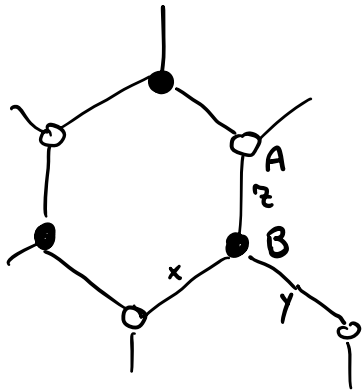


KITAEV MODEL.

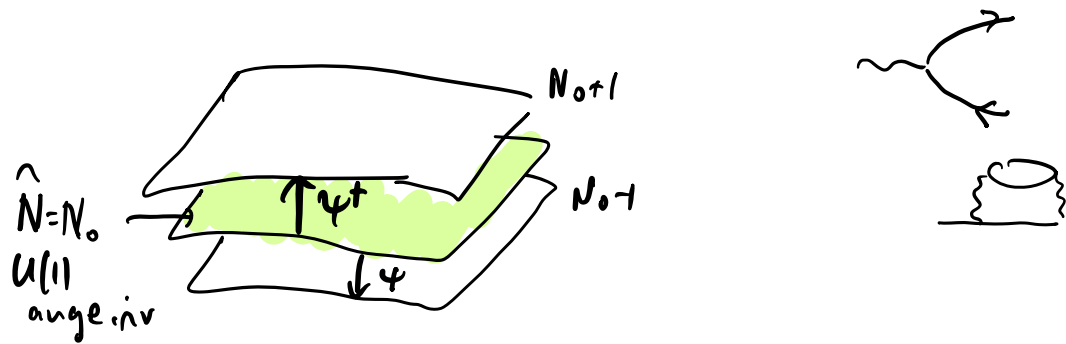
Kitaev, Annals of Physics, 321,2 (2006).

Method 2: Ancillary Qubits



$$H = -\left(\frac{1}{2}\right) \sum_{\langle i,j \rangle} K_{\alpha_{ij}} \sigma_i^{\alpha_{ij}} \sigma_j^{\alpha_{ij}}$$

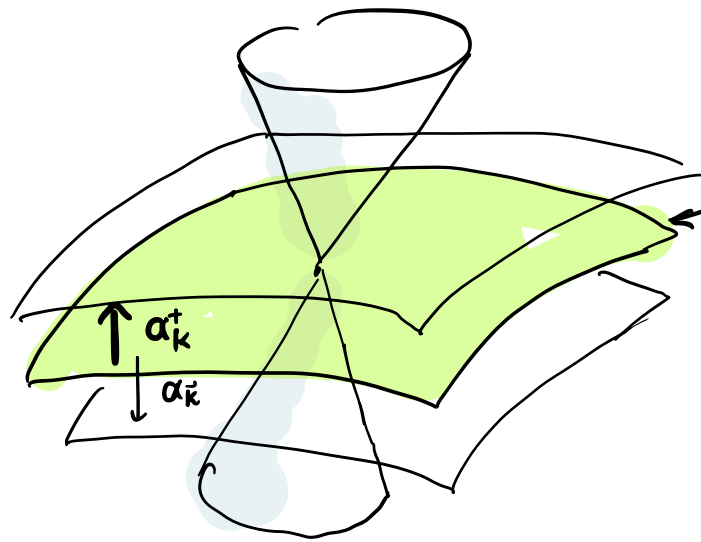
The Hilbert space of the spins can be considered to be a slice through the Fock-space of fractionalized excitations.



The description of many body physics in terms of field theory relies on expanding a Hilbert space of definite particle number to a Fock space, in which the conservation of charge is protected by gauge invariance

$$A_\mu \rightarrow A_\mu + \partial_\mu \phi$$

$$0 = \int \frac{\delta S}{\delta A_\mu} \delta_\mu \phi = \int J^\mu \delta_\mu \phi = - \int \phi \partial_\mu J^\mu$$



Physical Hilbert Space

$$|\Psi_{\text{phys}}\rangle = \prod_j \frac{1}{2} (1 + D_j) |\Psi\rangle$$

$$D_j = -2i\alpha_j\beta_j$$

To each site, we add an ancillary qubit $f^+ = \frac{\alpha - i\beta}{\sqrt{2}}$, $f = \frac{\alpha + i\beta}{\sqrt{2}}$,

$$f^+ f - \frac{1}{2} = -i\alpha\beta \quad D_j = -2i\alpha_j\beta_j = 1 \text{ (choice)}$$

$$\beta_j \vec{\sigma}_j = \vec{b}_j$$

$$\Rightarrow \vec{\sigma}_j = 2\beta_j \vec{b}_j = -i\vec{b}_j \times \vec{b}_j$$

$$\{b_i^\alpha, b_i^\delta\} = \delta^{\alpha\delta} \delta_{je}$$

$$\beta_j = -2i b_j^1 b_j^2 b_j^3$$

$$\text{(check } -2i b_j^1 b_j^2 b_j^3 = -2i \beta_j^3 \overbrace{\sigma_j^x \sigma_j^y \sigma_j^z}^i = \beta_j \text{)}$$

Unfortunately β_j is not independent of the b_j^α . However in the physical case where $D_j = 1$, $-i\alpha_j = \beta_j$, i.e

$$\vec{\sigma}_j = -2i\alpha_j \vec{b}_j$$

$$\begin{aligned}
 H &\stackrel{!}{=} \sum_{\langle i,j \rangle} K^{\gamma_{ij}} \sigma_i^{\gamma_{ij}} \sigma_j^{\gamma_{ij}} = +2 \sum_{\langle i,j \rangle} K^{\gamma} (\alpha_i b_i^{\gamma_{ij}}) (\alpha_j b_j^{\gamma_{ij}}) \\
 &\quad \delta_{ij} = \{x,y,z\} \\
 &= -2 \sum_{\langle i,j \rangle} K^{\gamma_{ij}} (\alpha_i \alpha_j) b_i^{\gamma_{ij}} b_j^{\gamma_{ij}}
 \end{aligned}$$

$$\begin{aligned}
 H &= \sum_{\langle i,j \rangle} K^{\gamma_{ij}} (i\alpha_i \alpha_j u_{ij}) \\
 u_{ij} &= +2i b_i^{\gamma_{ij}} b_j^{\gamma_{ij}} = \pm 1 \text{ GAUGE FIELD}
 \end{aligned}$$

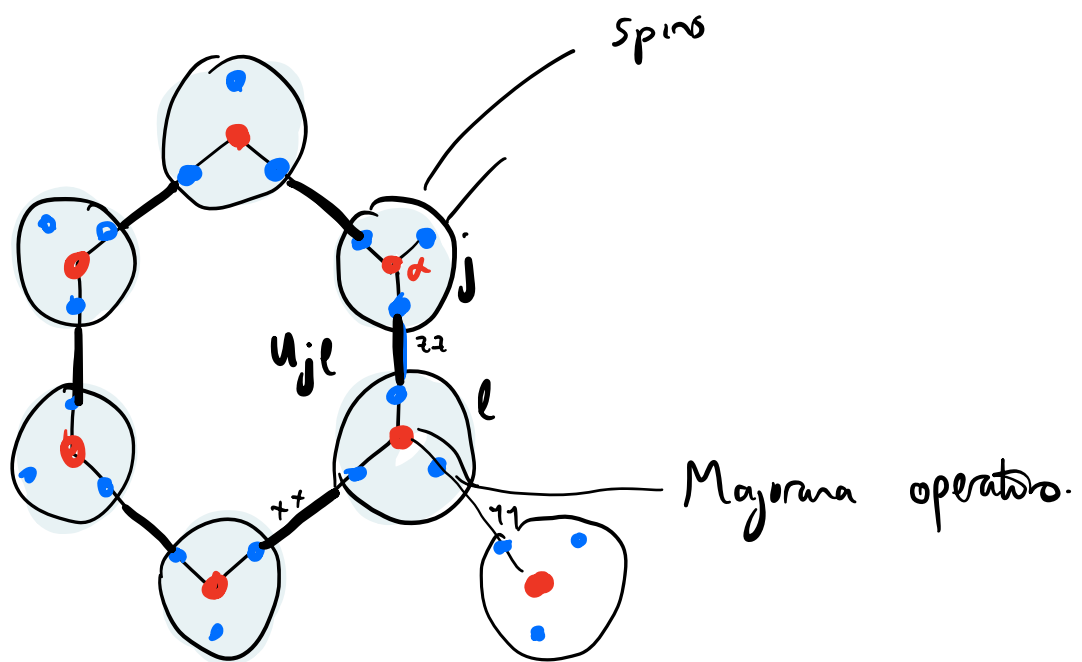
$$D_j = -2i\alpha_j \beta_j = -4\alpha_j b_j^1 b_j^2 b_j^3 = 1$$

$$\left. \begin{aligned}
 a_i &\rightarrow z_i a_i \\
 u_{ij} &\rightarrow z_i u_{ij} z_j
 \end{aligned} \right\}$$

$$z_i = \pm 1$$

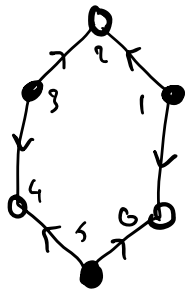
$$z_i z_j$$

GAUGE INVARIANCE.



The product

$$\hat{W}_P = \prod_{(e+1, e) \in P} (\hat{U}_{e+1, e})$$



$$= U_{12} U_{23} U_{34} U_{45} U_{56} U_{61} \quad -2i b_6^x b_1^x$$

$$= -U_{12} U_{23} U_{34} U_{45} U_{56} U_{61}$$

$$= 2^6 b_1^y b_2^y b_2^x b_3^x b_3^z b_4^z b_4^y b_5^y b_5^x b_6^x b_6^z b_1^z$$

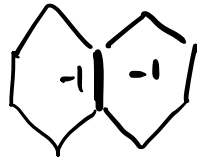
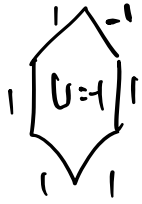
$$= -2^6 (b_1^z b_1^y) (b_2^y b_2^x) (b_3^x b_3^z) (b_4^y b_4^y) (b_5^x b_5^x) (b_6^z b_1^z)$$

$$-2^6 \left(\frac{-i\sigma_1^x}{2} \right) \left(\frac{-i\sigma_2^z}{2} \right) \left(\frac{-i\sigma_3^x}{2} \right) \left(\frac{-i\sigma_4^x}{2} \right) \left(\frac{-i\sigma_5^z}{2} \right) \left(\frac{-i\sigma_6^y}{2} \right)$$

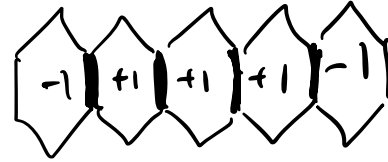
$$= \prod_{j \in P} \sigma_j^{a_j}$$

$$[W_P, H] = 0$$

Ground state energy is purely a functional of the W_P , infact $W_P = +1$ defines the GS.

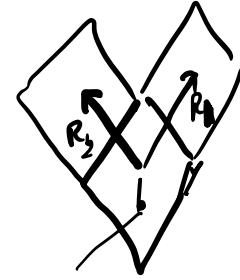


Two Visions



Two visions, separated.

In the GS all $W_p = +1$.



$$R_2 = \left(-\frac{\sqrt{3}}{2}, \frac{3}{2} \right)$$

$$R_1 = \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right)$$

$$H = \frac{1}{2} \sum_k \alpha_{ij} \left(\gamma(R_i - R_j) \alpha_A^{(i)} \alpha_B^{(j)} + h.c. \right)$$

$$\gamma(R) = i \left(k^z \delta_{R,0} + k^x \delta_{R,R_1} + k^y \delta_{R,R_2} \right)$$

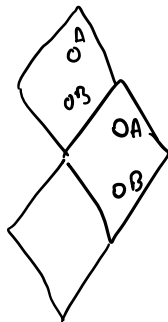
$$\begin{aligned} \langle k | R \rangle \gamma(R) \langle 0 | k \rangle \\ = \sum_R \gamma(R) e^{+ik \cdot R} \end{aligned}$$

$$\alpha_{\mathbf{k}} = \frac{1}{\sqrt{N}} \sum_j \begin{pmatrix} \alpha_A(j) \\ \alpha_B(j) \end{pmatrix} e^{-i\vec{k} \cdot \vec{R}_j} = \begin{pmatrix} \alpha_{\mathbf{k}A} \\ \alpha_{\mathbf{k}B} \end{pmatrix}$$

$$\begin{pmatrix} \alpha_A(j) \\ \alpha_B(j) \end{pmatrix} = \frac{1}{\sqrt{N}} \sum_{\mathbf{k} \in \text{BZ}} \alpha_{\mathbf{k}} e^{i\vec{k} \cdot \vec{R}_j}$$

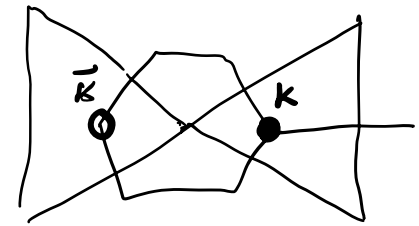
$$= \frac{1}{\sqrt{N}} \sum_{\mathbf{k} \in \Delta} \left(\alpha_{\mathbf{k}} e^{i\vec{k} \cdot \vec{R}_j} + \alpha_{\mathbf{k}}^{\dagger} e^{-i\vec{k} \cdot \vec{R}_j} \right)$$

$$H = \sum_{\mathbf{k} \in \text{BZ}} \frac{k}{2} \alpha_{\mathbf{k}}^{\dagger} \begin{pmatrix} 0 & v_{\mathbf{k}} \\ v_{\mathbf{k}}^{\dagger} & 0 \end{pmatrix} \alpha_{\mathbf{k}}$$



$$\langle \mathbf{k} | \mathcal{R} \rangle \delta(\mathcal{R}) \langle 0 | \mathbf{k} \rangle = \sum_{\mathcal{R}} \delta(\mathcal{R}) e^{-i\vec{k} \cdot \vec{R}_j}$$

$$\begin{aligned}
 H &= \frac{1}{2} \sum_{\mathbf{k} \in \square} \alpha_{\mathbf{k}}^{\dagger} \begin{pmatrix} 0 & \gamma_{\mathbf{k}} \\ \gamma_{\mathbf{k}}^* & 0 \end{pmatrix} \alpha_{\mathbf{k}} \\
 &= \sum_{\mathbf{k} \in \triangle} \alpha_{\mathbf{k}}^{\dagger} \begin{pmatrix} 0 & \delta_{\mathbf{k}} \\ \gamma_{\mathbf{k}}^* & 0 \end{pmatrix} \alpha_{\mathbf{k}} \\
 &= \sum_{\mathbf{k} \in \triangle} \alpha_{\mathbf{k}}^{\dagger} (\vec{\gamma}_{\mathbf{k}} \cdot \vec{\tau}) \alpha_{\mathbf{k}}
 \end{aligned}$$



$$\begin{aligned}
 \gamma_{\mathbf{k}} &= i \left[k_z + k_y e^{ik \cdot \mathbf{r}_2} + k_x e^{ik \cdot \mathbf{r}_3} \right] \\
 &= (\text{Re} \gamma_{\mathbf{k}}, -\text{Im} \gamma_{\mathbf{k}}, 0)
 \end{aligned}$$

$$(\gamma_{\mathbf{k}} \cdot \vec{\tau}) \begin{pmatrix} u & -v^* \\ v & u^* \end{pmatrix} = \begin{pmatrix} u & -v^* \\ v & u^* \end{pmatrix} \begin{pmatrix} |\gamma_{\mathbf{k}}| & \\ & -|\gamma_{\mathbf{k}}| \end{pmatrix} \quad \begin{aligned} \delta v &= |\delta| u \\ \delta u^* &= |\delta| v \end{aligned}$$

$$\begin{aligned}
 v_{\mathbf{k}} &= \delta^* / \sqrt{2} |\delta_{\mathbf{k}}| \\
 u_{\mathbf{k}} &= \frac{1}{\sqrt{2}}
 \end{aligned}
 \quad \begin{pmatrix} \gamma & \\ \gamma^* & \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{\delta^*}{|\delta| \sqrt{2}} \end{pmatrix} = \begin{pmatrix} \frac{|\delta|}{|\delta| \sqrt{2}} \\ \frac{\delta^*}{\sqrt{2}} \end{pmatrix} = |\delta| \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{\gamma^*}{|\delta| \sqrt{2}} \end{pmatrix}$$

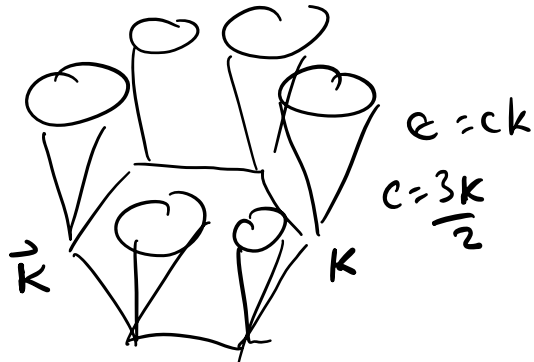
$$\alpha_{\mathbf{k}} = \overbrace{\begin{pmatrix} u & -v^* \\ v & u^* \end{pmatrix}}^u \begin{pmatrix} n_{k1} \\ n_{k2} \end{pmatrix}$$

$$\alpha = u n \quad n_{\mathbf{k}}^{\dagger} = \alpha^{\dagger} u_{\mathbf{k}}$$

$$H = \sum_{\mathbf{k} \in \triangle} \epsilon(\mathbf{k}) \left[n_{k1}^{\dagger} n_{k1} - n_{k2}^{\dagger} n_{k2} \right]$$

$$|\psi_{gs}\rangle = \prod_j \left(\frac{1+D_j}{2} \right) \prod_{k \in \Delta} n_{kz}^+ |\phi\rangle$$

$$k^x = k^y = k^z = k$$



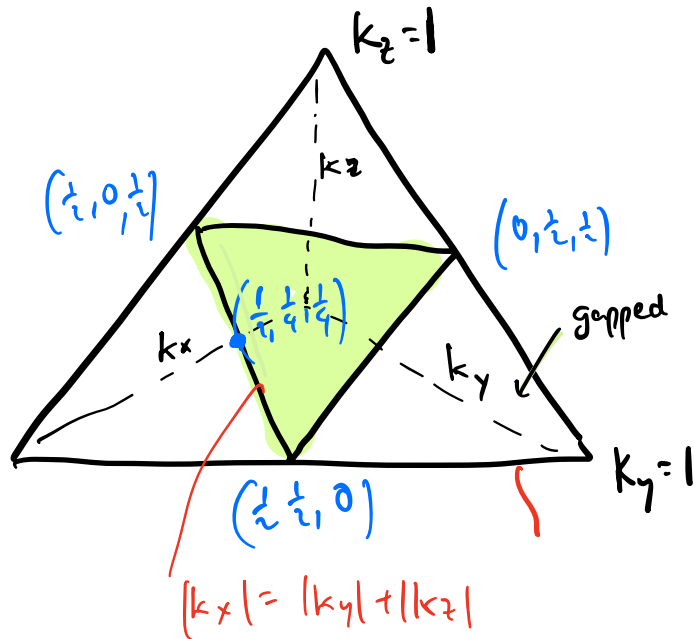
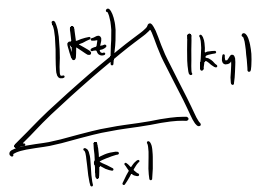
$$e(k) = k \left| 1 + e^{i\vec{k} \cdot \vec{R}_1} + e^{i\vec{k} \cdot \vec{R}_2} \right|$$

$$\left[k_z + k_x e^{ik \cdot R_1} + k_y e^{ik \cdot R_2} \right] = 0$$

$$\text{I } |k_x| \leq |k_y| + |k_z|$$

$$|k_y| \leq |k_x| + |k_z|$$

$$|k_z| \leq |k_x| + |k_y|$$



$$k_x + k_y + k_z = 1$$

$$k_x = 1 - 2k < 2k$$

$$\begin{matrix} 4k > 1 \\ k > 1/4 \end{matrix}$$

$$k_x = k_y = k_z = k$$

$$\chi(k) = ik(1 + e^{i\vec{k}\cdot\vec{R}_1} + e^{i\vec{k}\cdot\vec{R}_2})$$

$$k_x = k_y = k_z = k$$

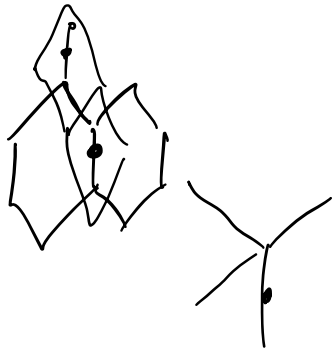
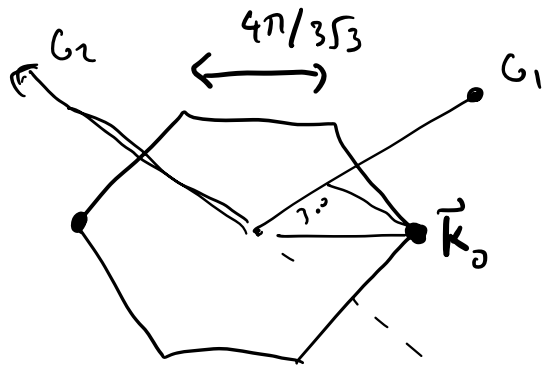
$$\vec{k} = k_1 \vec{G}_1 + k_2 \vec{G}_2$$

$$\vec{k}\cdot\vec{R}_1 = 2\pi k_1 = 2\pi/3$$

$$\vec{k}\cdot\vec{R}_2 = 2\pi k_2 = -2\pi/3$$

$$\vec{k} = \frac{2}{3} \vec{G}_2 + \frac{1}{3} \vec{G}_1$$

$$-\vec{k} = -\frac{2}{3} \vec{G}_2 - \frac{1}{3} \vec{G}_1 = \frac{1}{3} \vec{G}_2 + \frac{2}{3} \vec{G}_1$$



$$n_1 = \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right)$$

$$n_2 = \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right)$$

$$\vec{G}_1 = \frac{2\pi}{3} (\sqrt{3}, 1)$$

$$\vec{G}_2 = \frac{2\pi}{3} (-\sqrt{3}, 1)$$

$$\propto \frac{2\sqrt{3}}{2} = 2\pi$$

$$\vec{k}_0 = \frac{2}{3} (\vec{G}_1 - \vec{G}_2)$$

$$= \frac{4\pi}{3\sqrt{3}} (1, 1, 0)$$

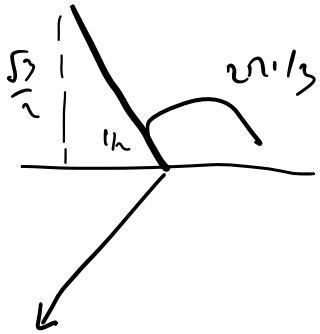
$$\vec{G}_1 \cdot n_1 = \frac{2\pi}{3} \left(\frac{3}{2} + \frac{3}{2} \right) = \frac{2\pi}{3} \cdot 3 = 2\pi \checkmark$$

$$\frac{2\pi}{2 \cdot 3} (3 + 3) = 2\pi$$

Velocity

$$\gamma(\vec{k}_0 + \delta k) = \nabla_{\vec{k}} \gamma \cdot \delta \vec{k}$$

$$\begin{aligned} \nabla_{\vec{k}} \gamma &= ik \left(i\vec{R}_2 e^{i\vec{k}_0 \cdot \vec{R}_1} + i\vec{R}_2 e^{i\vec{k}_0 \cdot \vec{R}_2} \right) \\ &= -k \left(e^{i\vec{k}_0 \cdot \vec{R}_1} \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right) + e^{i\vec{k}_0 \cdot \vec{R}_2} \left(-\frac{\sqrt{3}}{2}, \frac{3}{2} \right) \right) \\ &= -k \left(e^{2\pi i/3} \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right) + e^{-2\pi i/3} \left(-\frac{\sqrt{3}}{2}, \frac{3}{2} \right) \right) \\ &= -k \left(\left(-\frac{1}{2} + \frac{\sqrt{3}i}{2} \right) \left(\frac{\sqrt{3}}{2}, \frac{3}{2} \right) + \left(-\frac{1}{2} - \frac{\sqrt{3}i}{2} \right) \left(-\frac{\sqrt{3}}{2}, \frac{3}{2} \right) \right) \\ &= \frac{3}{2} k (-i, 1) \end{aligned}$$



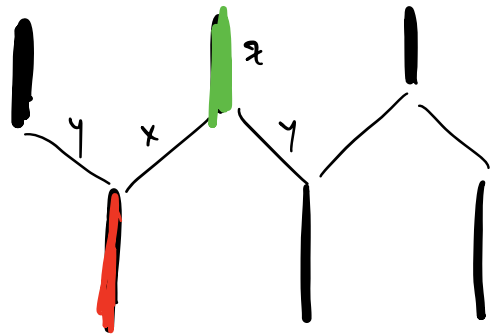
$$\gamma(k) \approx \frac{3k}{2} (-i\delta k_x + \delta k_y)$$

$$E(k) = \frac{3k}{2} \sqrt{\delta k_x^2 + \delta k_y^2}$$

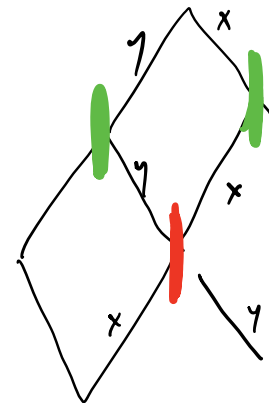
$$\Rightarrow \begin{cases} c = \frac{3k}{2} \\ e = c \delta k \end{cases}$$

GAPPED PHASE = Anyons

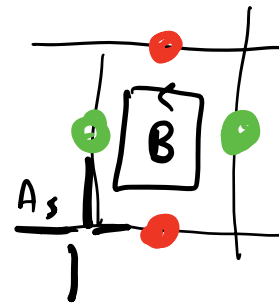
$$|k_x| + |k_y| \leq |k_z|$$



(TT) or (LL)



$$H_{\text{eff}} = -\mathcal{J}_{\text{eff}} \left(\sum_{\text{vertices}} A_s + \sum_{\text{plaquettes}} B_p \right)$$



$$A_s = \prod_{\text{sites}} \sigma_j^x$$

$$B_p = \prod_{\text{bonds}} \sigma_j^z$$

Toric code

$$Z_{\text{eff}} = -Z_x^2 Z_y^2 / 16(Z_z)^3.$$