\bar{1}} t_0 - t_0 s_1) + \tilde{\mathcal{J}}_{12} (s_{\bar{2}} s_{\bar{1}} - s_1 s_2). \end{aligned} \quad (2)$$

Eq. (1) is equivalent to Eq. (2) of the main text, except that here we introduce the condensed notation $t_j = s_j^2$. The layers are labeled according to the sequence $[\dots \bar{3}.\bar{2}.\bar{1}.0.1.2.3 \dots]$ where numbers represent AO layers, periods represent TiO₂ layers, and the central AO layer 0 is the one whose layer polarization is being expanded. The coefficients $\mathcal{J}(D)$, $\mathcal{J}_0(D)$, ... are odd functions of D ; the $\tilde{\mathcal{J}}_1(D)$, $\tilde{\mathcal{J}}'_1(D)$, ... are even functions of D ; and the combinations of spins appearing in Eqs. (1-2) are the symmetry-appropriate ones. The D -dependence of the coefficients has been suppressed for simplicity of notation.

The corresponding expansion for the layer polarization of a TiO₂ layer is

$$\begin{aligned} p_{\text{TiO}_2}^{(-)} &= \mathcal{K} + \mathcal{K}_1 (s_{\bar{1}} + s_1) + \mathcal{K}'_1 (t_{\bar{2}} + t_2) + \mathcal{K}_2 (s_{\bar{2}} + s_2) \\ &+ \mathcal{K}'_2 (t_{\bar{1}} + t_1) + \mathcal{K}_{12} (s_{\bar{2}} s_{\bar{1}} + s_1 s_2) + \mathcal{K}_{\bar{1}1} s_{\bar{1}} s_1 \\ &+ \mathcal{K}_{\bar{2}1} (s_{\bar{2}} s_1 + s_{\bar{1}} s_2) \end{aligned} \quad (3)$$

and

$$\begin{aligned} p_{\text{TiO}_2}^{(+)} &= \tilde{\mathcal{K}}_1 (s_{\bar{1}} - s_1) + \tilde{\mathcal{K}}'_1 (t_{\bar{2}} - t_2) + \tilde{\mathcal{K}}_2 (s_{\bar{2}} - s_2) \\ &+ \tilde{\mathcal{K}}_{12} (s_{\bar{2}} s_{\bar{1}} - s_1 s_2) + \tilde{\mathcal{K}}'_{12} (s_{\bar{1}} t_{\bar{2}} - t_2 s_1) \\ &+ \tilde{\mathcal{K}}'_{21} (s_{\bar{2}} t_{\bar{1}} - t_1 s_2) + \tilde{\mathcal{K}}_{\bar{1}1} (s_{\bar{1}} t_1 - t_1 s_1) \\ &+ \tilde{\mathcal{K}}'_{\bar{1}2} (s_{\bar{1}} t_2 - t_2 s_{\bar{1}}) + \tilde{\mathcal{K}}_{\bar{2}1} (s_{\bar{2}} s_1 - s_{\bar{1}} s_2) \\ &+ \tilde{\mathcal{K}}'_{21} (s_{\bar{2}} t_1 - t_1 s_2), \end{aligned} \quad (4)$$

where \mathcal{K} and $\tilde{\mathcal{K}}$ coefficients are odd and even in D respectively. Here the labeling scheme is $[\dots \bar{3}.\bar{2}.\bar{1}.1.2.3 \dots]$ with the central TiO₂ layer being the one whose layer polarization is being expanded.

The parameters $\mathcal{J}(D)$, $\tilde{\mathcal{J}}(D)$, $\mathcal{K}(D)$, and $\tilde{\mathcal{K}}(D)$ are expressed as fifth-order Taylor expansions in D , with the Taylor coefficients obtained from a least-squares fitting to the first-principles calculations of $p_j(D)$ for superlattices in the database. Also, recall that the supercell lattice constant $c(D)$ was expanded in terms of parameters $\mathcal{C}_q(D)$ ($q = 1, \dots, 4$) as specified in Eq. (3) of the main text. The $\mathcal{C}(D)$ were similarly fitted to the database, but now as fourth-order Taylor expansions in D .

The fitted parameter values are given in the following tables.

Tables

TABLE I: Fitted first-, third-, and fifth-order-in- D Taylor coefficients of effective cluster interactions for $p_{\text{AO}}^{(-)}(D)$ in Eq. (1), and zeroth-, second-, and fourth-order-in- D Taylor coefficients of effective cluster interactions for $p_{\text{AO}}^{(+)}(D)$ in Eq. (2).

	\mathcal{J}	\mathcal{J}_0	\mathcal{J}'_0	\mathcal{J}_1	\mathcal{J}'_1	\mathcal{J}_2	\mathcal{J}_{01}	\mathcal{J}_{02}	\mathcal{J}_{11}	\mathcal{J}_{12}
1st	2.2771	0.1113	0.0819	0.0034	0.0007	-0.0018	0.0197	0.0031	0.0026	0.0013
3rd	-0.0459	-0.2090	-0.0319	-0.0513	-0.0129	-0.0078	-0.0512	-0.0211	-0.0149	-0.0064
5th	0.2793	0.3690	0.1236	0.1287	0.0533	0.0293	0.0554	0.0339	0.0071	0.0209
	$\tilde{\mathcal{J}}_1$	$\tilde{\mathcal{J}}'_1$	$\tilde{\mathcal{J}}_{01}$	$\tilde{\mathcal{J}}'_{01}$	$\tilde{\mathcal{J}}'_{10}$	$\tilde{\mathcal{J}}_{12}$				
0th	-0.0210	-0.0044	-0.0015	0.0001	-0.0002	-0.0005				
2nd	0.0657	0.0014	0.0047	0.0037	0.0029	-0.0017				
4th	-0.0569	-0.0035	0.0010	-0.0139	-0.0164	0.0081				

TABLE II: Fitted first-, third-, and fifth-order-in- D Taylor coefficients of effective cluster interactions for $p_{\text{TiO}_2}^{(-)}(D)$ in Eq. (3), and zeroth-, second-, and fourth-order Taylor coefficients of effective cluster interactions for $p_{\text{TiO}_2}^{(+)}(D)$ in Eq. (4).

	\mathcal{K}	\mathcal{K}_1	\mathcal{K}'_1	\mathcal{K}_2	\mathcal{K}'_2	\mathcal{K}_{12}	\mathcal{K}_{11}	\mathcal{K}_{21}		
1st	1.5292	-0.0097	-0.0060	0.0075	-0.0013	-0.0076	0.0033	-0.0018		
3rd	0.0415	0.1203	0.0376	0.0035	0.0074	0.0148	0.0025	0.0091		
5th	-0.0468	-0.1727	-0.0702	-0.0421	-0.0068	-0.0084	-0.0095	-0.0219		
	$\tilde{\mathcal{K}}_1$	$\tilde{\mathcal{K}}'_1$	$\tilde{\mathcal{K}}_2$	$\tilde{\mathcal{K}}_{12}$	$\tilde{\mathcal{K}}'_{12}$	$\tilde{\mathcal{K}}'_{21}$	$\tilde{\mathcal{K}}'_{11}$	$\tilde{\mathcal{K}}'_{12}$	$\tilde{\mathcal{K}}_{21}$	$\tilde{\mathcal{K}}'_{21}$
0th	0.0634	-0.0057	0.0027	0.0023	0.0004	0.0001	0.0010	-0.0001	-0.0003	0.0000
2nd	-0.1468	0.0101	-0.0200	0.0106	-0.0057	-0.0169	0.0089	-0.0018	0.0014	0.0063
4th	0.0522	0.0059	0.0045	-0.0470	0.0247	0.0679	-0.0447	0.0147	-0.0030	-0.0372

TABLE III: Fitted zeroth-, second-, and fourth-order-in- D Taylor coefficients of effective cluster interactions for $c(D)$ in Eq. (3) of main text.

	\mathcal{C}_1	\mathcal{C}_2	\mathcal{C}_3	\mathcal{C}_4
0th	7.2305	0.1968	0.0543	0.0097
2nd	0.4430	0.3293	0.1675	0.0832
4th	-0.1926	-0.6367	-0.1546	-0.3744