

## INTRODUCTION TO MANY BODY PHYSICS: 620. Fall 2025

### Questions 6. Finite Temperature and superfluidity (Due Fri, 12th Dec. )

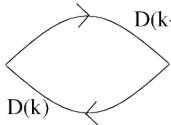
1. Use the method of complex contour integration to carry out the Matsubara sums in the following:

- (i) Derive the density of a spinless Bose Gas at finite temperature from the boson propagator  $D(k) \equiv D(\mathbf{k}, i\nu_n) = [i\nu_n - \omega_{\mathbf{k}}]^{-1}$ , where  $\omega_{\mathbf{k}} = E_{\mathbf{k}} - \mu$  is the energy of a boson, measured relative to the chemical potential.

$$\rho(T) = \frac{N}{V} = V^{-1} \sum_{\mathbf{k}} \langle T b_{\mathbf{k}}(0^-) b_{\mathbf{k}}^\dagger(0) \rangle = -(\beta V)^{-1} \sum_{i\nu_n, \mathbf{k}} D(k) e^{i\nu_n 0^+}. \quad (1)$$

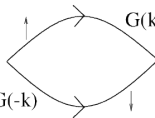
How do you need to modify your answer to take account of Bose Einstein condensation?

- (ii) The dynamic charge-susceptibility of a free Bose gas, i.e

$$\chi_c(q, i\nu_n) = \text{Diagram} = T \sum_{i\nu_r} \int \frac{d^3k}{(2\pi)^3} D(q+k) D(k). \quad (2)$$


where  $\nu_r$  is the Bose Matsubara frequency of the internal loop. Please analytically extend your final answer to real frequencies ( $i\nu_n \rightarrow \nu$ ).

“pair-susceptibility” of a spin-1/2 free Fermi gas, i.e.

$$\chi_P(q, i\nu_n) = \text{Diagram} = T \sum_{i\omega_r} \int \frac{d^3k}{(2\pi)^3} G(q+k) G(-k) \quad (3)$$


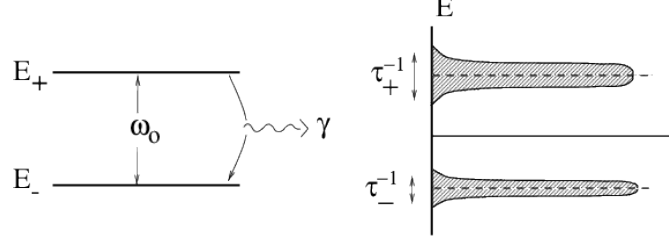
where  $G(k) \equiv G(\mathbf{k}, i\omega_n) = [i\omega_n - \epsilon_{\mathbf{k}}]^{-1}$ . (Note the direction of the arrows: why is there no minus sign for the Fermion loop?) Show that the static pair susceptibility,  $\chi_P(0)$  is given by

$$\chi_P = \int \frac{d^3k}{(2\pi)^3} \frac{\tanh[\beta\epsilon_{\mathbf{k}}/2]}{2\epsilon_{\mathbf{k}}} \quad (4)$$

Can you see that this quantity diverges at low temperatures? How does it diverge, and why? (Hint: replace the integration near the Fermi surface by an integration over the density of states, which you may treat as a constant to obtain the divergence)

2. A simple model an atom with two atomic levels coupled to a radiation field is described by the Hamiltonian

$$H = H_o + H_I + H_{\text{photon}}, \quad (5)$$



where

$$H_o = \tilde{E}_- c_-^\dagger c_- + \tilde{E}_+ c_+^\dagger c_+ \quad (6)$$

describes the atom, treating it as a *fermion*

$$H_I = V^{-1/2} \sum_{\vec{q}} g(\omega_{\vec{q}}) \left( c_+^\dagger c_- + c_-^\dagger c_+ \right) \left[ a_{\vec{q}}^\dagger + a_{-\vec{q}} \right] \quad (7)$$

describes the coupling to the radiation field ( $V$  is the volume of the box enclosing the radiation) and

$$H_{\text{photon}} = \sum_{\vec{q}} \omega_{\vec{q}} a_{\vec{q}}^\dagger a_{\vec{q}}, \quad (\omega_{\vec{q}} = cq) \quad (8)$$

is the Hamiltonian for the electromagnetic field. The “dipole” matrix element  $g(\omega)$  is weak enough to be treated by second order perturbation theory and the polarization of the photon is ignored.

- (i) Calculate the self-energy  $\Sigma_+(\omega)$  and  $\Sigma_-(\omega)$  for an atom in the  $+$  and  $-$  states.
- (ii) Use the self-energy obtained above to calculate the life-times  $\tau_\pm$  of the atomic states, i.e.

$$\tau_\pm^{-1} = 2\text{Im}\Sigma_\pm(\tilde{E}_\pm - i\delta). \quad (9)$$

If the gas of atoms is non-degenerate, i.e the Fermi functions are all small compared with unity,  $f(E_\pm) \sim 0$  show that

$$\begin{aligned} \tau_+^{-1} &= 2\pi |g(\omega_o)|^2 F(\omega_o) [1 + n(\omega_o)] \\ \tau_-^{-1} &= 2\pi |g(\omega_o)|^2 F(\omega_o) n(\omega_o), \end{aligned} \quad (10)$$

where  $\omega_o = \tilde{E}_+ - \tilde{E}_-$  is the separation of the atomic levels and

$$F(\omega) = \int \frac{d^3q}{(2\pi)^3} \delta(\omega - \omega_q) = \frac{\omega^2}{2\pi c^3} \quad (11)$$

is the density of state of the photons at energy  $\omega$ . What do these results have to do with stimulated emission? Do your final results depend on the initial assumption that the atoms were fermions?

- (iii) Why is the decay rate of the upper state larger than the decay rate of the lower state by the factor  $[1 + n(\omega_0)]/n(\omega_0)$ ? (Hint: appeal to detailed balance)

3. A one dimensional superconductor of *spinless* electrons is described by a p-wave BCS model

$$H = \sum_k \epsilon_k c_k^\dagger c_k - \frac{g}{L} \sum_{k,k'} \gamma_k \gamma_{k'} (c_k^\dagger c_{-k}^\dagger) (c_{-k'} c_{k'}) \quad (12)$$

where  $\epsilon_k = -2t \cos k - \mu$  is the dispersion of the electrons, while  $\gamma_k = \sin k$  is the (odd-parity) form factor of the pairing. Consider periodic boundary conditions on a chain with  $L$  sites, so that the momenta take the discrete values  $k = \frac{2\pi n}{L}$  where  $n$  is an integer and  $k \in [-\pi, \pi]$ .

- (i) Why must  $\gamma_k$  be an odd-parity function of momentum, i.e  $\gamma_k = -\gamma_{-k}$ ?  
(ii) Show that the BCS mean-field Hamiltonian for this model takes the form

$$H_{BCS} = \sum_{k \in [0, \pi]} \psi_k^\dagger \begin{pmatrix} \epsilon_k & 2\Delta\gamma_k \\ 2\bar{\Delta}\gamma_k & -\epsilon_k \end{pmatrix} \psi_k + L \frac{\bar{\Delta}\Delta}{g} - L\mu$$

where  $\bar{\Delta} = \Delta^*$  and

$$\psi_k = \begin{pmatrix} c_k \\ c_{-k}^\dagger \end{pmatrix}, \quad \psi_k^\dagger = (c_k^\dagger, c_{-k})$$

are Nambu spinors. (Key question: why does the summation run over half the Brillouin zone?)

- (iii) What is the mean-field relationship between  $\Delta$  and the fermion pair density  $\langle c_{-k} c_k \rangle$ ?  
(iv) Calculate the Nambu Greens function  $\mathcal{G}_{\alpha\beta}(k, i\omega_n) = -\int_0^\beta d\tau \langle T \psi_{k\alpha}(\tau) \psi_{k\beta}^\dagger(0) \rangle$  using the above mean-field Hamiltonian. How does the anomalous (Gor'kov) propagator

$$F(k, i\omega_n) = -\int_0^\beta \langle T c_k(\tau) c_{-k}(0) \rangle e^{i\omega_n \tau}$$

differ from an s-wave superconductor?

- (v) Calculate the quasiparticle excitation spectrum of this superconductor.  
(vi) Calculate the free energy as a function of  $\Delta = |\Delta|e^{i\phi}$ . Plot or sketch your result, illustrating the phase *and* amplitude dependence of the Free energy. One way to do this is to sum the Free energy of the quasiparticles, adding in the term  $L|\Delta|^2/g$ . At what value of  $\Delta$  is the free energy minimized?  
(vii) Rewrite the mean-field Hamiltonian of 3(b) in real space for the special case where  $\Delta = it$  and  $\mu = 0$ . What happens to the excitation spectrum when the link between the first and last sites is cut?