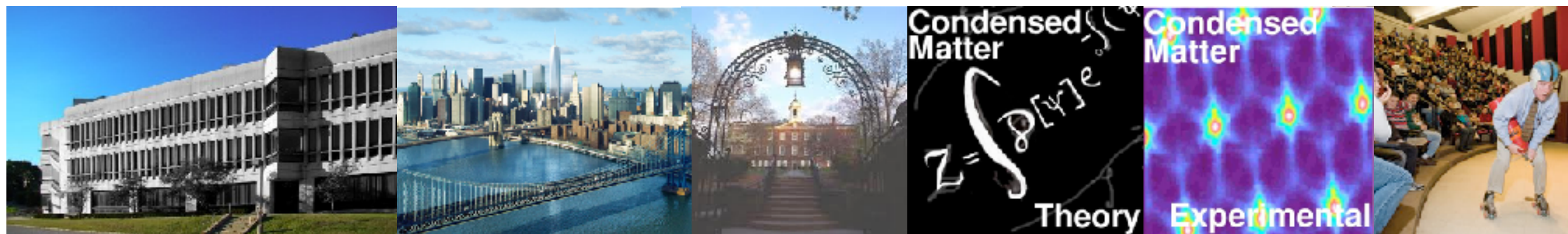


Heavy Fermion Physics: a 21st Century perspective

Piers Coleman: Rutgers Center for Materials Theory, USA

Julich,
21 Sept, 2015



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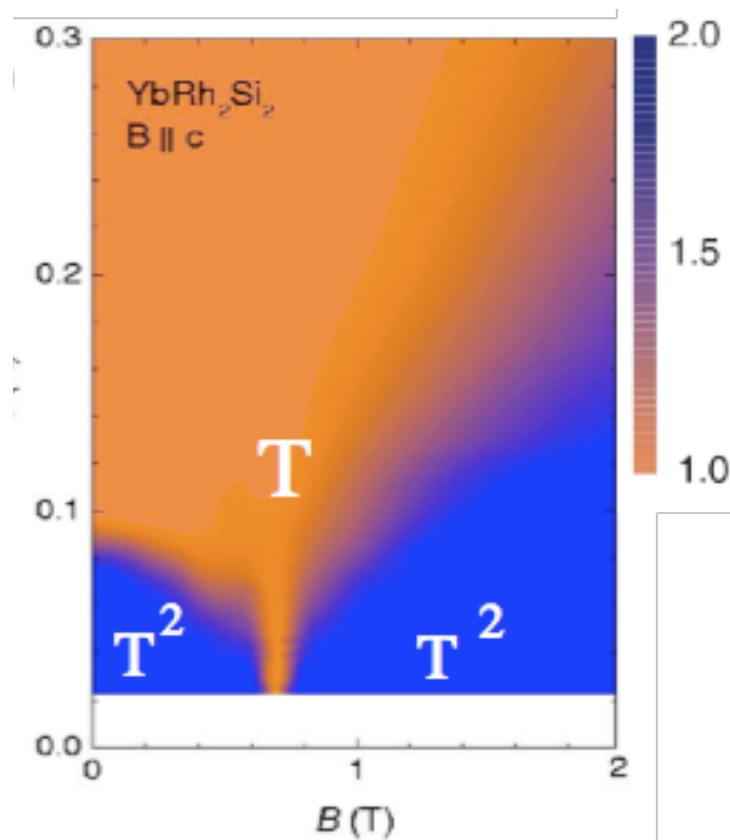


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Quantum Criticality & Strange Metals

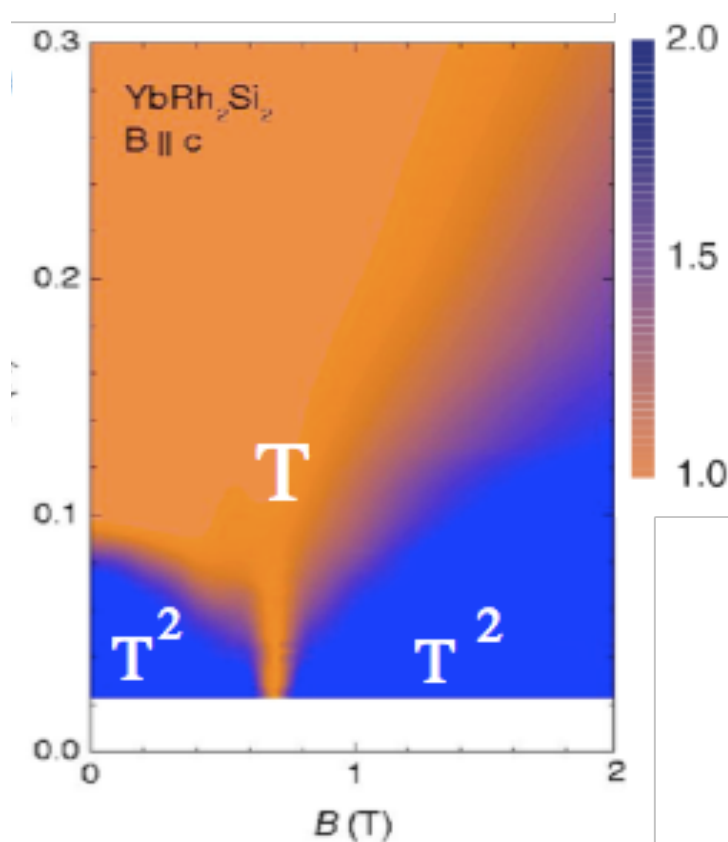


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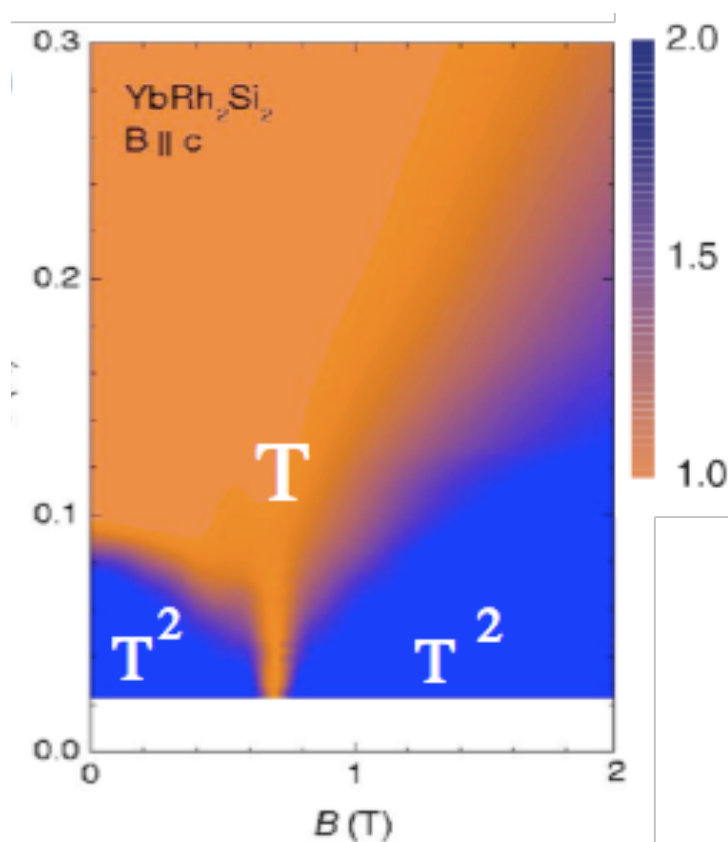


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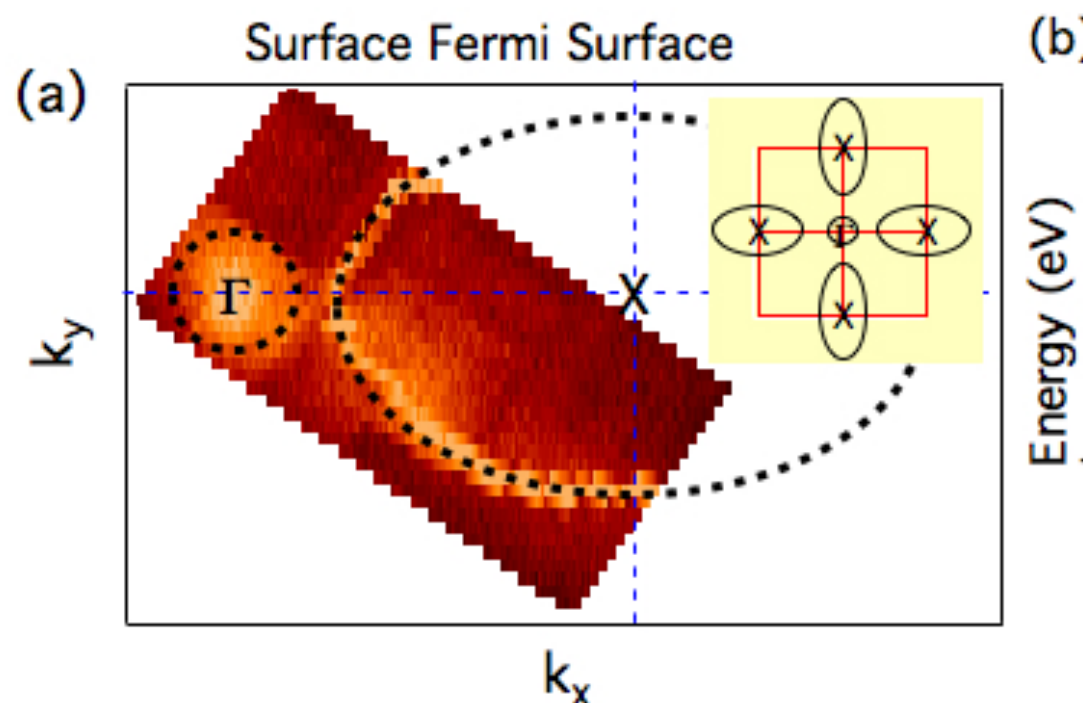
Quantum Criticality & Strange Metals



Heavy Fermion Superconductivity



Topological Kondo Insulators

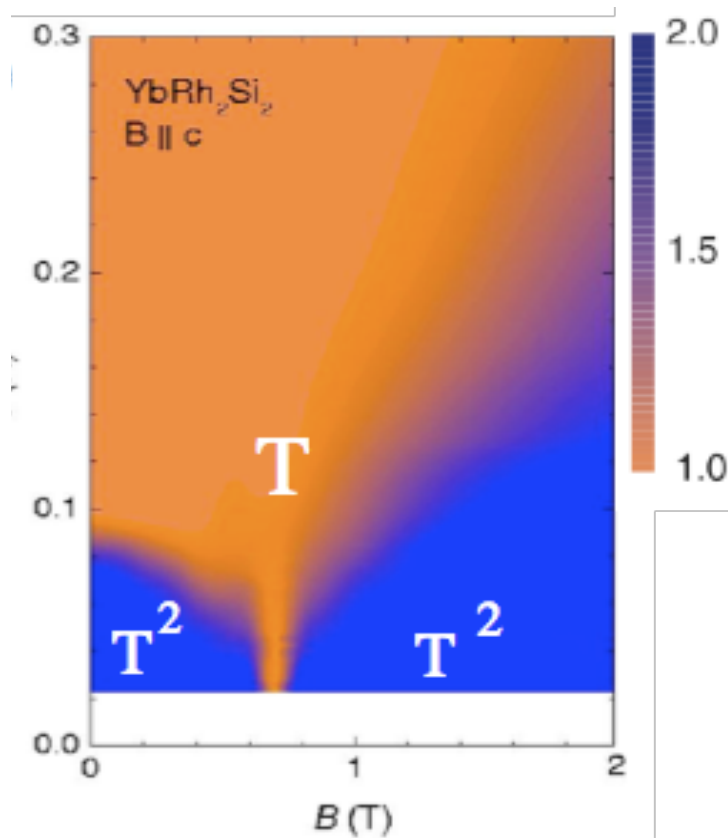


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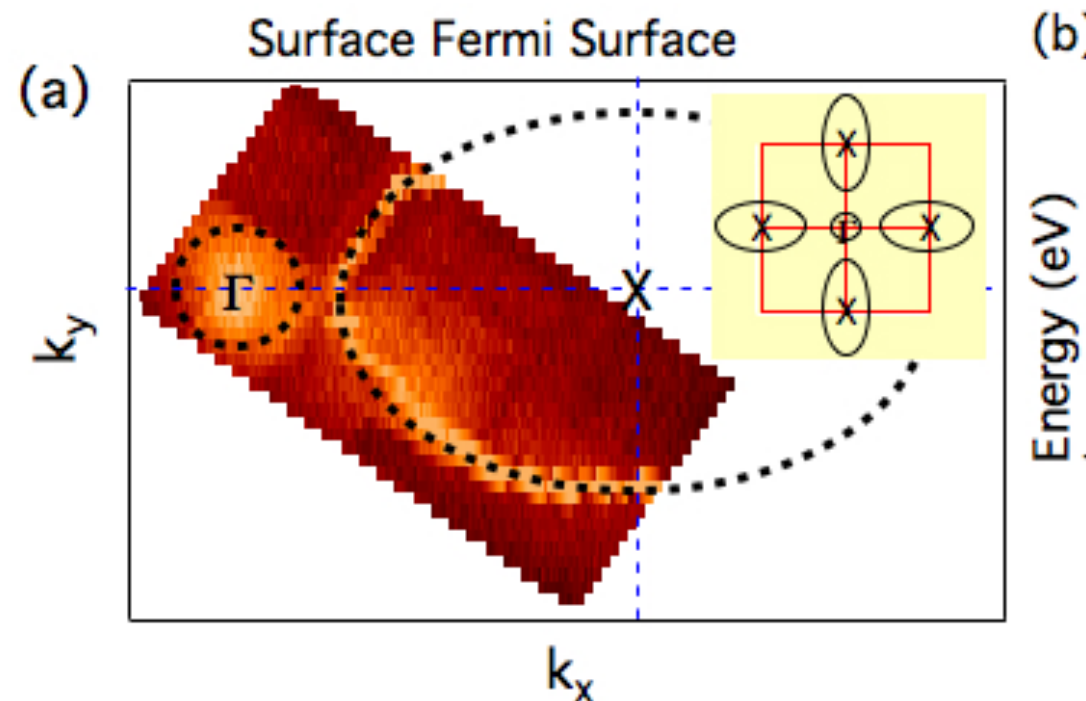
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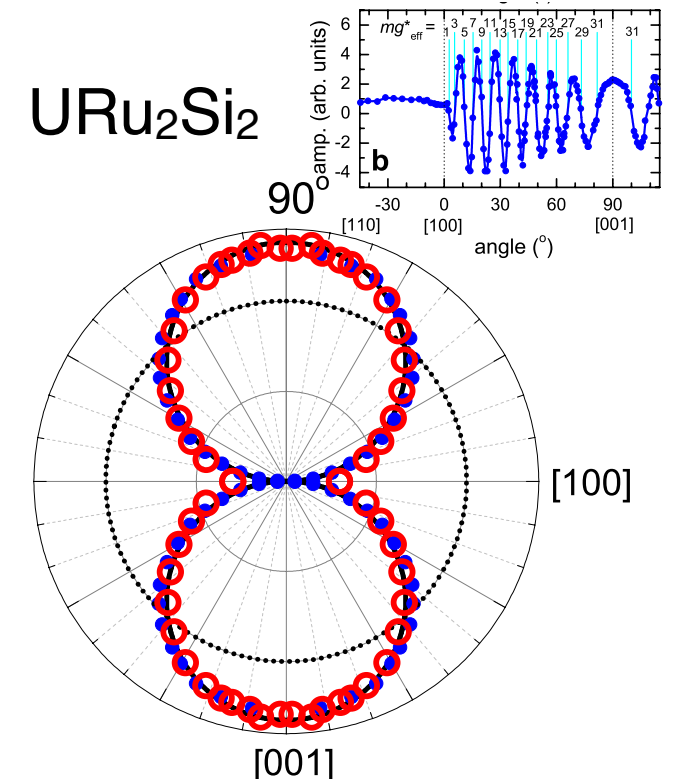
Heavy Fermion Superconductivity



Topological Kondo Insulators



Hidden Order



Collaborators.

Y Komijani	U. Cincinatti
Q. Si	Rice
R. Ramazashvili	CNRS, Toulouse
C. Pepin	CEA, Saclay
Aline Ramires	PSI, Zurich
Jerome Rech	CNRS, Marseille
Rebecca Flint	Iowa State
Premi Chandra	Rutgers
Andriy Nevidomskyy	Rice
Alexei Tselik	Brookhaven NL
Hai-Young Kee	U. Toronto
Natan Andrei	CMT, Rutgers
Onur Erten	ASU, Tempe
Tamaghna Hazra	CMT, Rutgers
Eduardo Miranda	Campinas
Maxim Dzero	Kent State
Victor Galitski	U. Maryland
Kai Sun	U. Michigan
Gergely Zarand	TU, Budapest

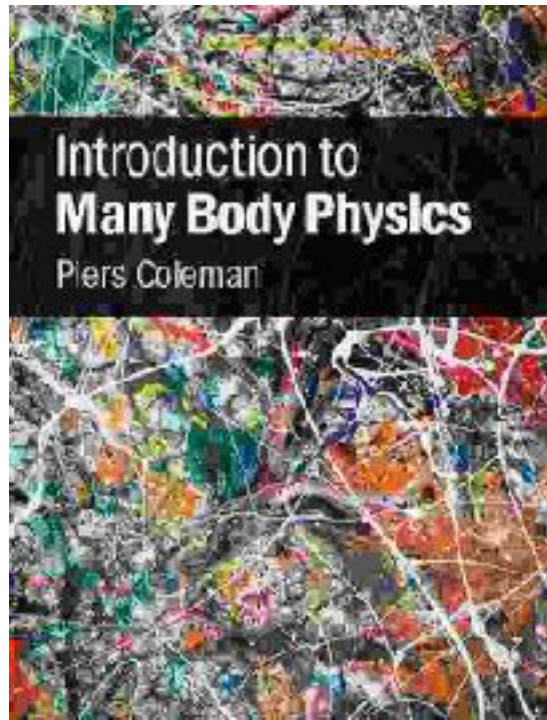
Experimentalists:

H. von Lohneysen	Karlsruhe
G. Aeppli	ETH, Zurich
A. Schröder	Kent State
S. Nakatsuji	ISSP
G. Lonzarich	Cambridge
S. Paschen	Vienna
J. Thompson	Los Alamos
J. Allen	U. Michigan
Z. Fisk	UC Irvine
F. Steglich	Dresden/Zhejiang
H. Yuan	Zhejiang



Notes:

"Introduction to Many Body Physics", Ch 8,15-16", PC, CUP to be published (2015).



"Heavy Fermions: electrons at the edge of magnetism." Wiley encyclopedia of magnetism. PC. cond-mat/0612006.

"I2CAM-FAPERJ Lectures on Heavy Fermion Physics", (X=I, II, III)
http://physics.rutgers.edu/~coleman/talks/RIO13_X.pdf

"Julich lectures: Heavy Fermion Physics: A 21st Century Perspective"
arXiv:1509.05769

General reading:

A. Hewson, "Kondo effect to heavy fermions", CUP, (1993).

"The Theory of Quantum Liquids", Nozieres and Pines (Perseus 1999).

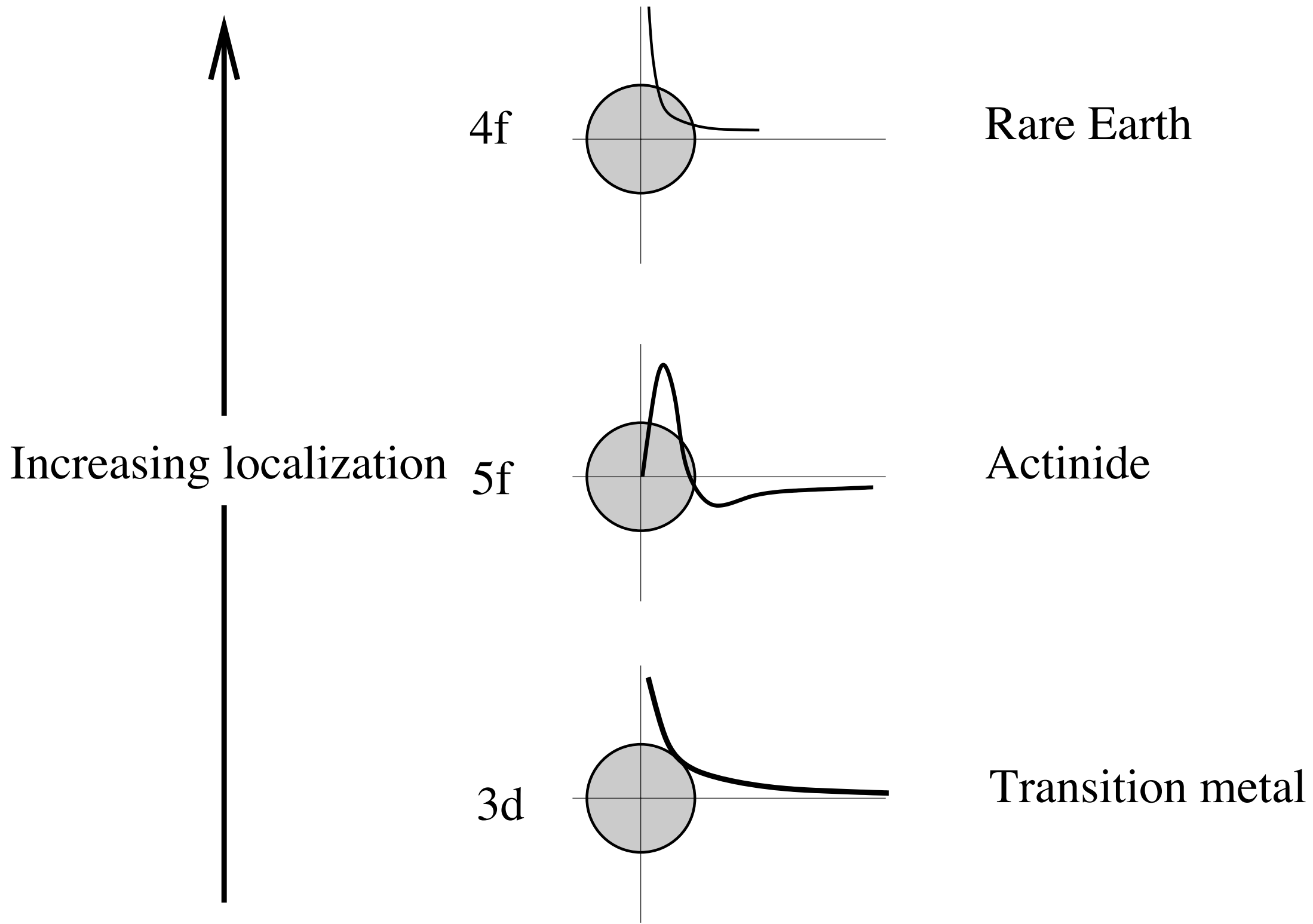
Outline of the Topics

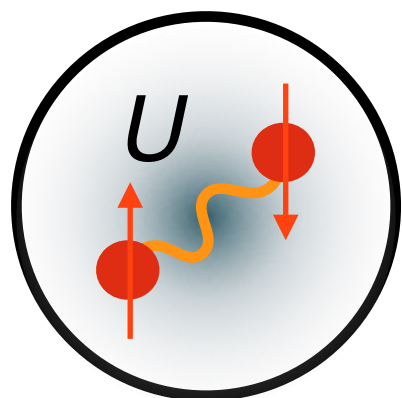
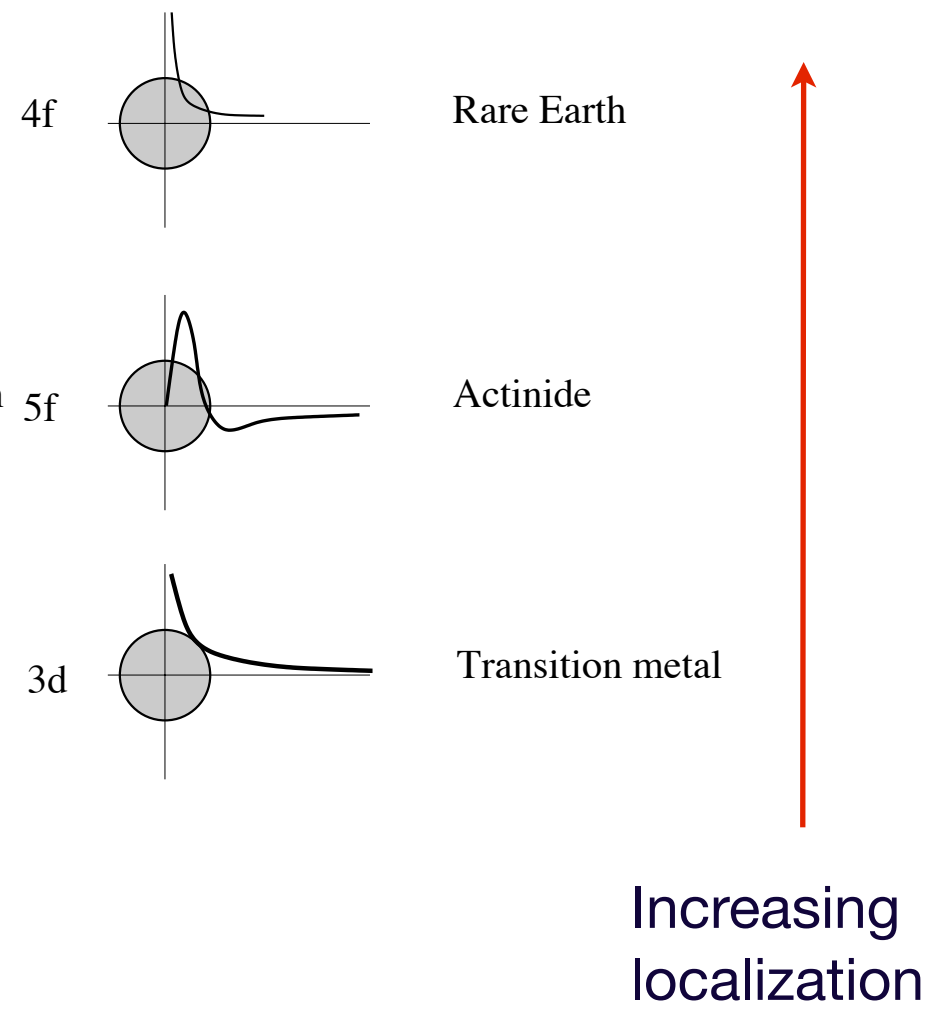
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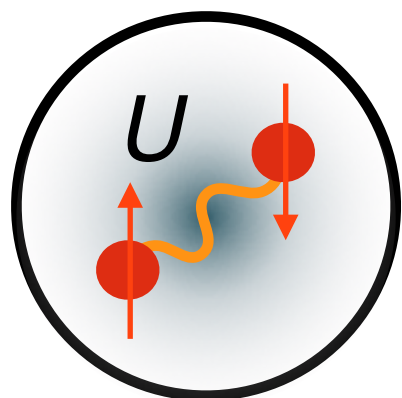
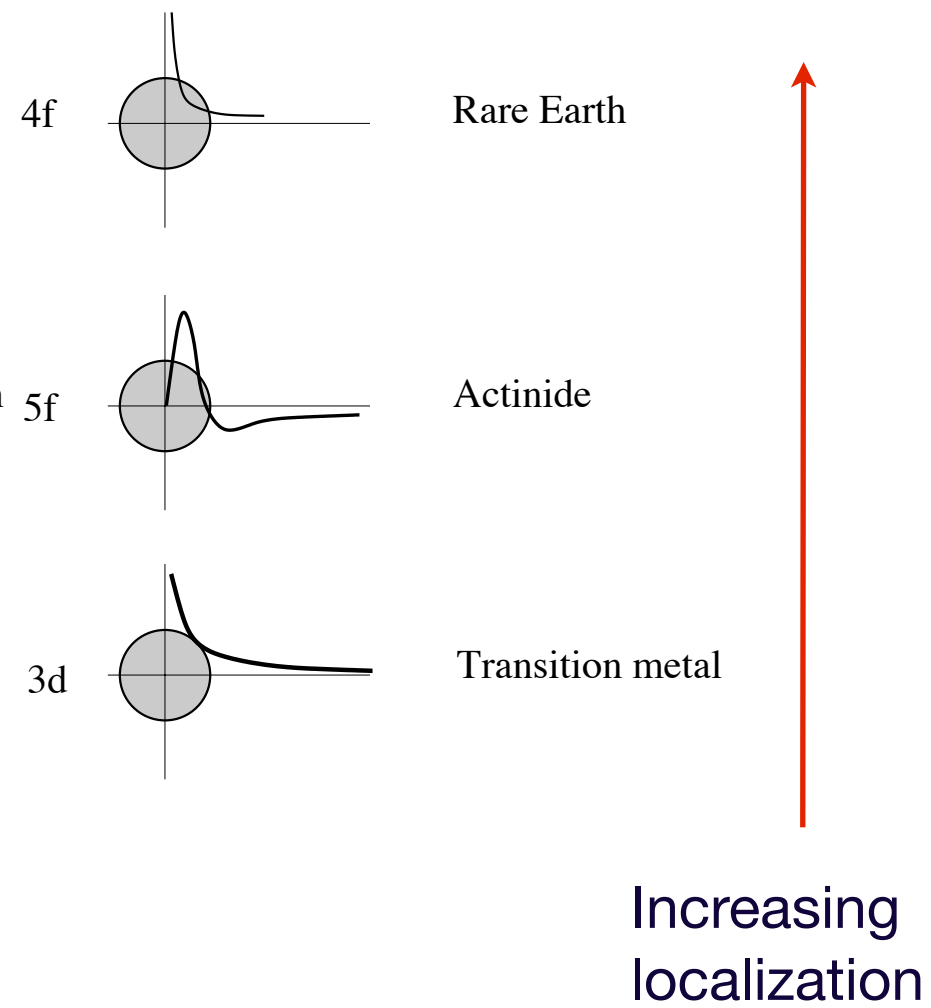
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Mott Mechanism.

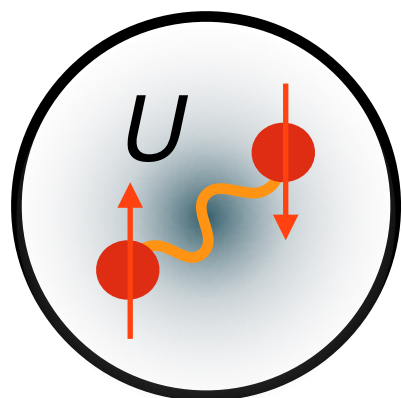
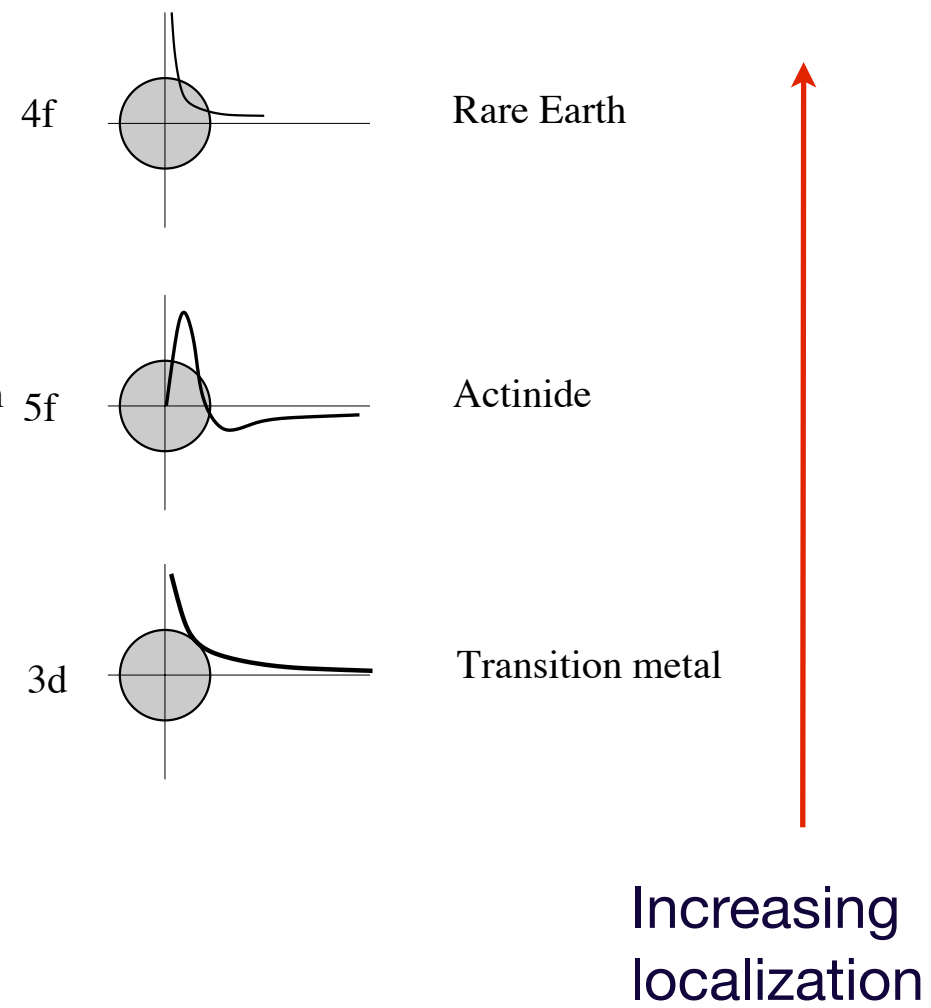
Anderson U (Anderson 1959)



- No double occupancy: strongly correlated

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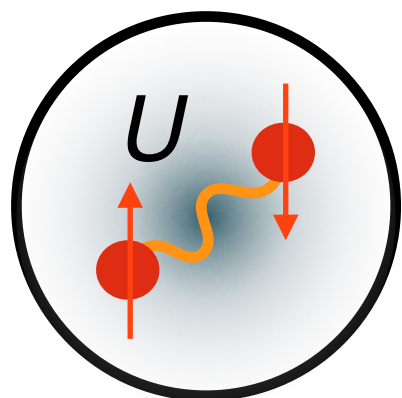
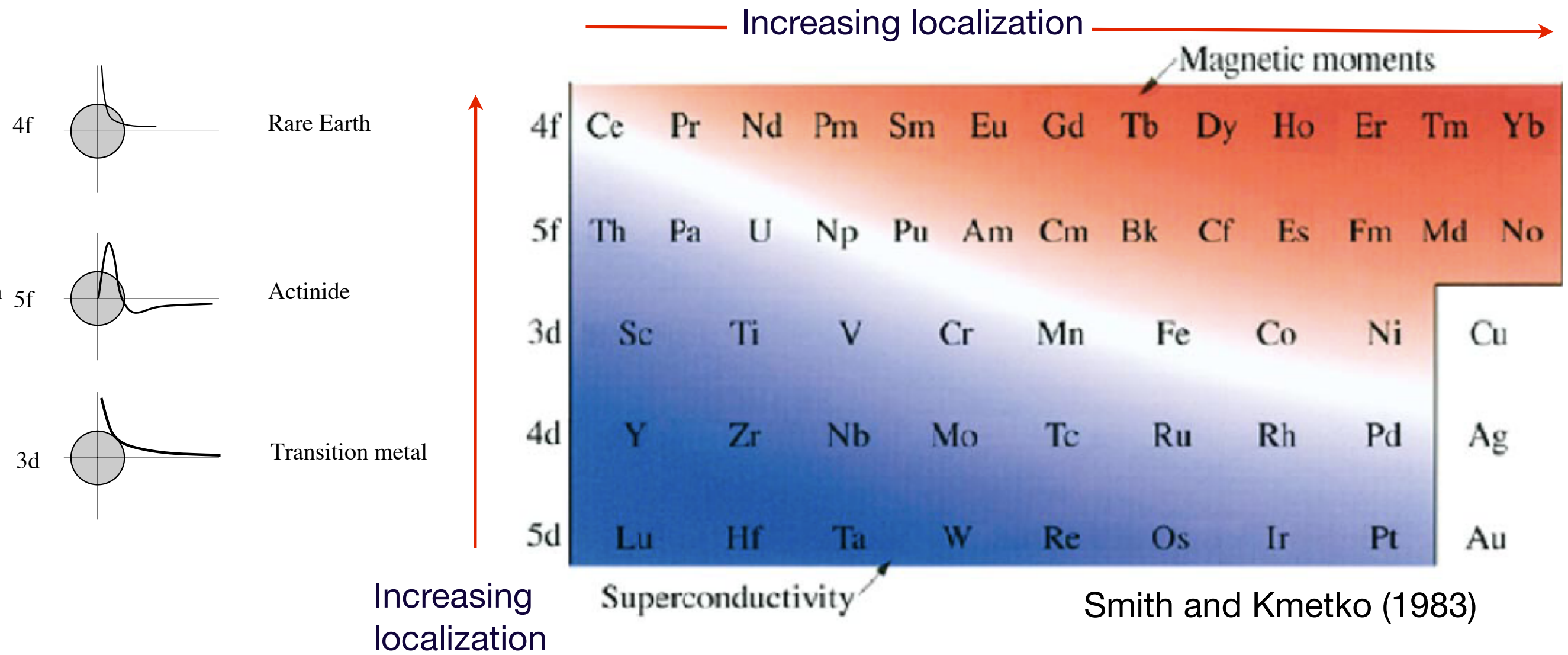
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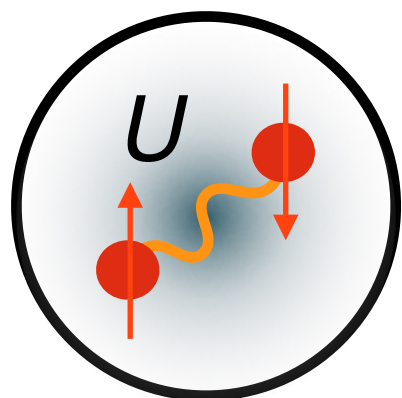
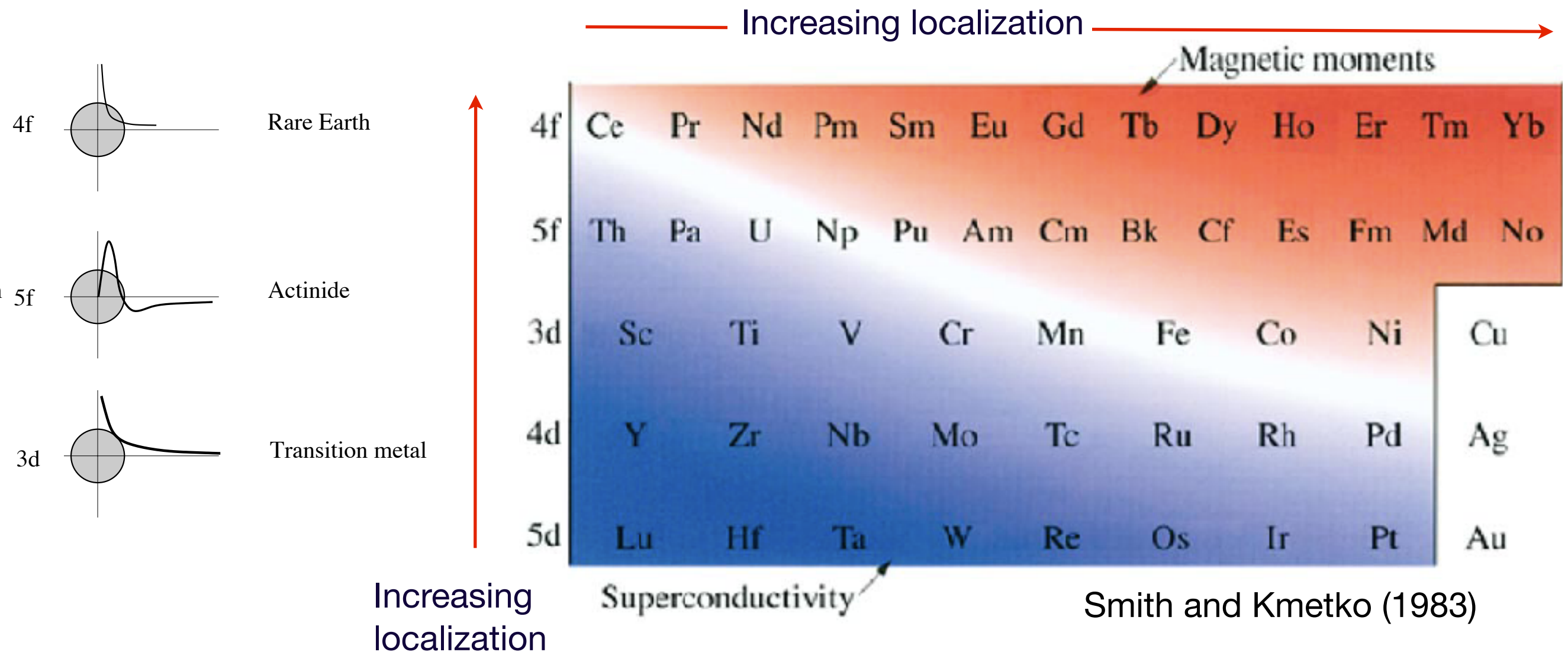
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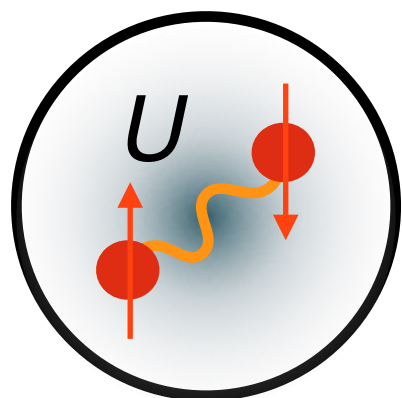
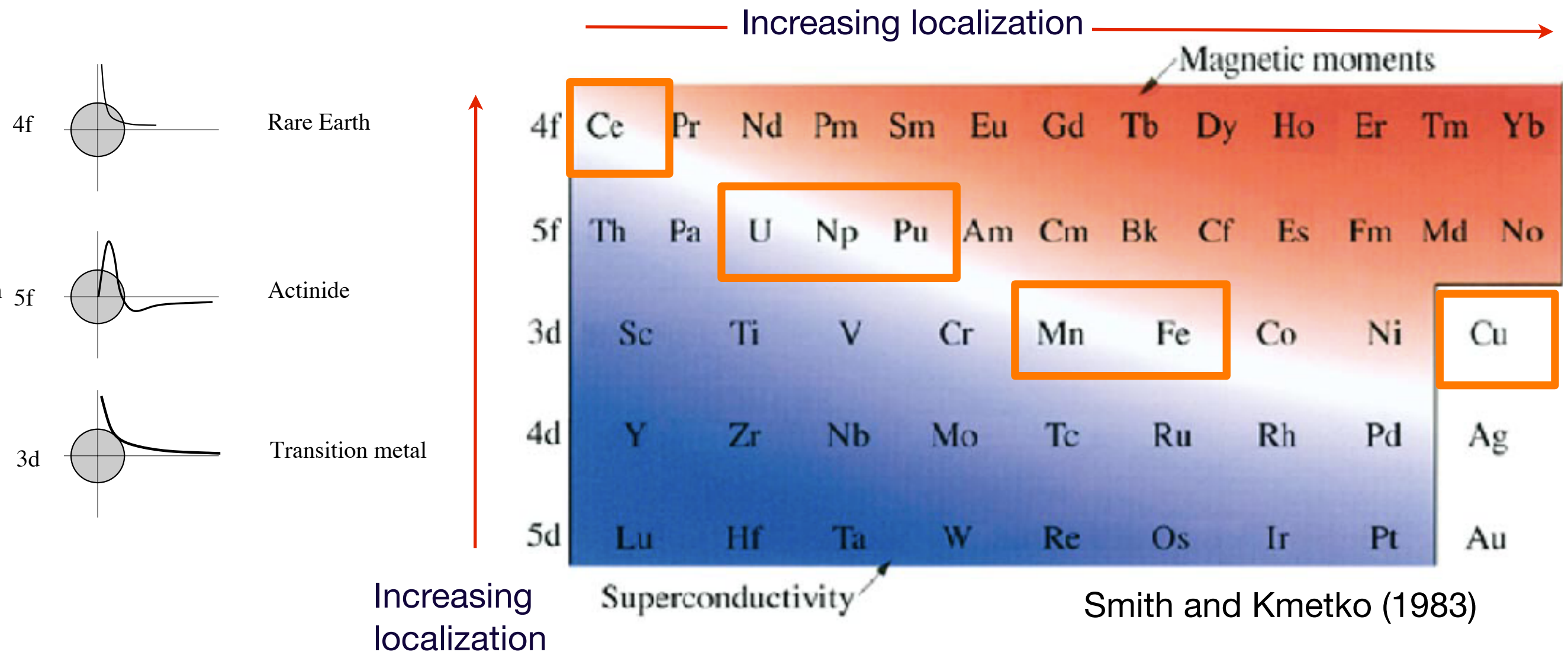
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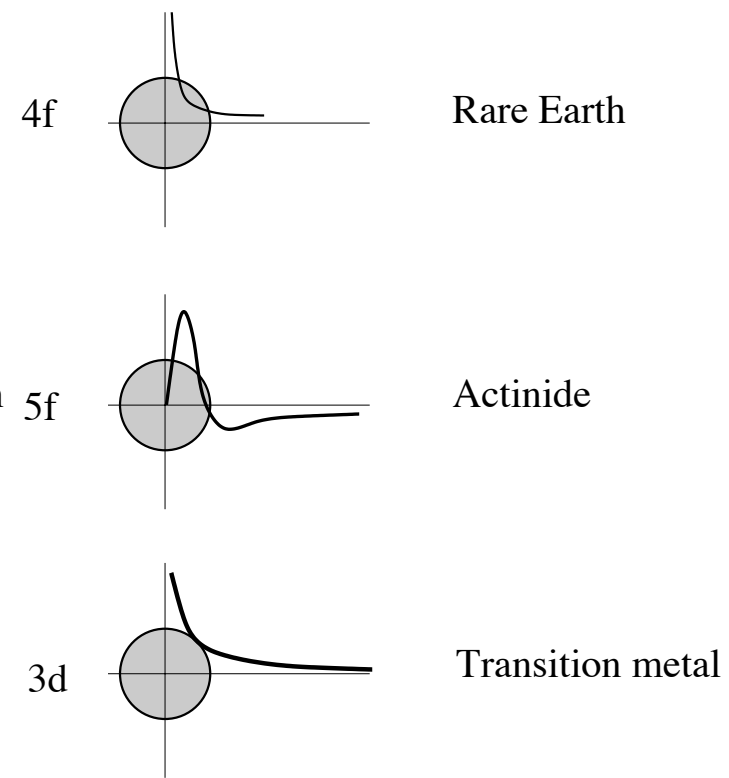
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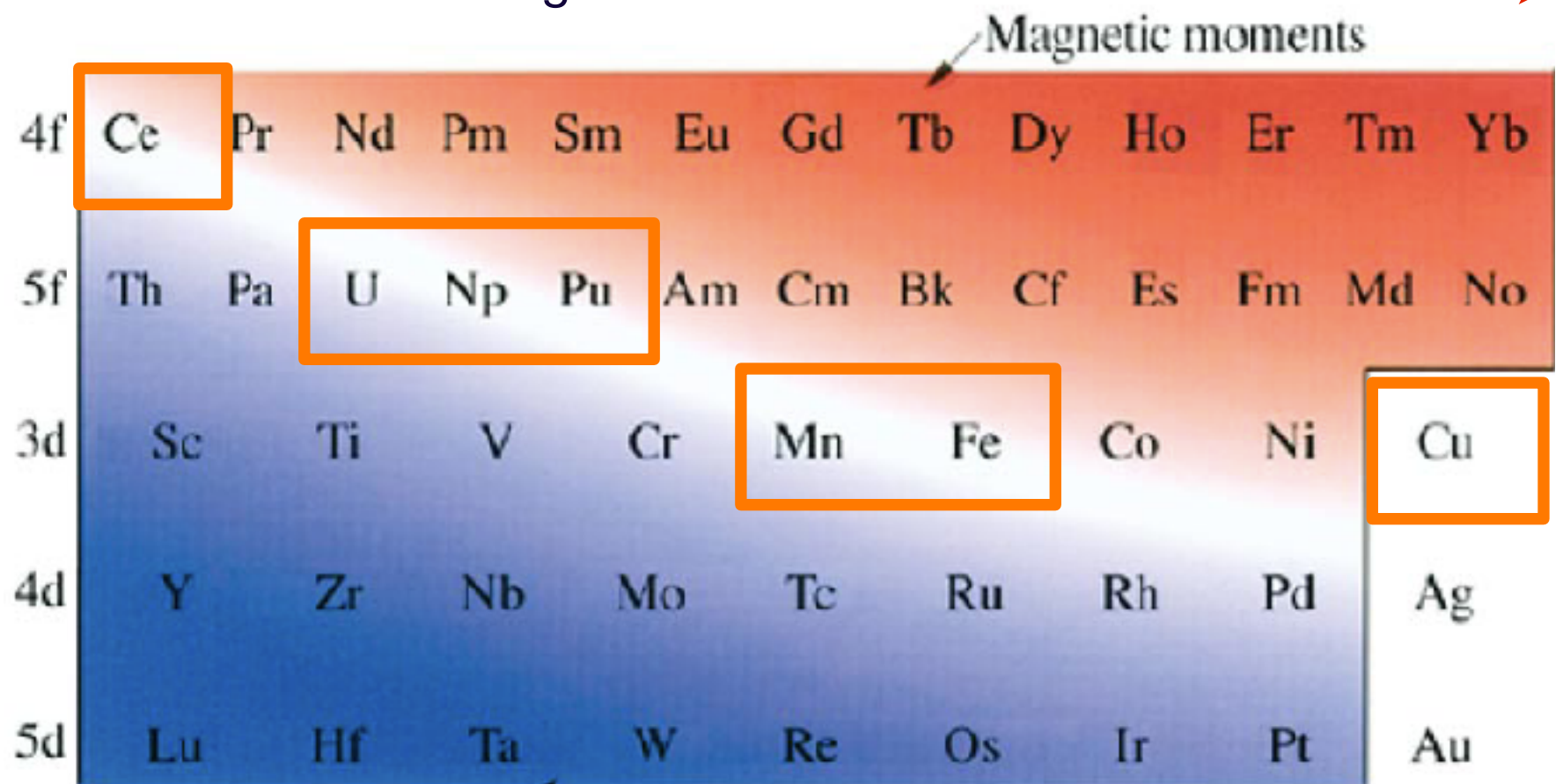
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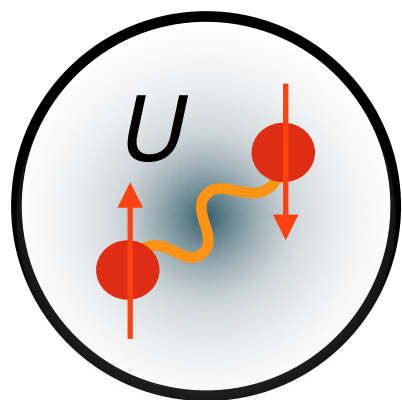


Increasing localization

Increasing localization



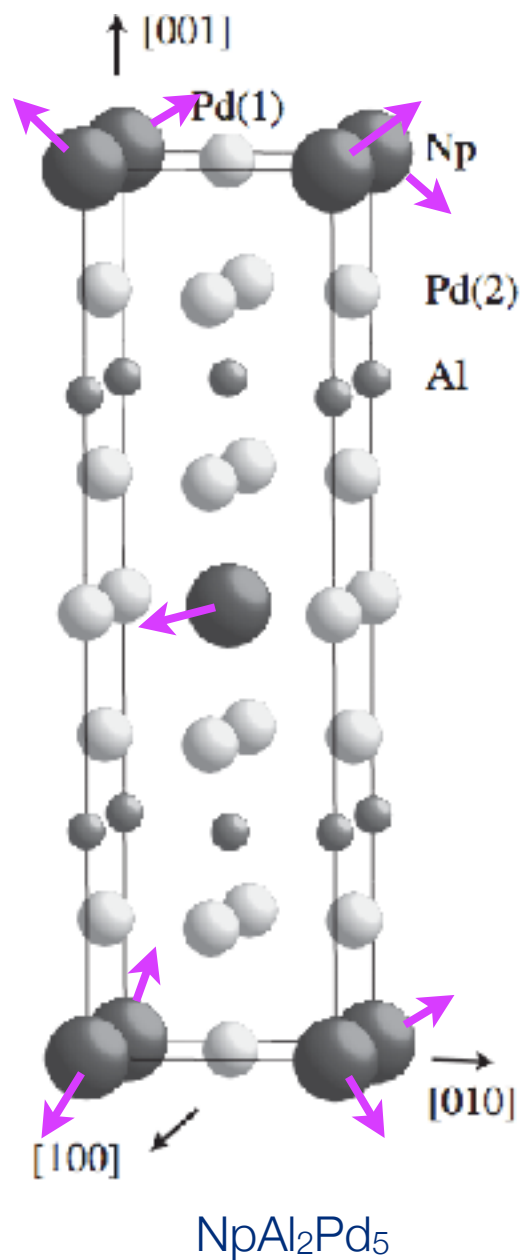
Smith and Kmetko (1983)



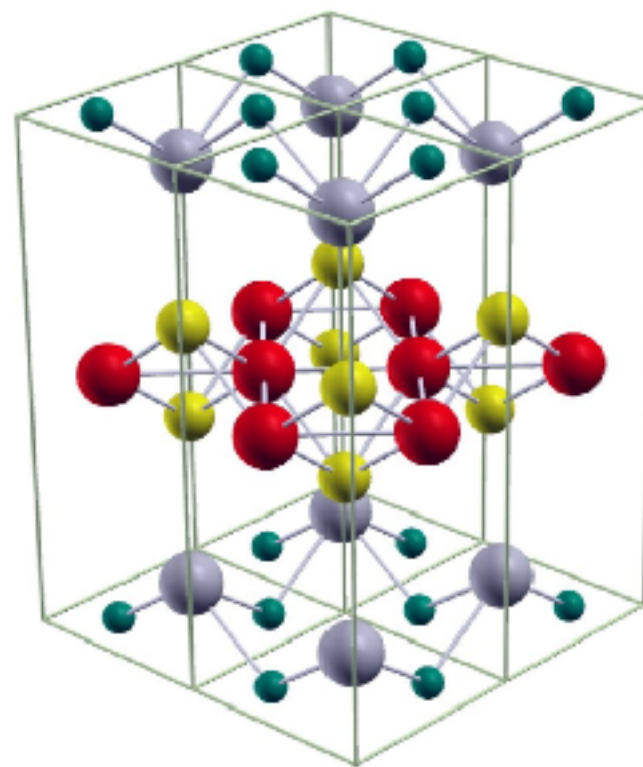
Many things are possible at the brink of magnetism.

Mott Mechanism.

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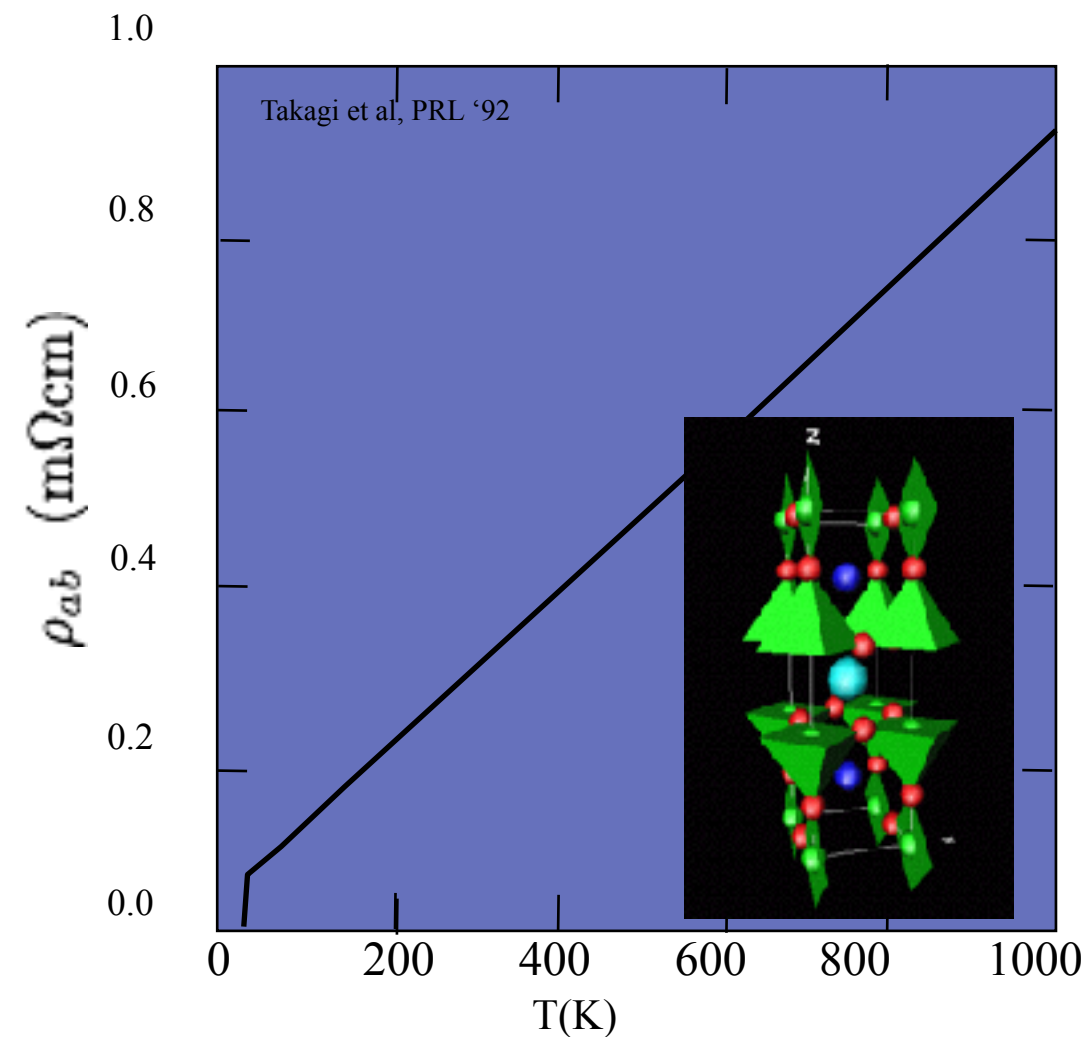


HF 115s
 $T_c = 0.2 - 18.5$ K



Z.A. Ren et.al, Beijing, (08)

Iron based sc
 $T_c = 6 - 53 ++ ?$ K

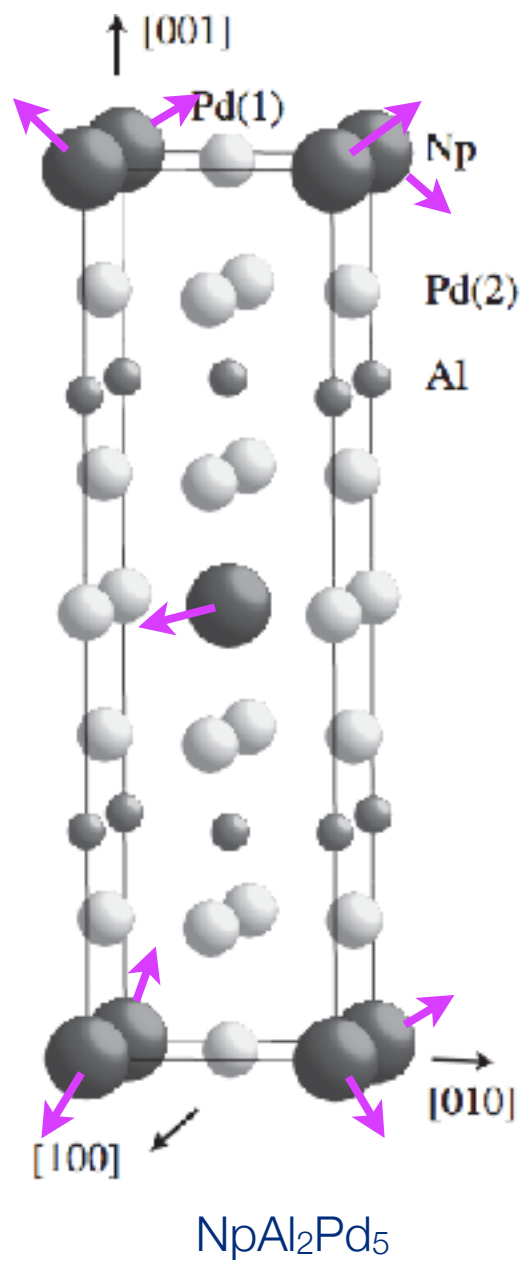


Cuprates $T_c = 11 - 92$ K

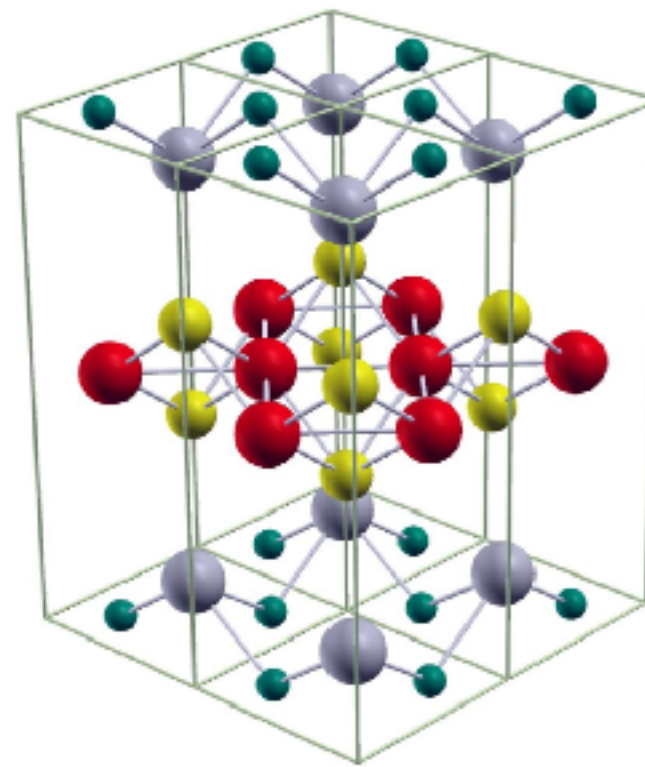
Diversity of new ground-states on the brink of localization.

f-electron systems: 4f Ce, Yb systems 5f U, Np, Pu systems.

d-electron systems: e.g. Pnictides, Cuprate SC.

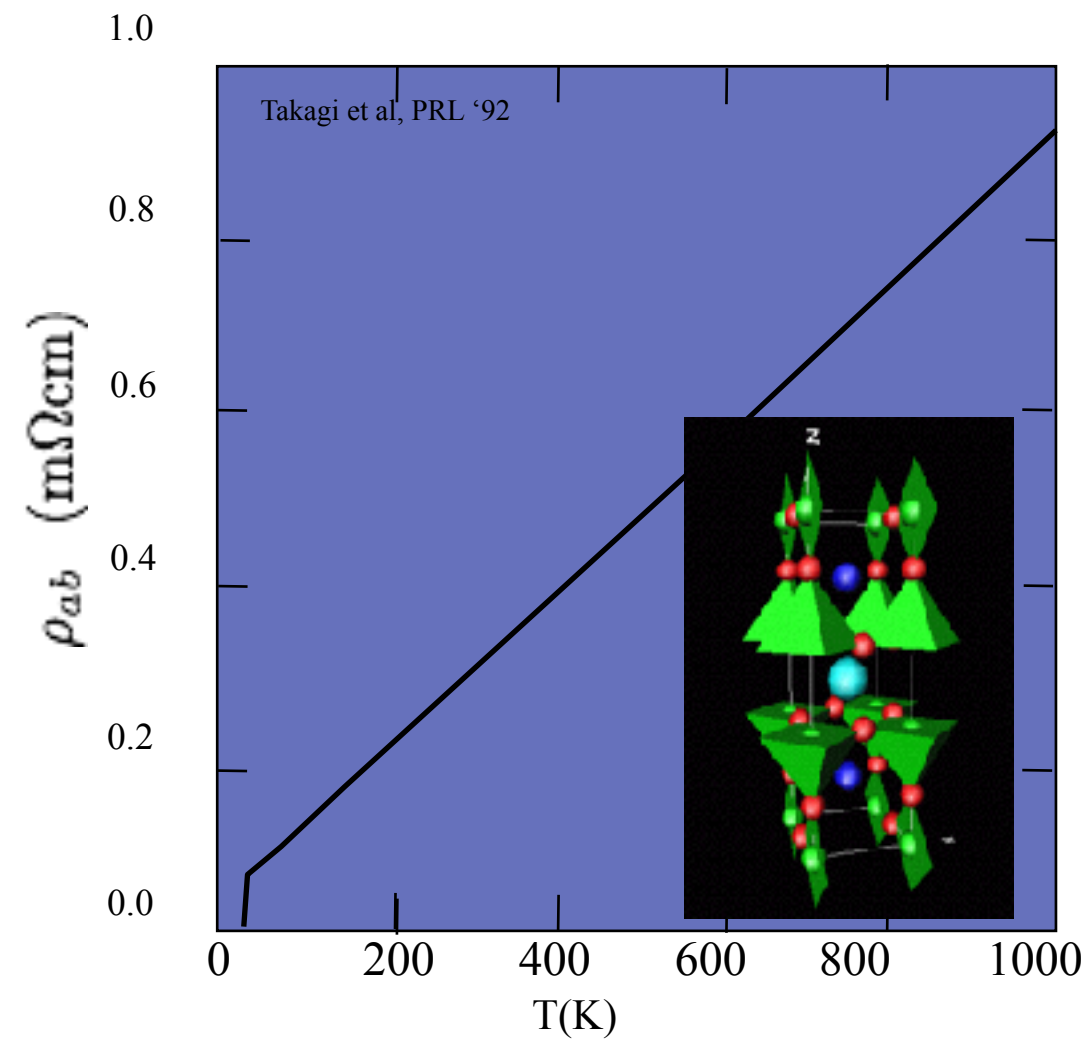


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Z.A. Ren et.al, Beijing, (08)

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Cuprates $T_c = 11 - 92$ K

A new era of mysteries

Outline of the Topics

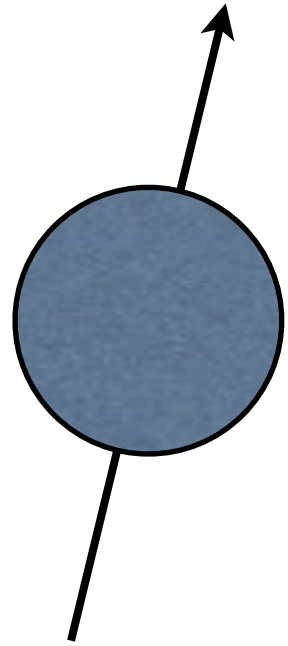
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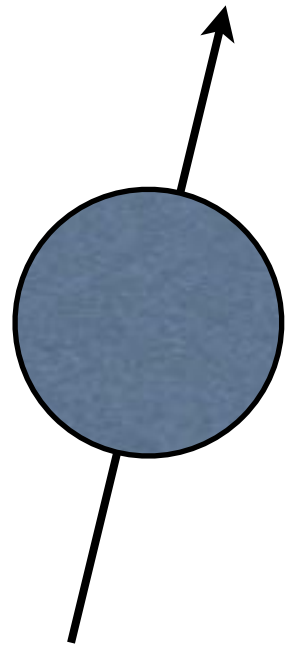
Please ask questions!

Heavy Fermions + Kondo

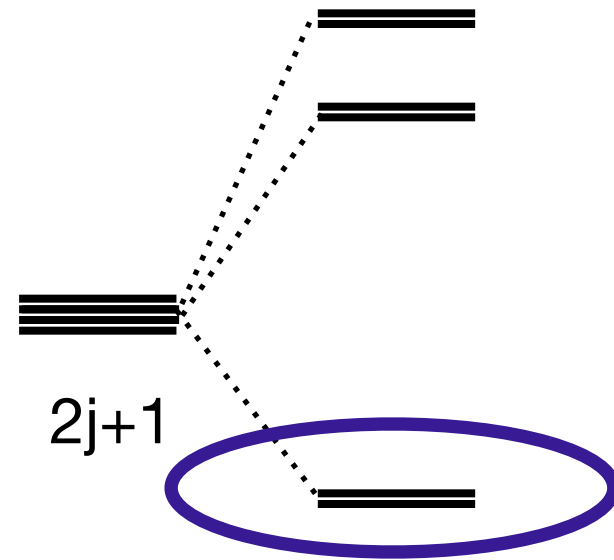


Spin (4f,5f):
“quark” of heavy
electron physics.

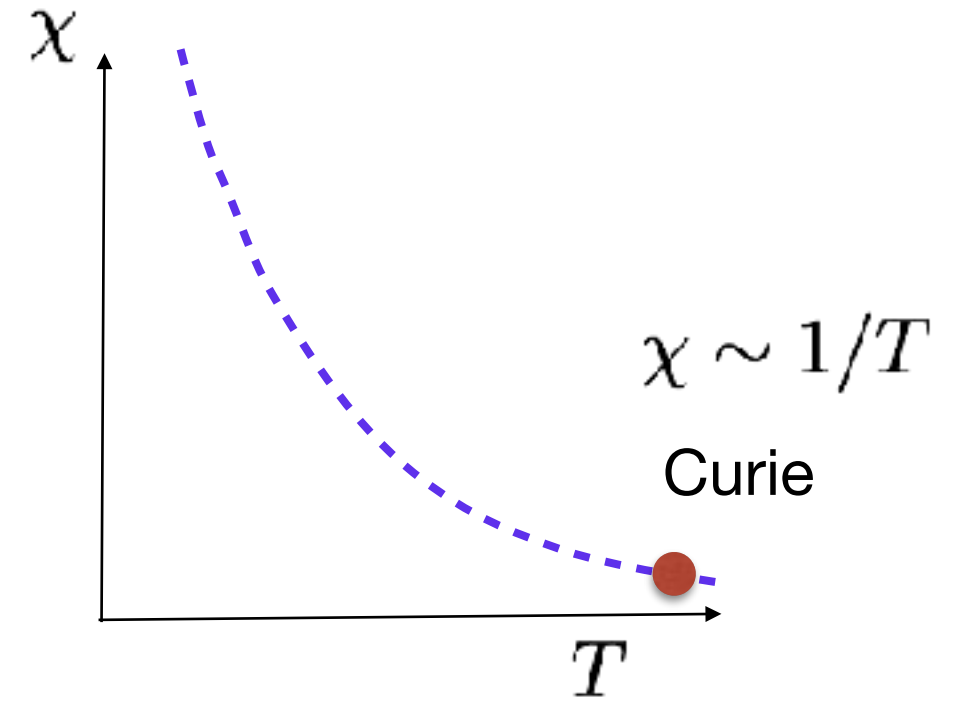
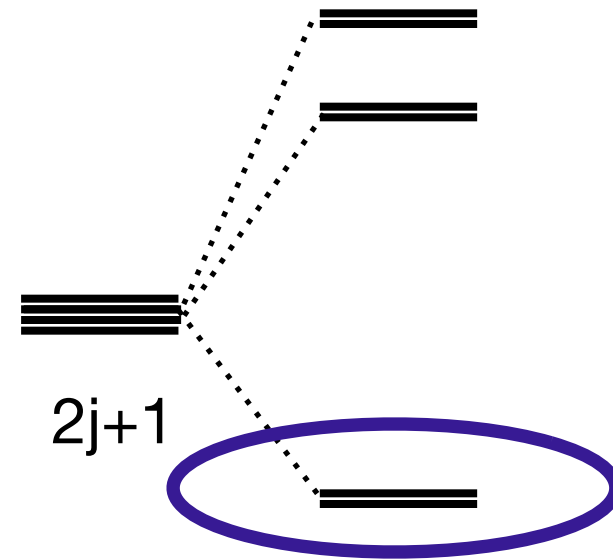
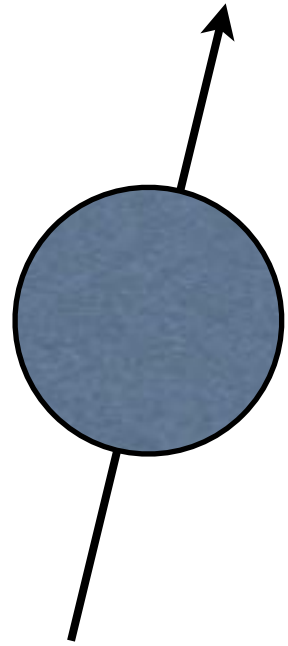
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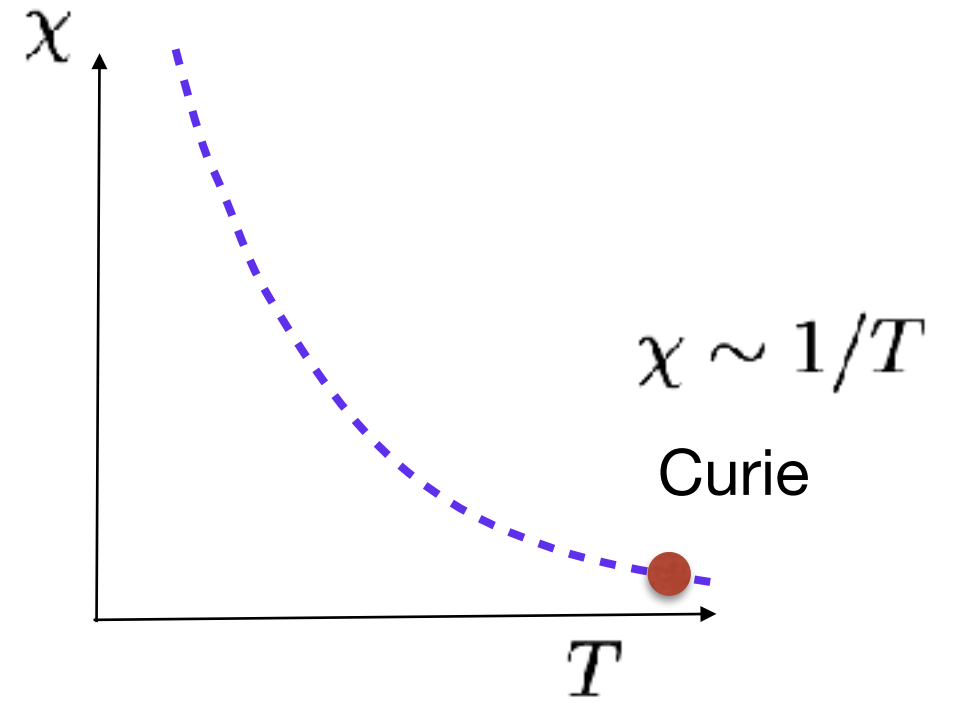
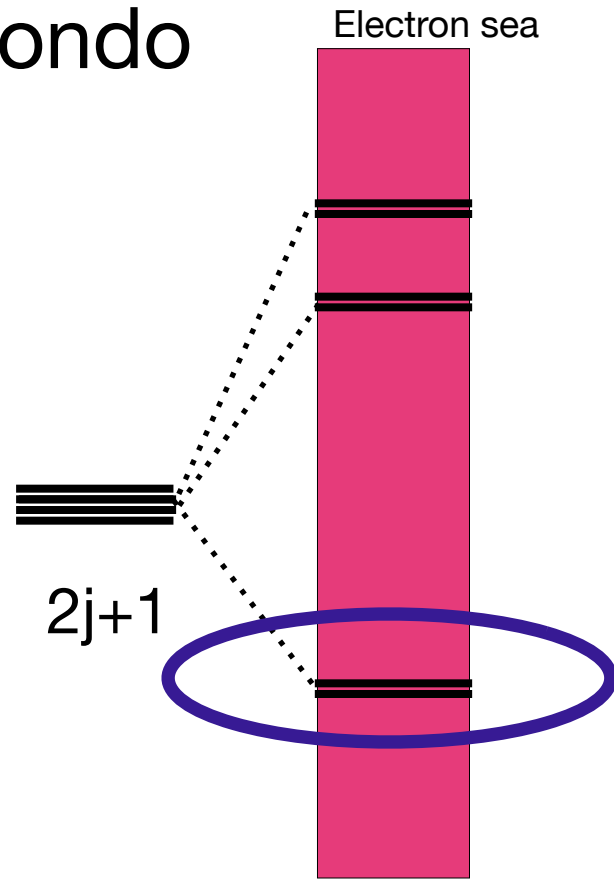
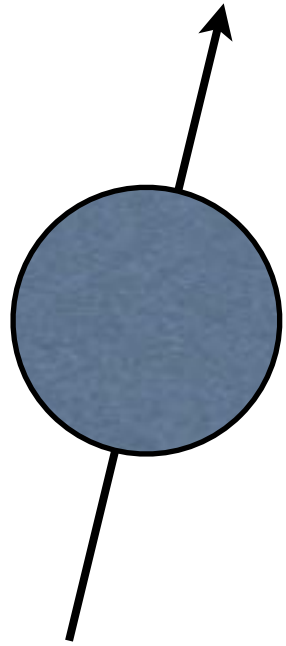


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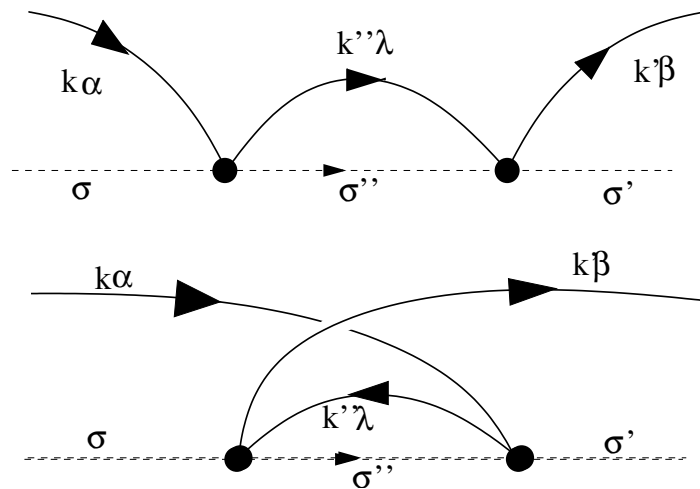


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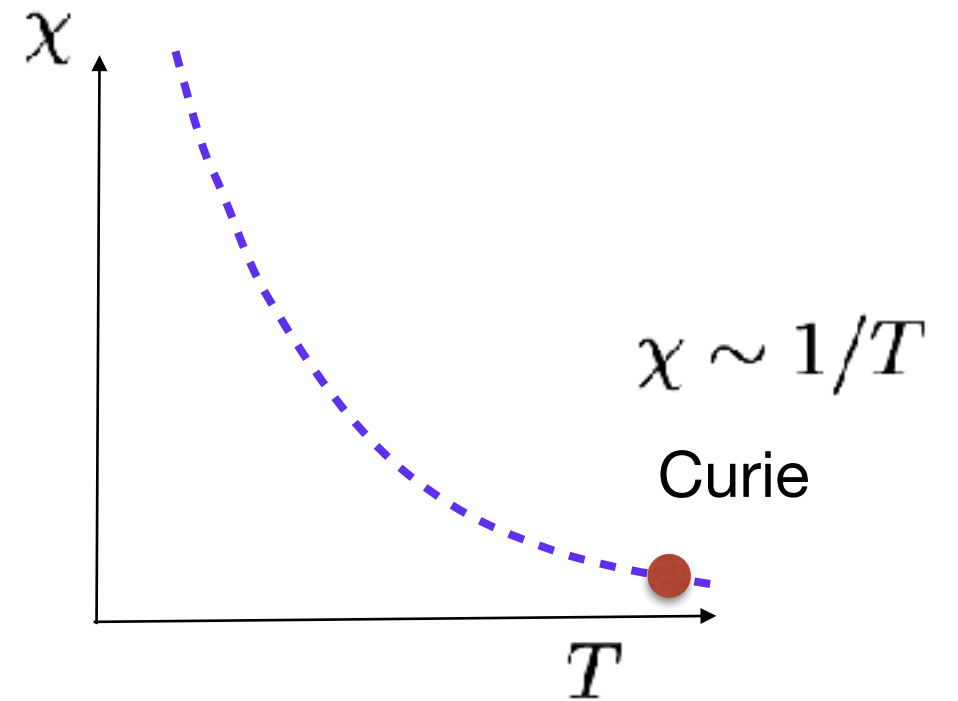
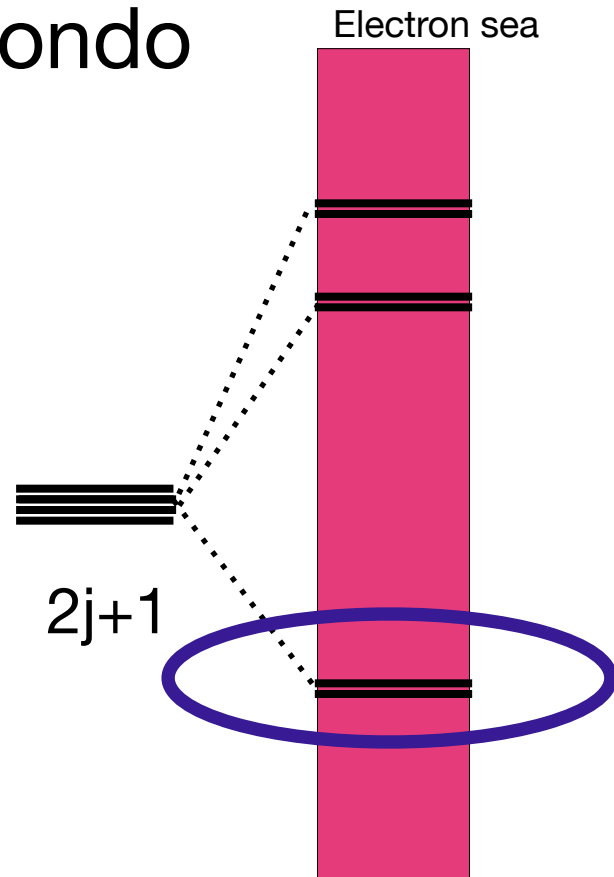
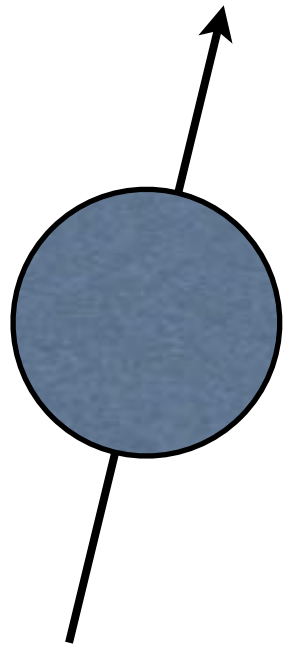
$$\approx -\frac{J^2 \rho \delta D}{D} (\sigma^b \sigma^a)_{\beta\alpha} (S^a S^b)_{\sigma'\sigma}$$

$$\approx \frac{J^2 \rho \delta D}{D} (\sigma^a \sigma^b)_{\beta\alpha} (S^a S^b)_{\sigma'\sigma}$$

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + J \vec{S} \cdot \vec{\sigma}(0)$$

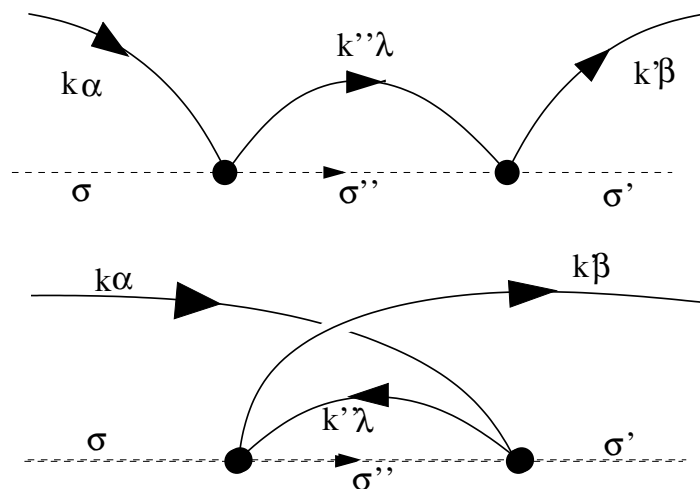
J. Kondo, 1962

Heavy Fermions + Kondo



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$$J \rightarrow J(T) = J + 2J^2\rho \ln \frac{D}{T}$$



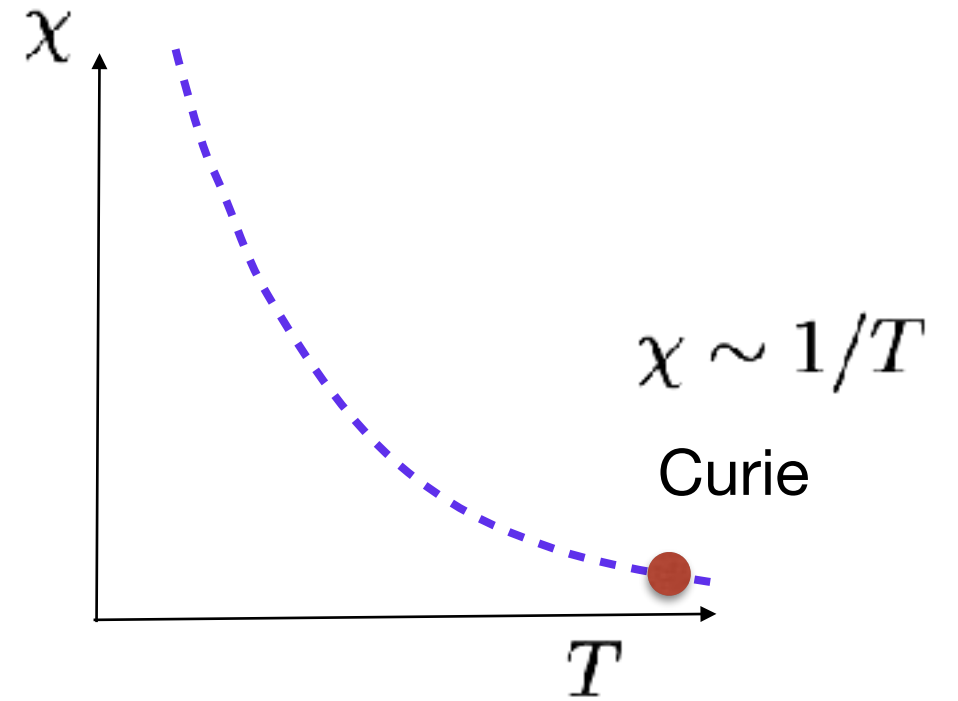
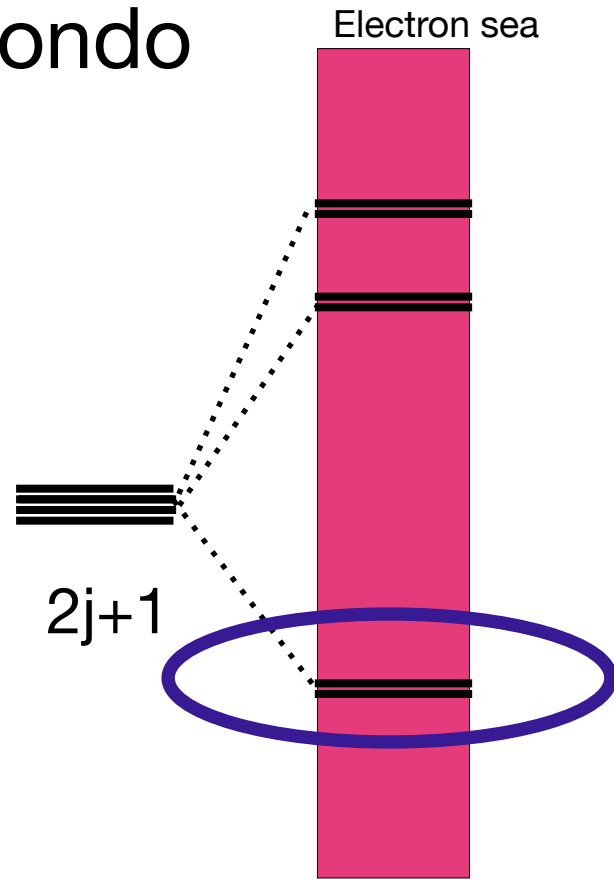
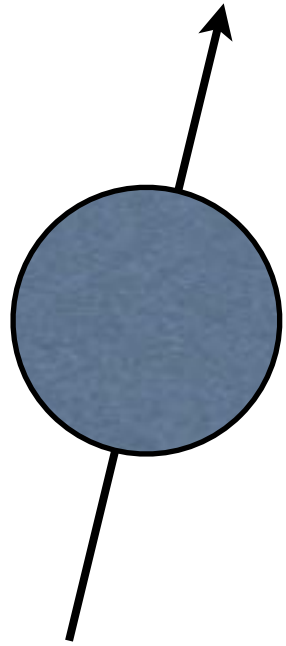
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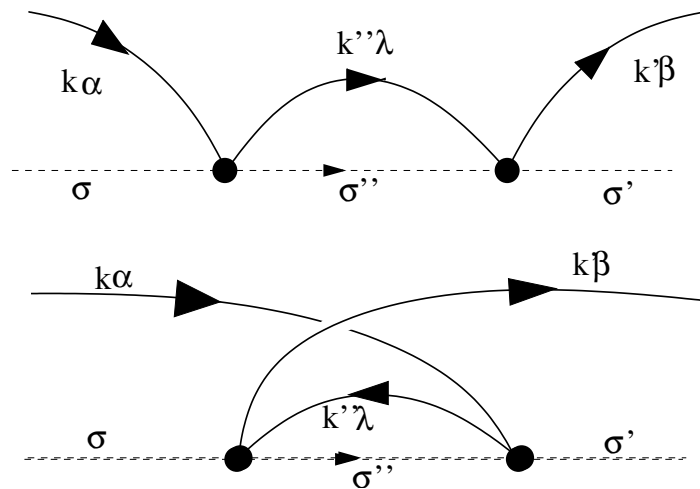
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“Scales to
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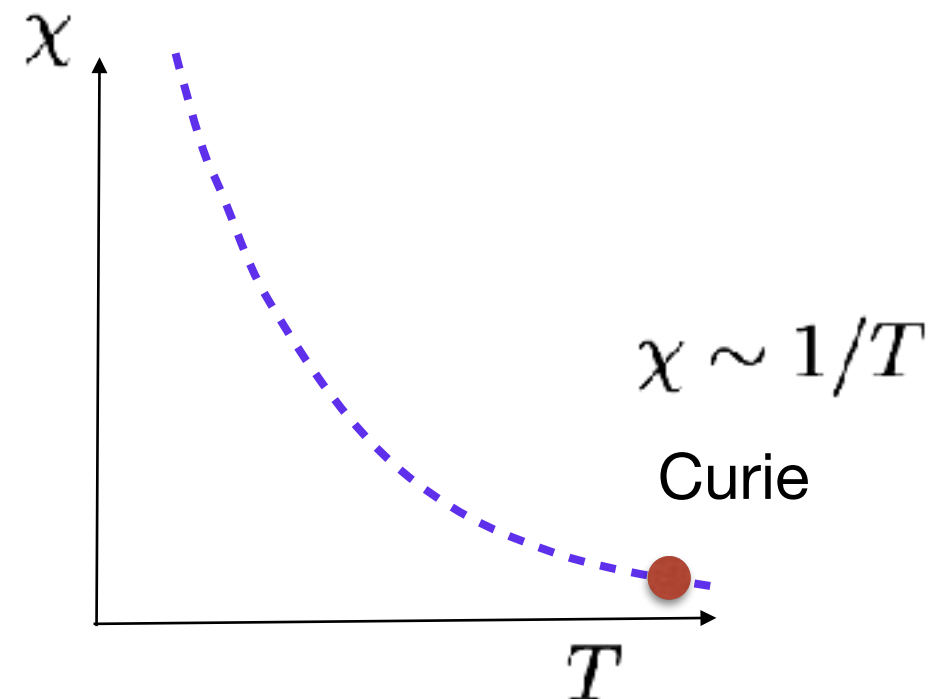
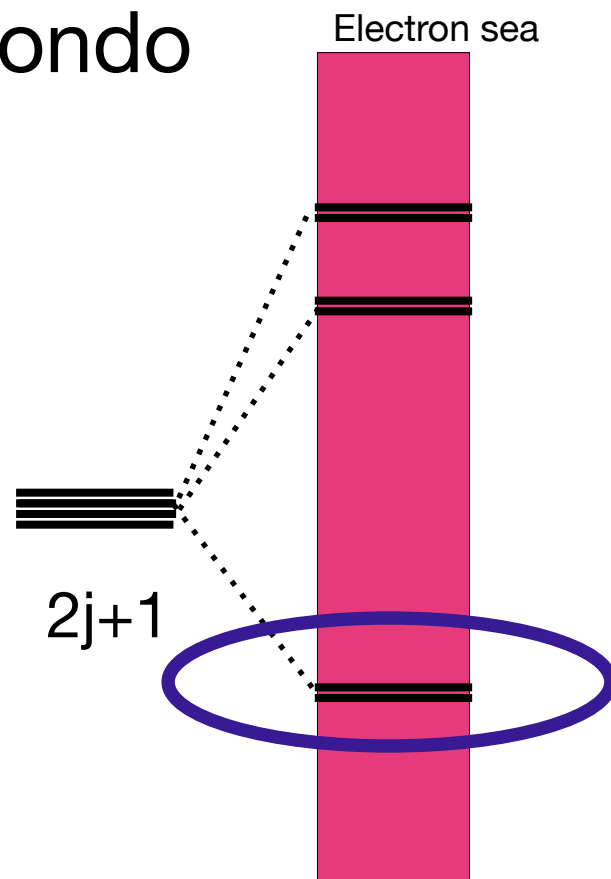
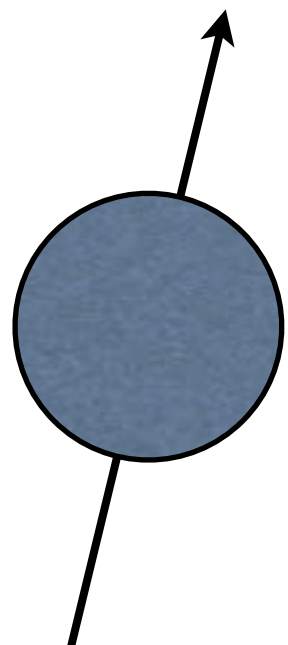
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Heavy Fermions + Kondo



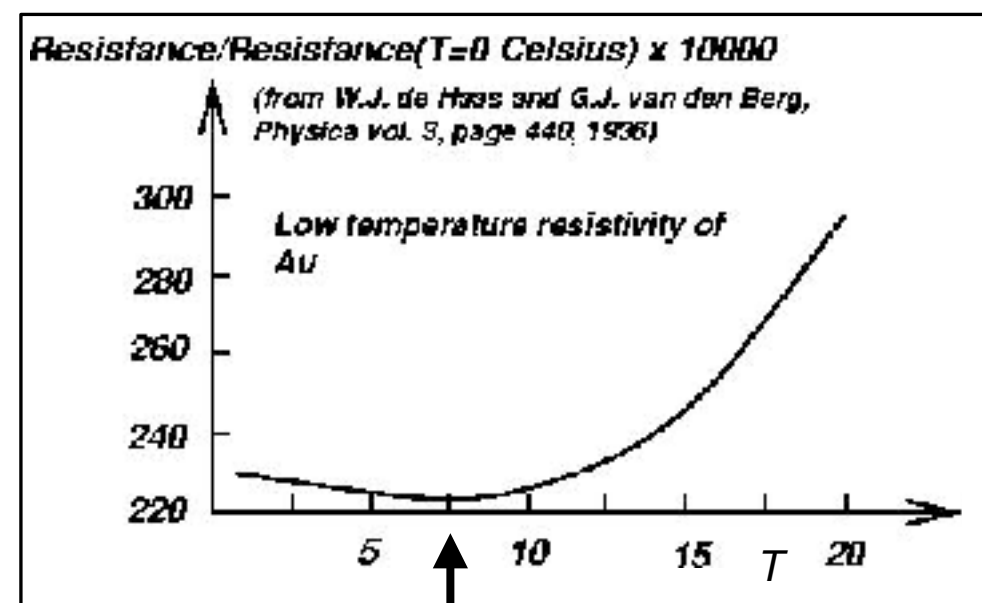
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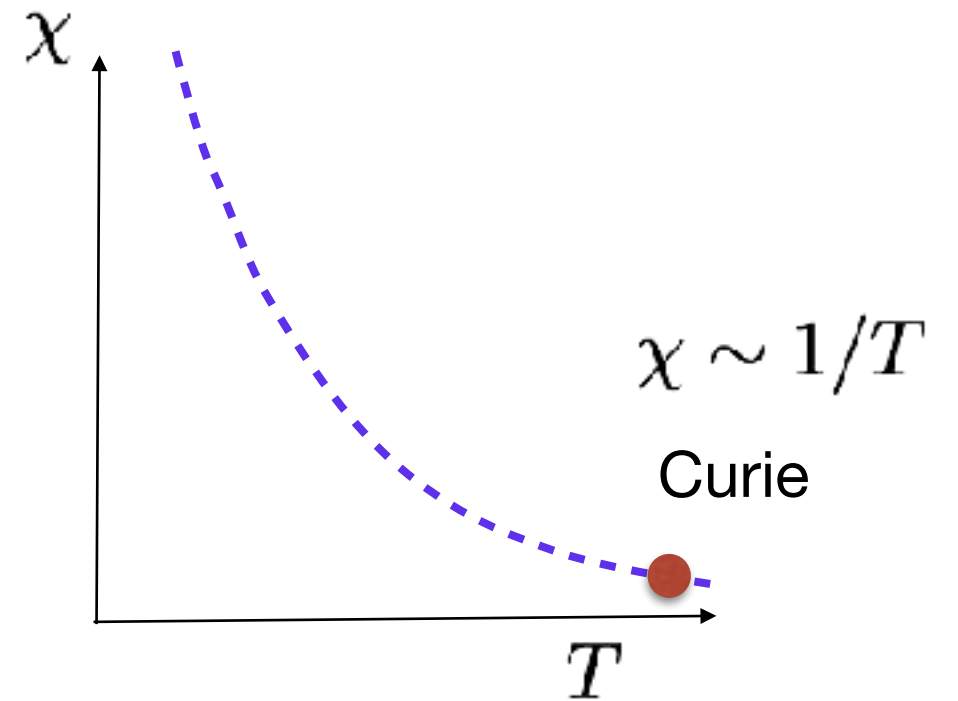
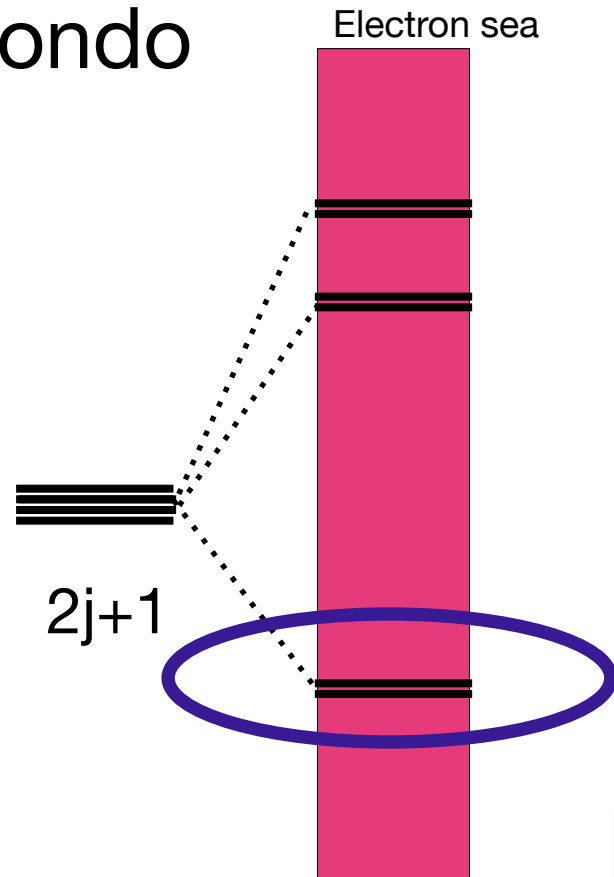
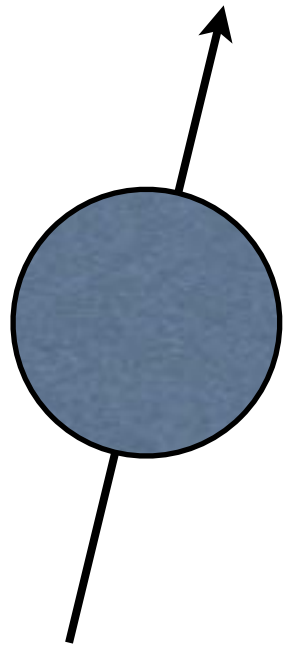
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J. Kondo, 1962



“Kondo Resistance Minimum”

Heavy Fermions + Kondo



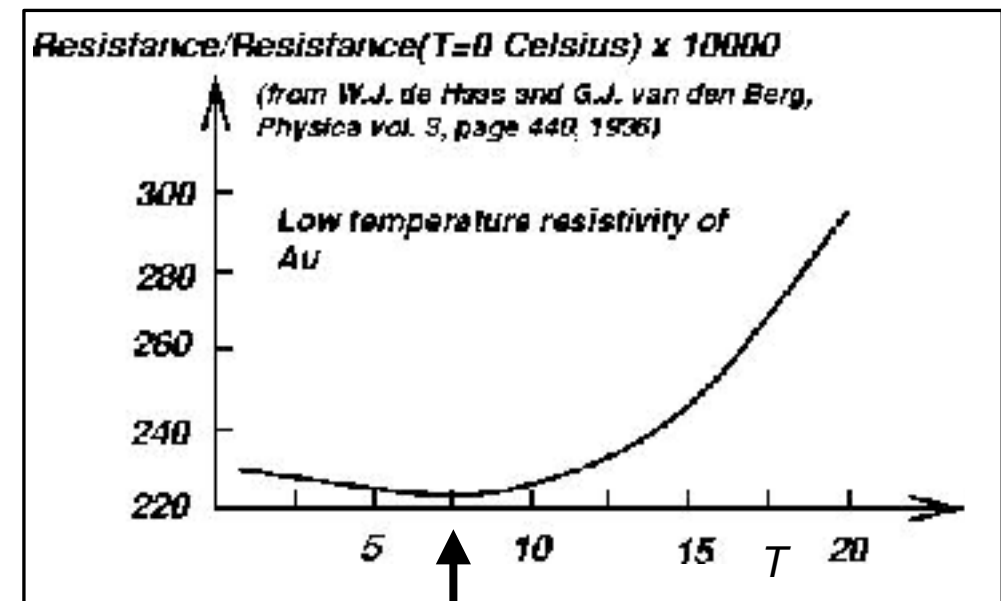
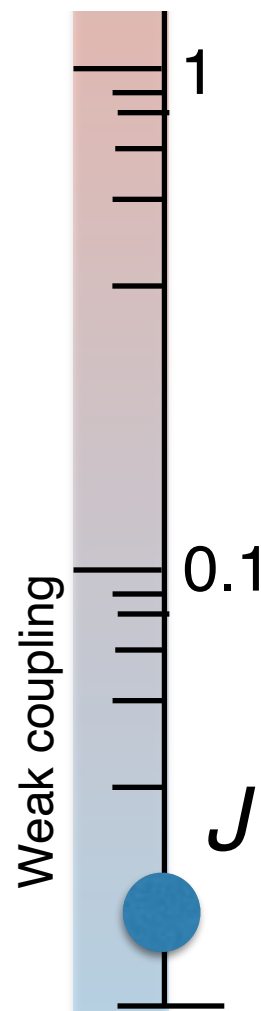
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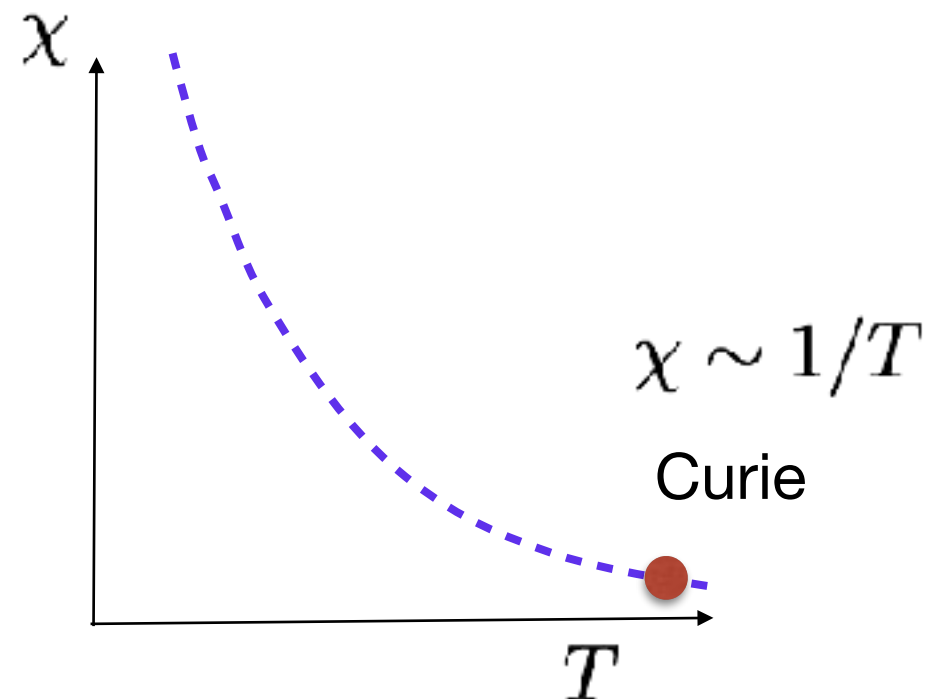
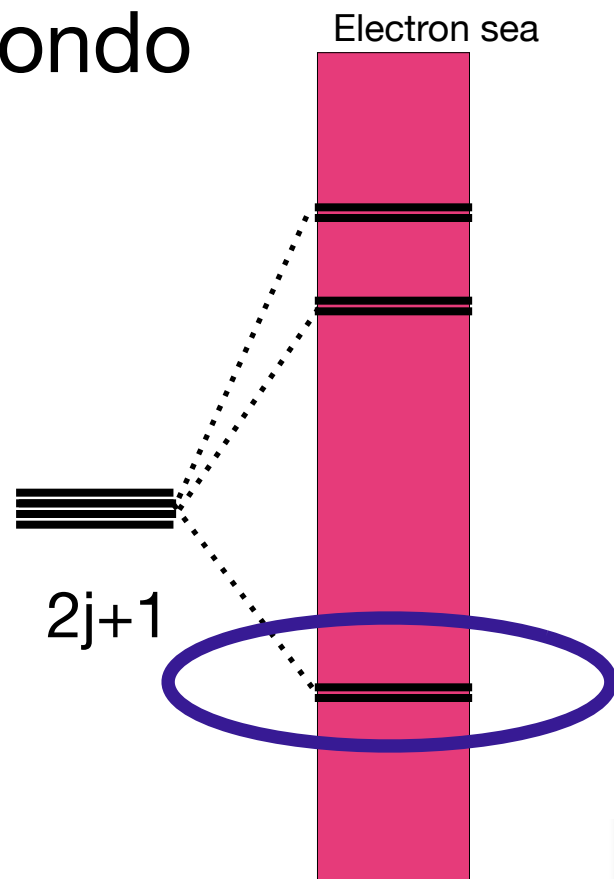
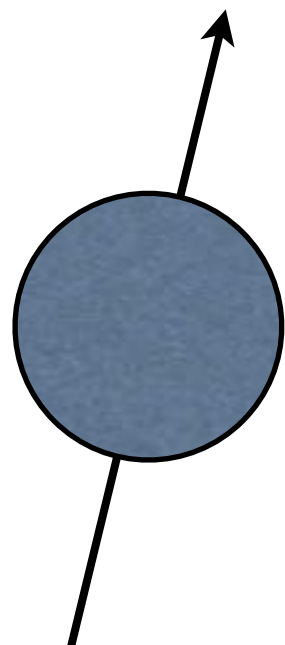
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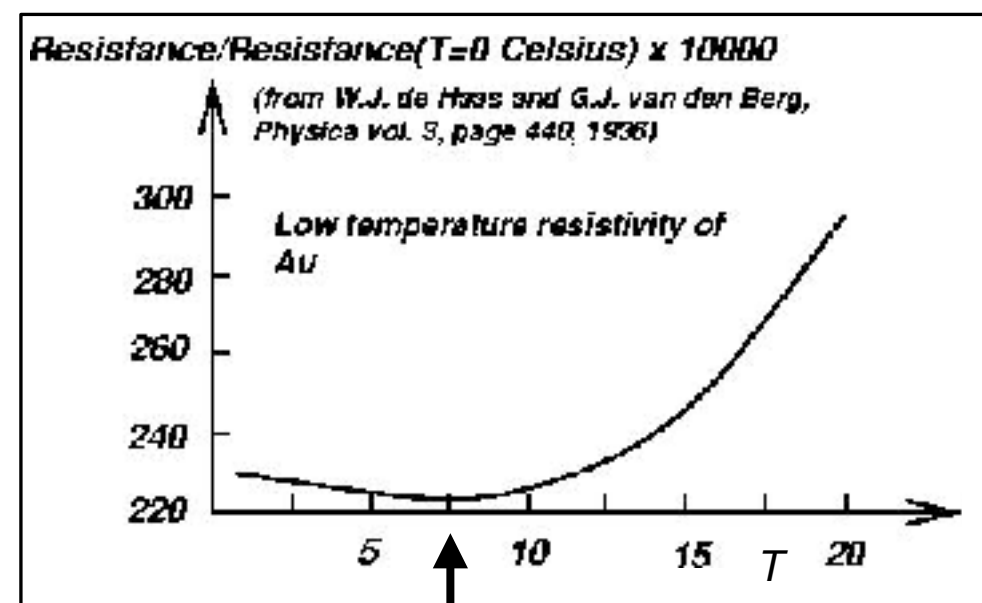
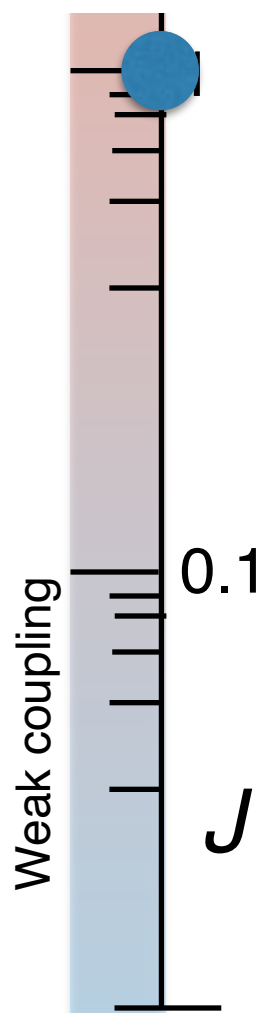
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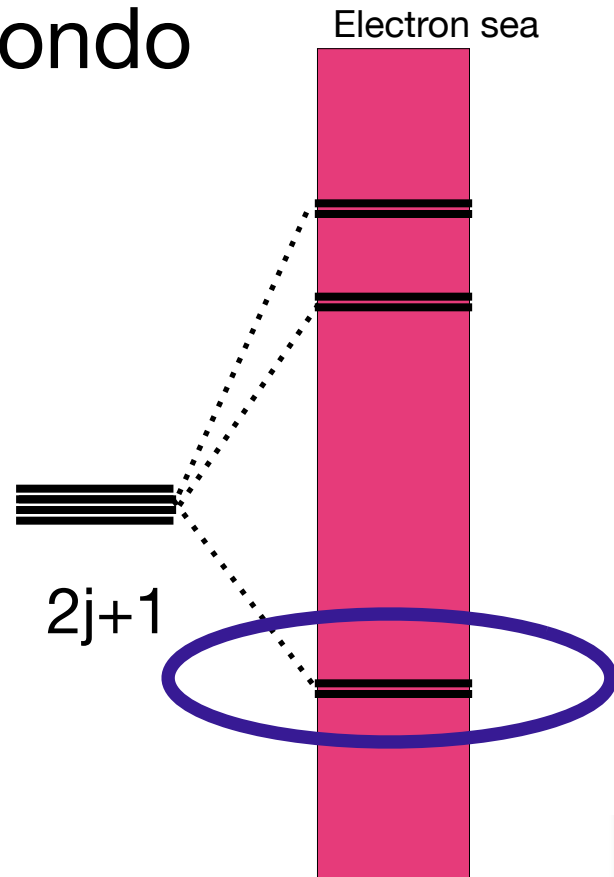
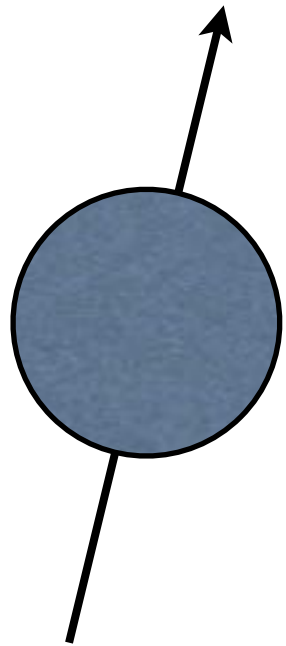
$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + \underbrace{J}_{\text{Kondo}} \vec{S} \cdot \vec{\sigma}(0)$$

J. Kondo, 1962



“Kondo Resistance Minimum”

Heavy Fermions + Kondo



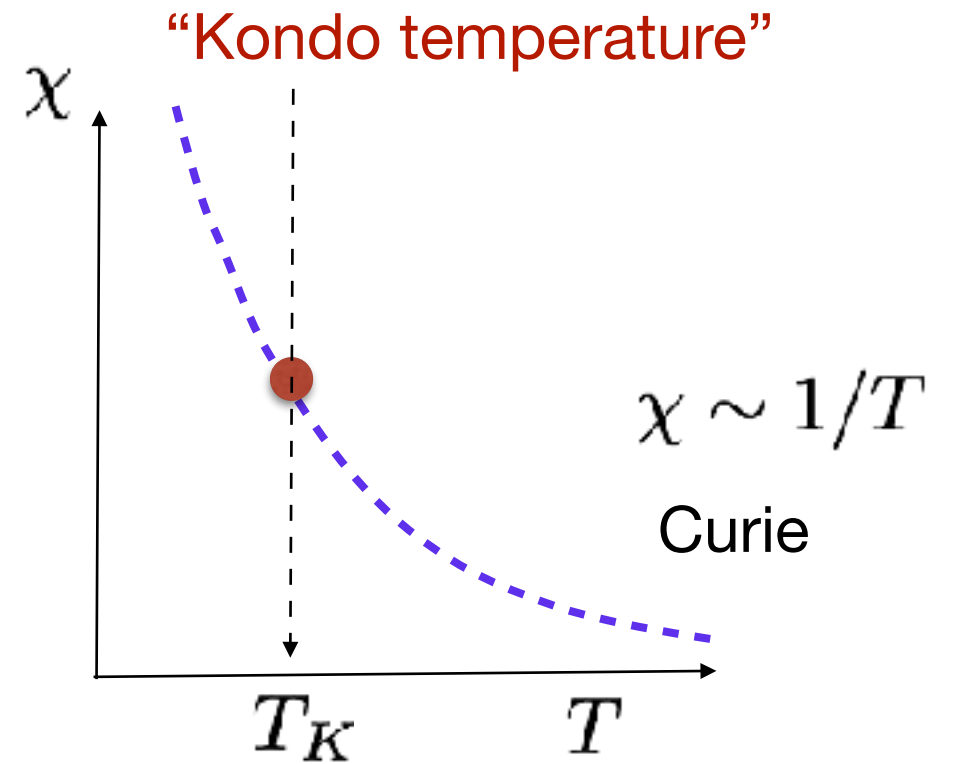
Spin (4f,5f):
“quark” of heavy
electron physics.

$$J \rightarrow J(T) = J + 2J^2\rho \ln \frac{D}{T}$$

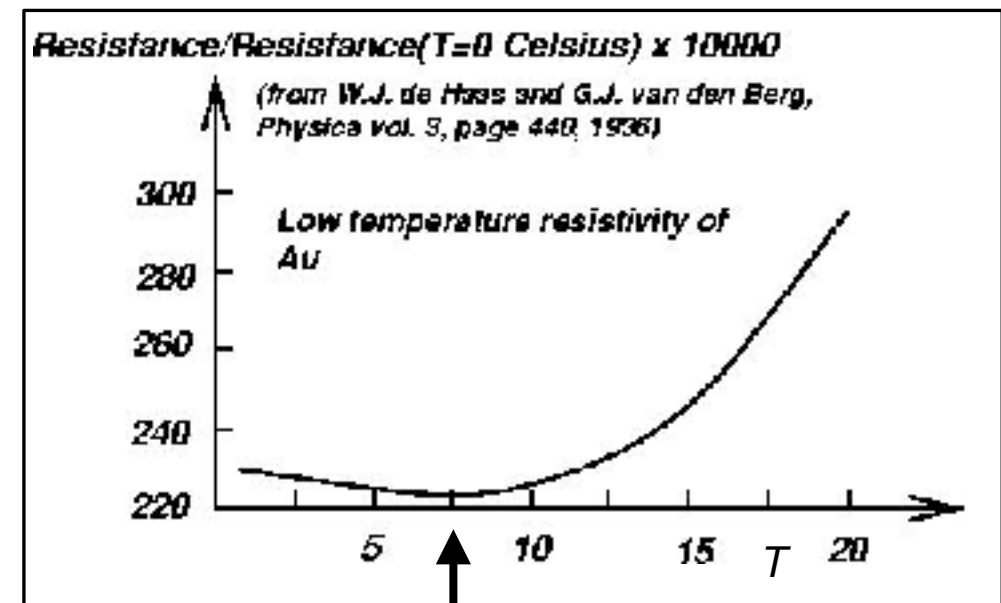
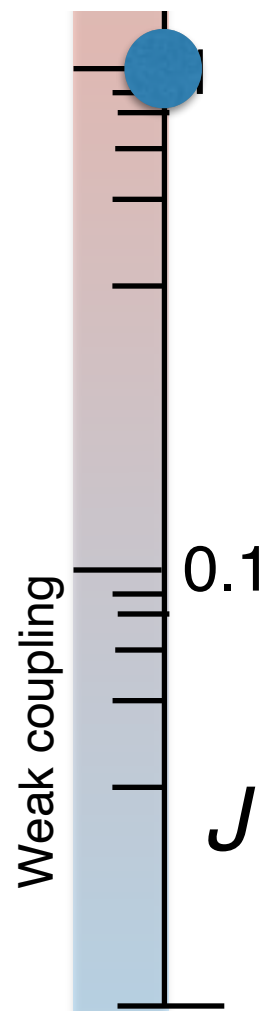
“Scales to
Strong Coupling”

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + J \vec{S} \cdot \vec{\sigma}(0)$$

J. Kondo, 1962

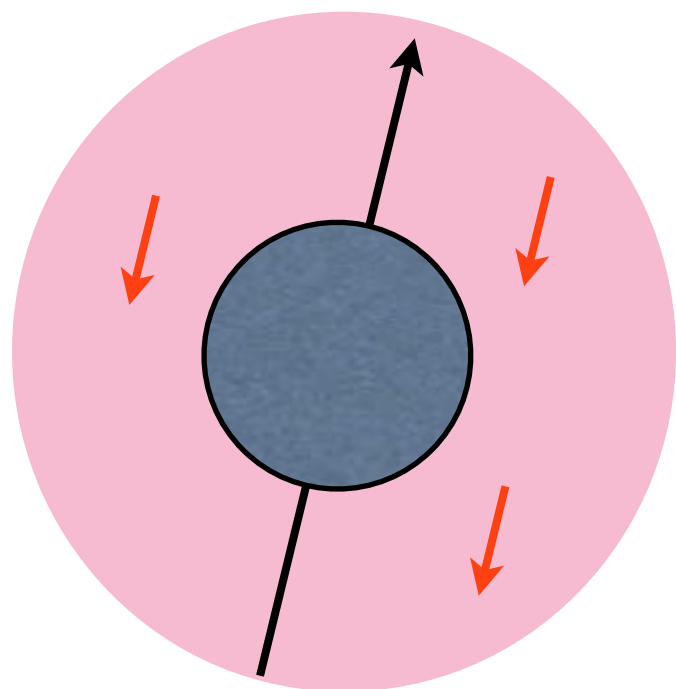


$$T_K = W \sqrt{J\rho} e^{-\frac{1}{2J\rho}}$$



“Kondo Resistance Minimum”

Heavy Fermions + Kondo



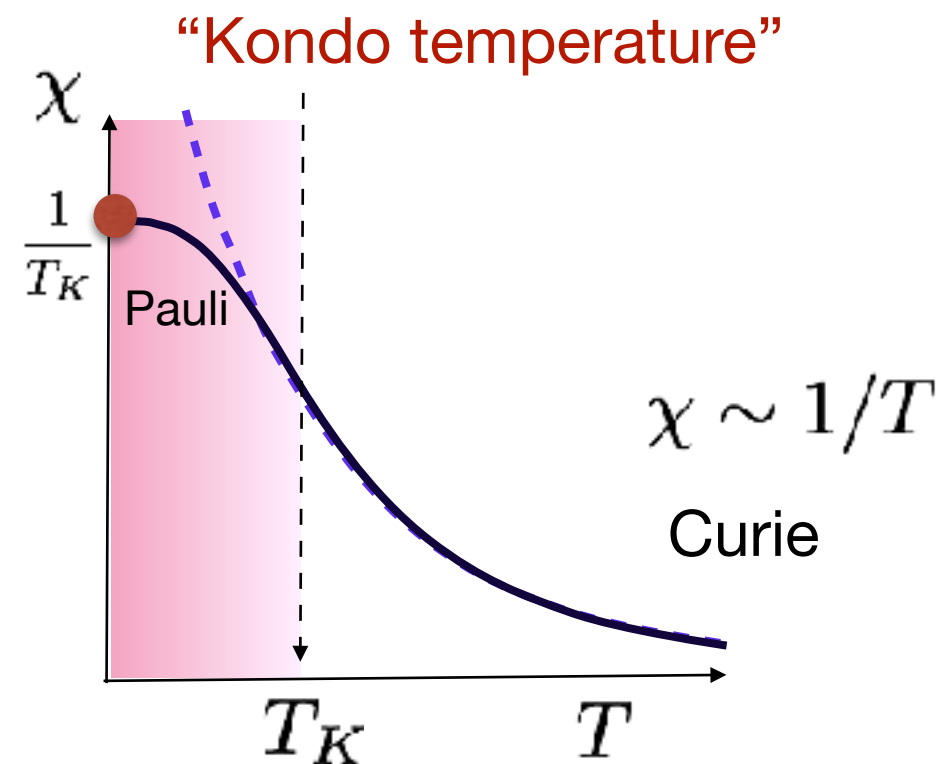
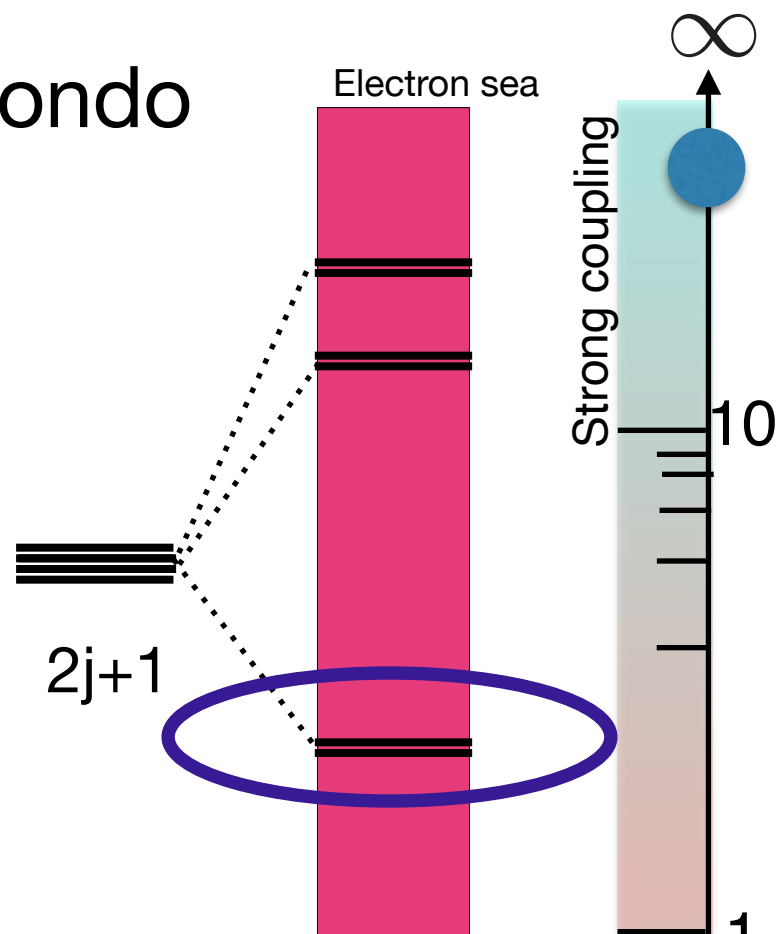
Spin screened by conduction electrons: **entangled**

$$J \rightarrow J(T) = J + 2J^2\rho \ln \frac{D}{T}$$

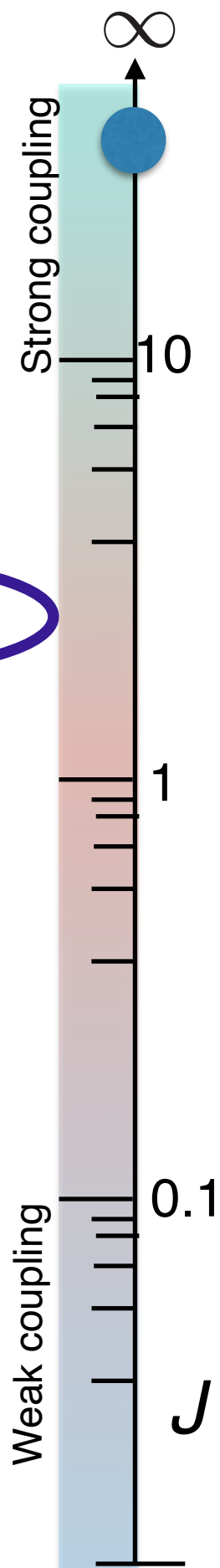
“Scales to Strong Coupling”

$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + \underbrace{J}_{\text{Kondo coupling}} \vec{S} \cdot \vec{\sigma}(0)$$

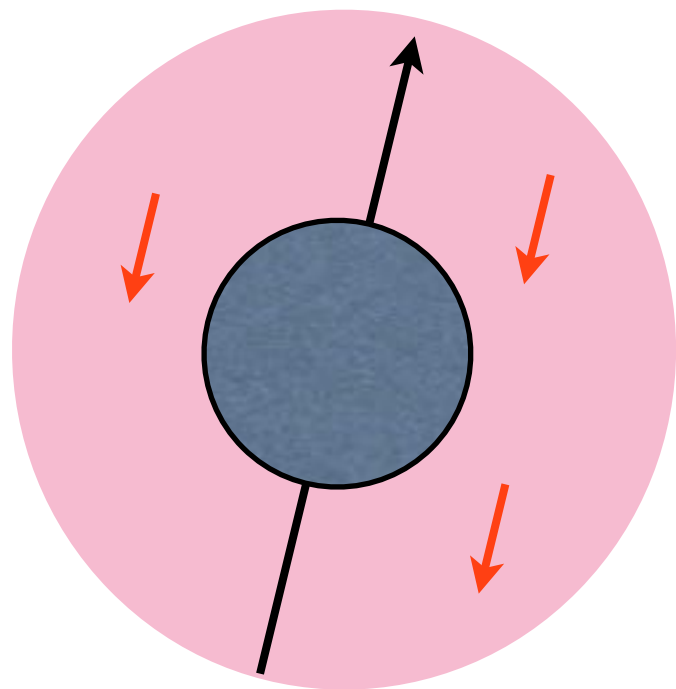
J. Kondo, 1962



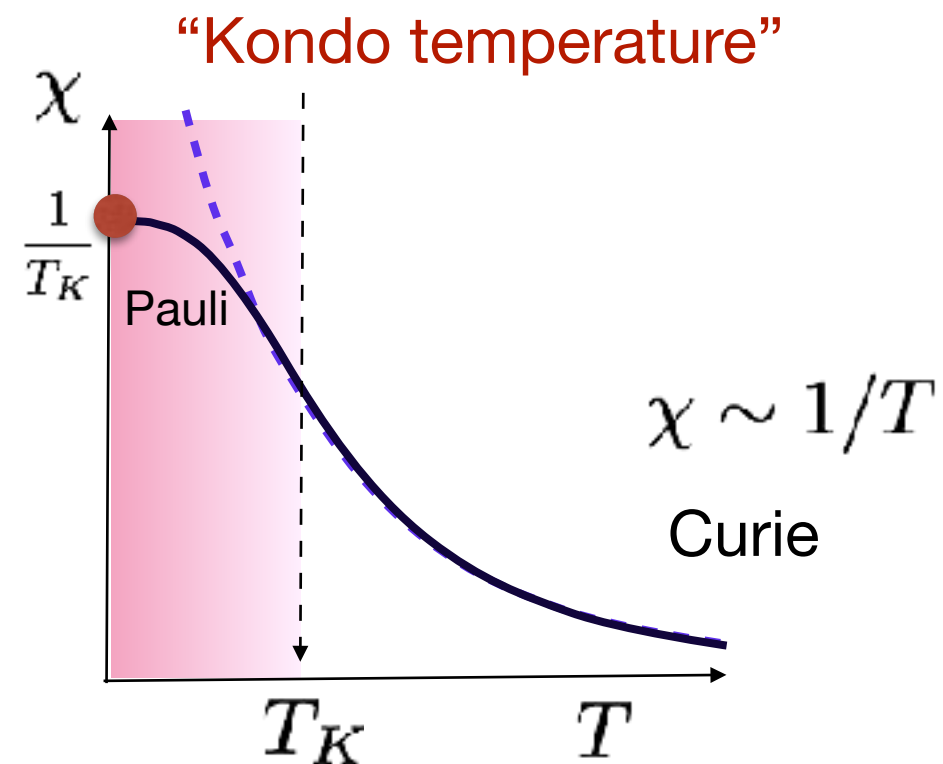
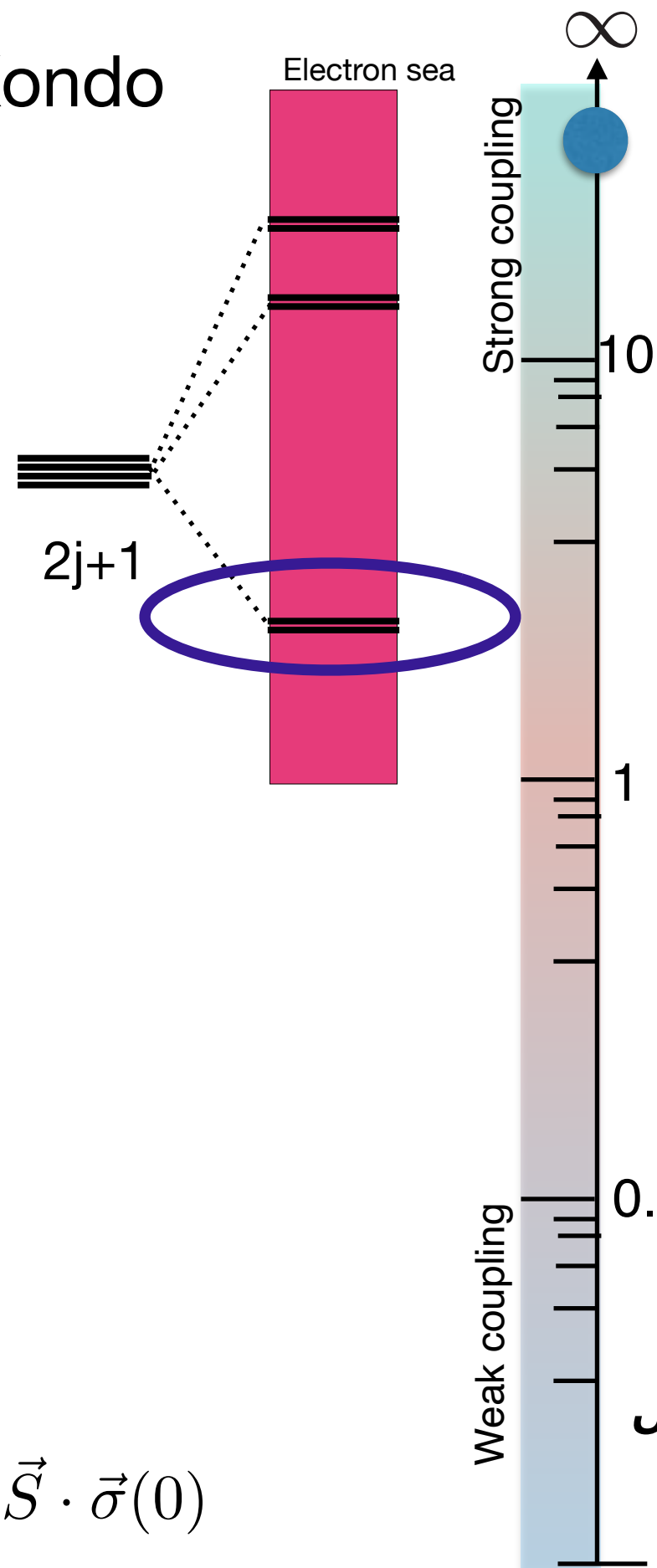
$$TK = W \sqrt{J\rho} e^{-\frac{1}{2J\rho}}$$



Heavy Fermions + Kondo



Spin screened by conduction electrons: entangled

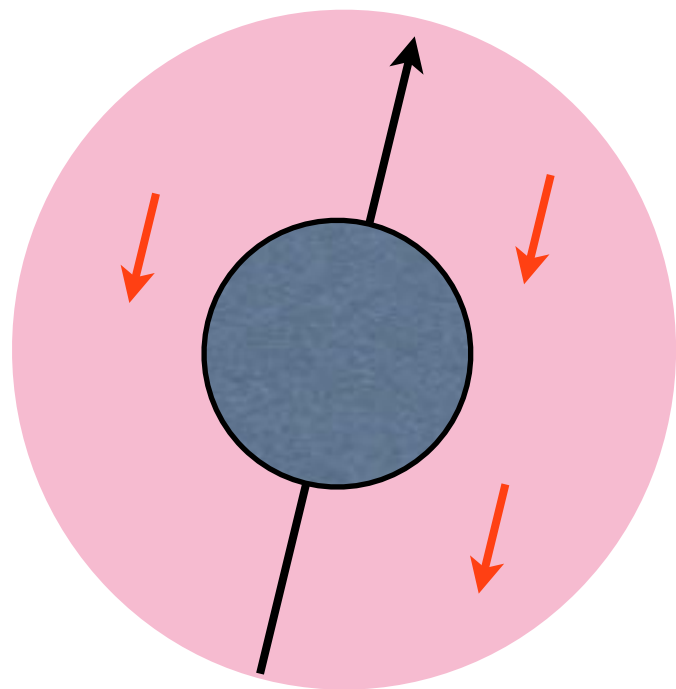


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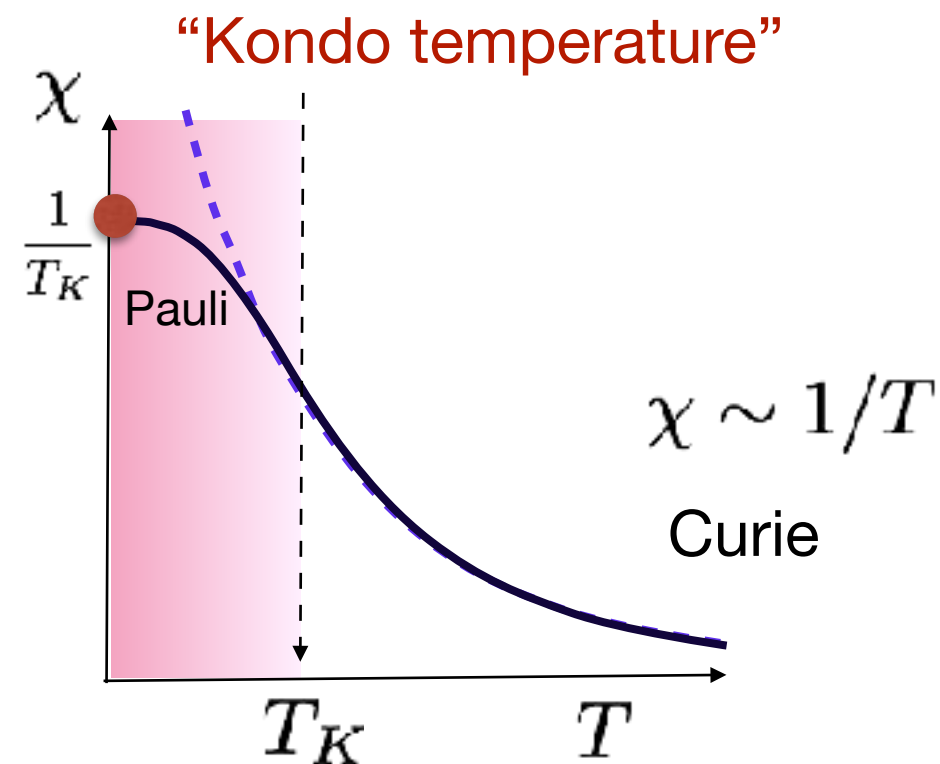
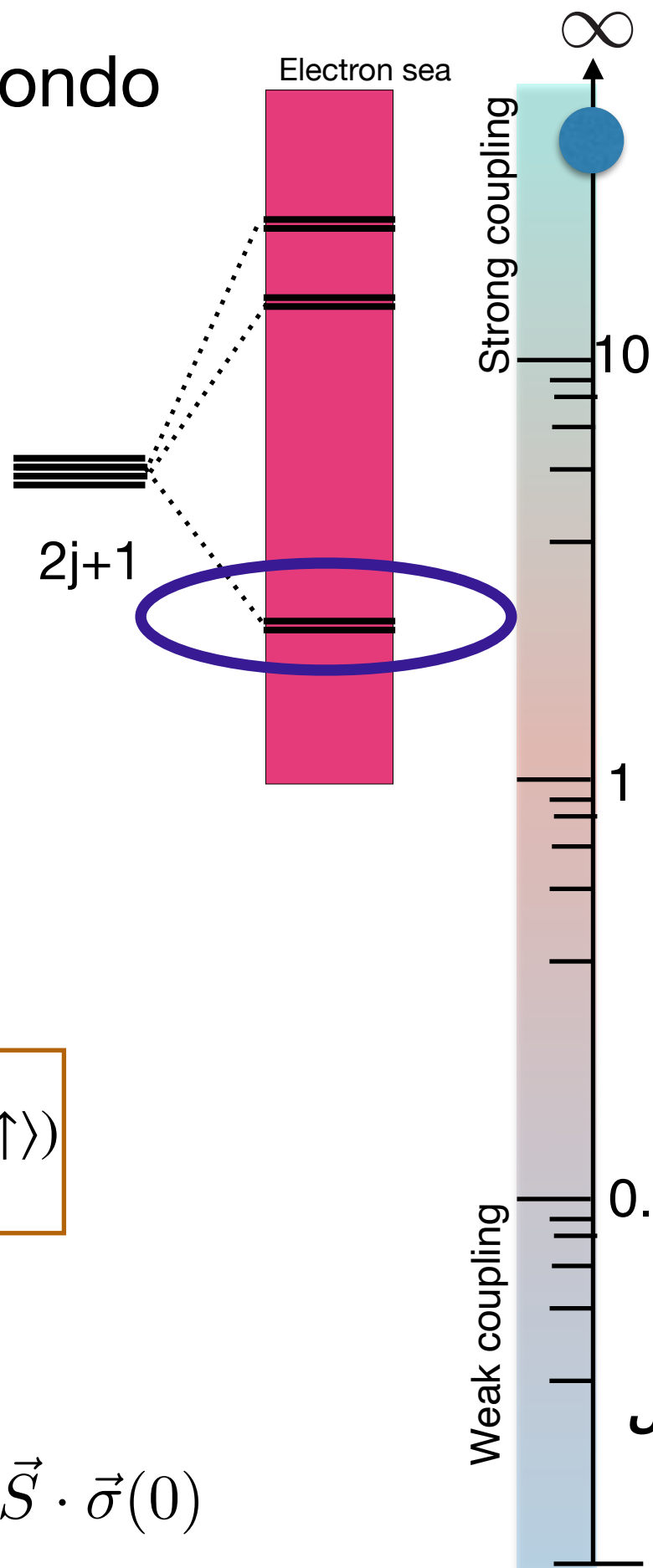
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J. Kondo, 1962

Heavy Fermions + Kondo



Spin screened by conduction electrons: entangled



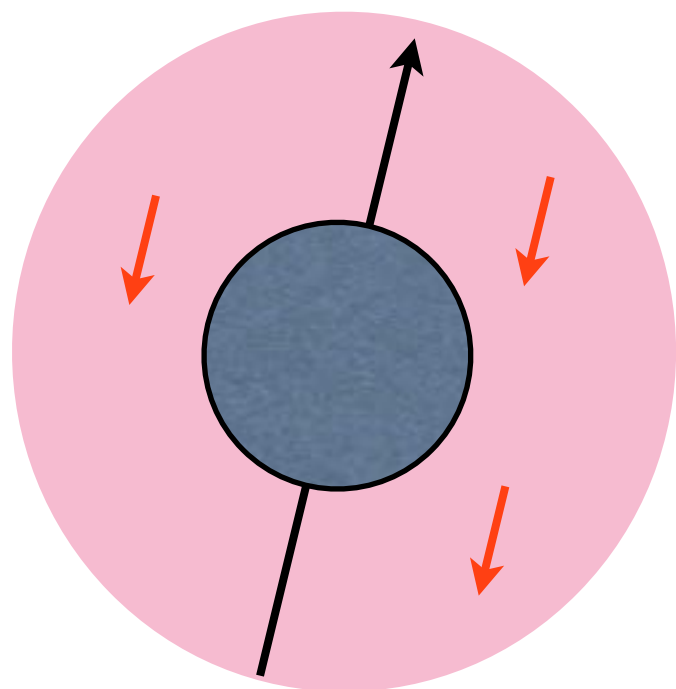
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$$|GS\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

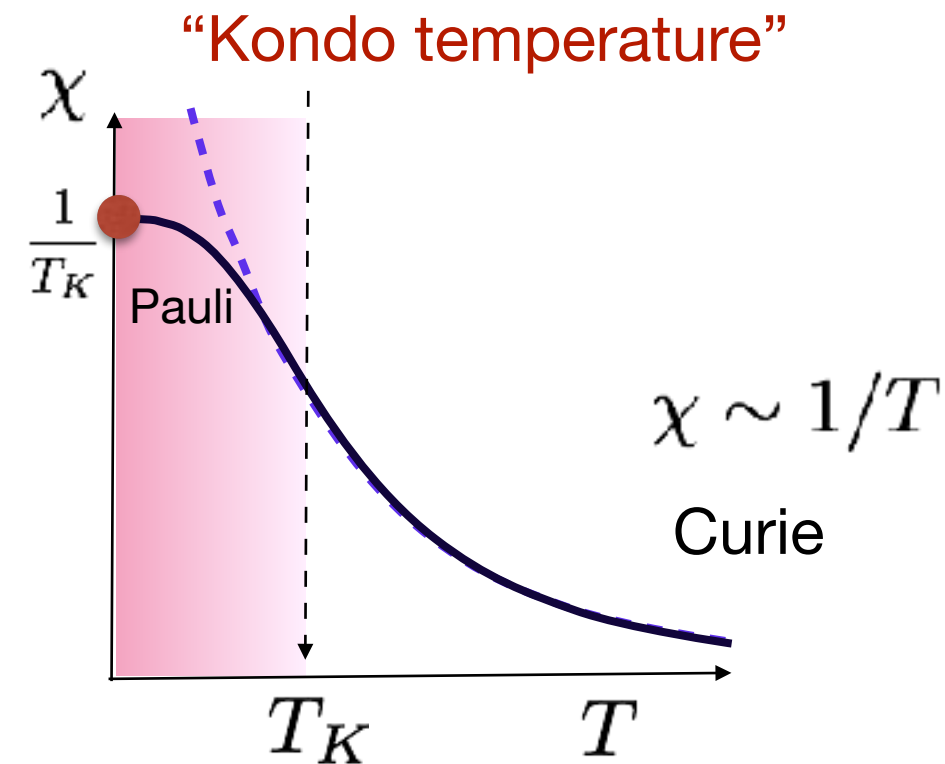
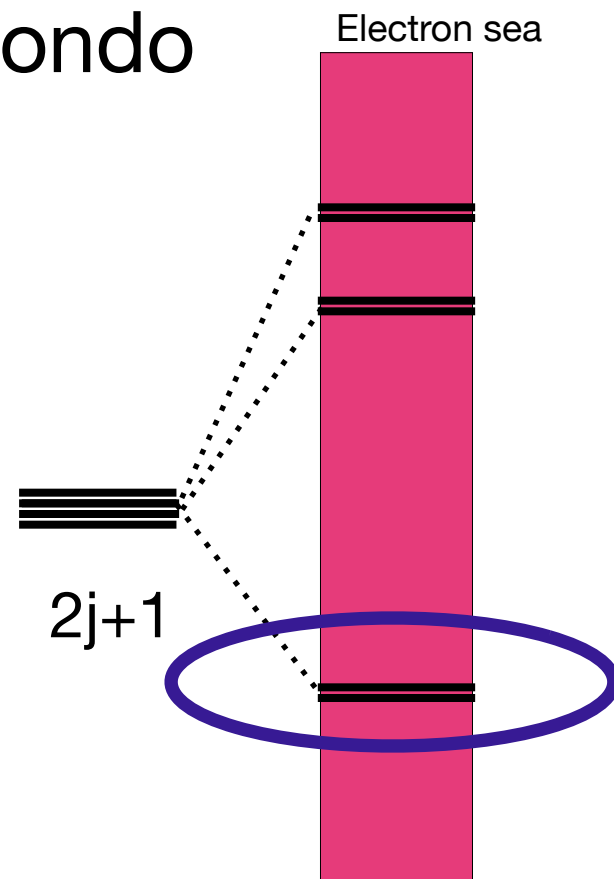
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Heavy Fermions + Kondo



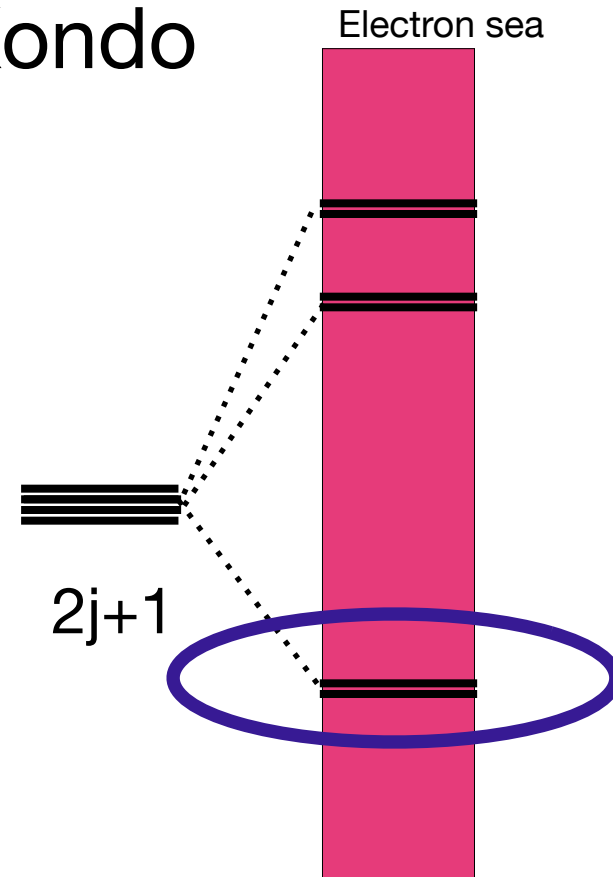
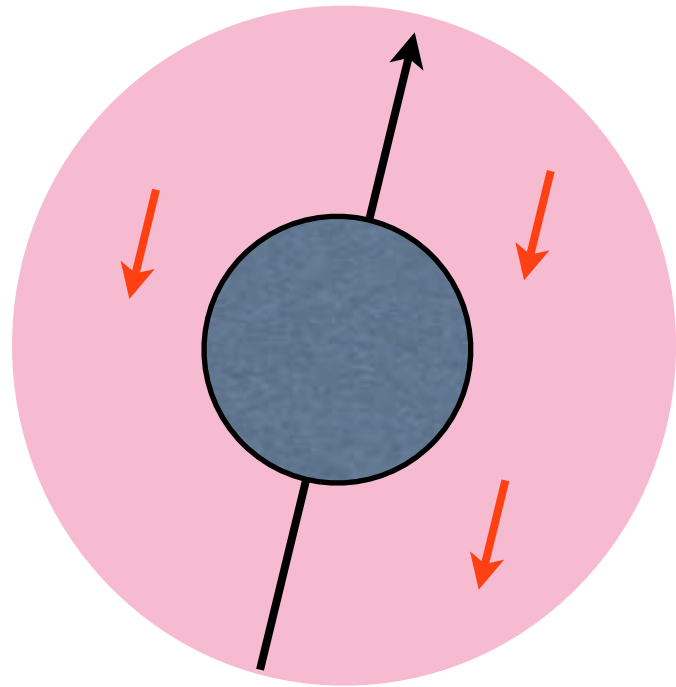
Spin screened by
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electrons: entangled



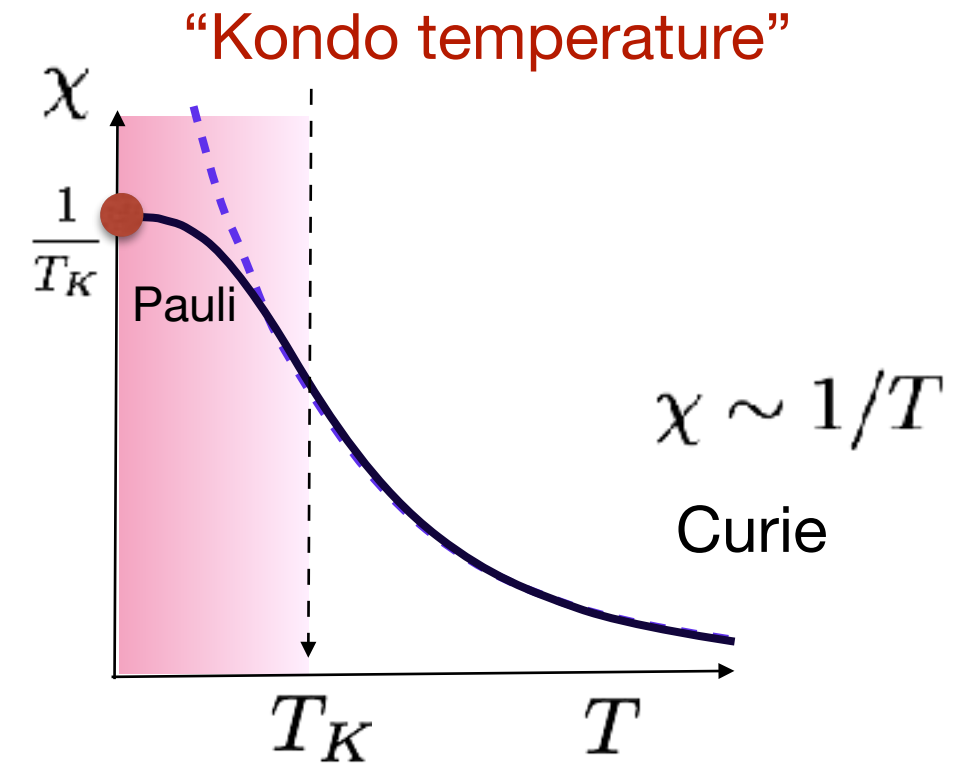
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Heavy Fermions + Kondo



Spin screened by conduction electrons: entangled



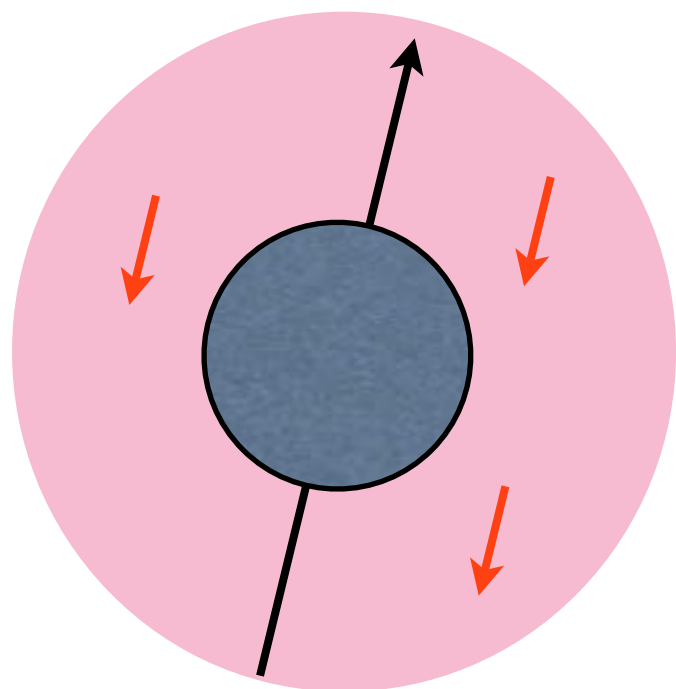
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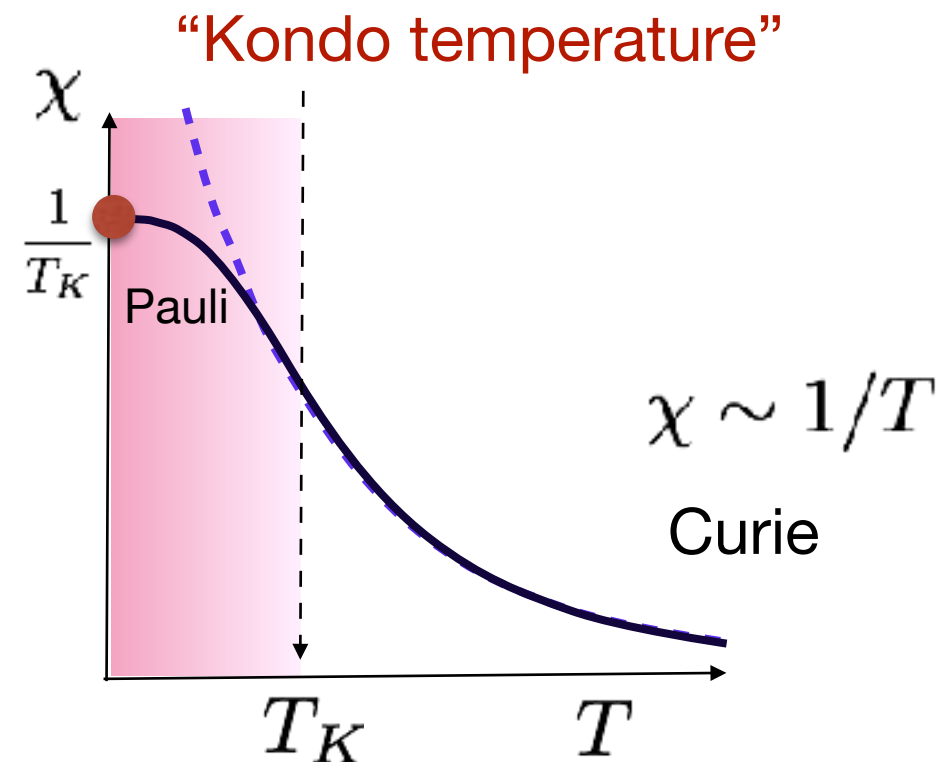
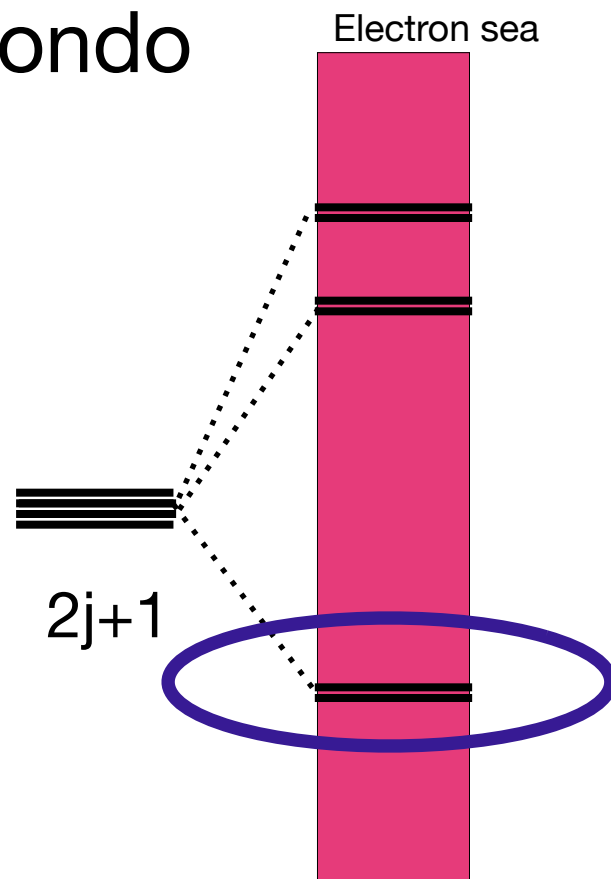
$$S(T) = \int_0^T \frac{C_V}{T'} dT'$$

Spin entanglement entropy

Heavy Fermions + Kondo



Spin screened by conduction electrons: entangled

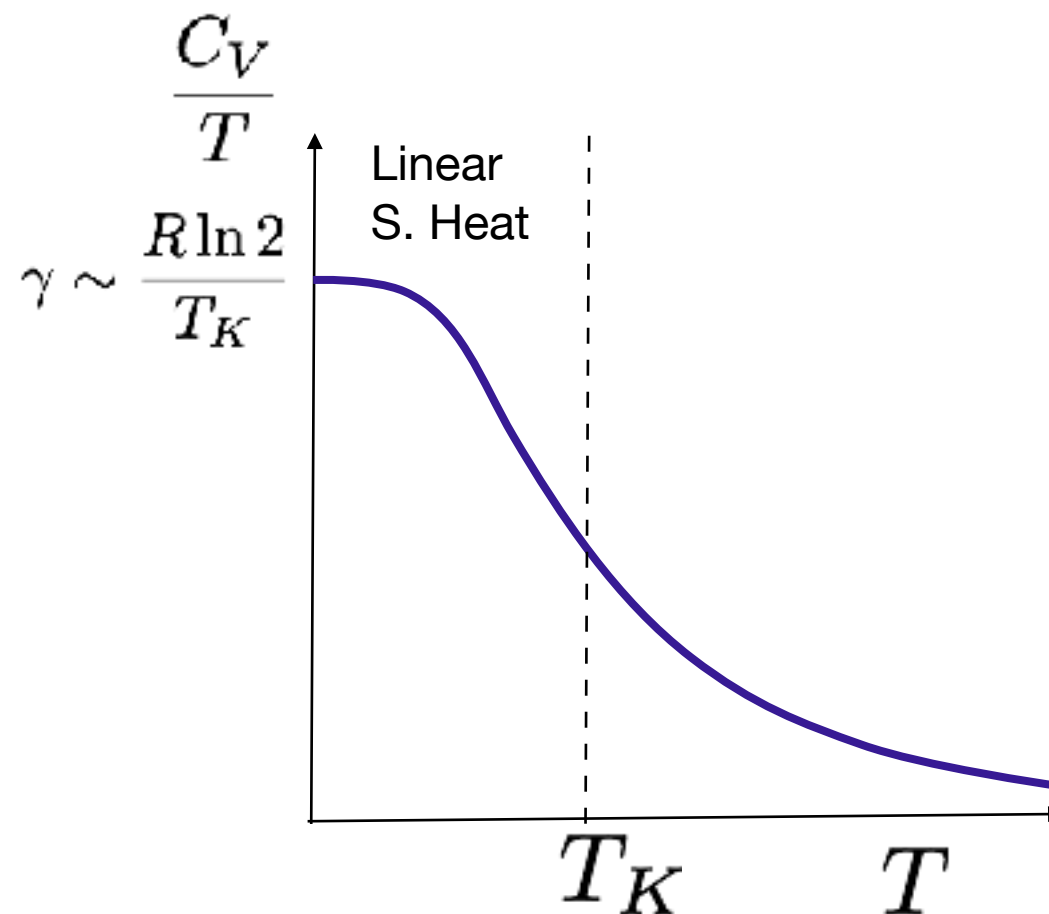


$$T_K = W \sqrt{J\rho} e^{-\frac{1}{2J\rho}}$$

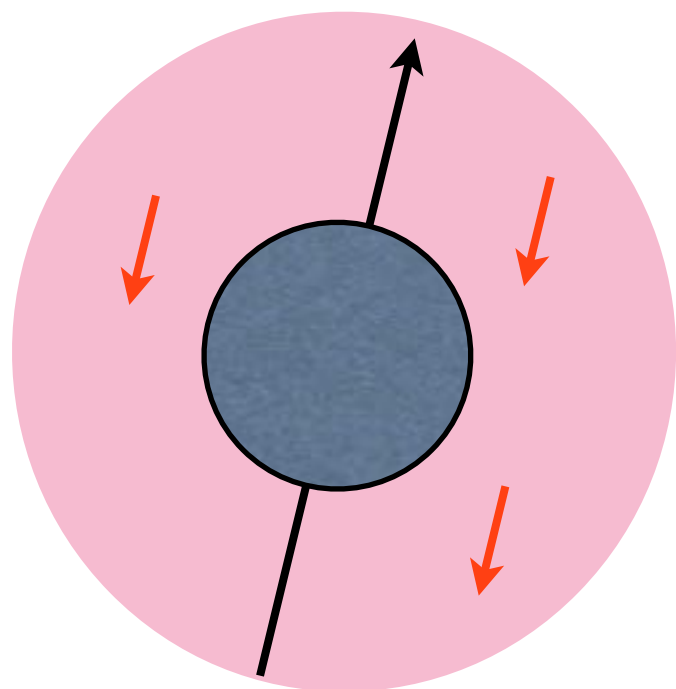
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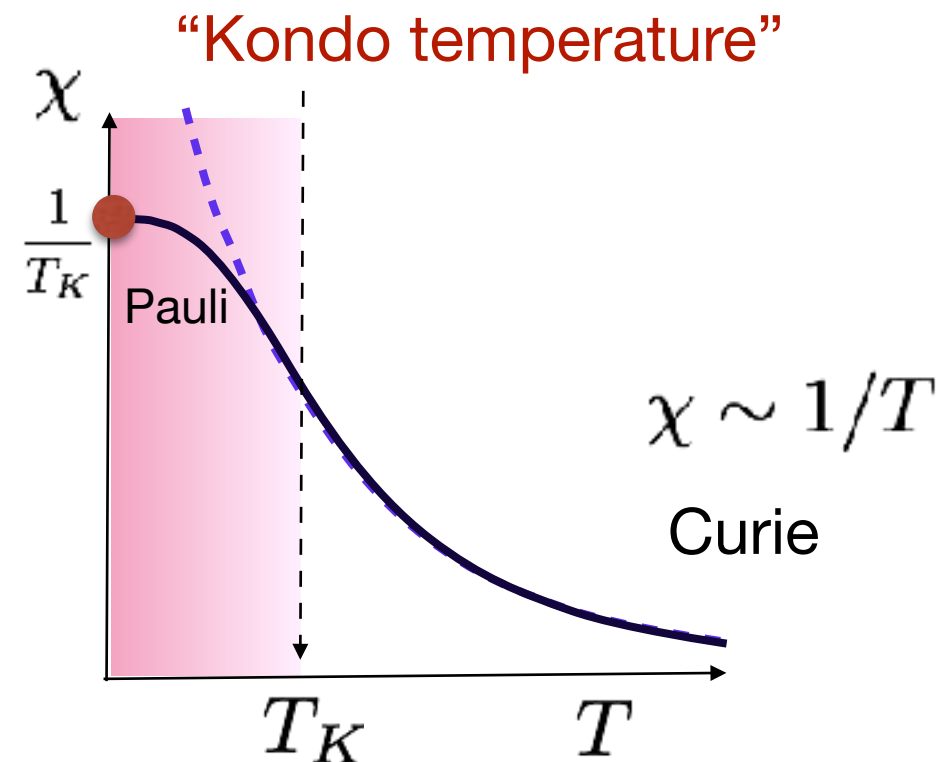
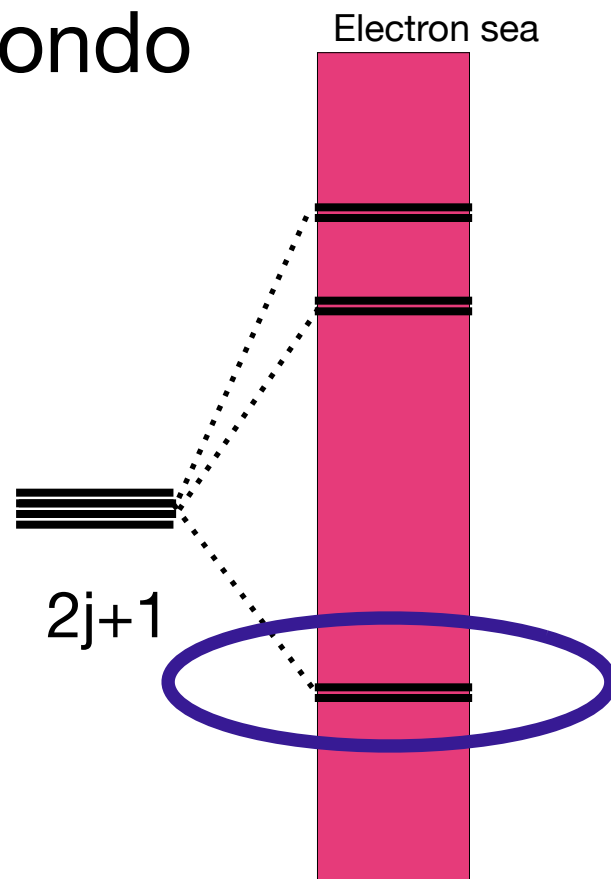
Spin entanglement entropy



Heavy Fermions + Kondo



Spin screened by conduction electrons: entangled

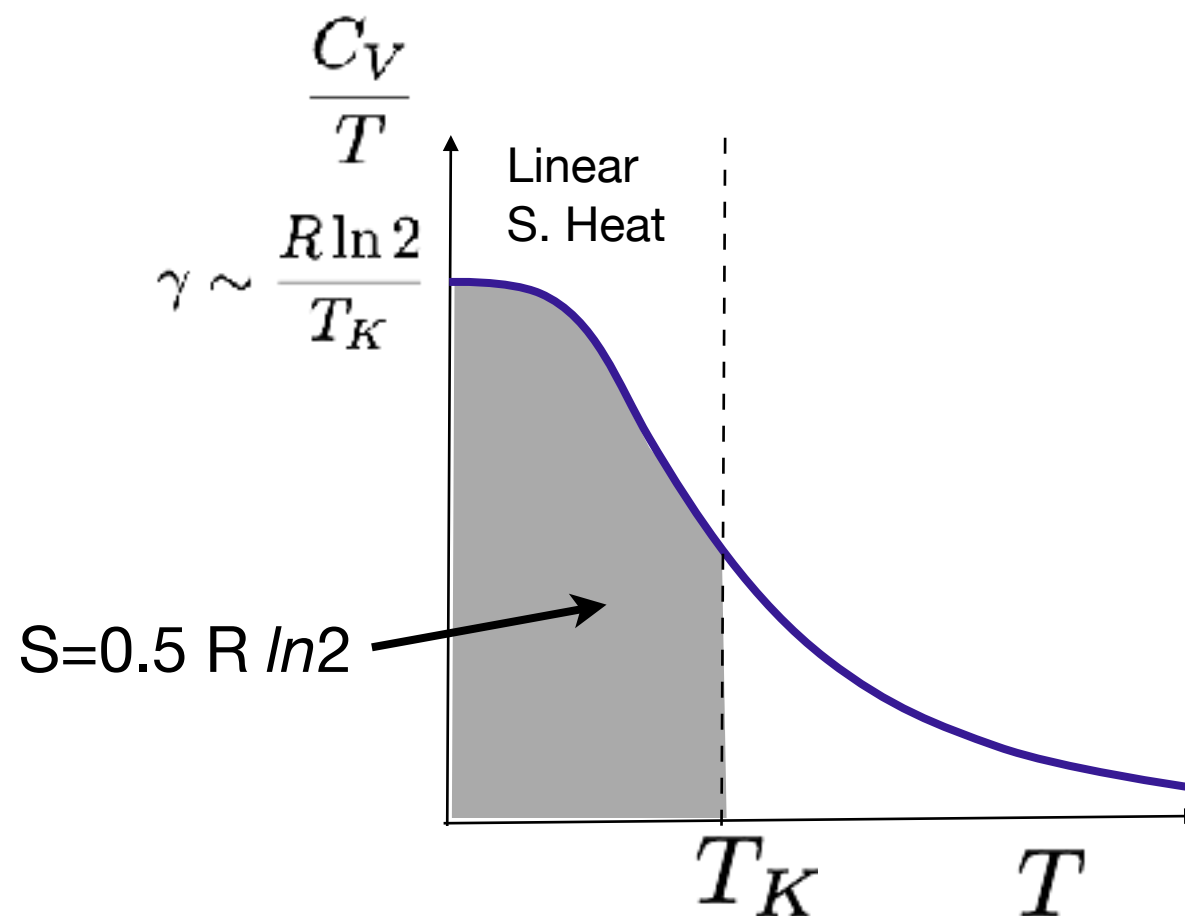


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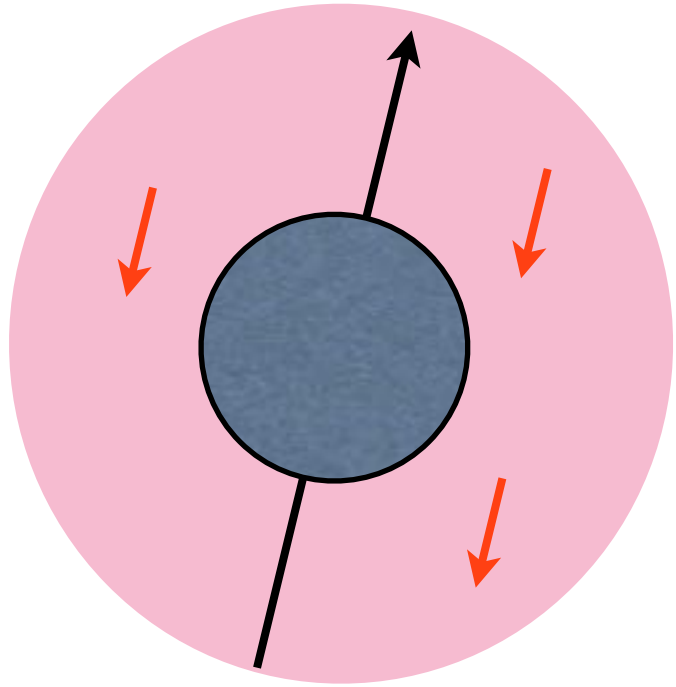
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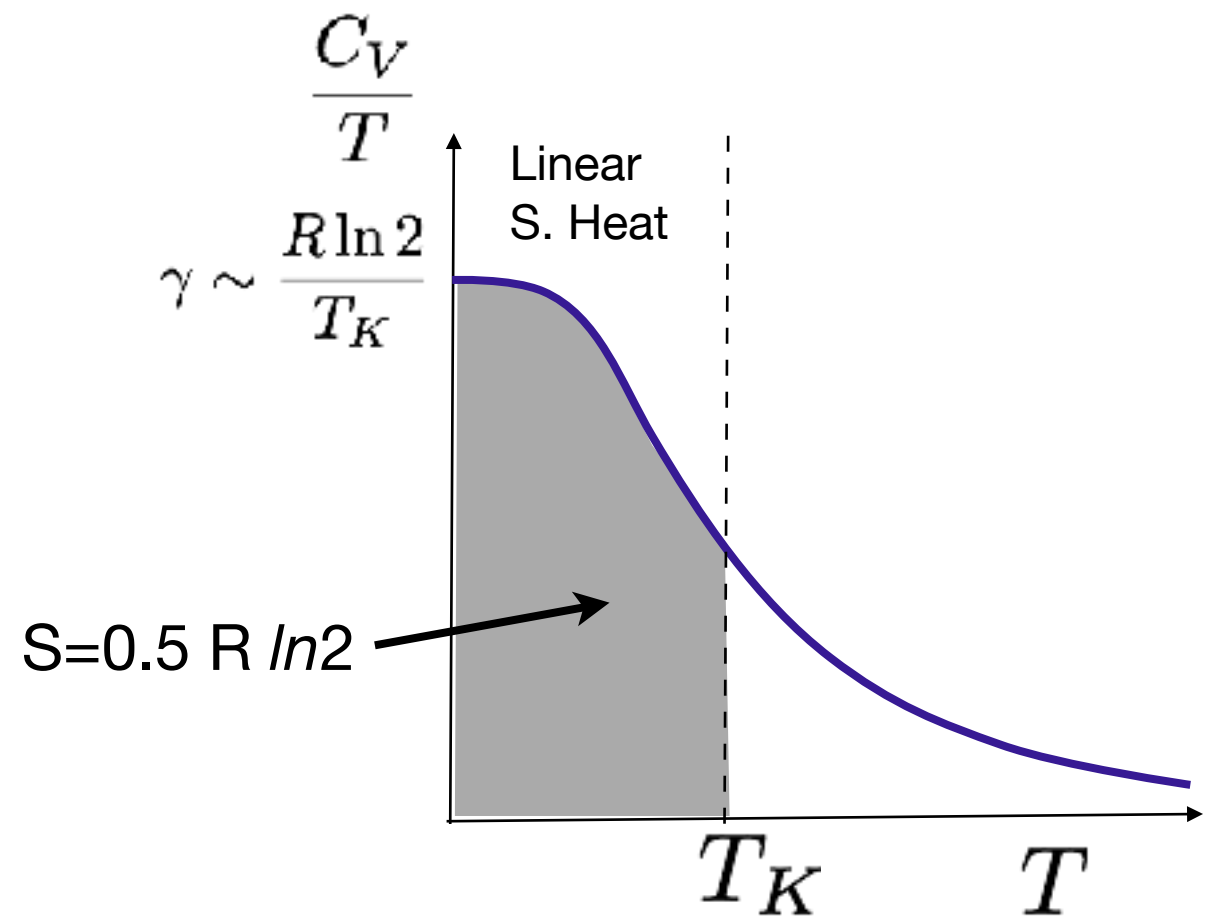


Heavy Fermion Primer

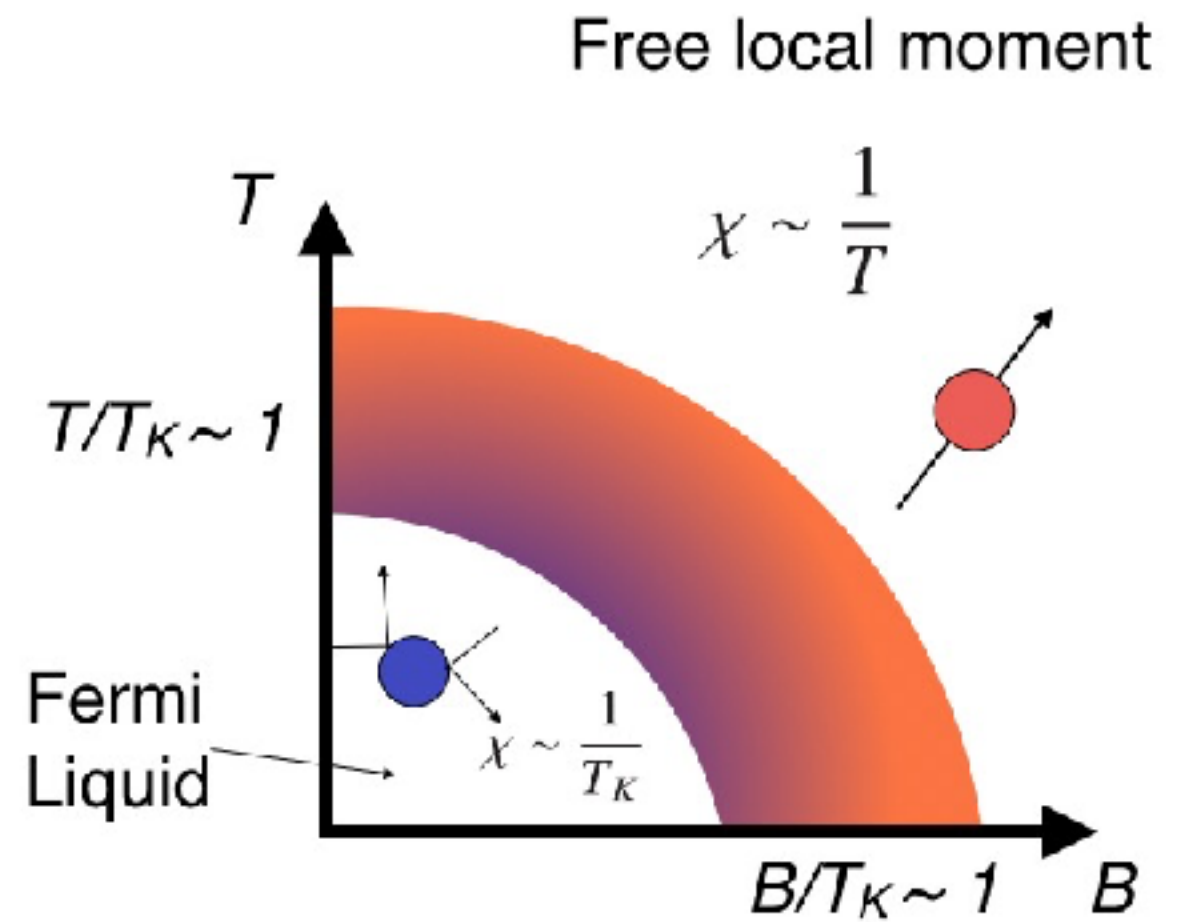
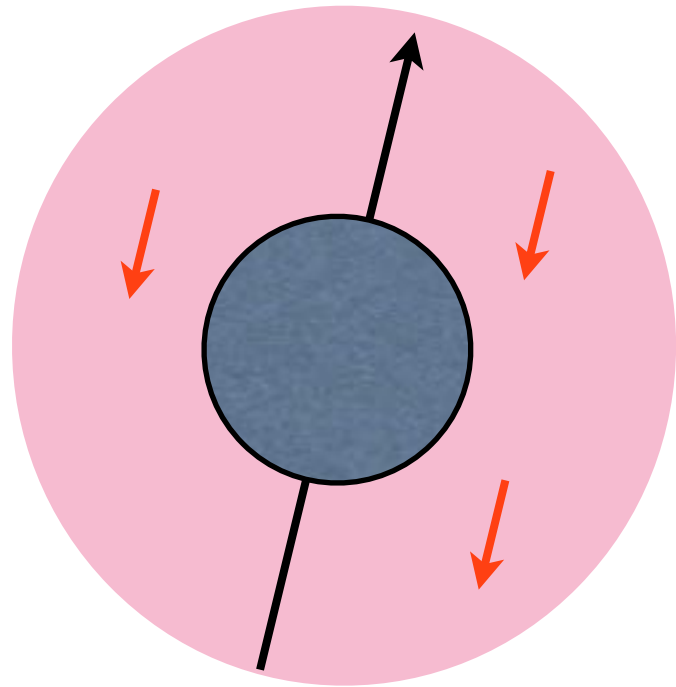


$$S(T) = \int_0^T \frac{C_V}{T'} dT'$$

Spin entanglement entropy

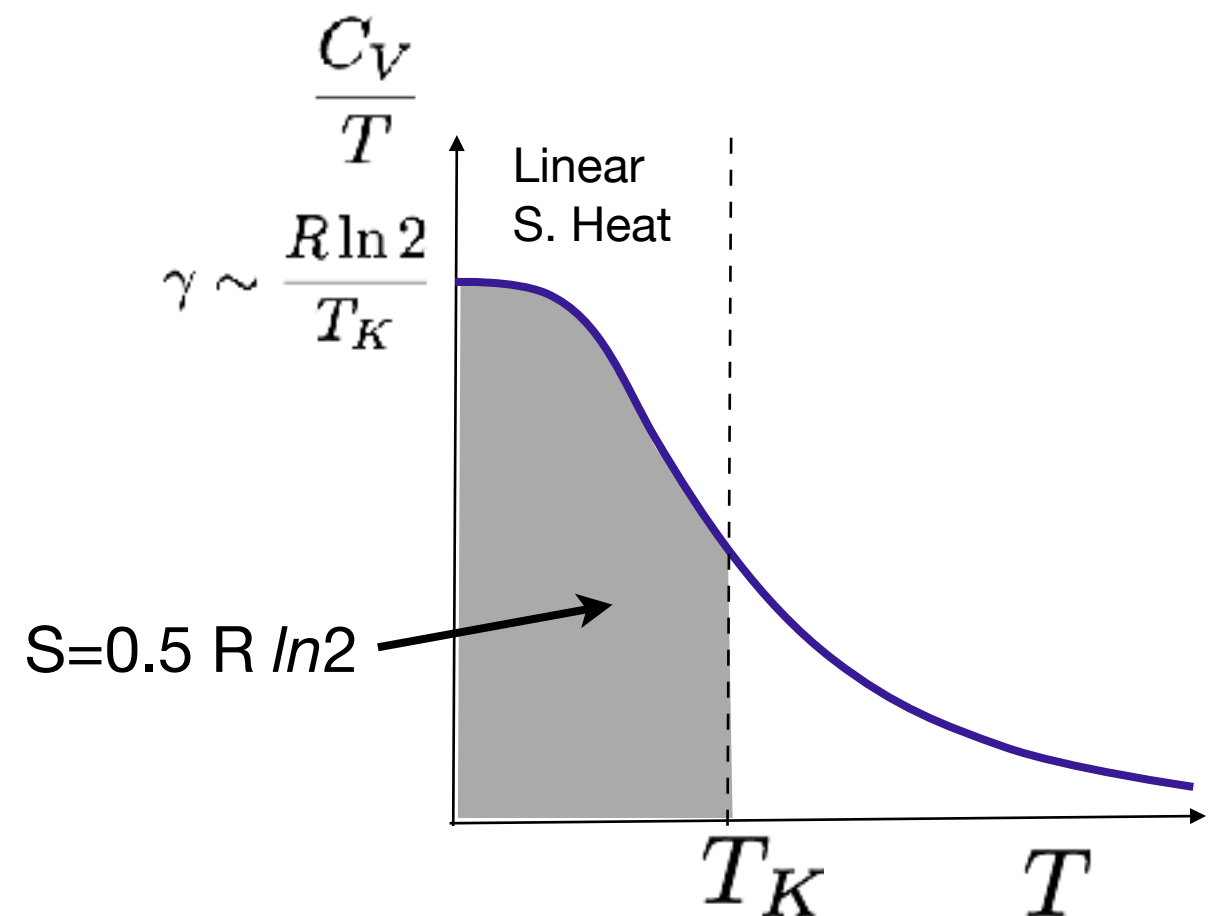


Heavy Fermion Primer

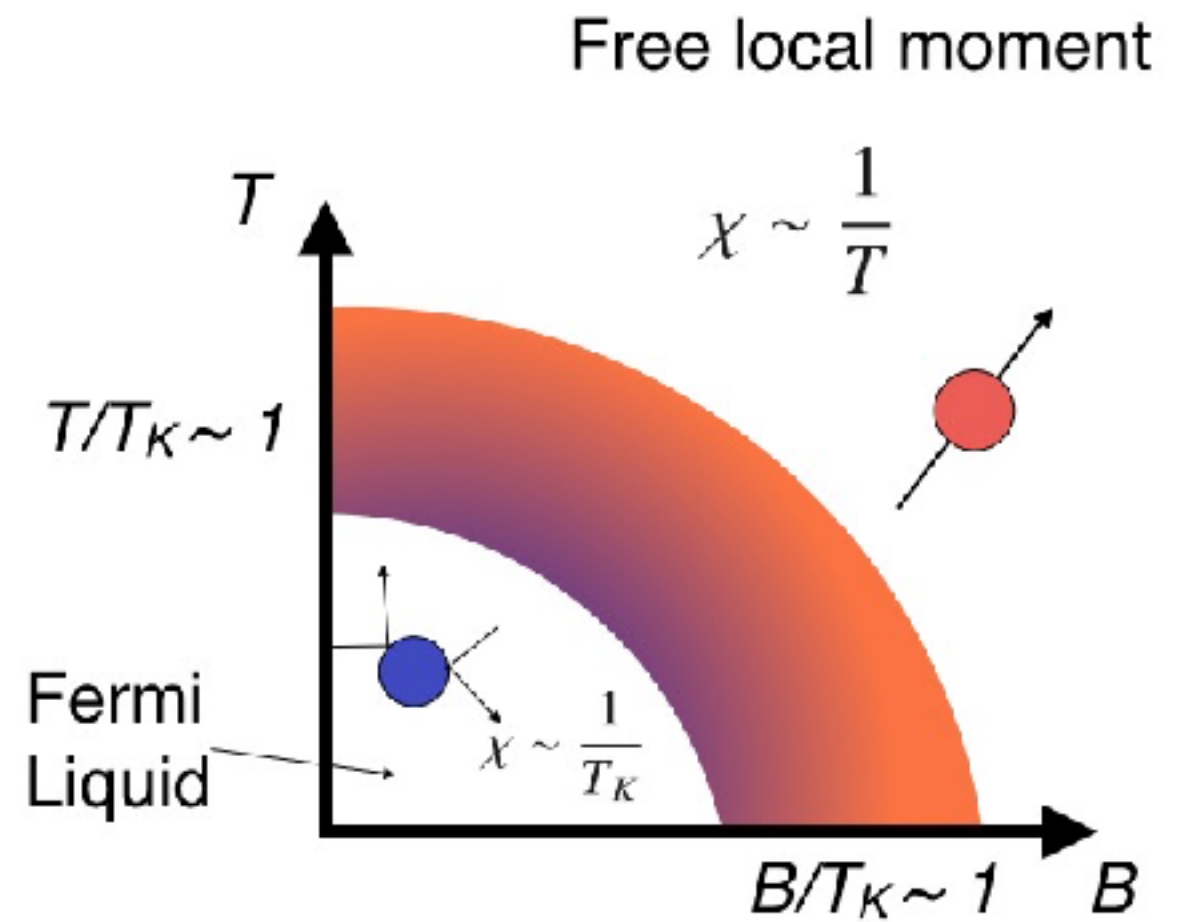
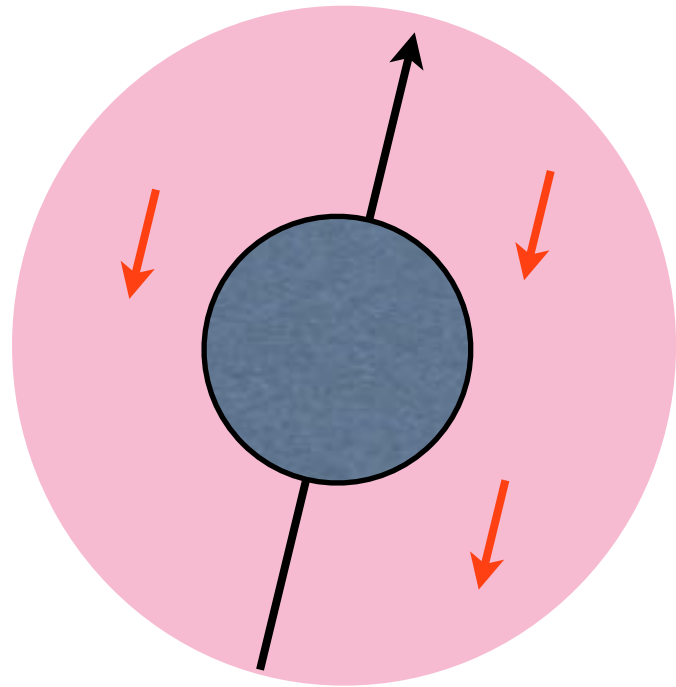


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Spin entanglement entropy

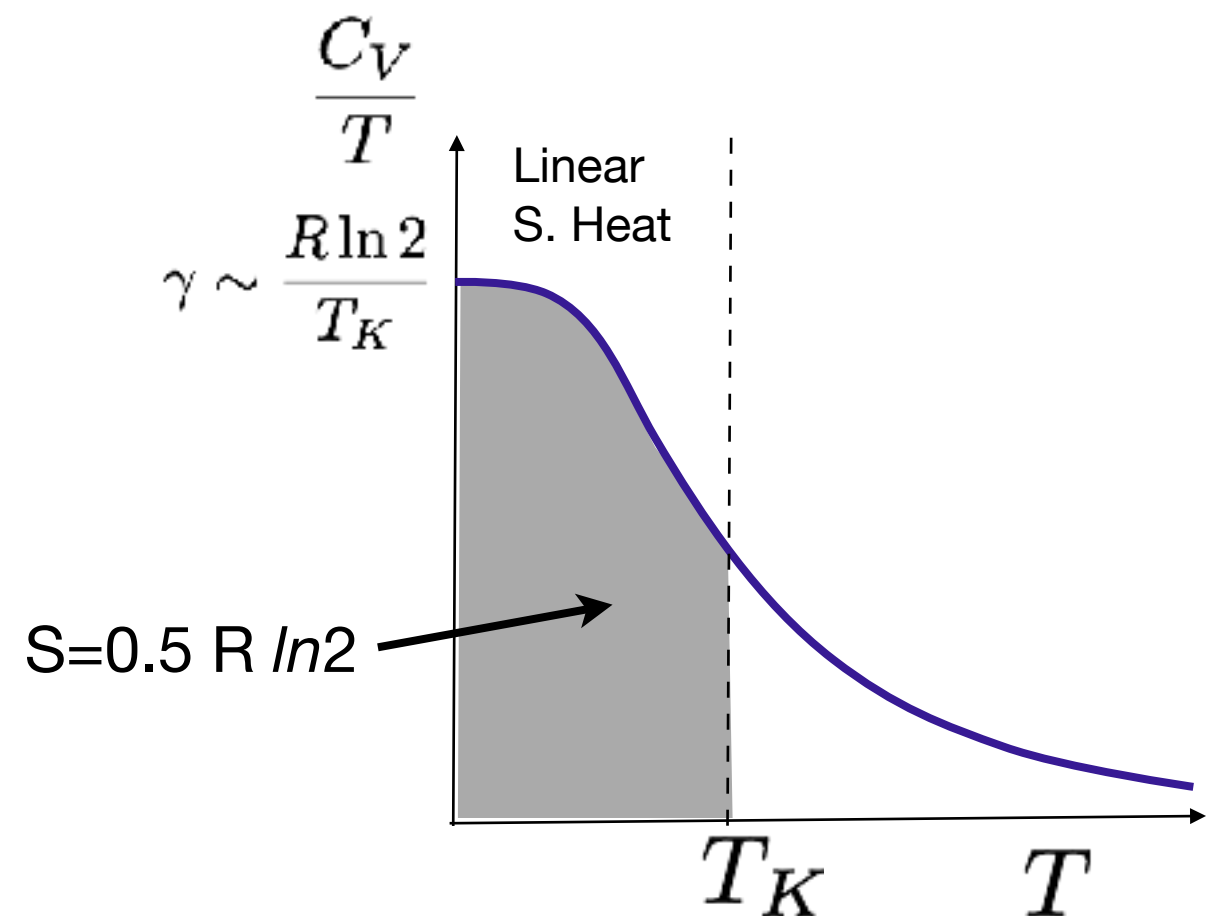


Heavy Fermion Primer

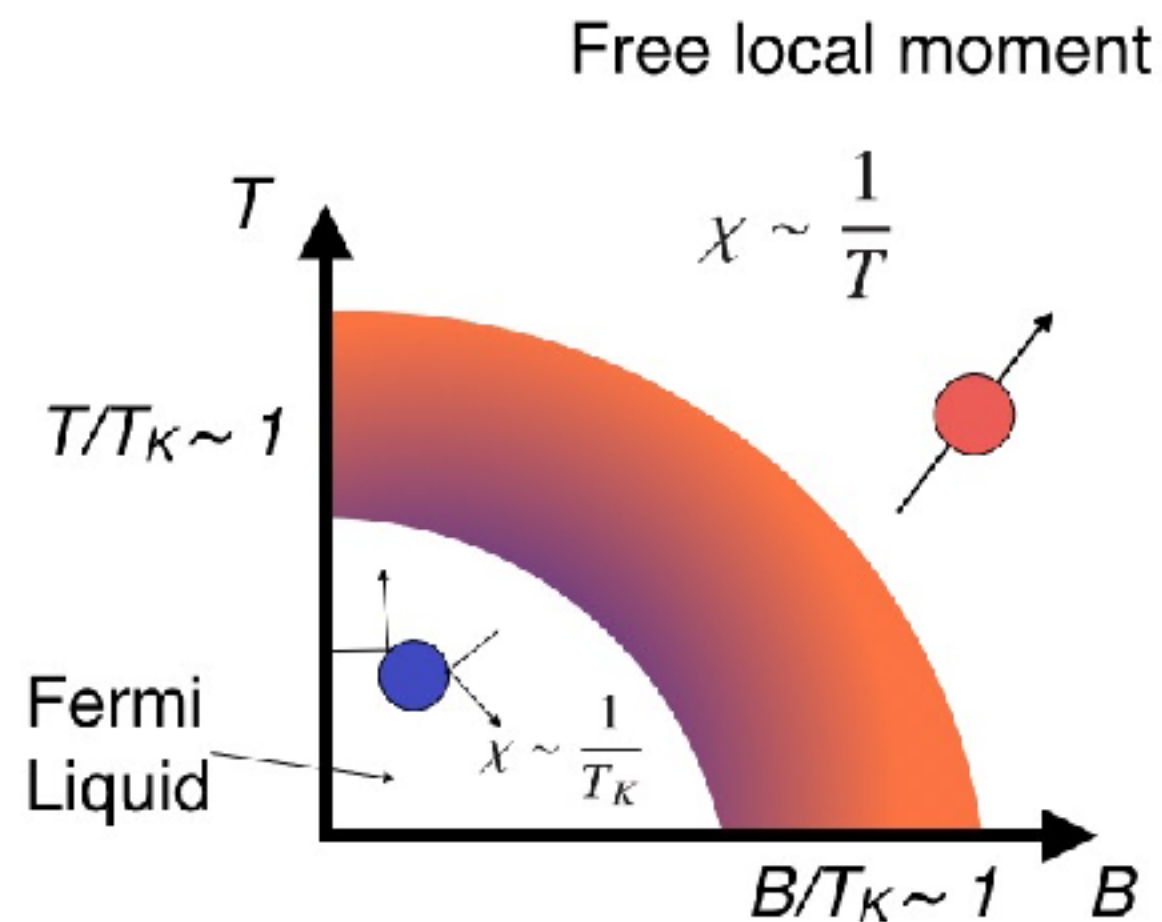
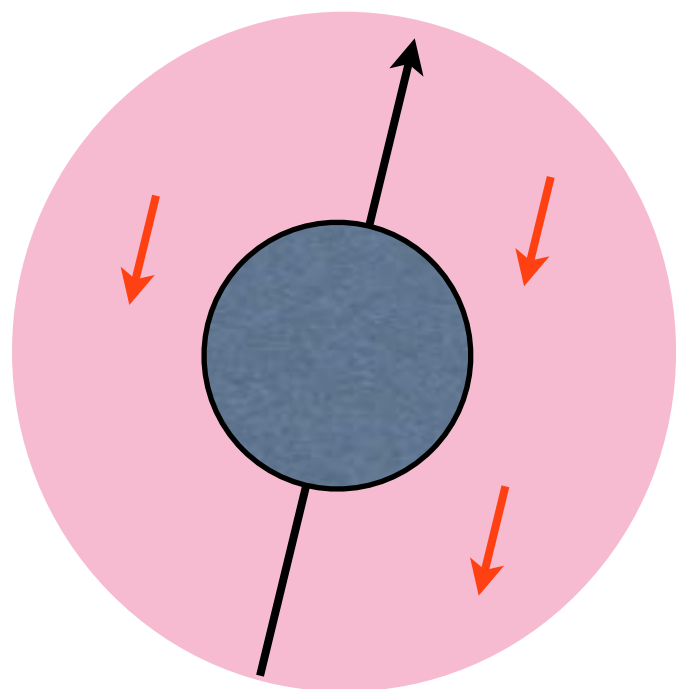


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Spin entanglement entropy



Heavy Fermion Primer

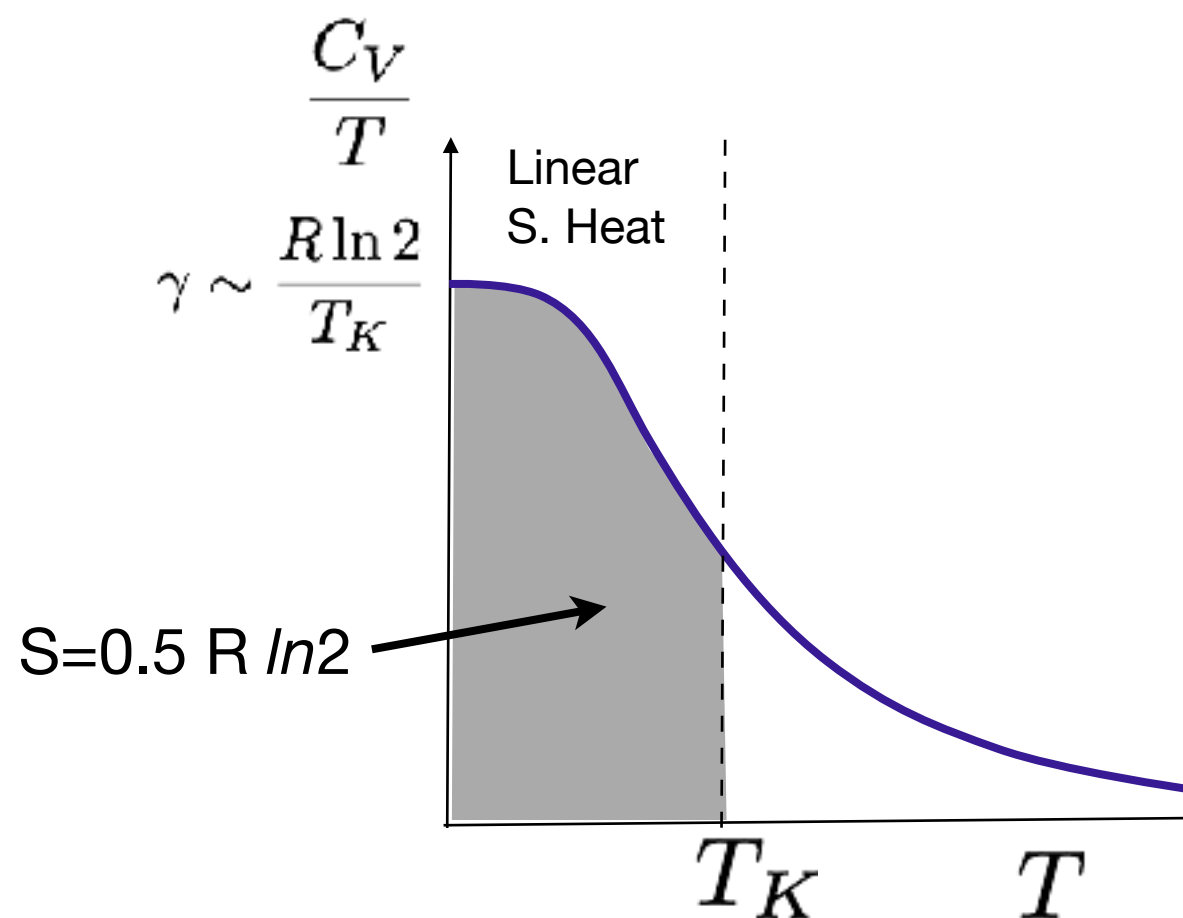


$$\frac{R(T)}{R_U} = n_i \Phi \left(\frac{T}{T_K} \right)$$

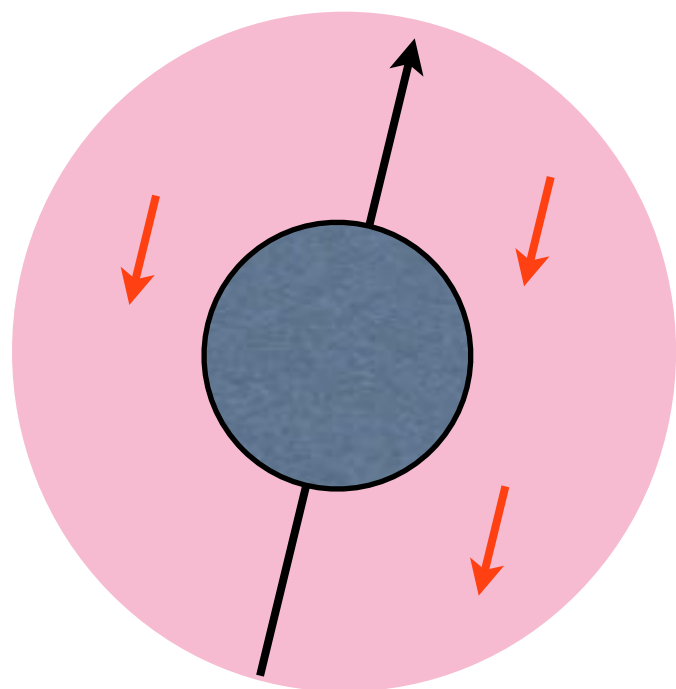
Universality

$$S(T) = \int_0^T \frac{C_V}{T'} dT'$$

Spin entanglement entropy



Heavy Fermion Primer

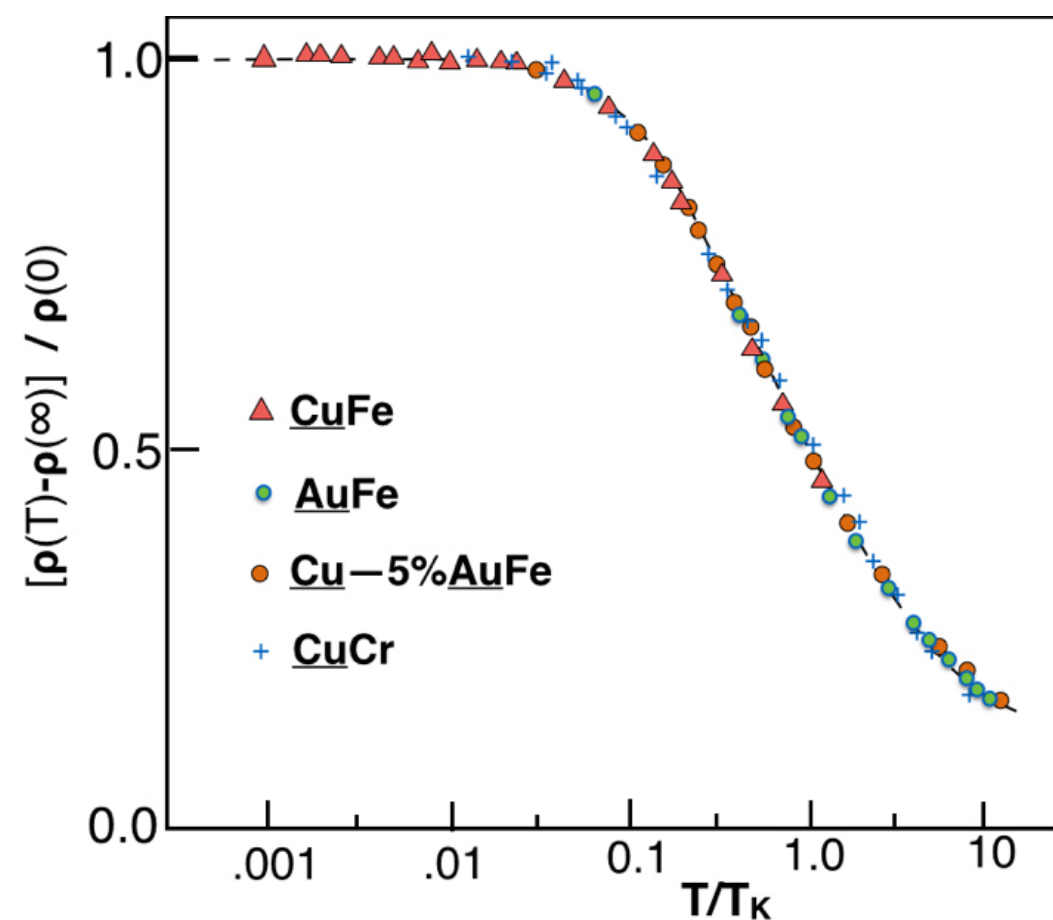
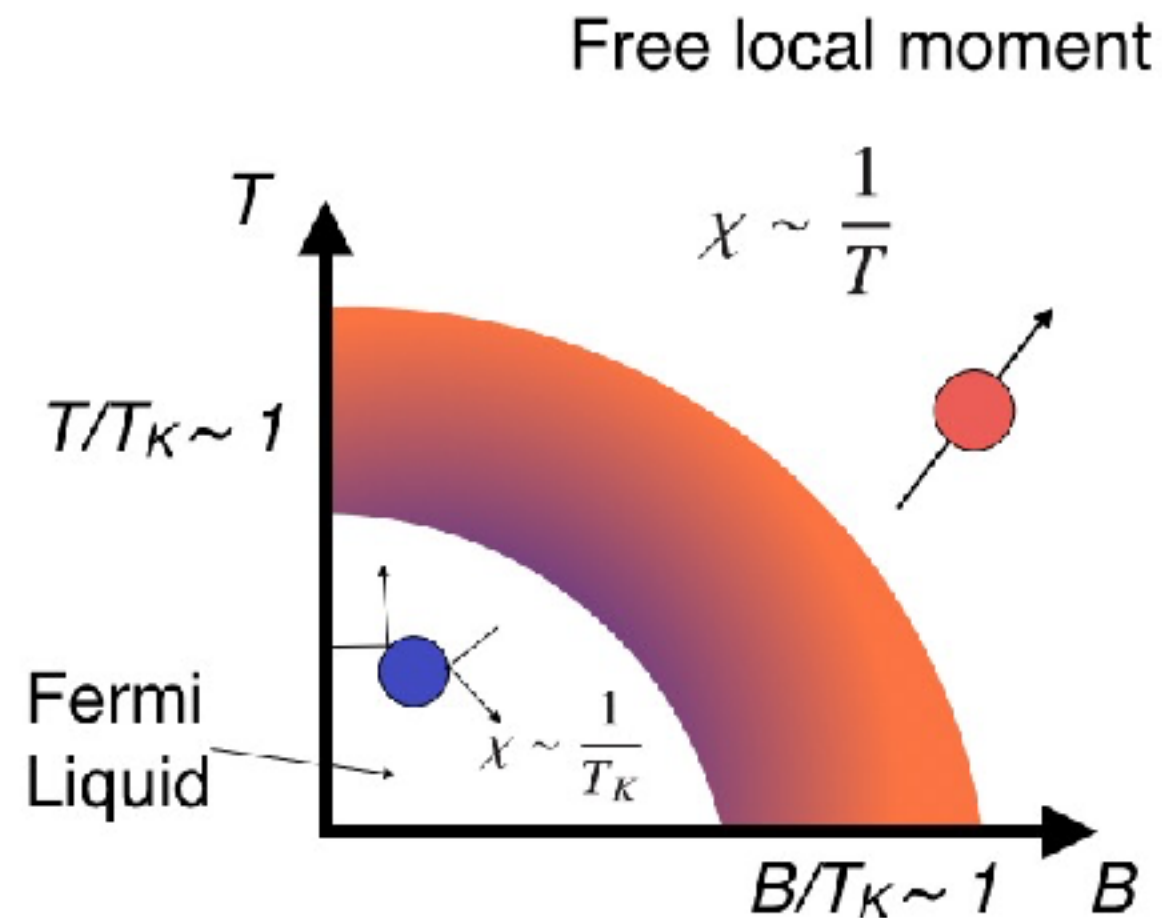


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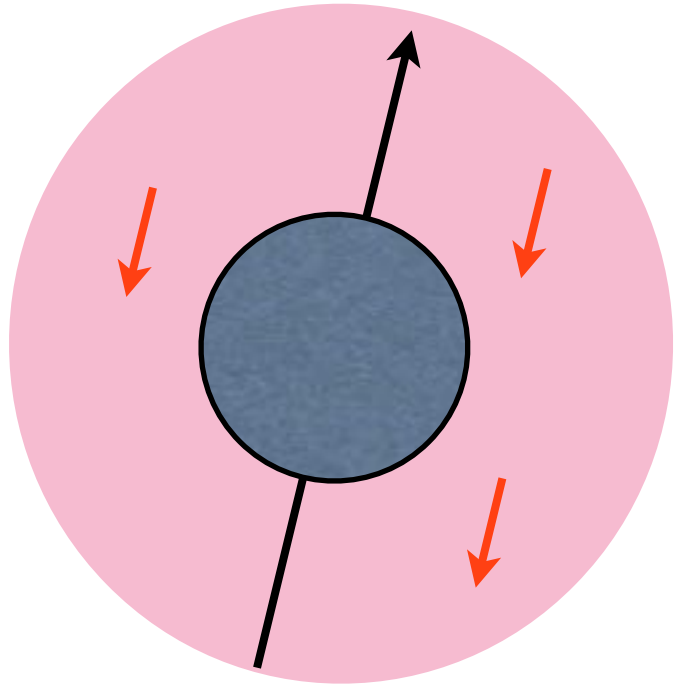
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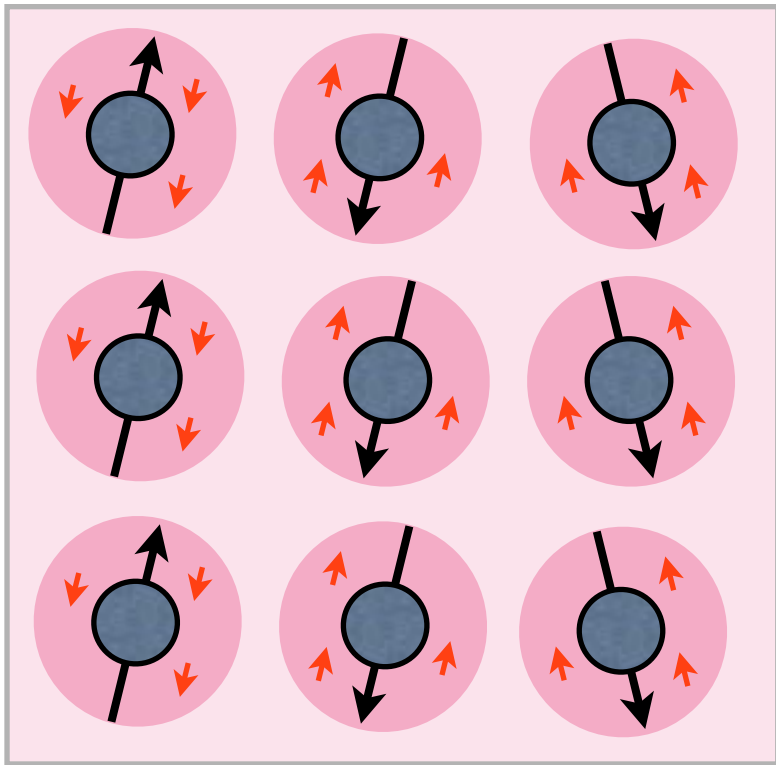
Spin entanglement entropy



Heavy Fermion Primer

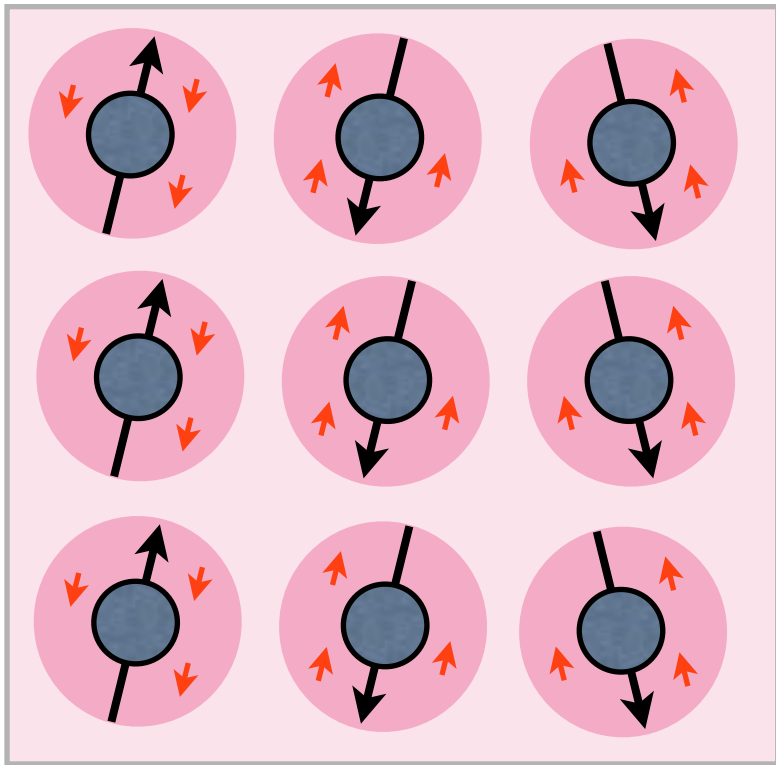


$$H = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} + J \sum_j \vec{S}_j \cdot \vec{\sigma}(j)$$



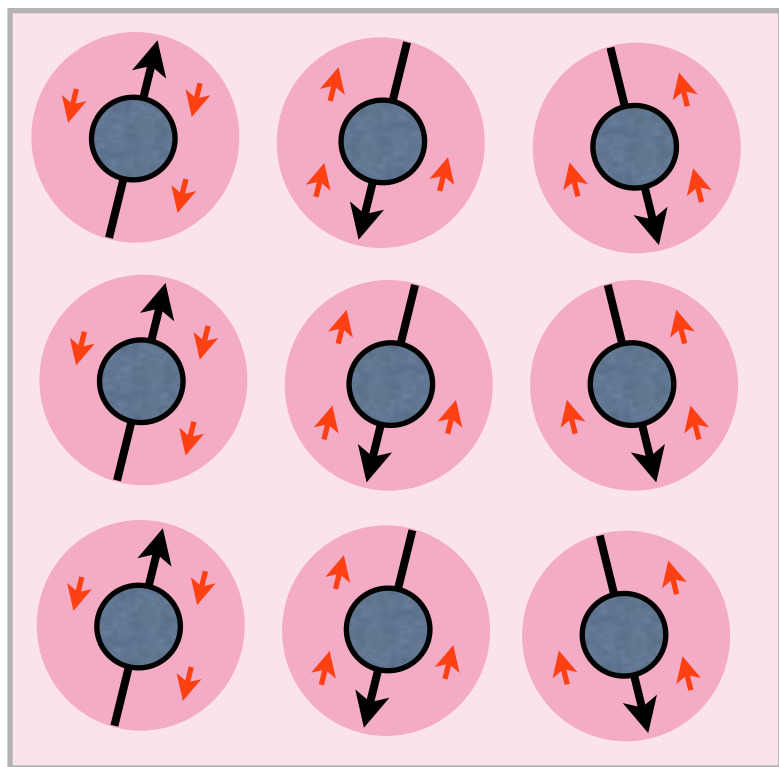
“Kondo Lattice”

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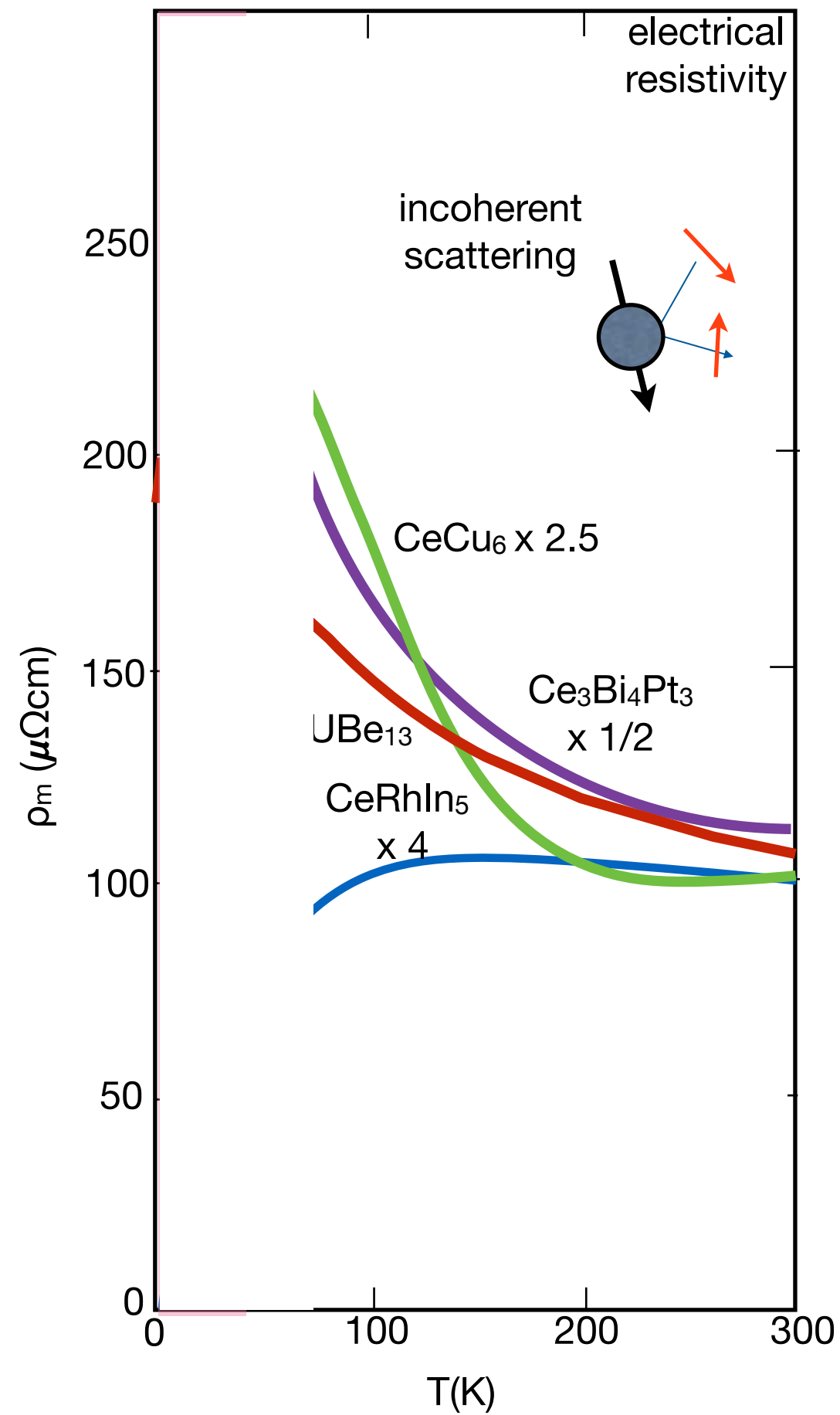


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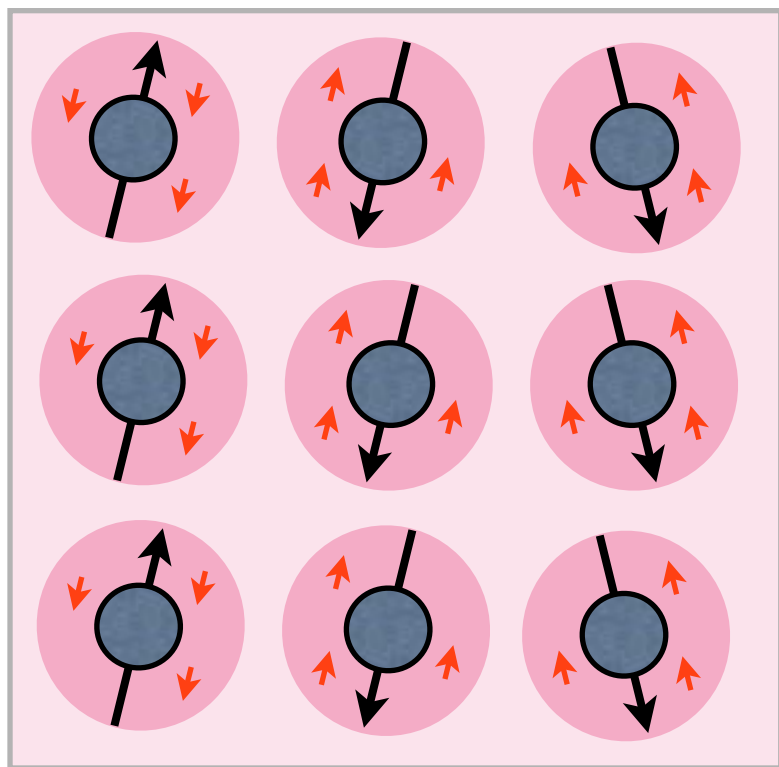
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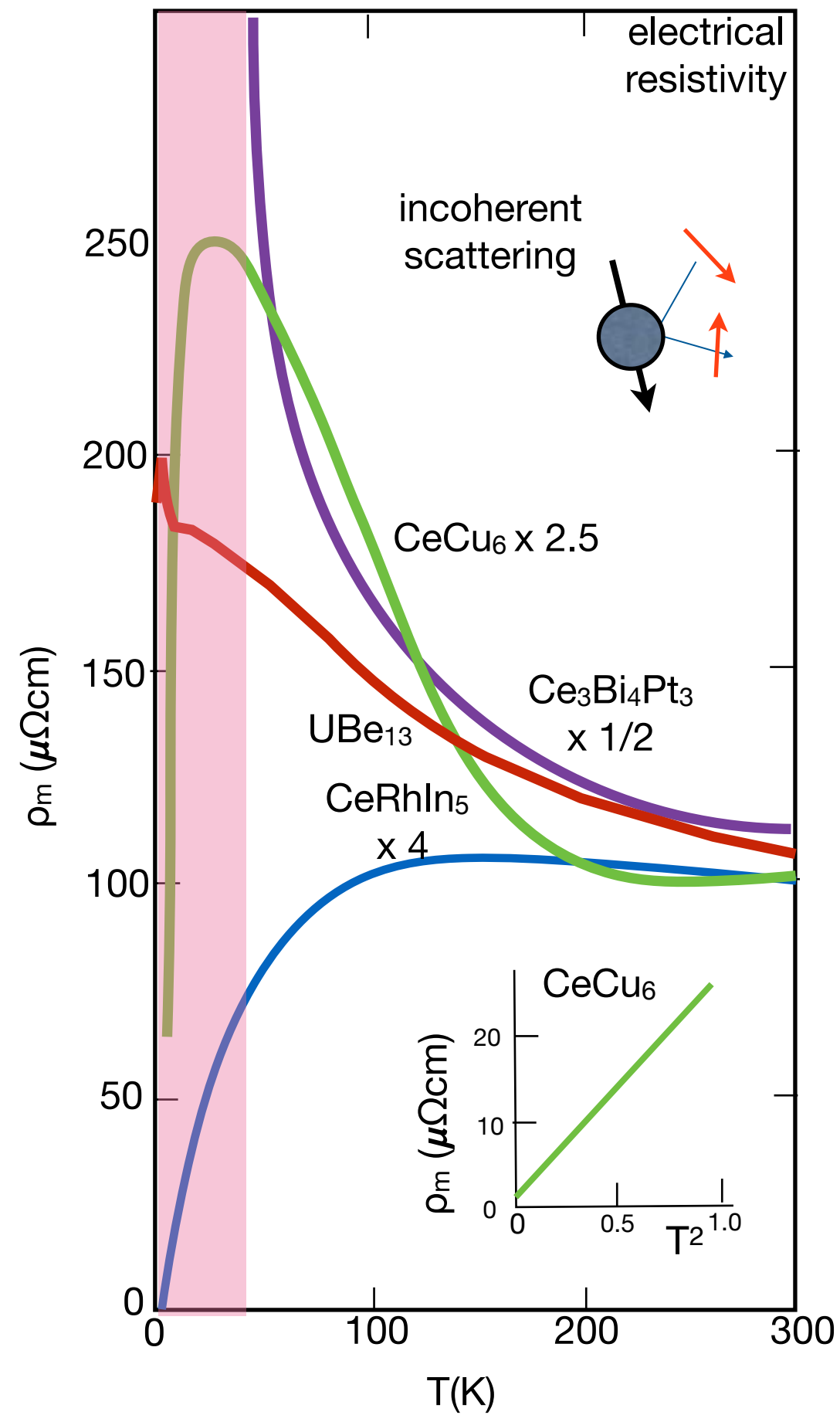
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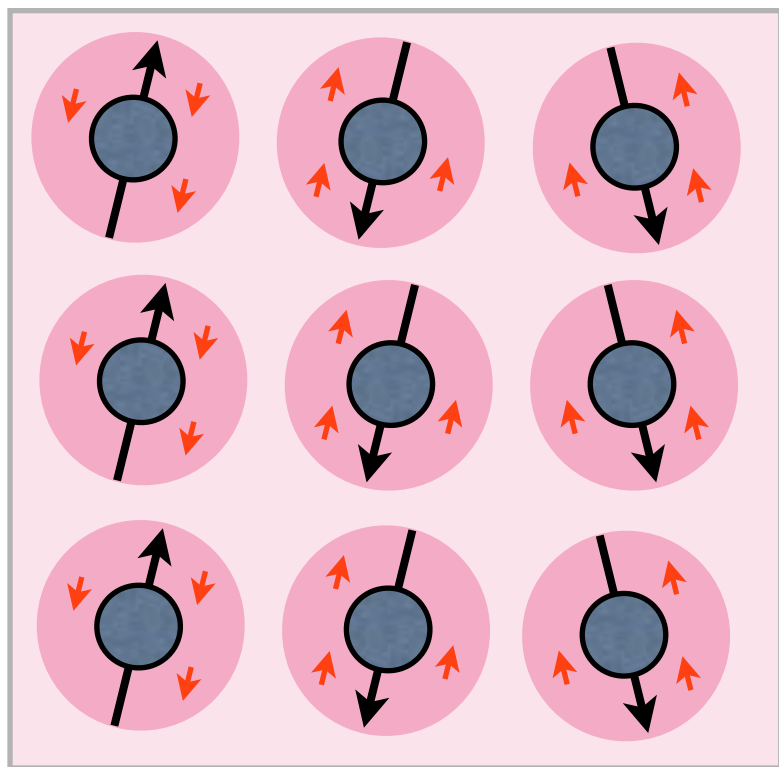
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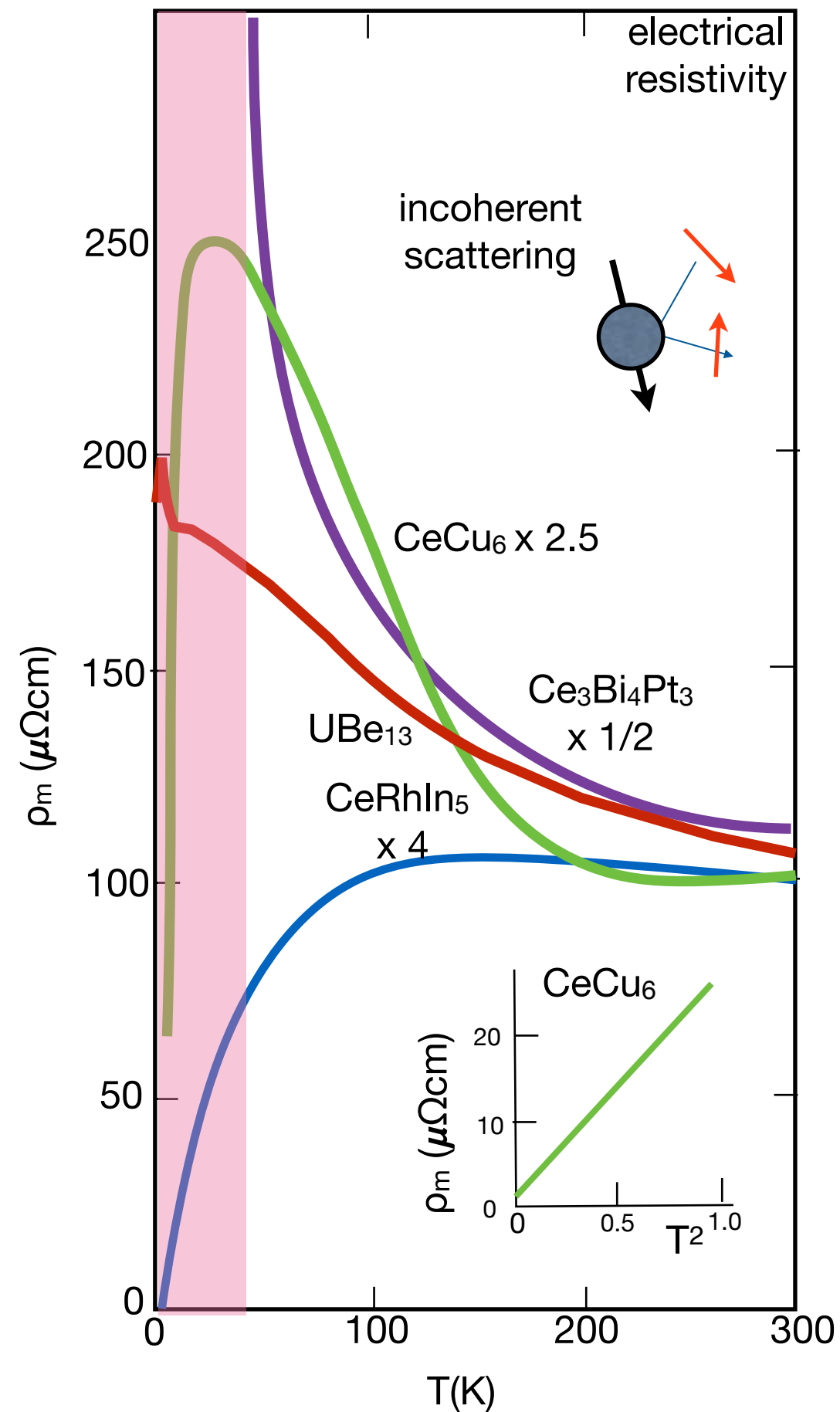
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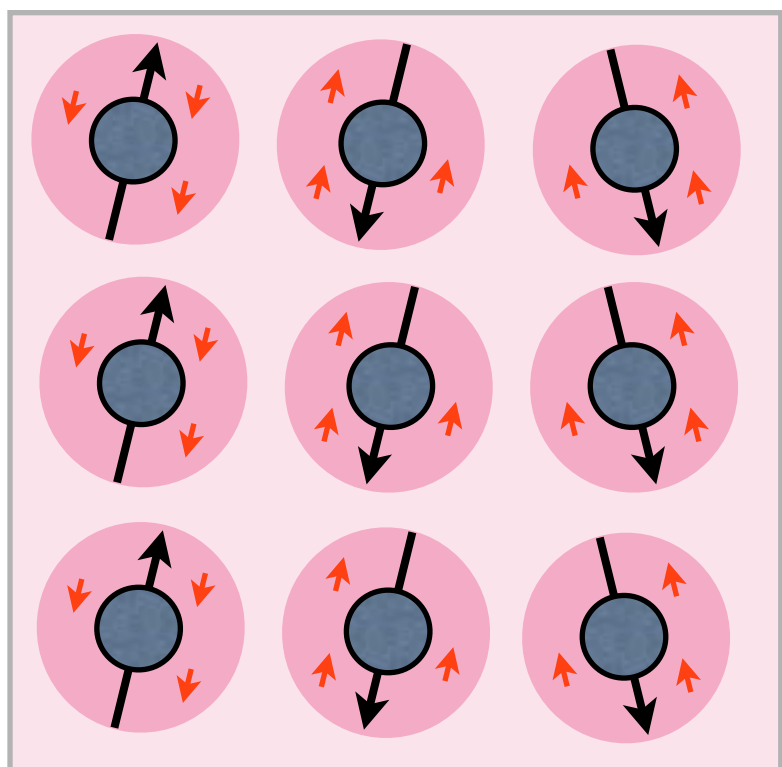
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Entangled spins and electrons

→ **Heavy Fermion Metals**



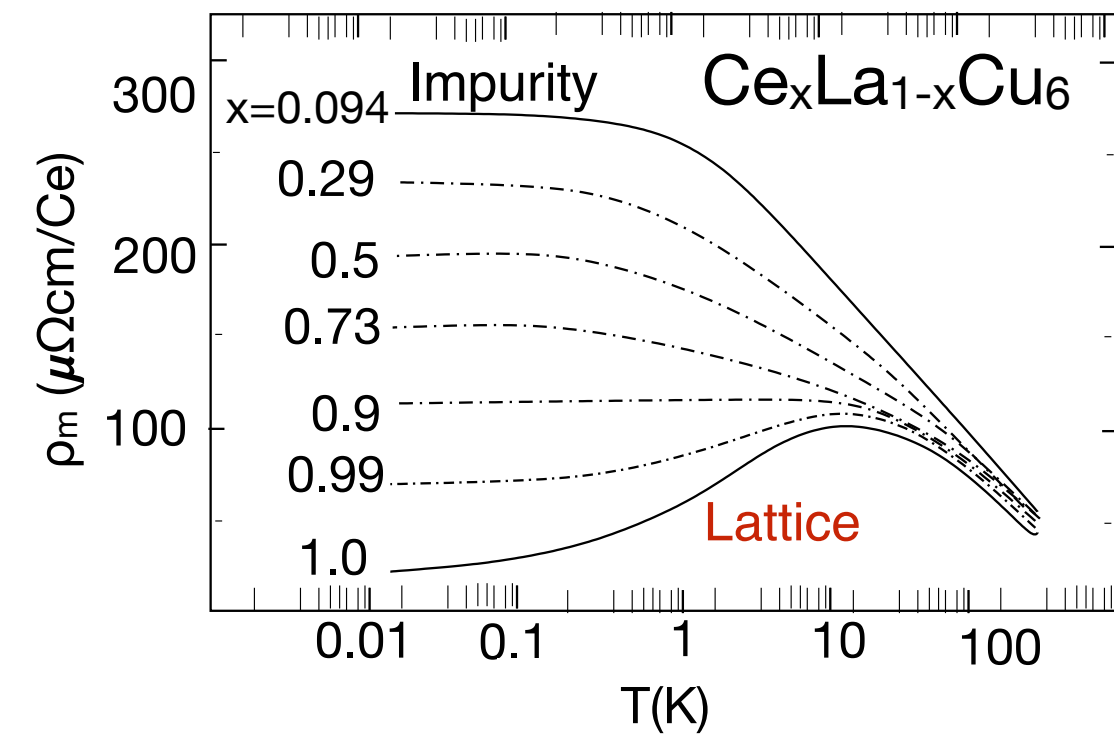
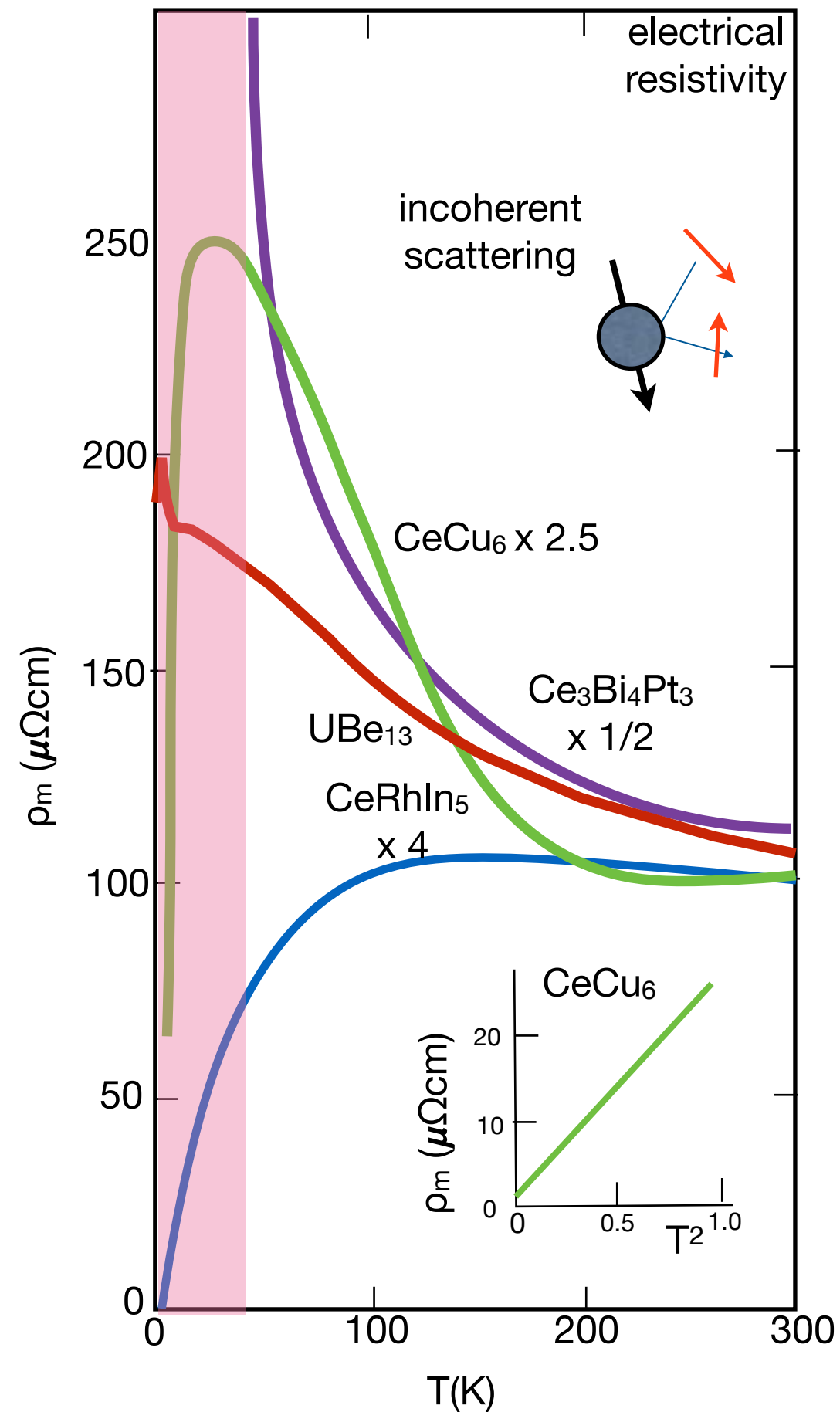
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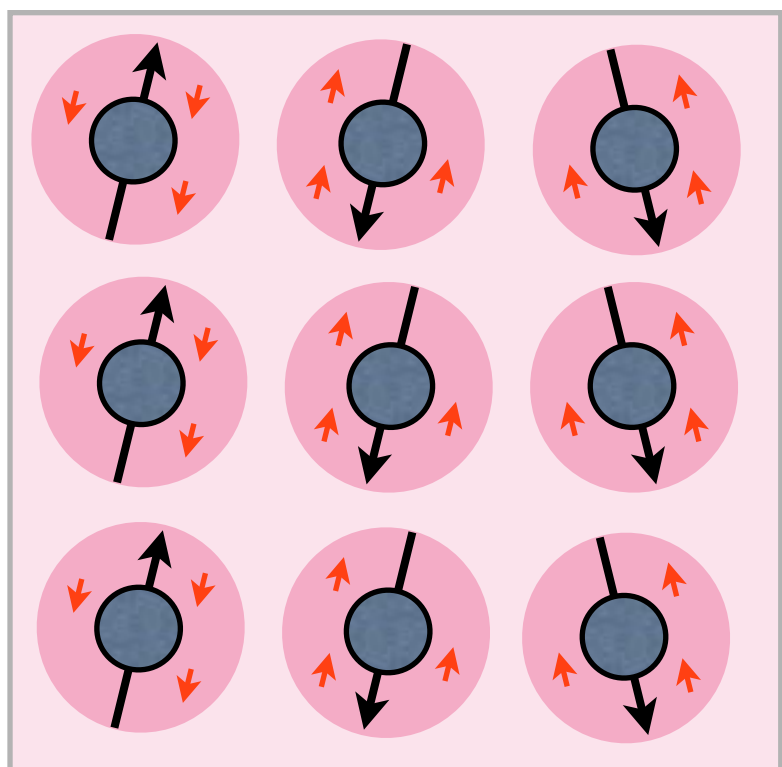
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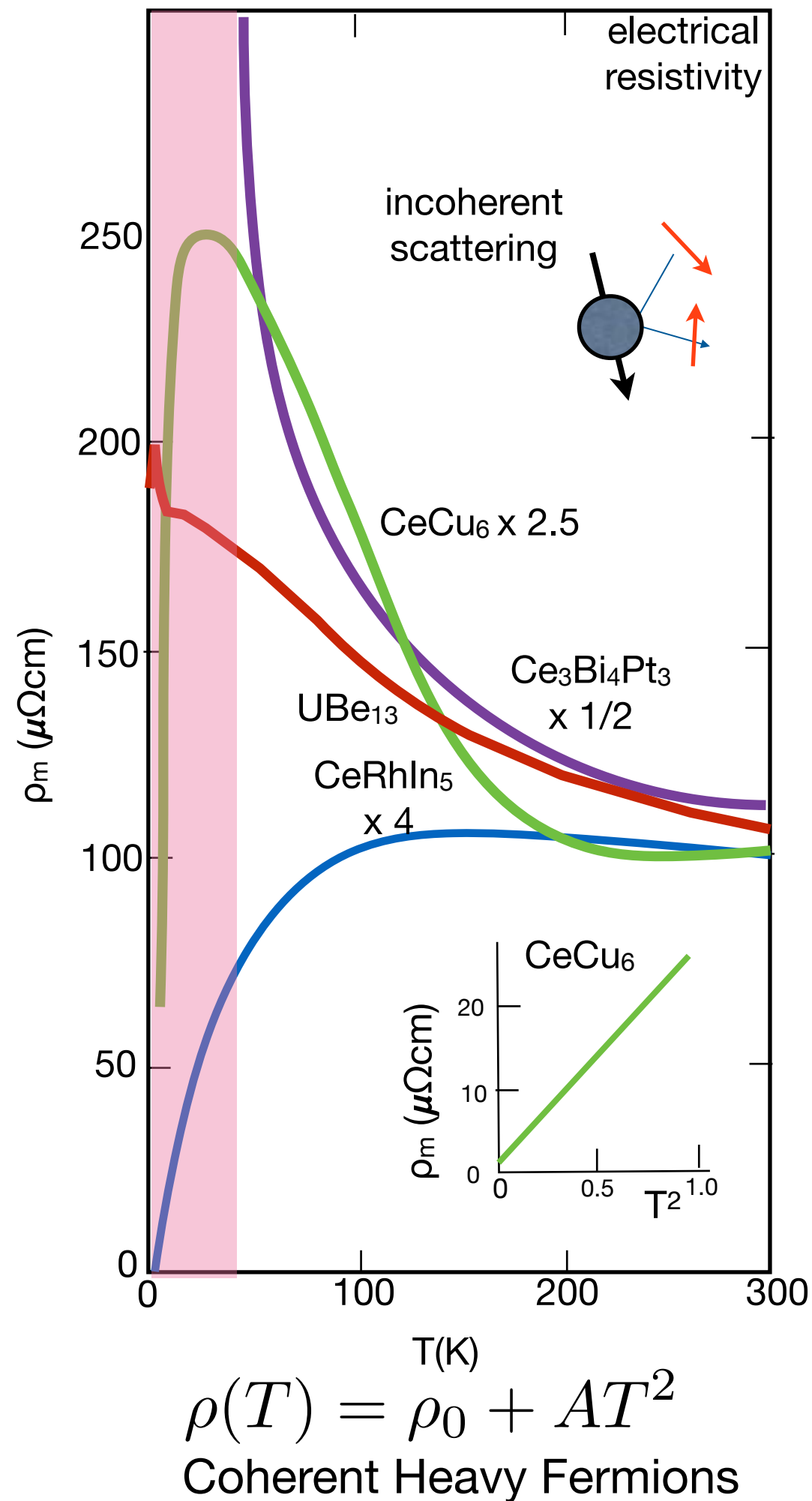
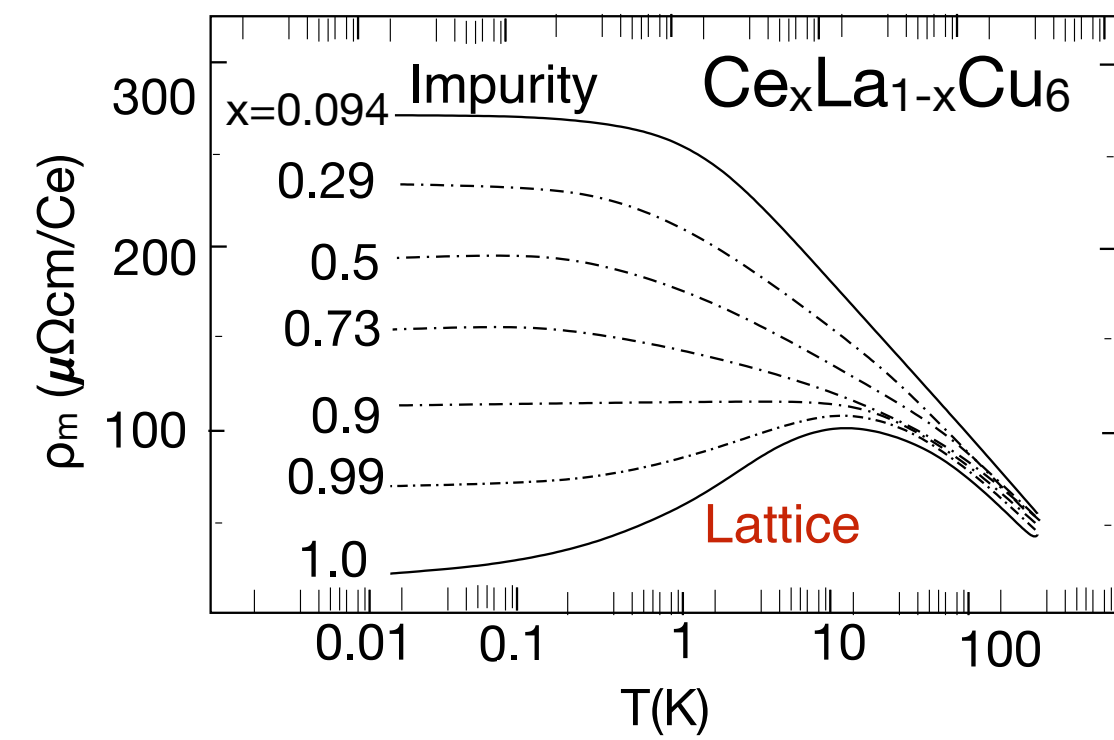
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“Kondo Lattice”

Entangled spins and electrons

→ **Heavy Fermion Metals**





DONIACH'S

Hypothesis.

Doniach (1977)

$$H = \sum_k \varepsilon_k c_{k\sigma}^\dagger c_{k\sigma} + J \sum_j (\psi_j^\dagger \vec{\sigma} \psi_j) \cdot \vec{S}_j$$

Kondo Lattice Model
(Kasuya, 1951)



DONIACH'S

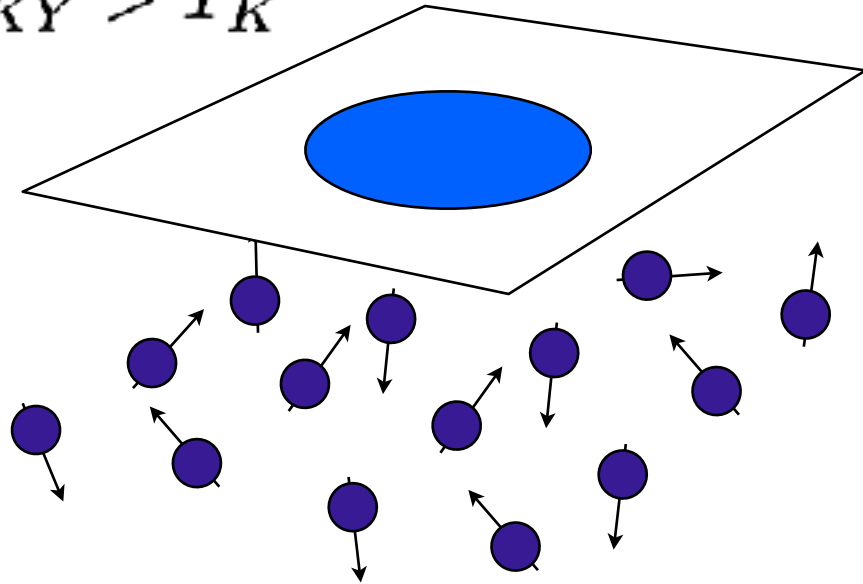
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Kondo Lattice Model
(Kasuya, 1951)

$T_{RKKY} > T_K$





DONIACH'S

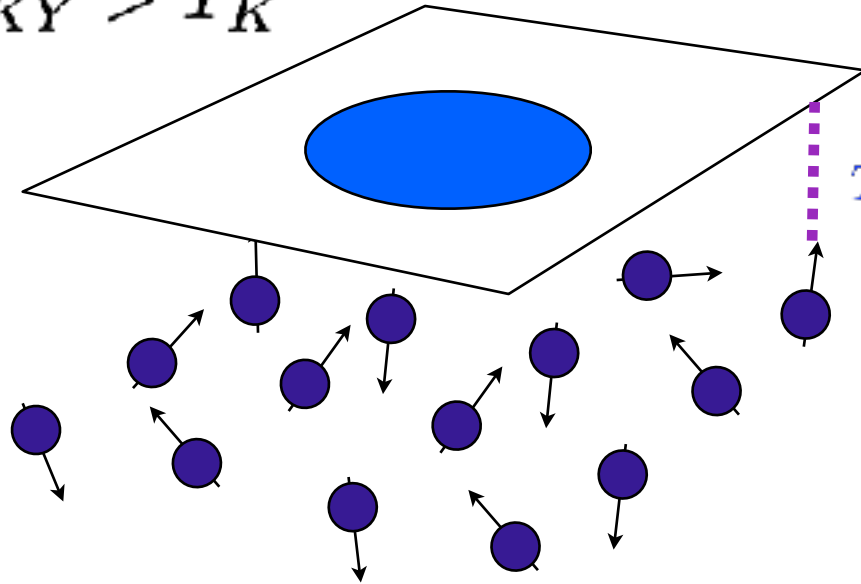
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Kondo Lattice Model
(Kasuya, 1951)

$$T_{RKKY} > T_K$$



$$T_K \sim D \exp \left[-\frac{1}{2J\rho} \right]$$



DONIACH'S

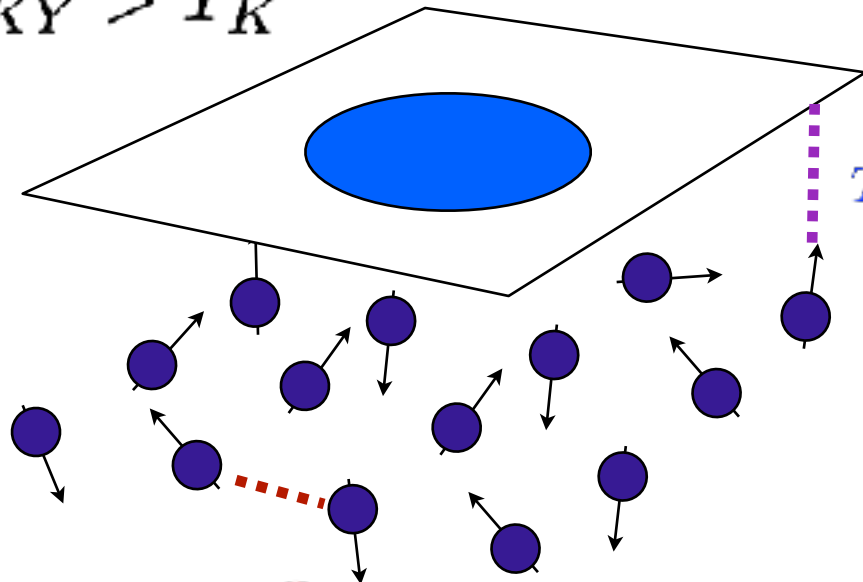
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DONIACH'S

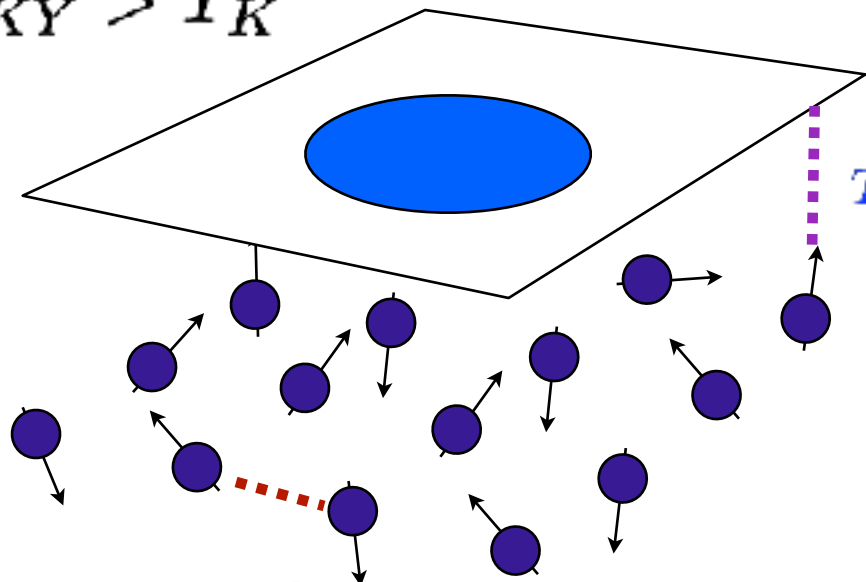
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DONIACH'S

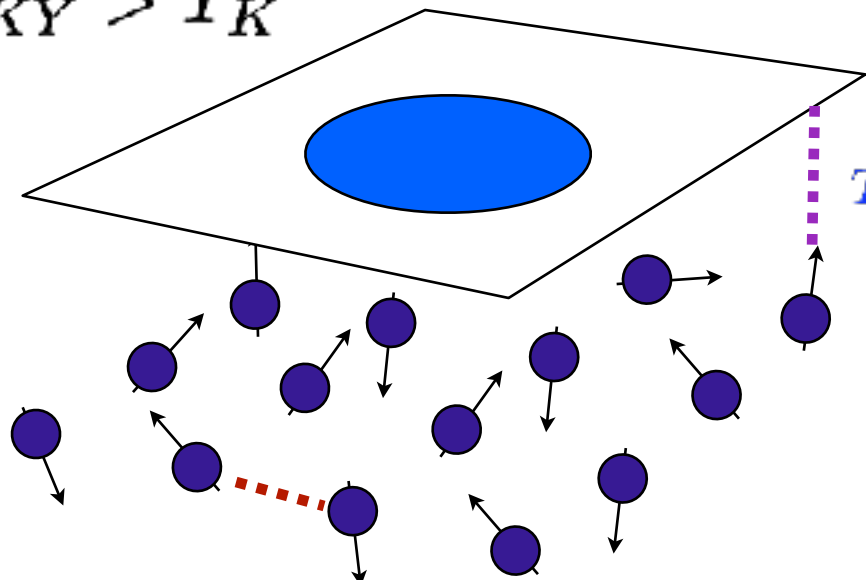
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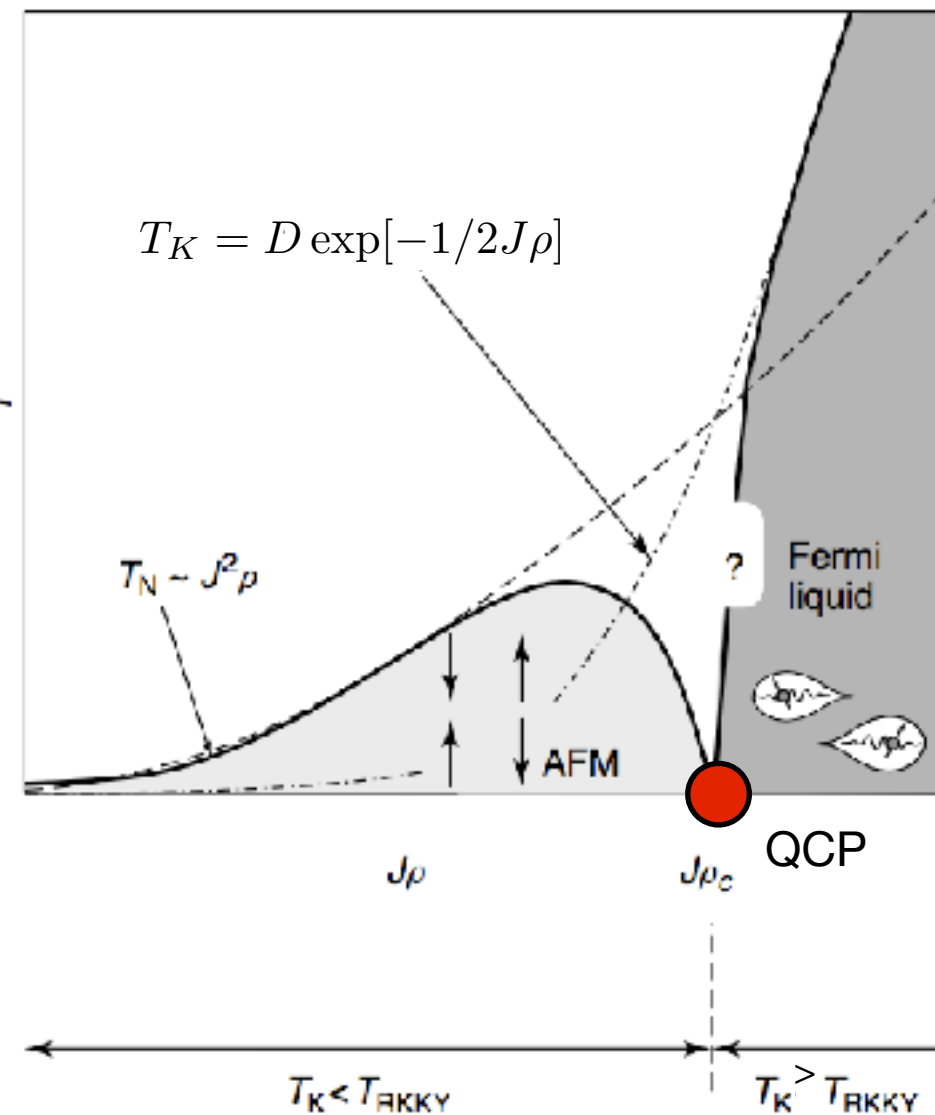
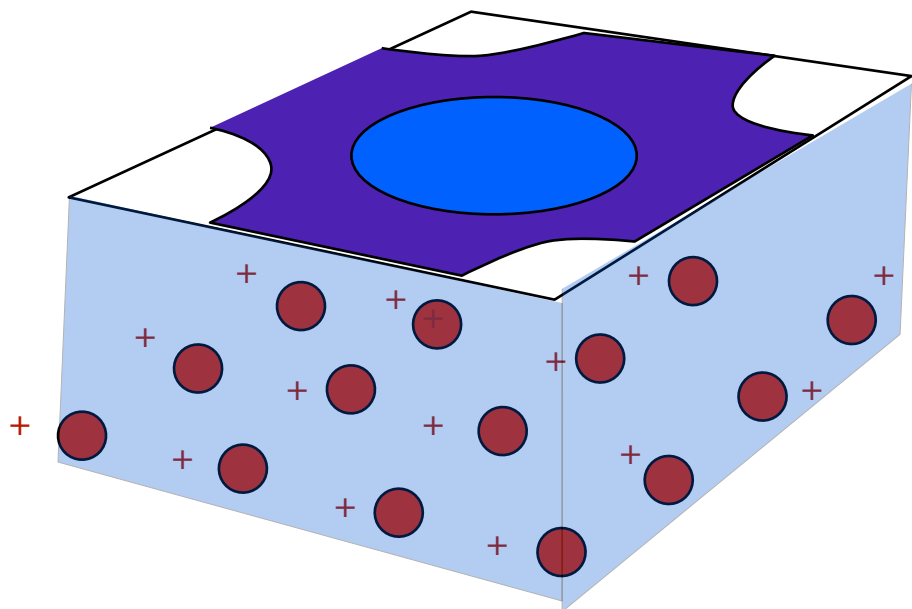


$$T_K \sim D \exp\left[-\frac{1}{2J\rho}\right]$$

$$T_{RKKY} \sim J^2 \rho$$

$$T_{RKKY} < T_K$$

Large Fermi surface of heavy Fermions





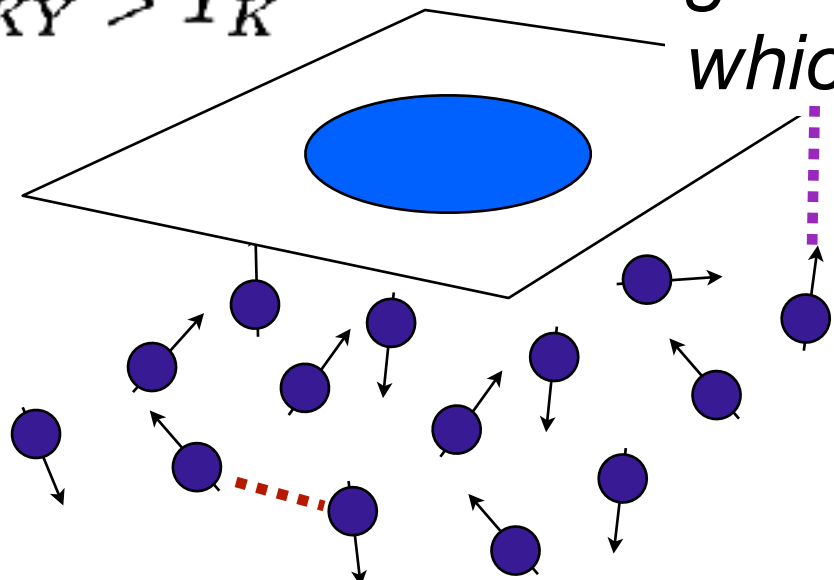
DONIACH'S Hypothesis.

Doniach (1977)

The main result ... is that there should be a second-order transition at zero temperature, as the exchange is varied, between an antiferromagnetic ground state for weak J and a Kondo-like state in which the local moments are quenched.

(Kusuya, 1971)

$$T_{RKKY} > T_K$$

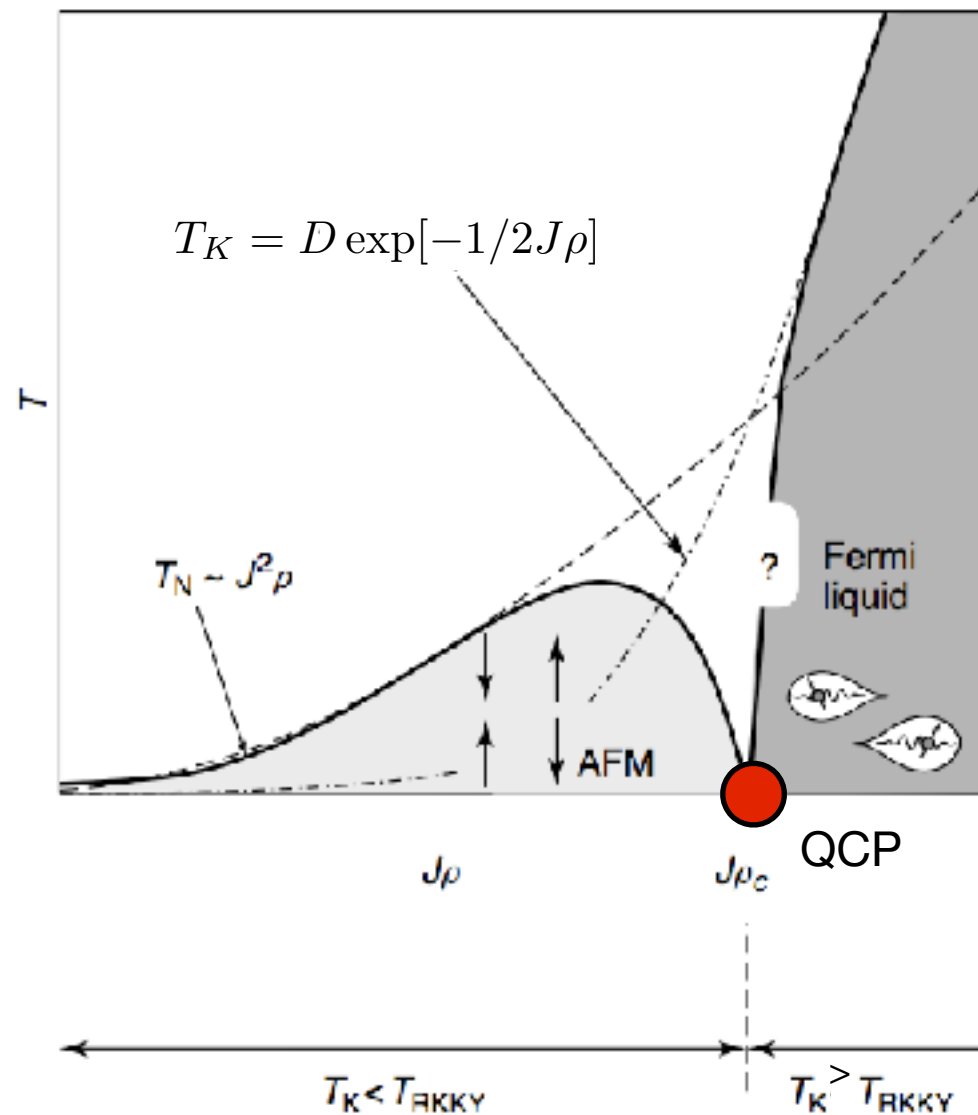
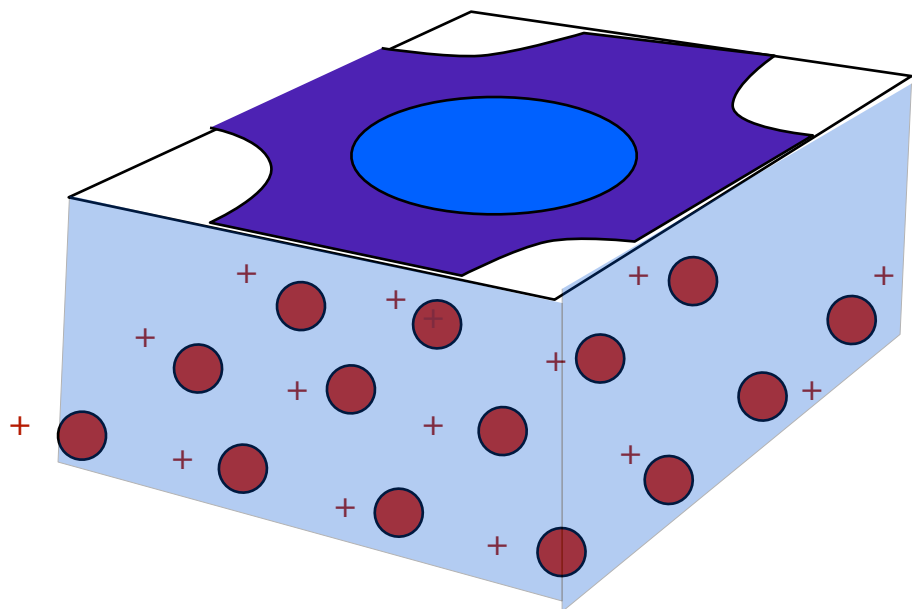


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Large Fermi surface of heavy Fermions





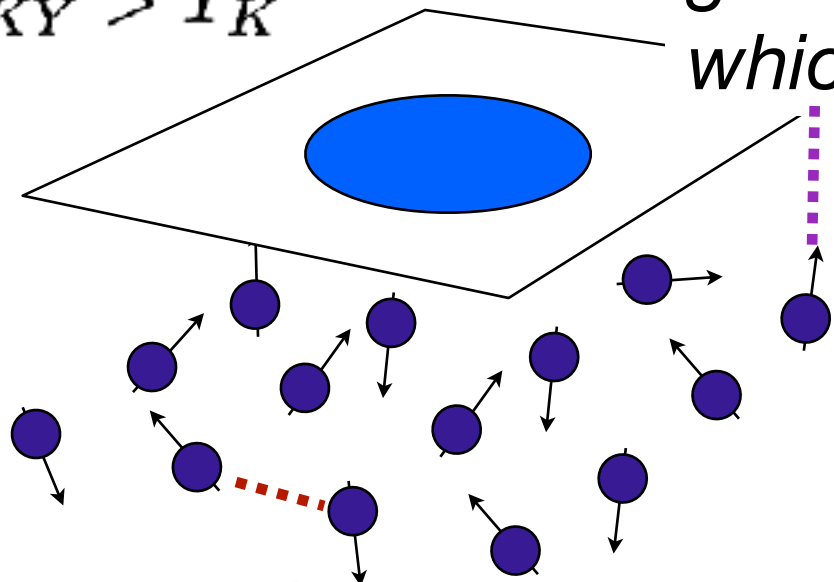
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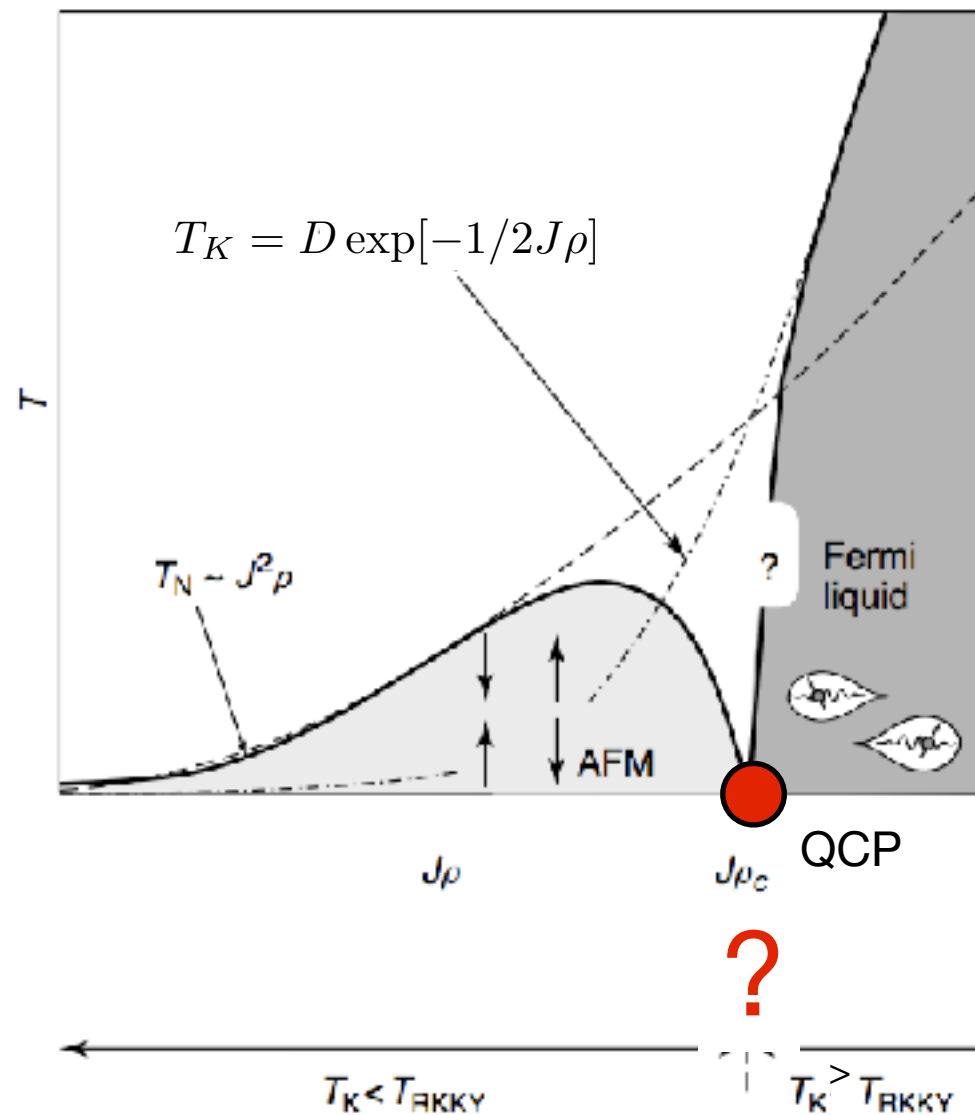
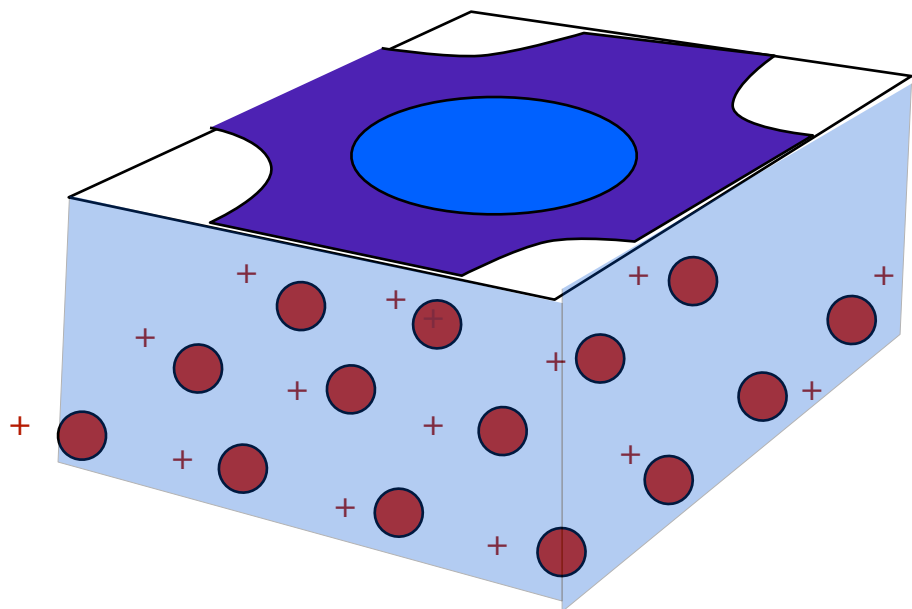


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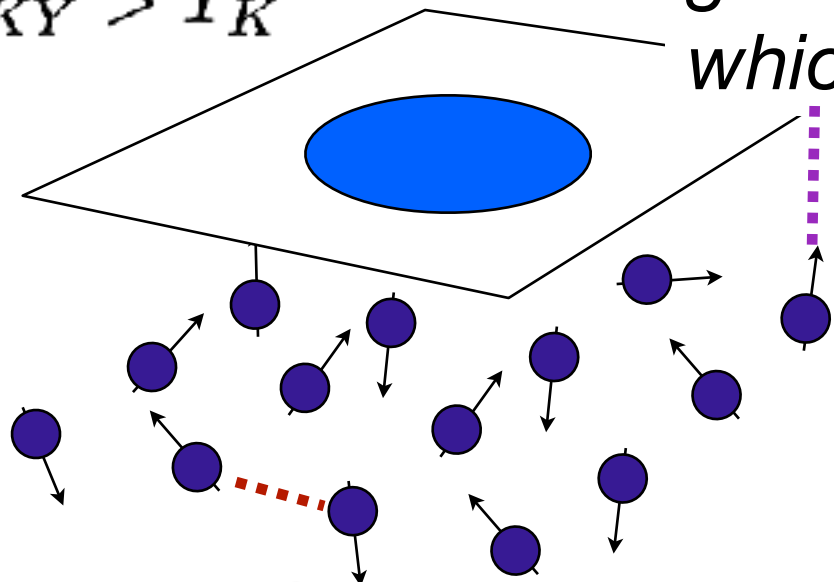
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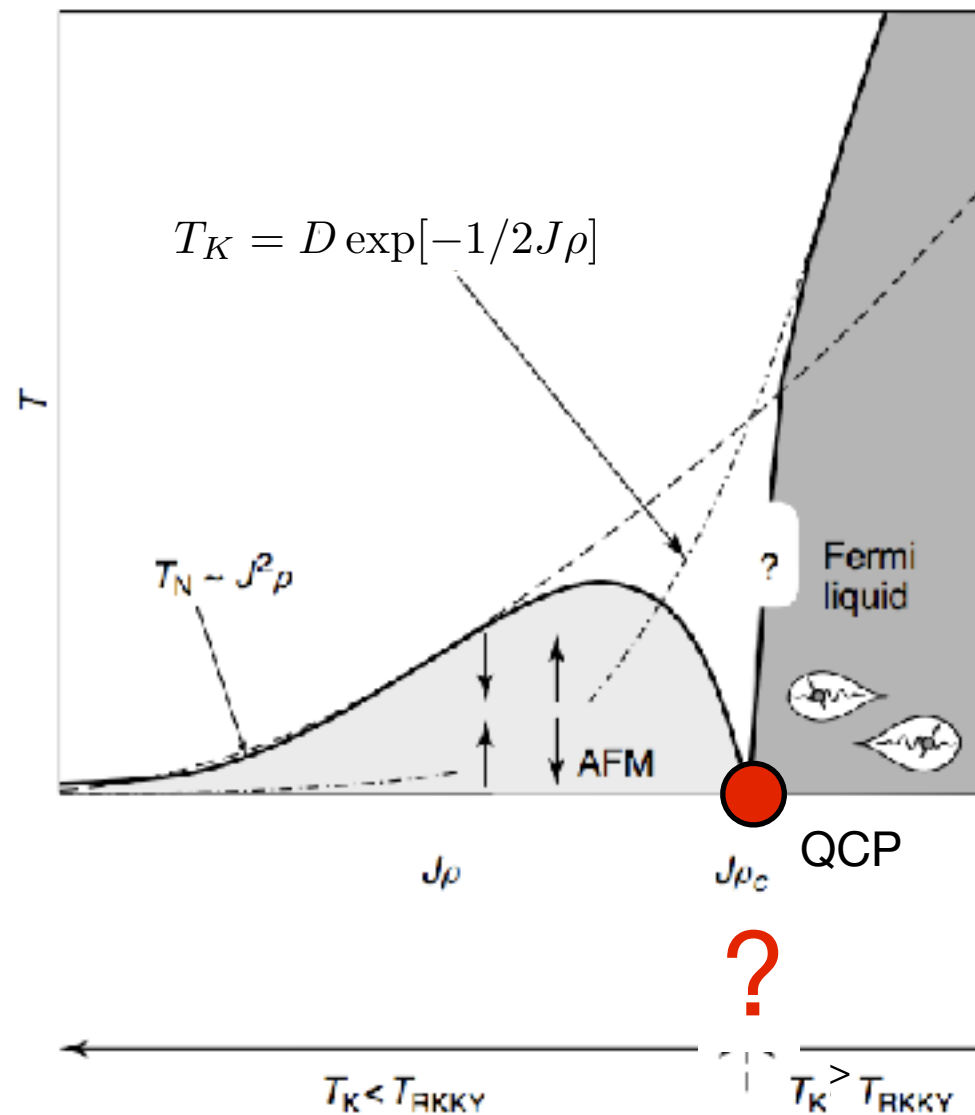
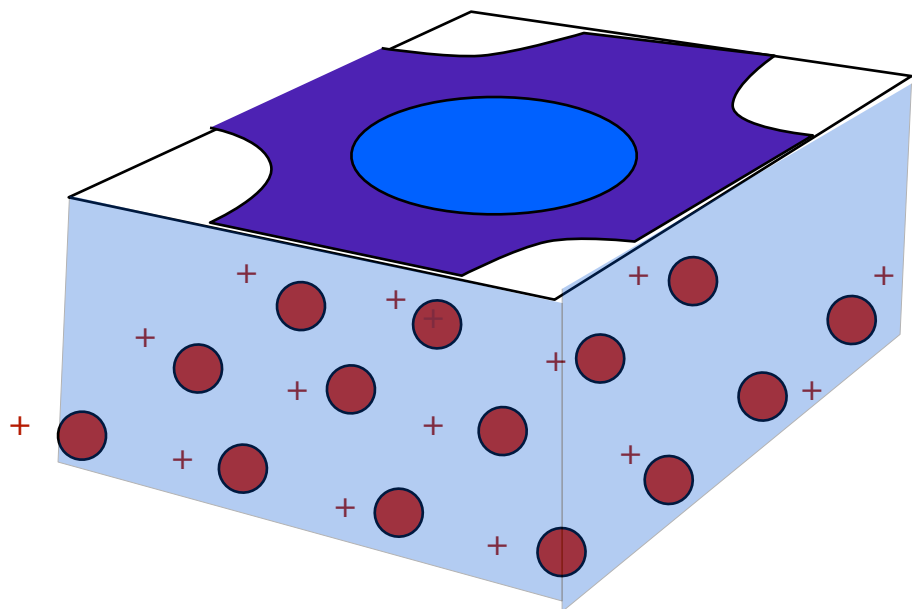


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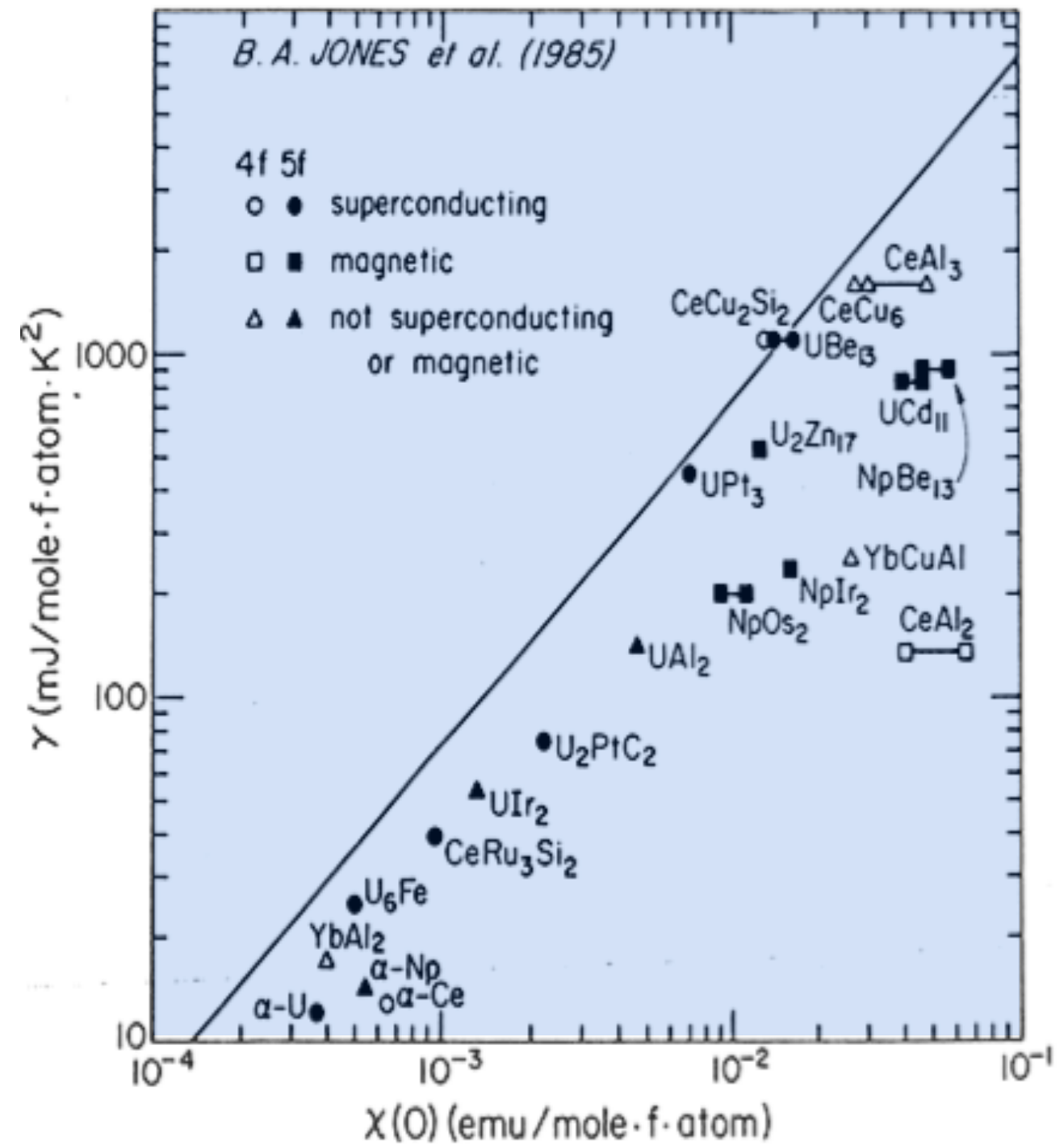
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FRACTIONALIZATION?

Heavy Fermions: magnetically polarizable Landau Fermi liquids.

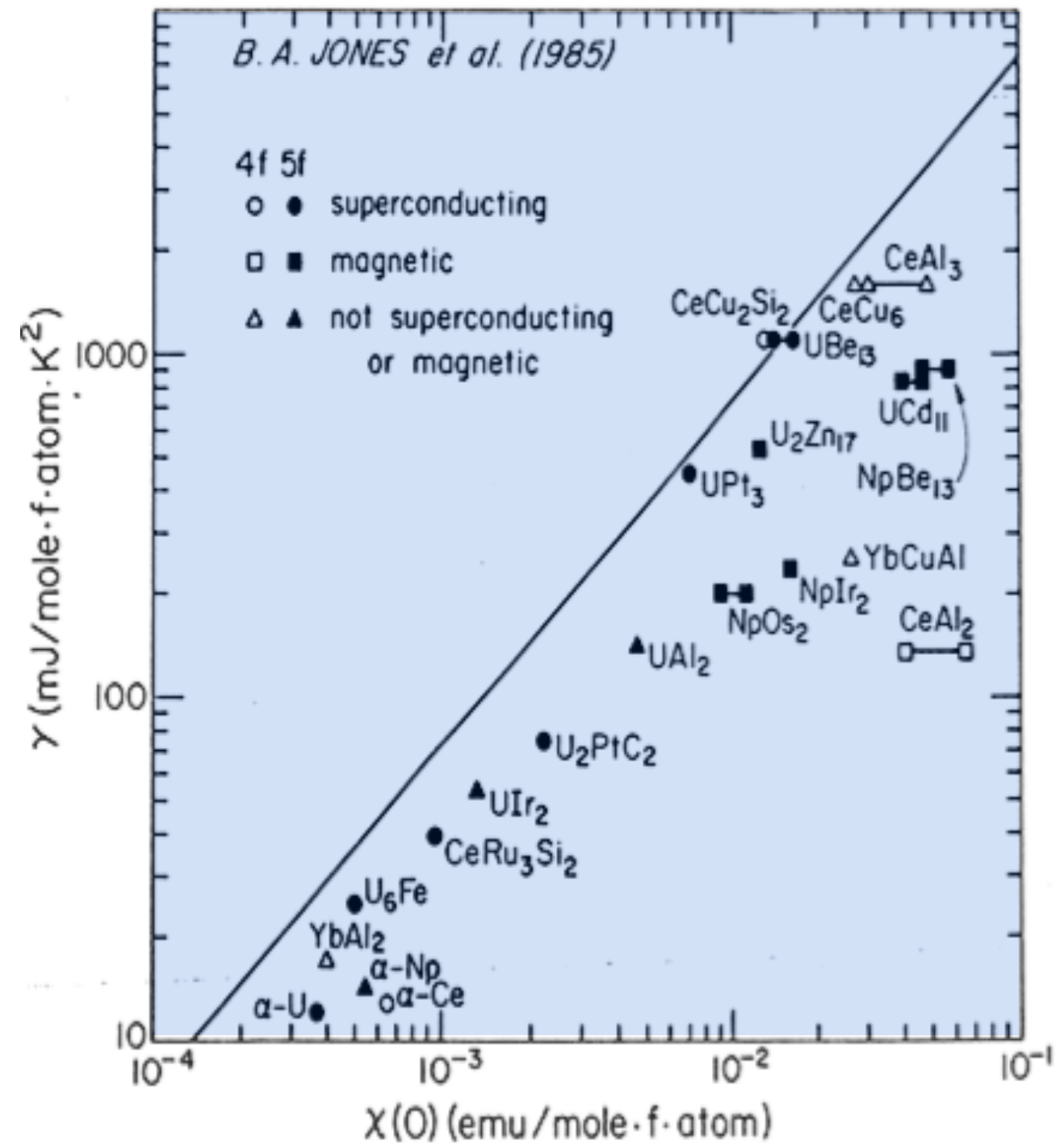
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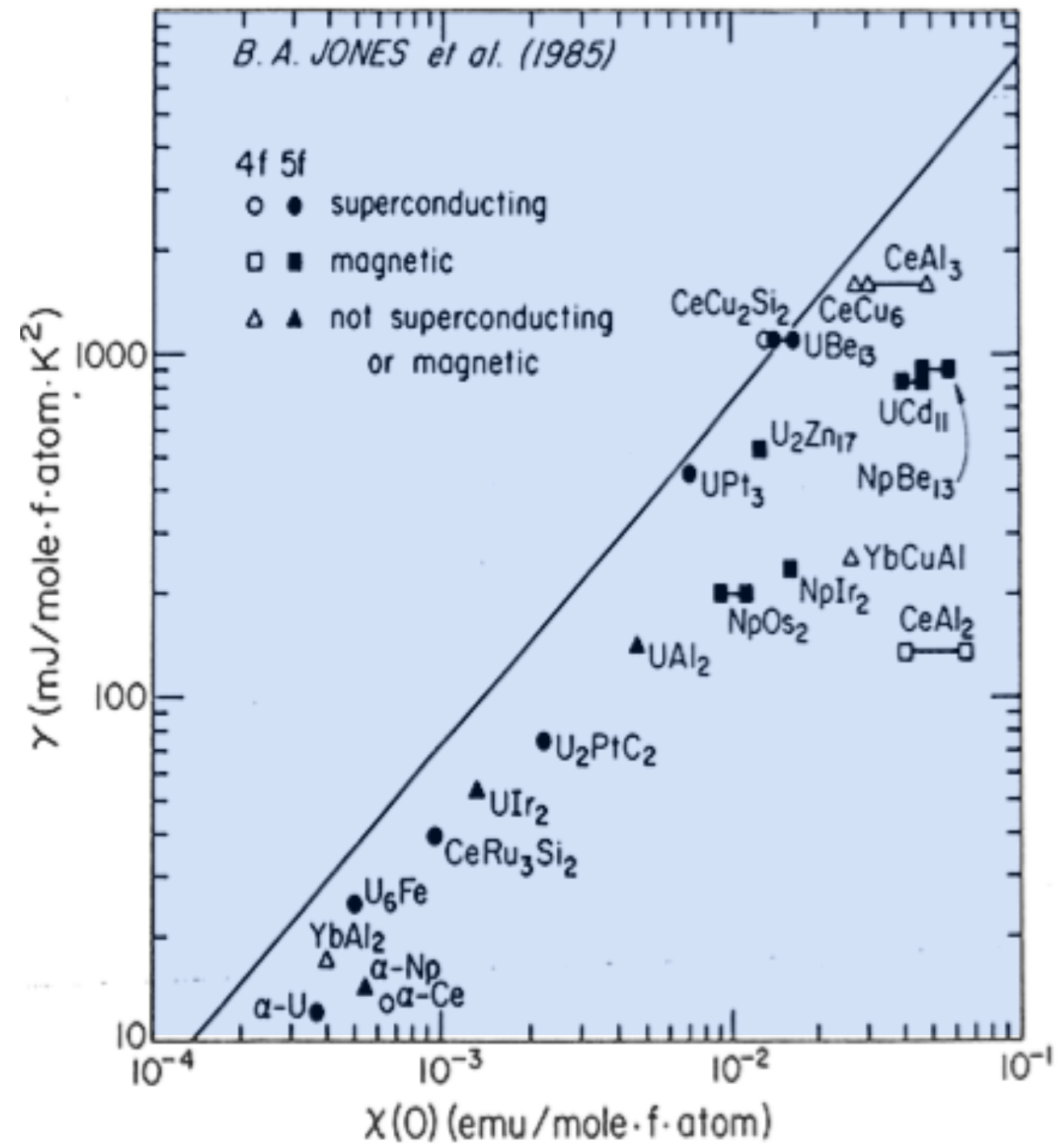


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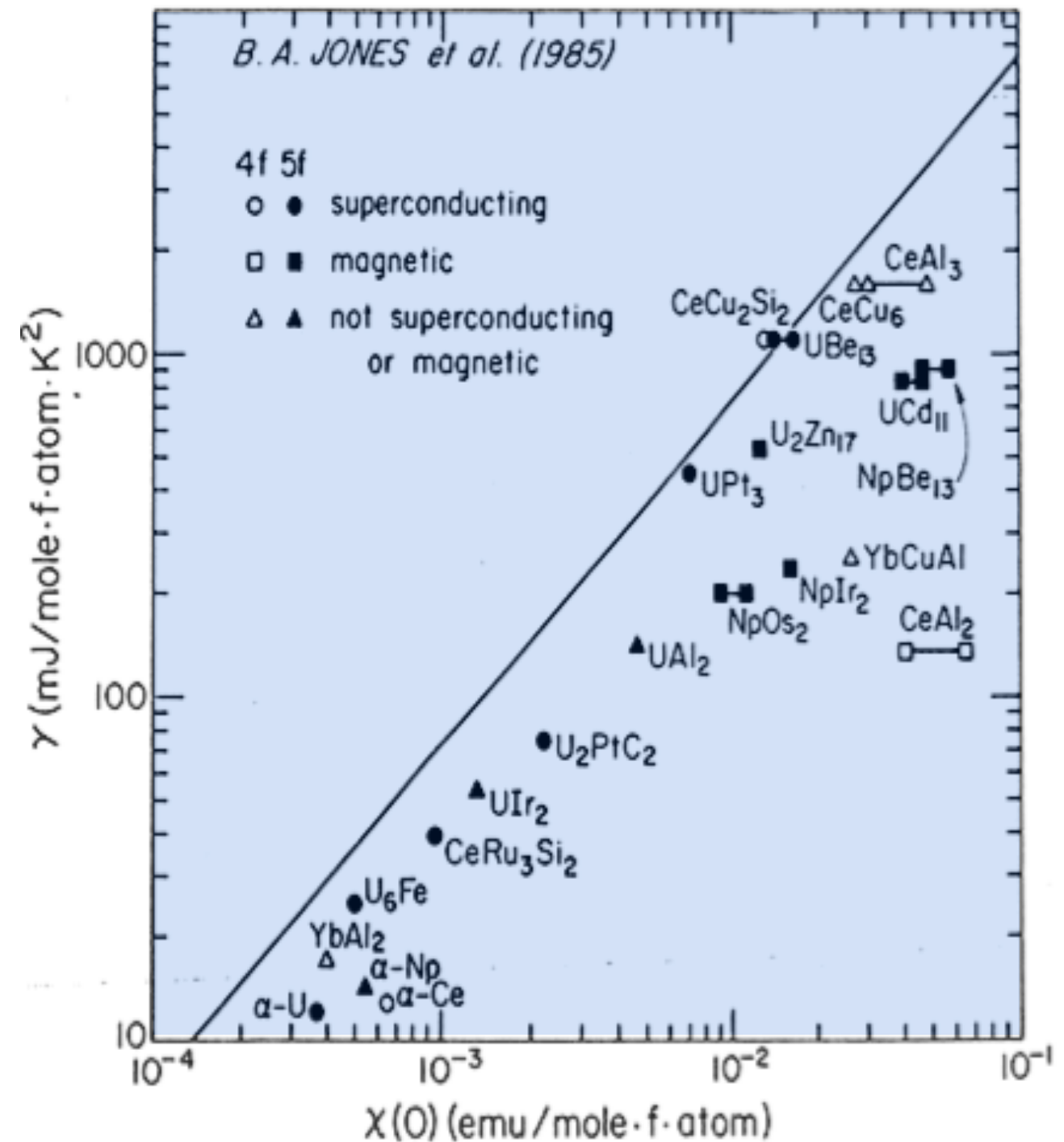
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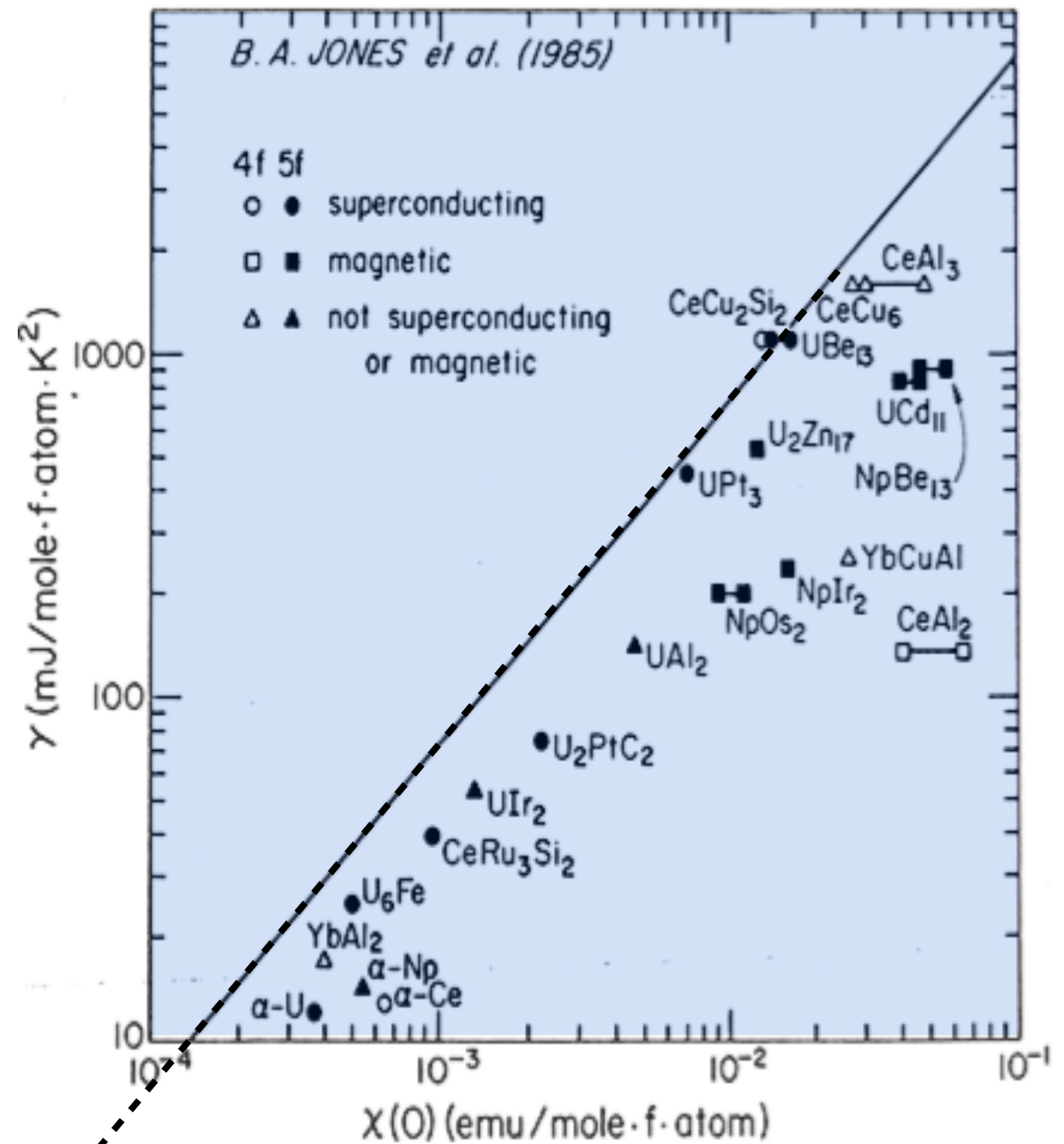
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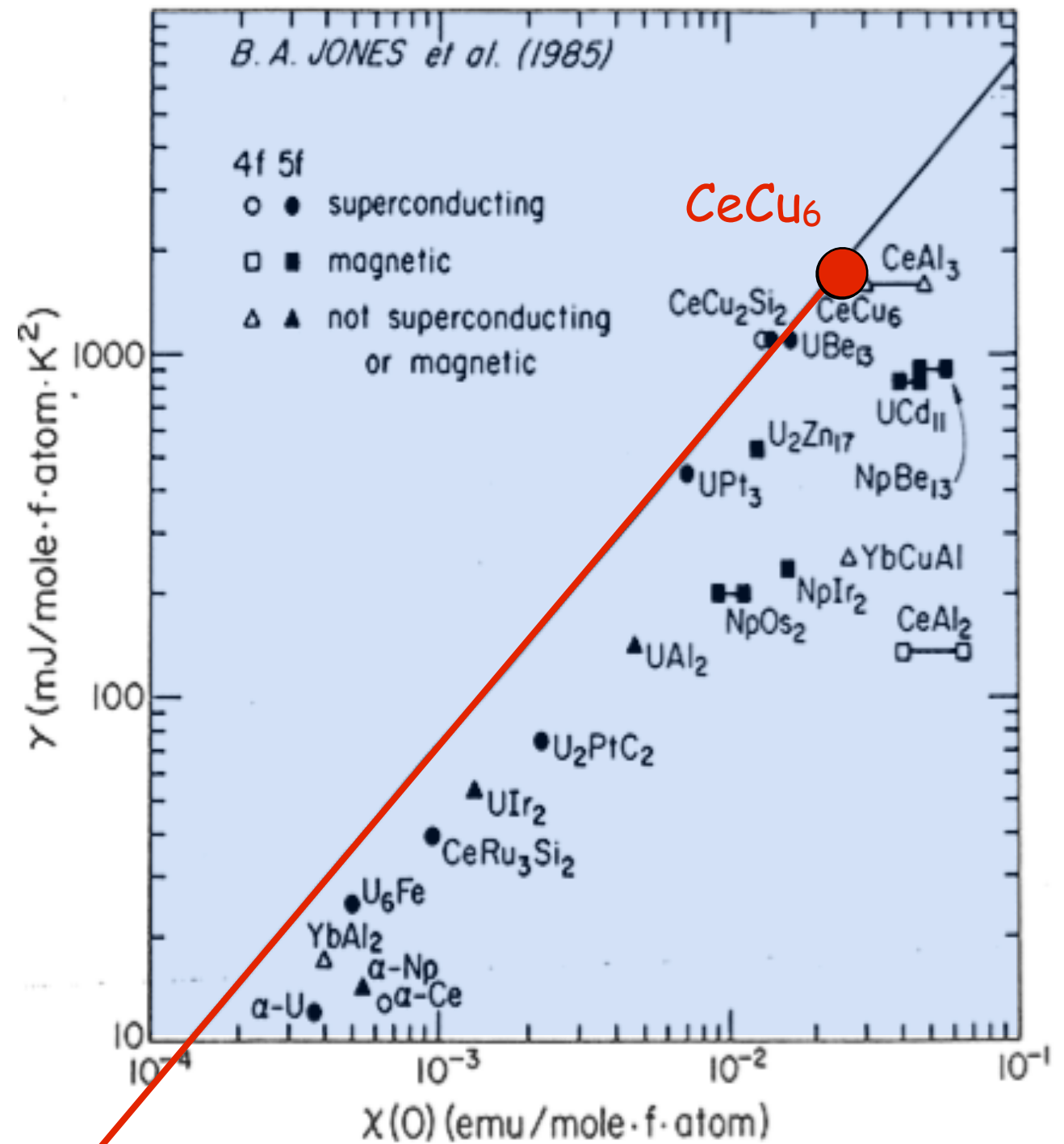
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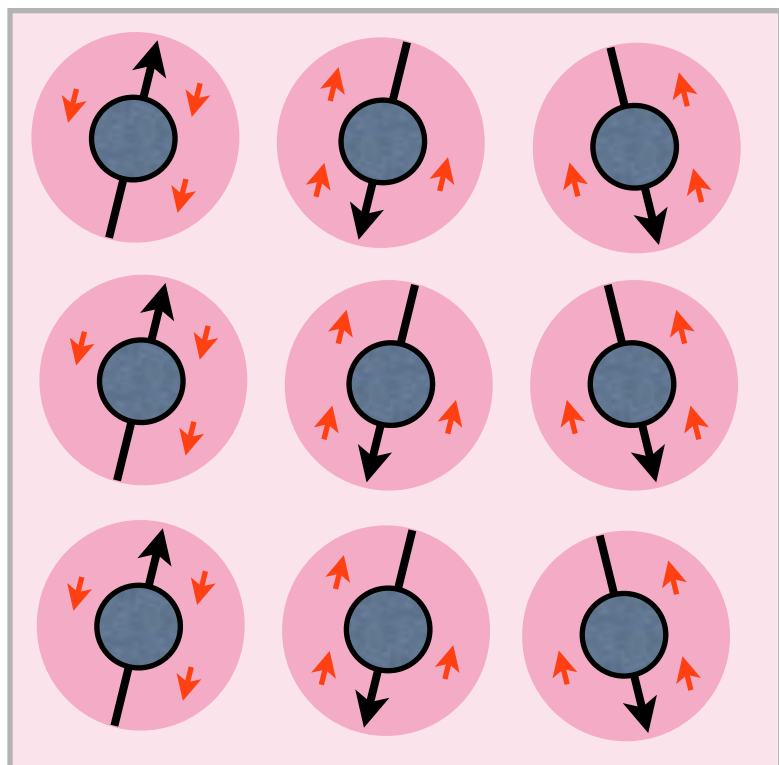
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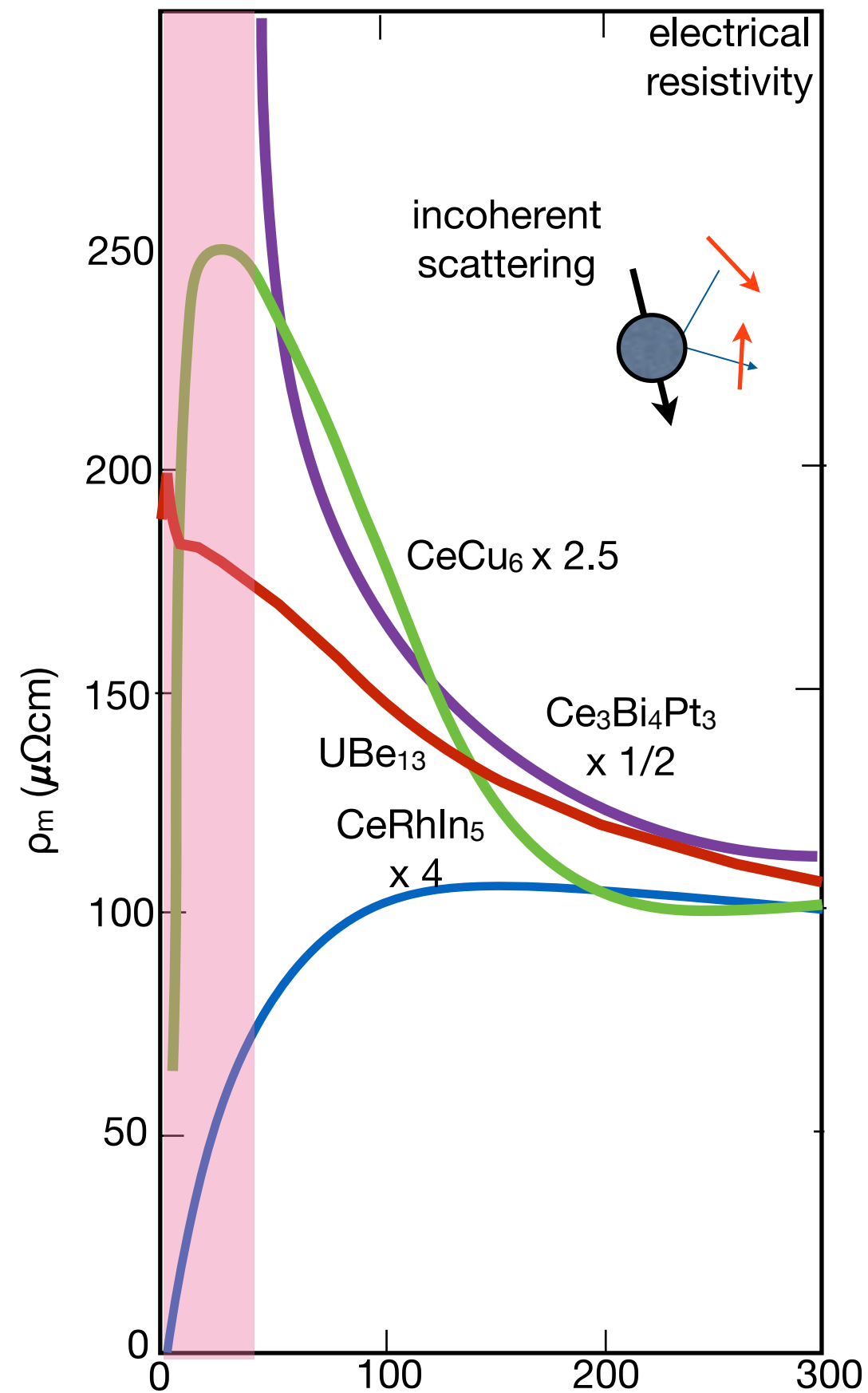


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“Kondo Lattice”

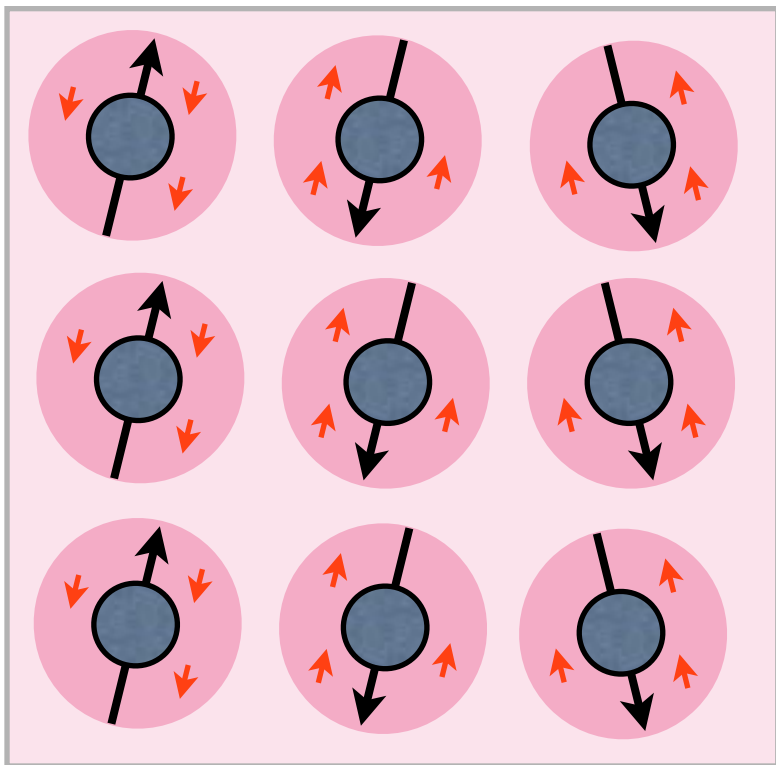
Entangled spins and electrons
 → **Heavy Fermion Metals**



$$\rho(T) = \rho_0 + AT^2$$

Coherent Heavy Fermions

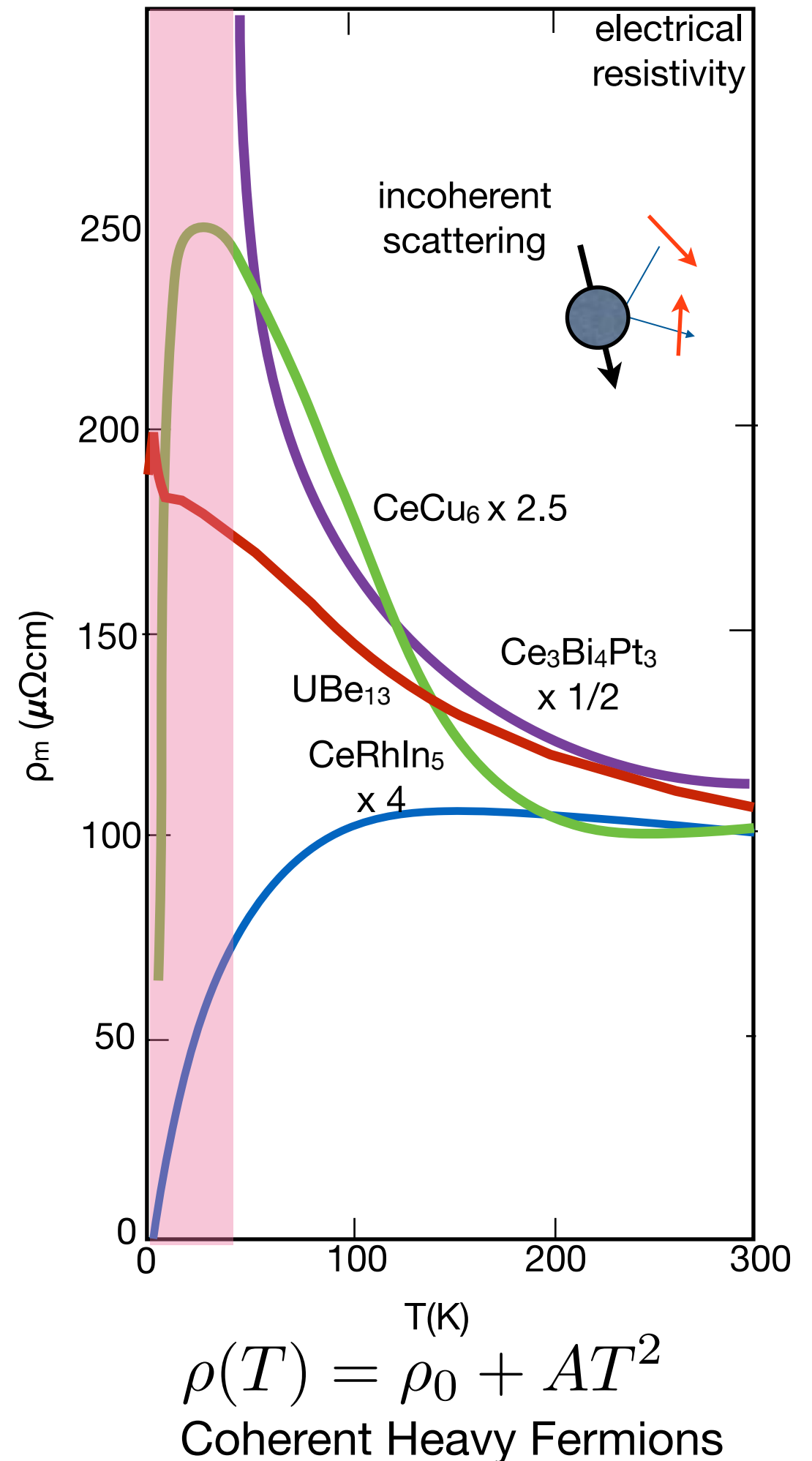
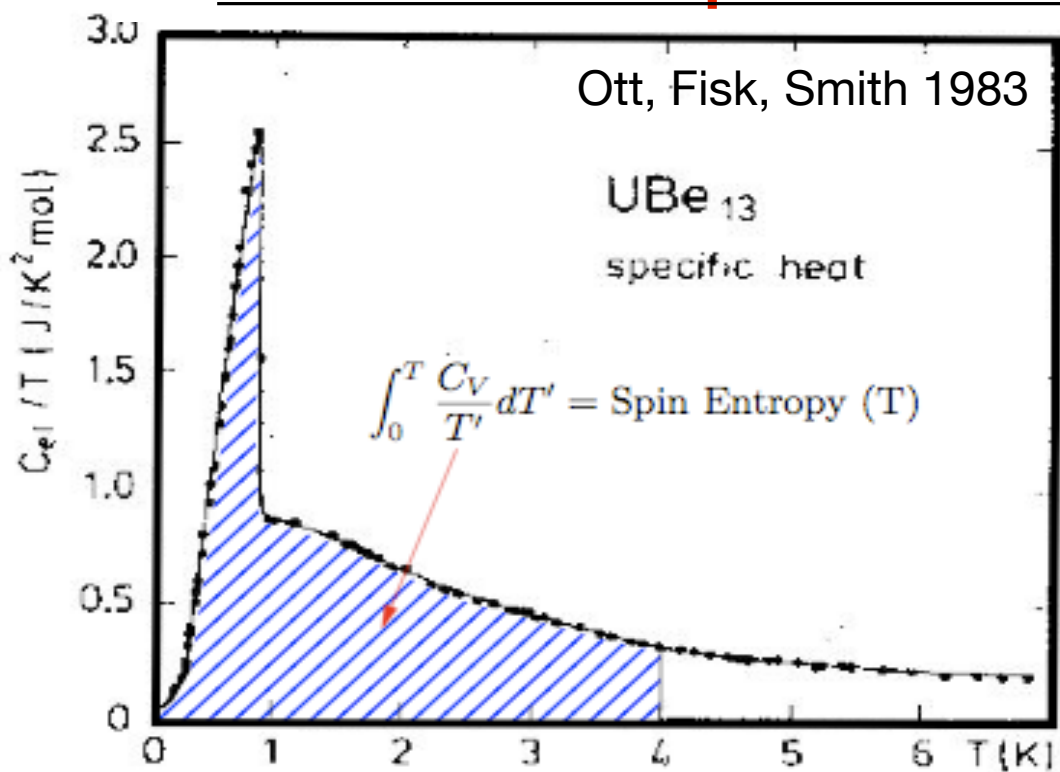
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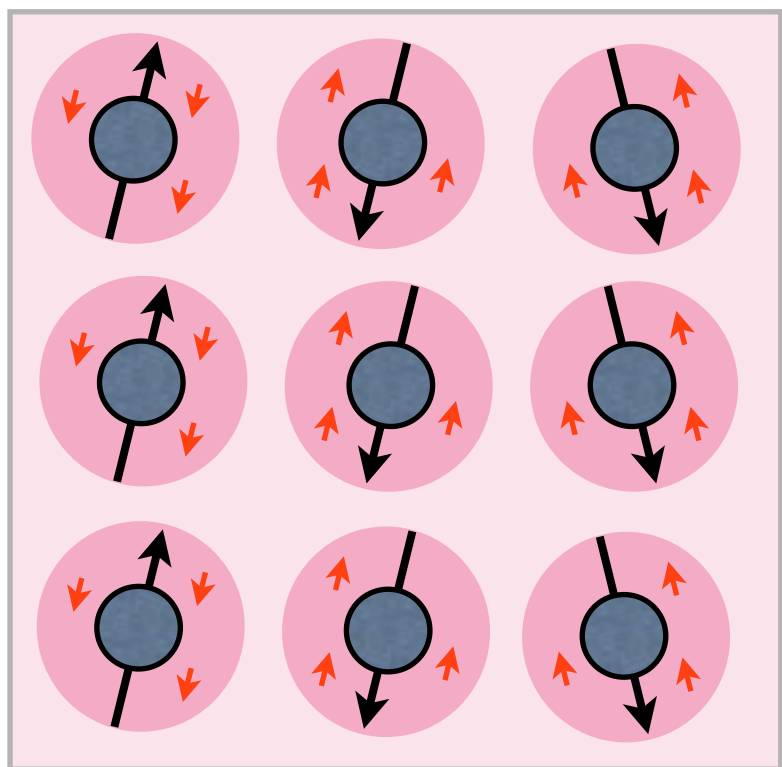
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Entangled spins and electrons

→ **New kinds of superconductor**

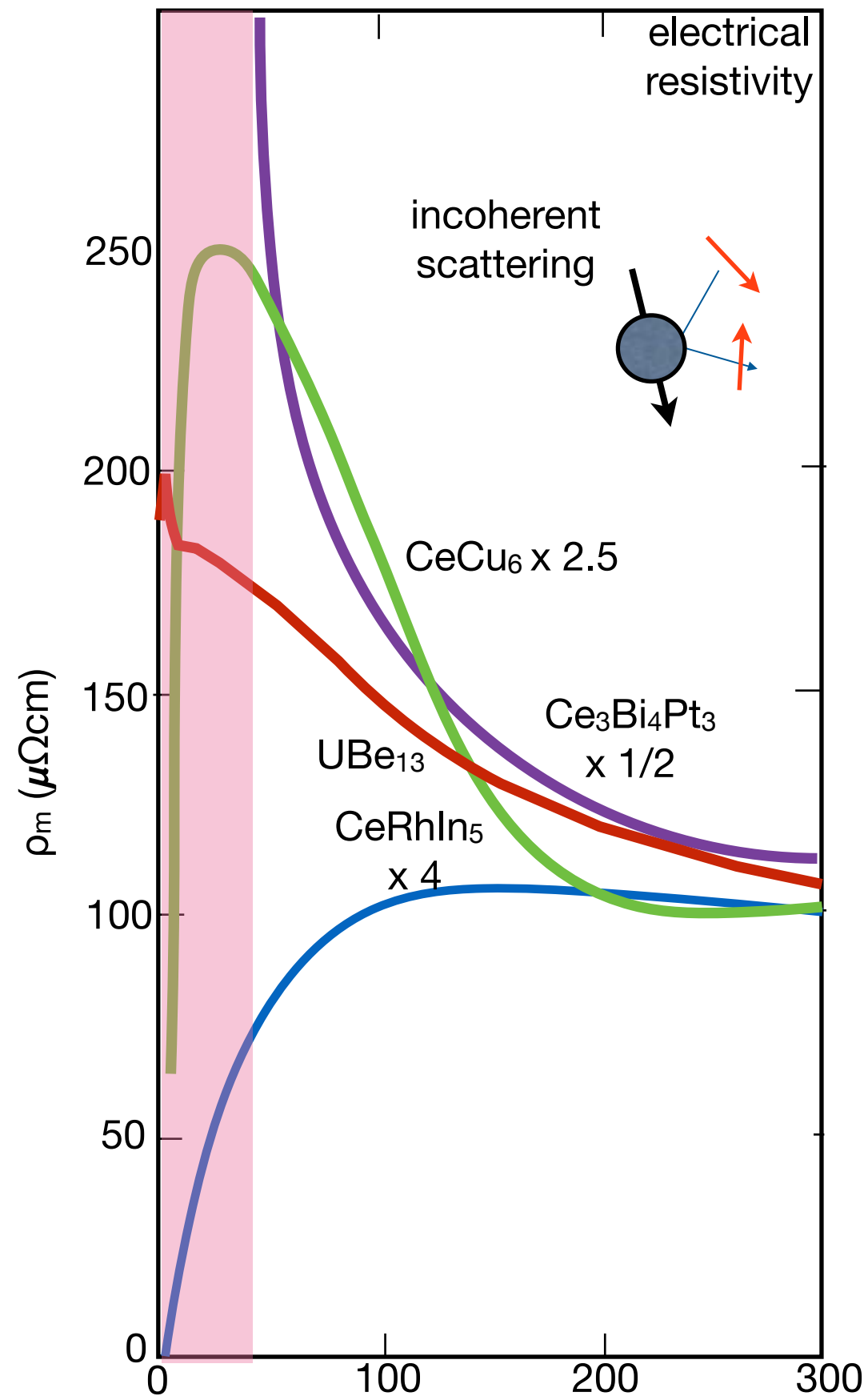


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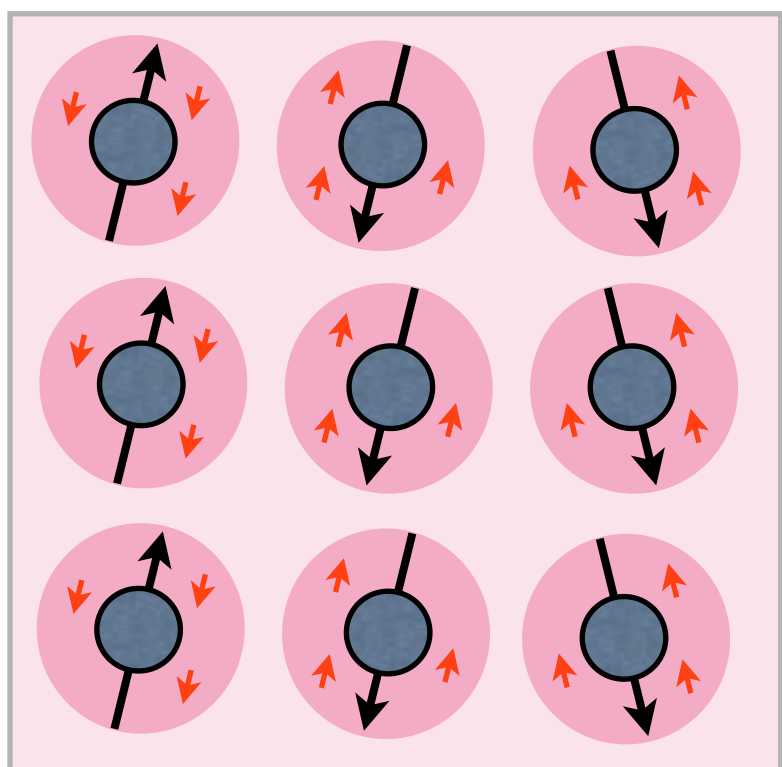
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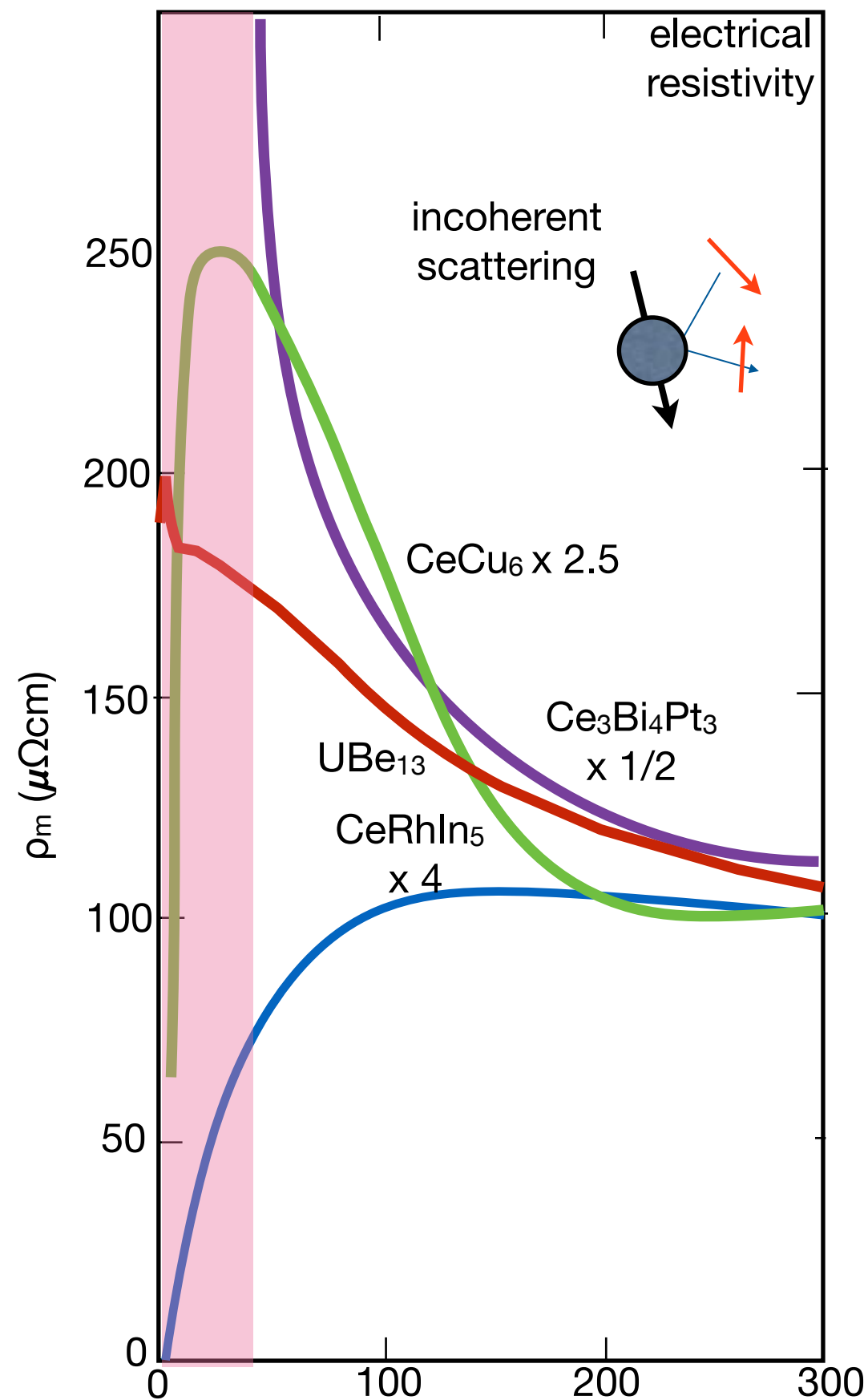
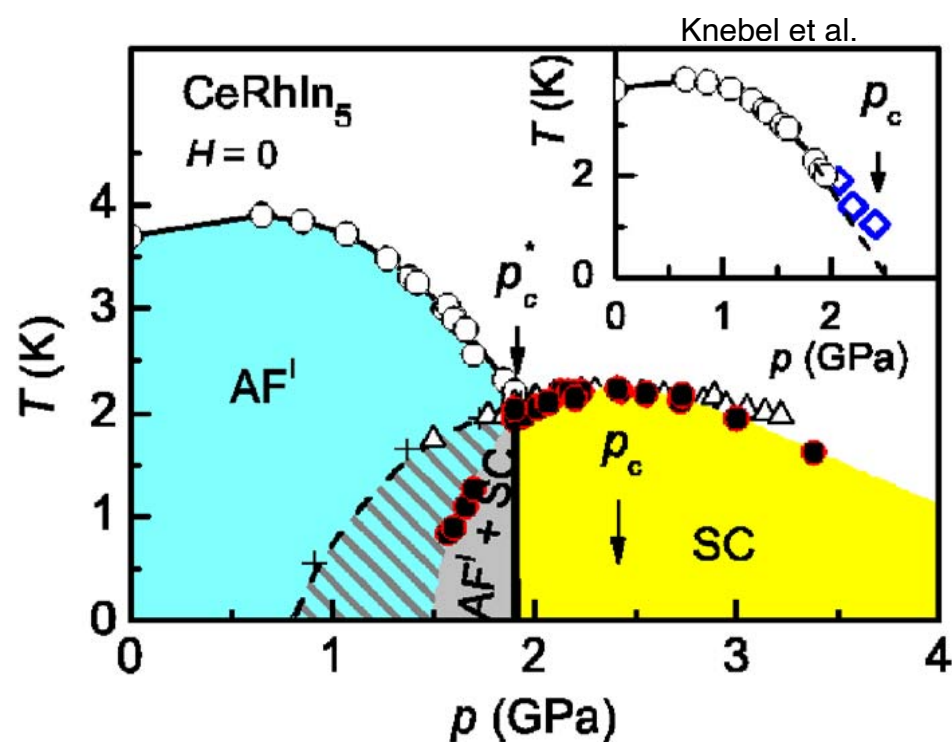
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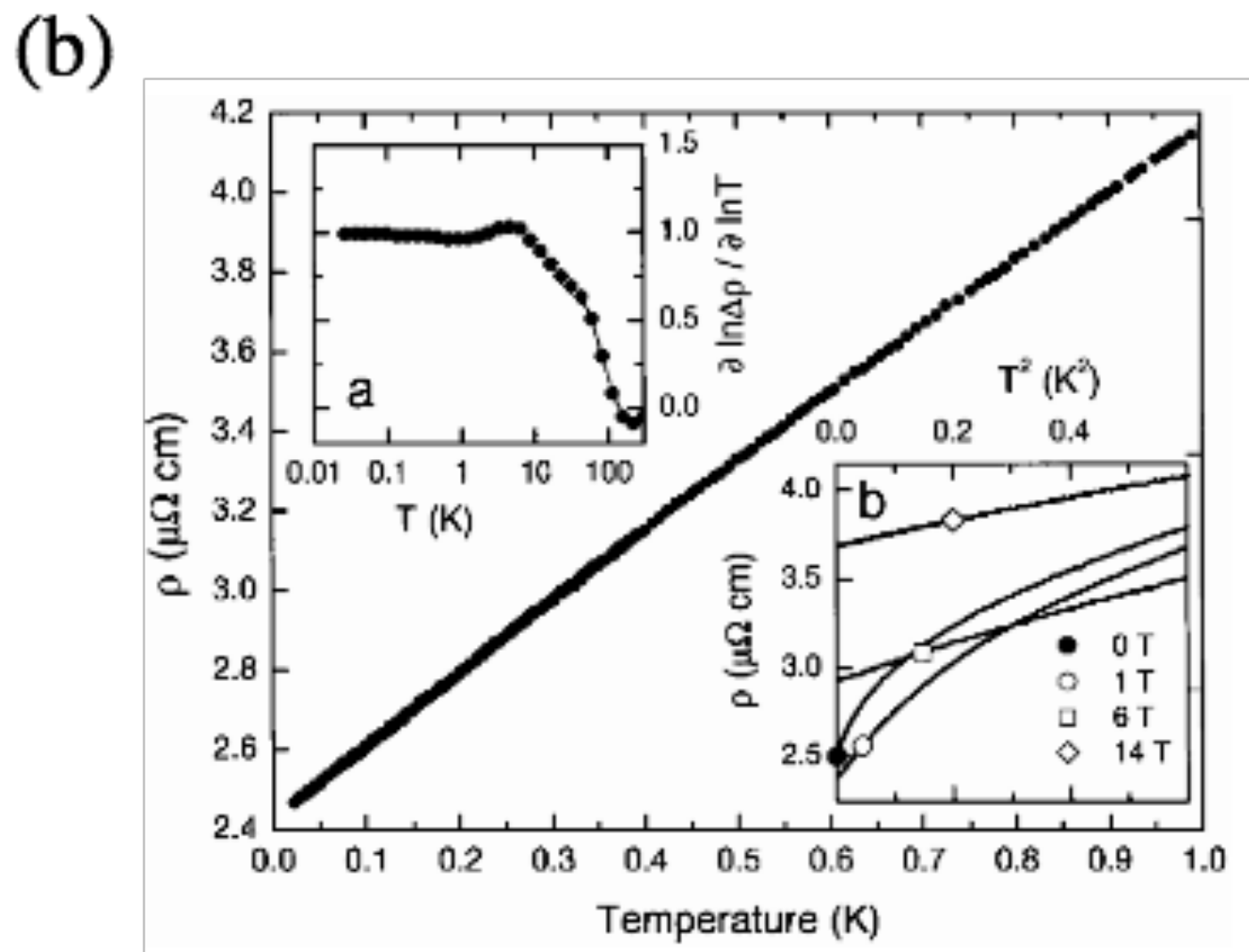
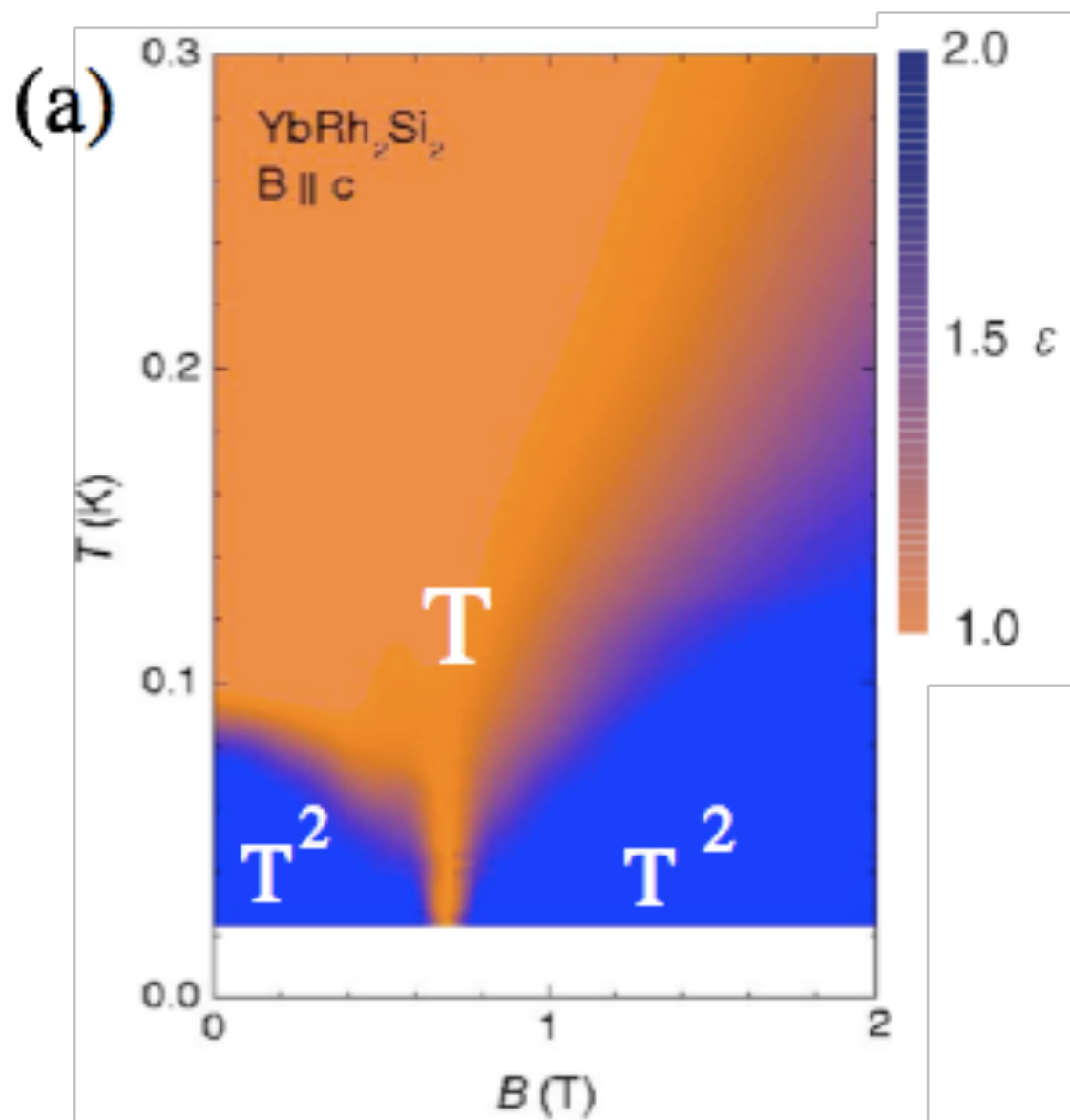
→ **AFM/Superconductivity**



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Coherent Heavy Fermions

YbRh₂Si₂ : Field tuned quantum criticality.



Custers et al, (2003)

Outline of the Topics

1. Trends in the periodic table.
2. Introduction: Heavy Fermions and the Kondo Lattice.
3. From Anderson to Kondo
4. Kondo Insulators: the simplest heavy fermions.
5. Oshikawa's Theorem.
6. Large N expansion for the Kondo Lattice
7. Heavy Fermion Superconductivity
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Please ask questions!

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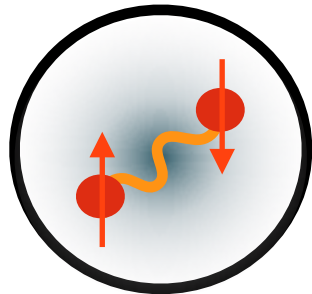
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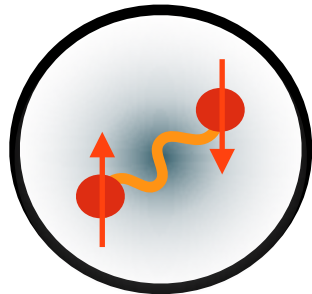
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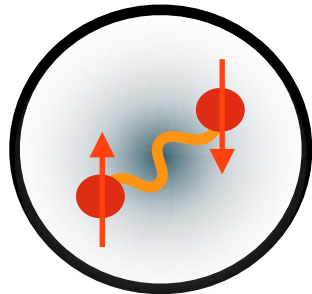
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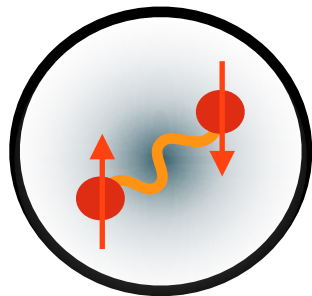


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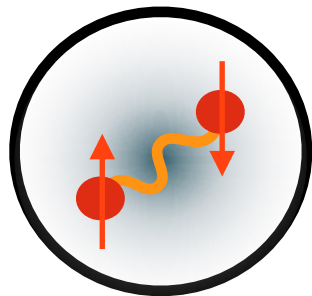
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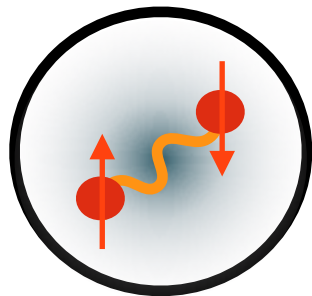
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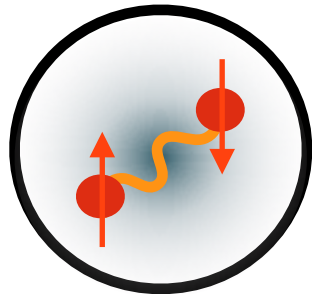
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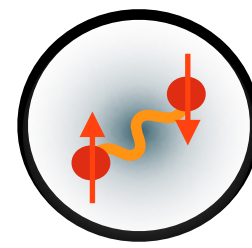
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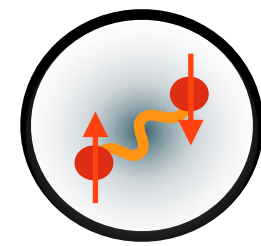
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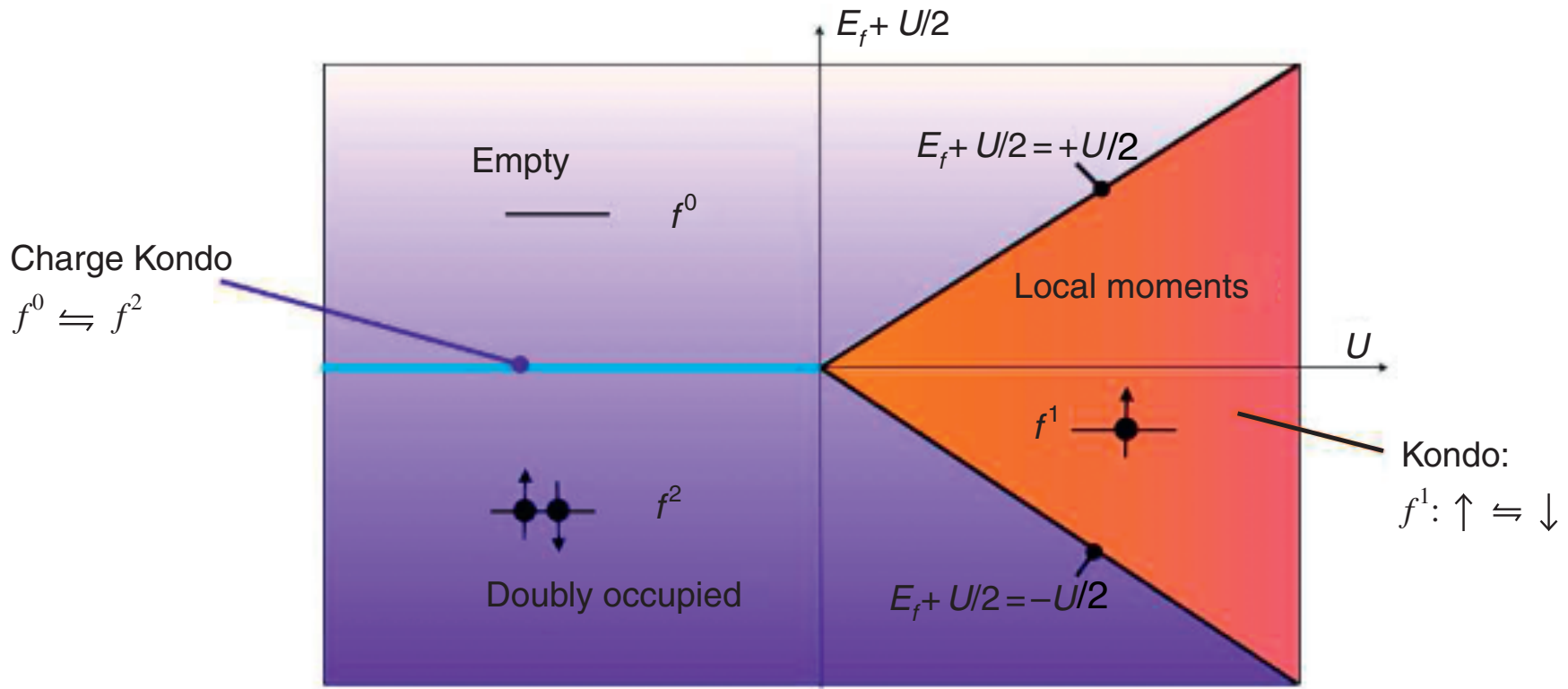
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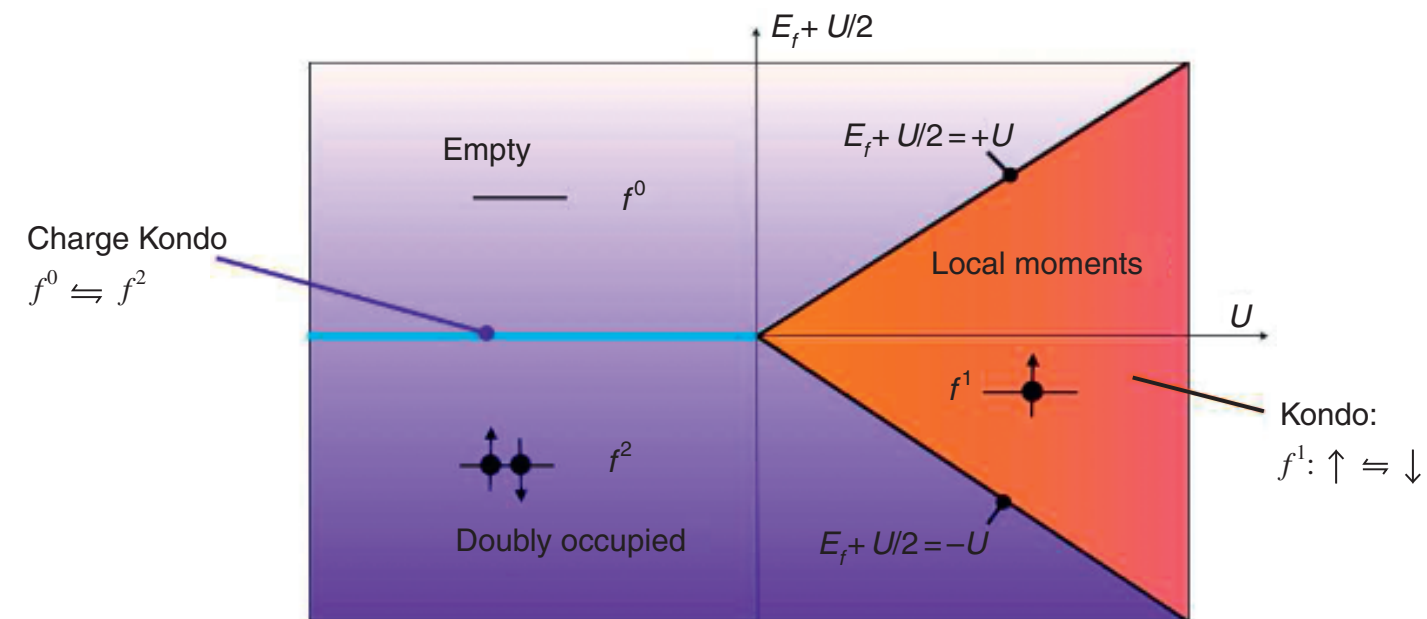
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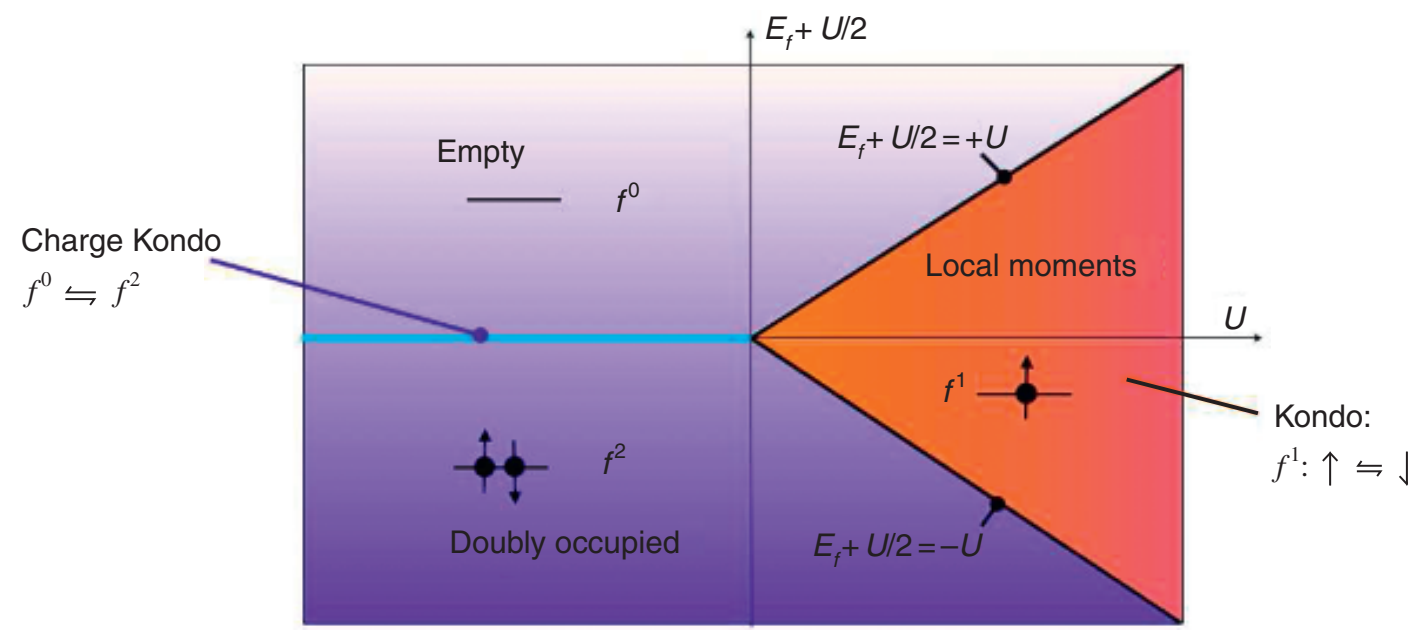
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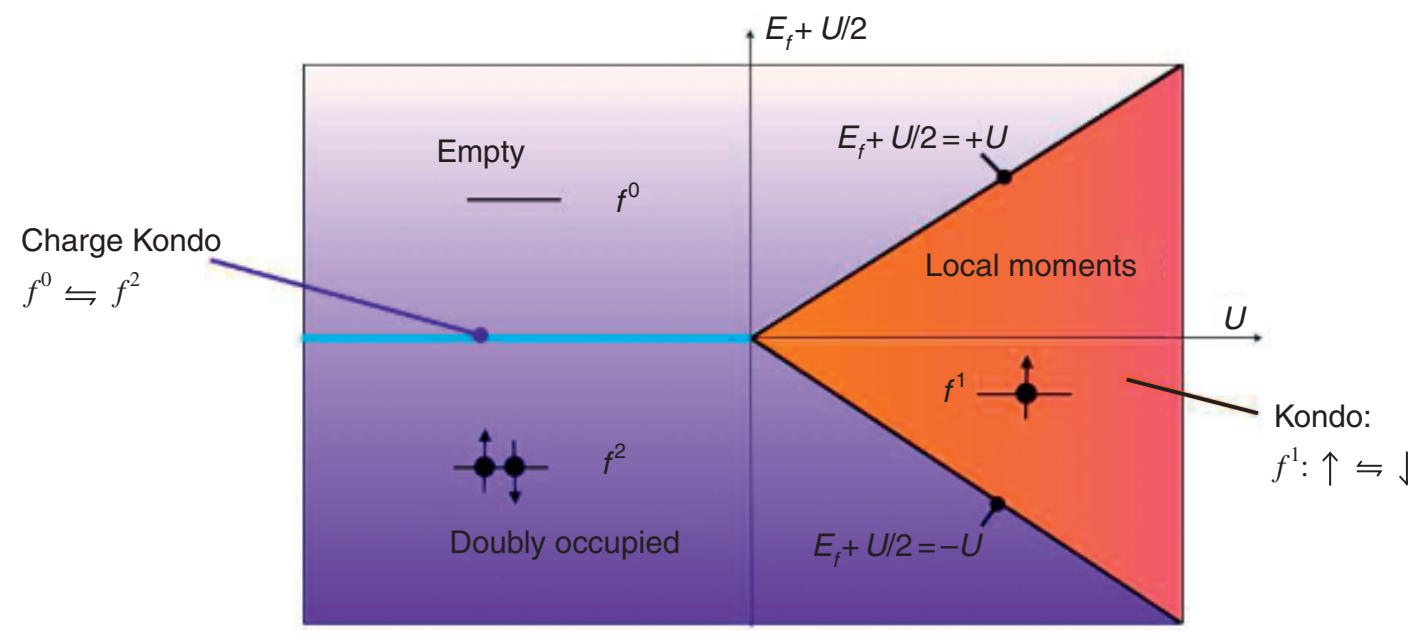
Anderson 1961



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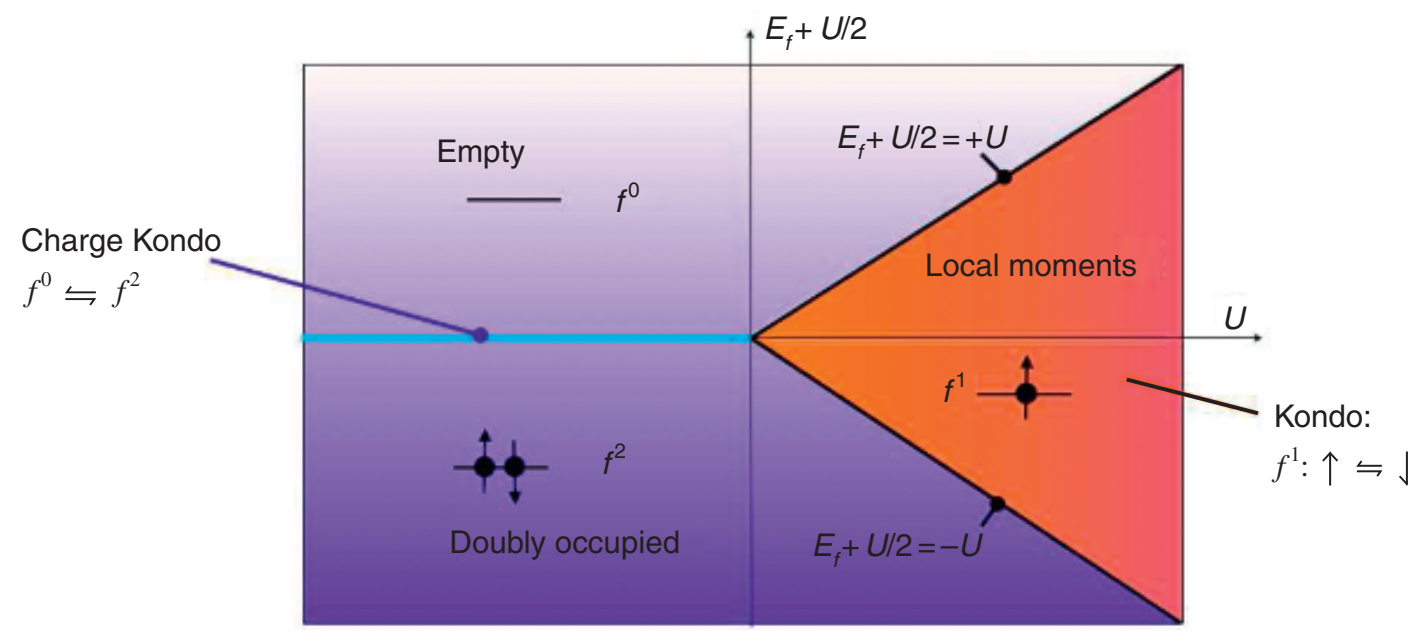
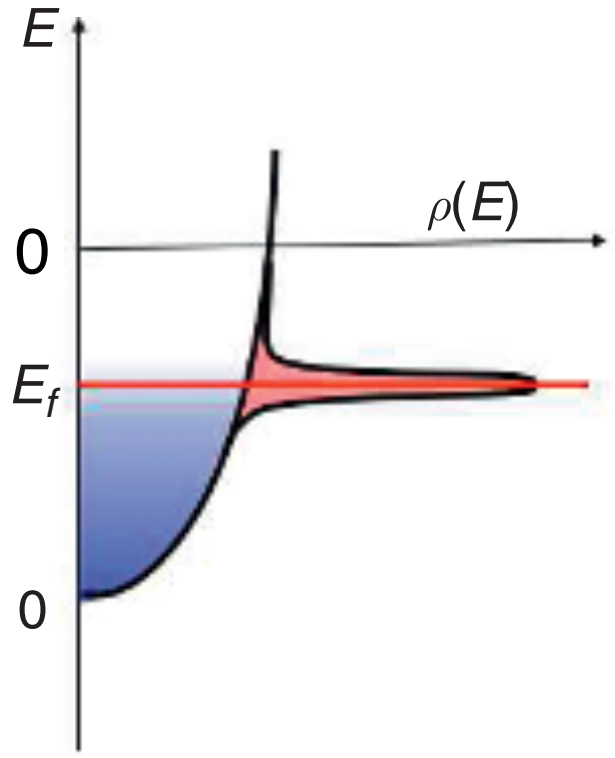
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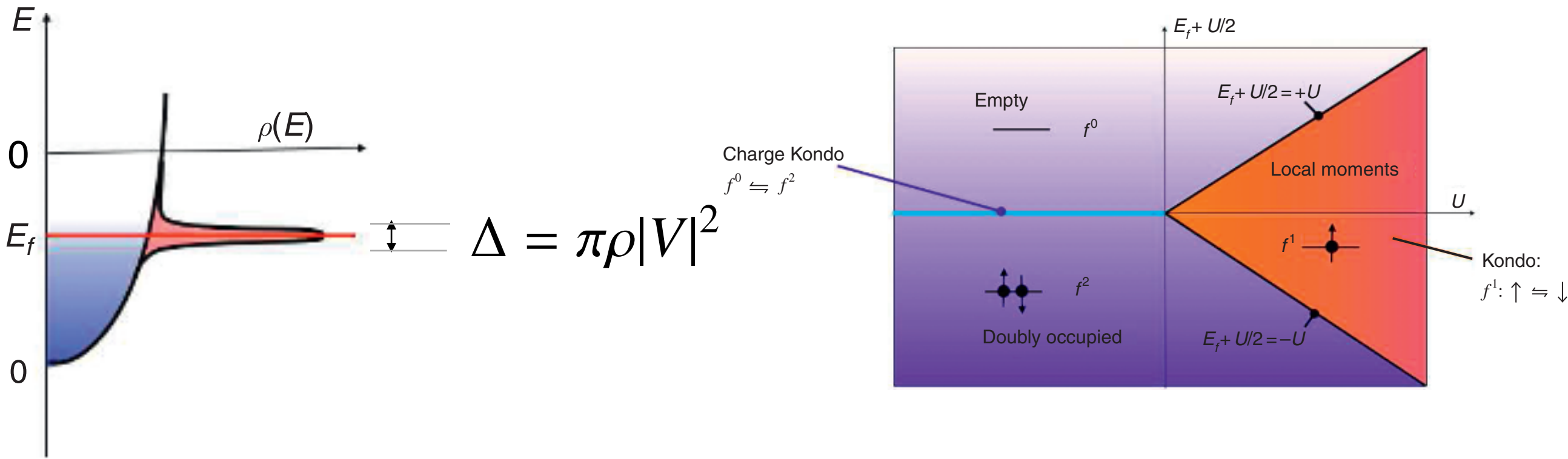
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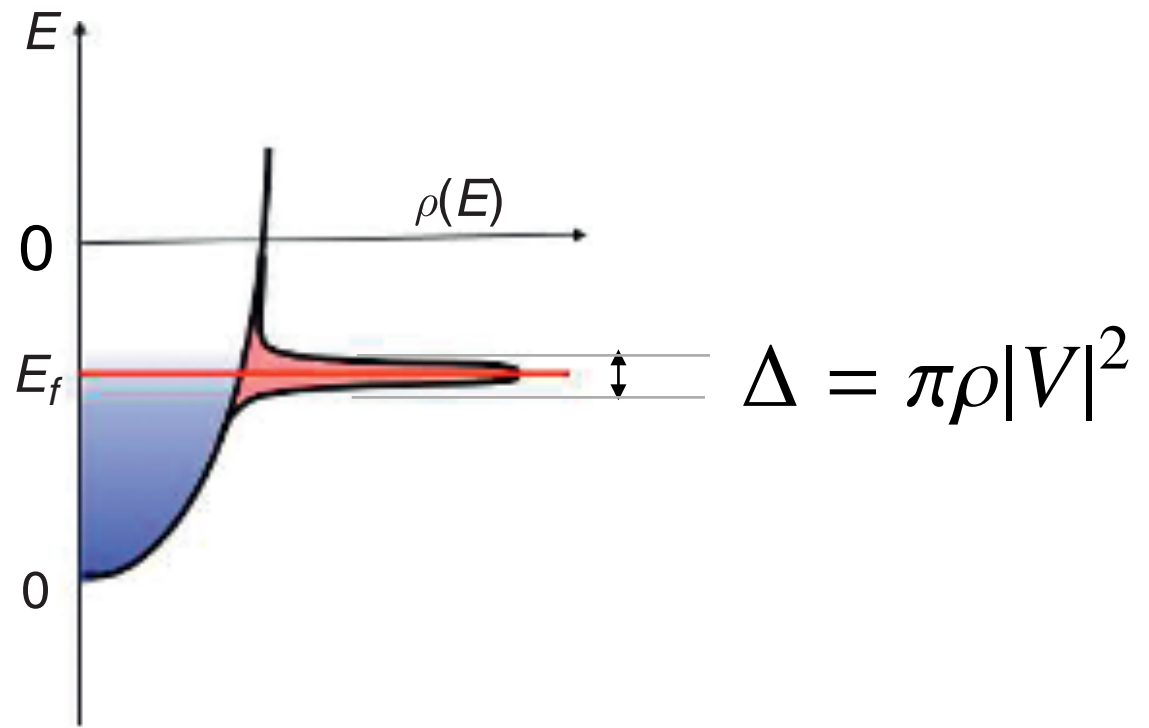
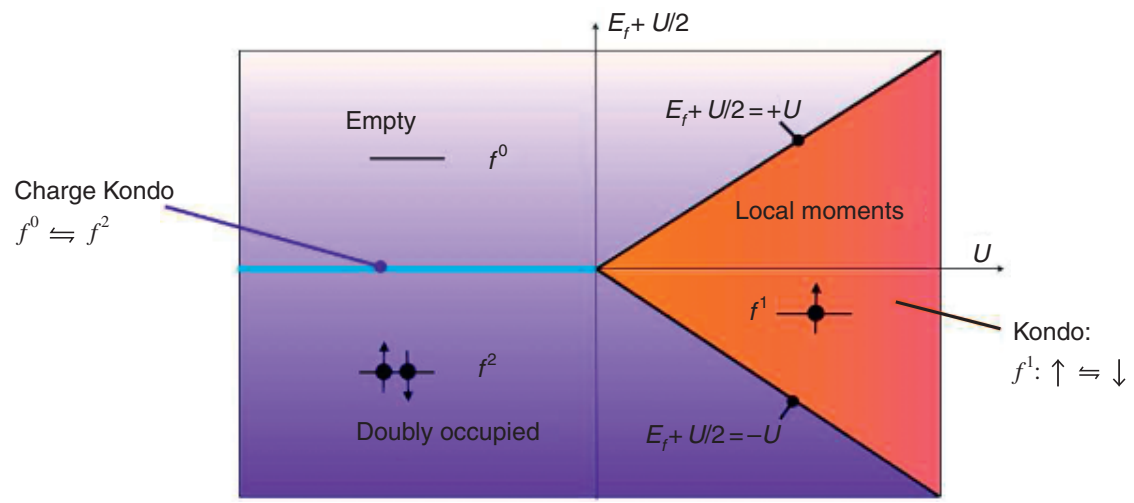
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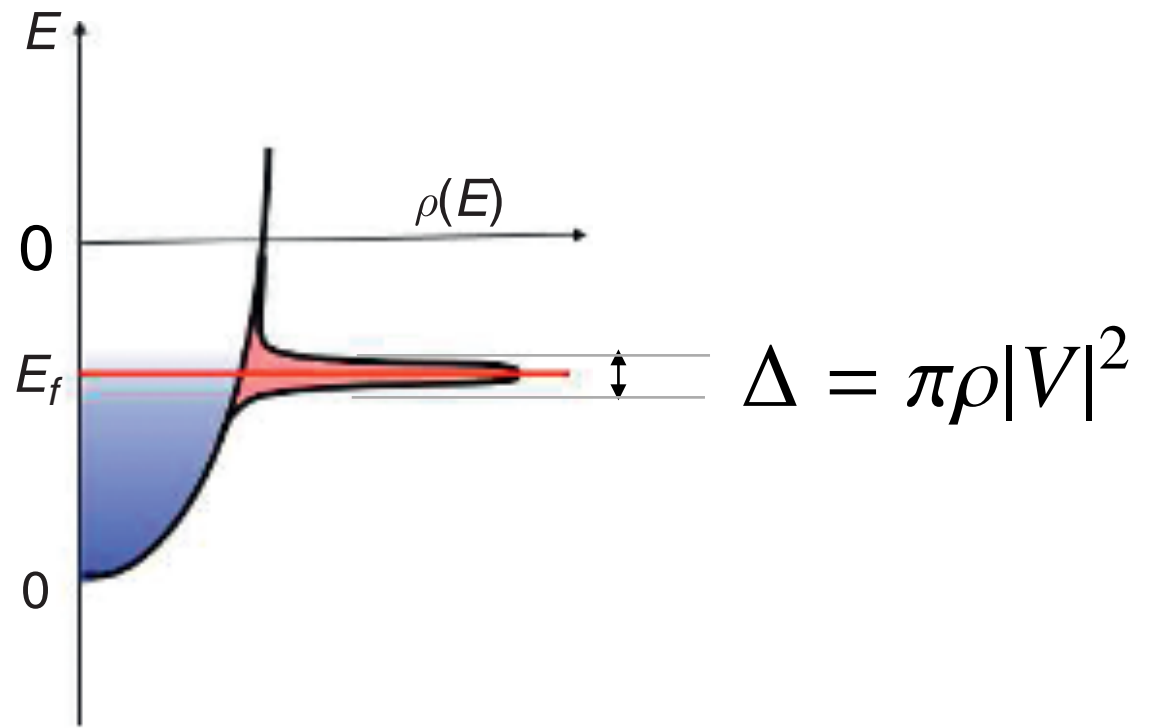
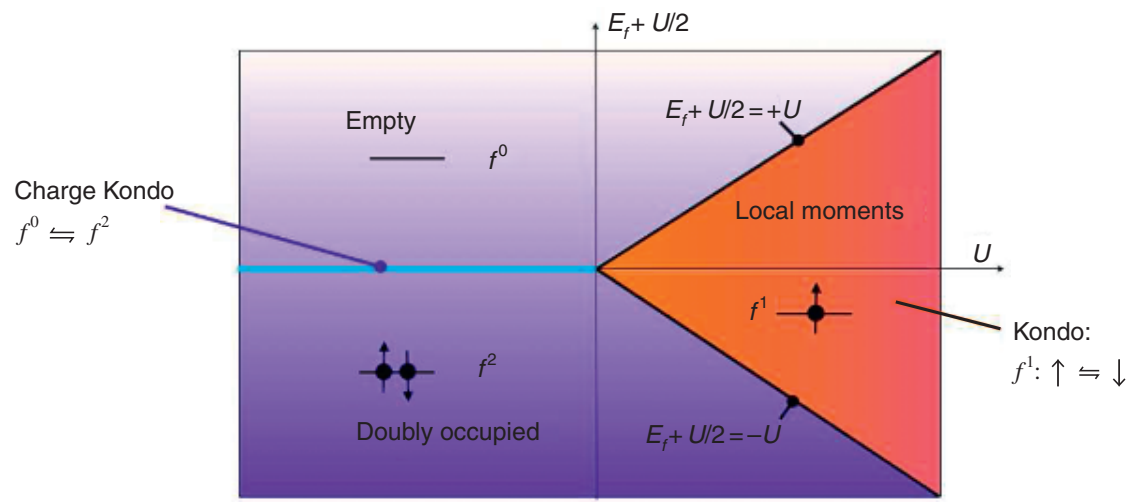
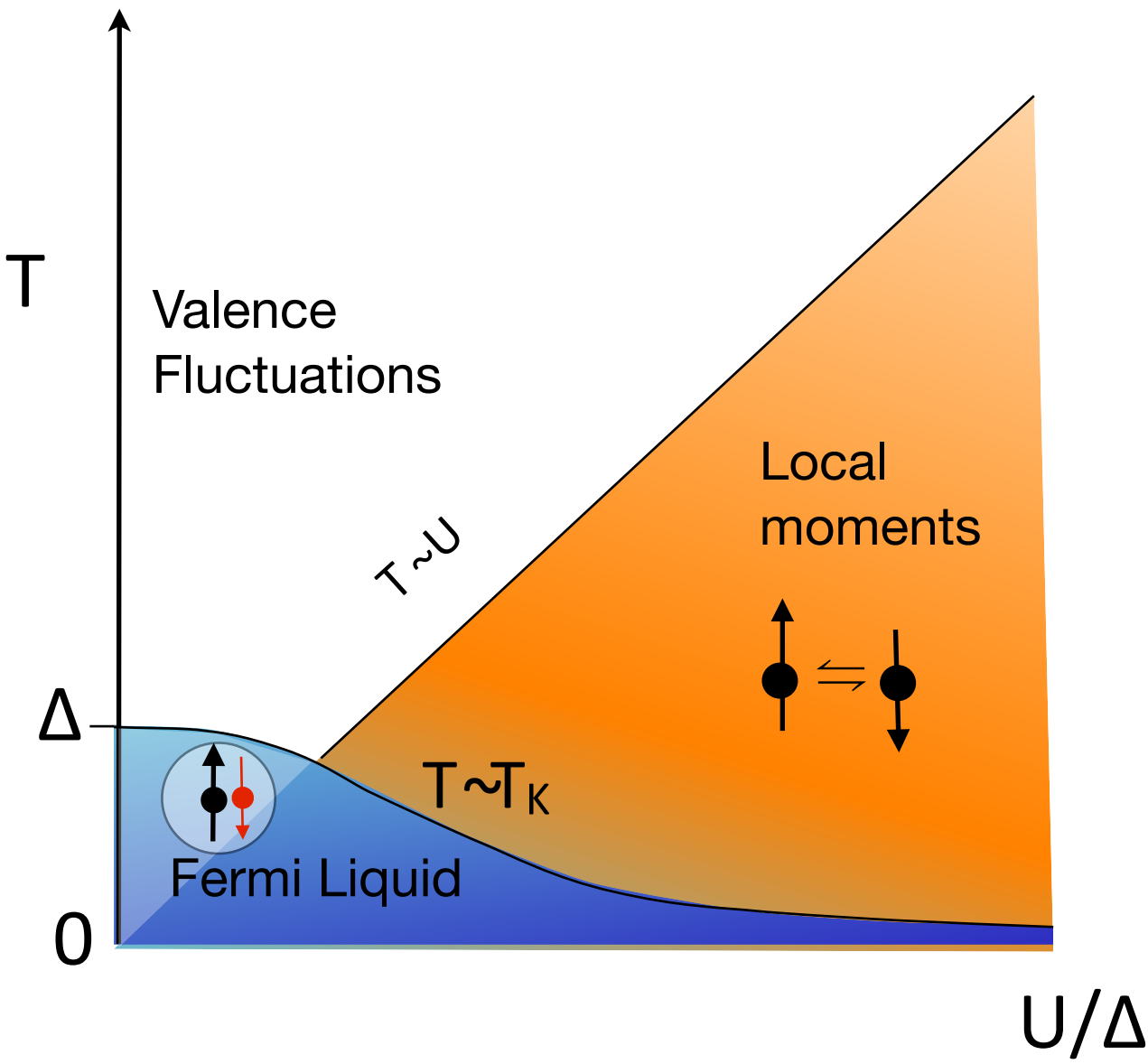
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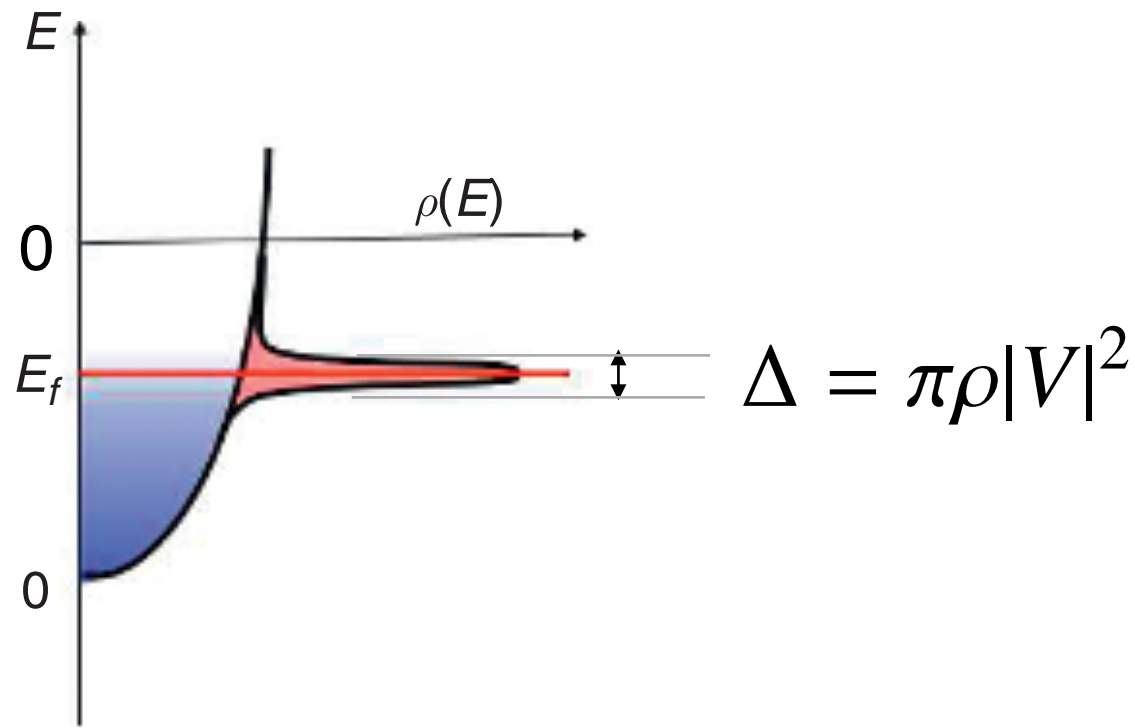
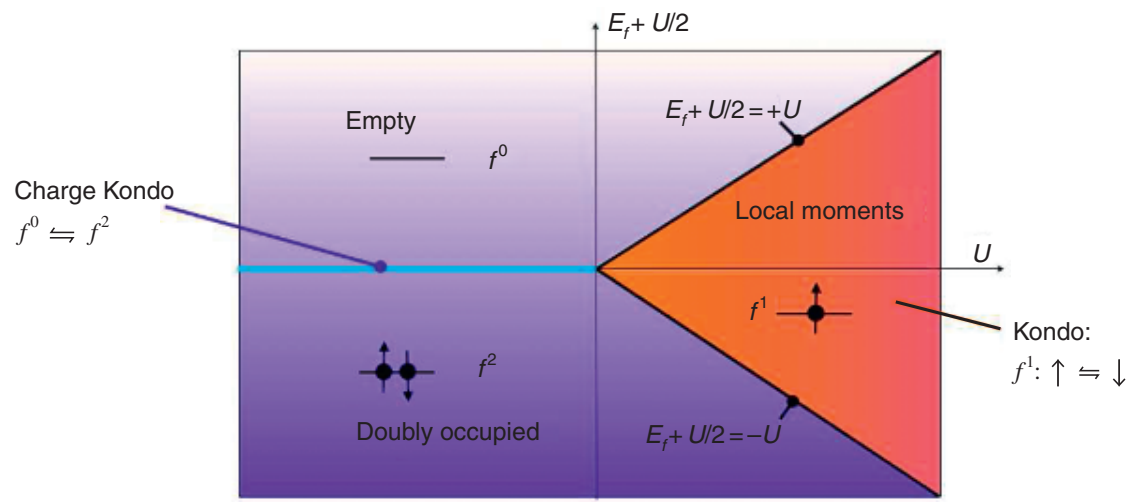
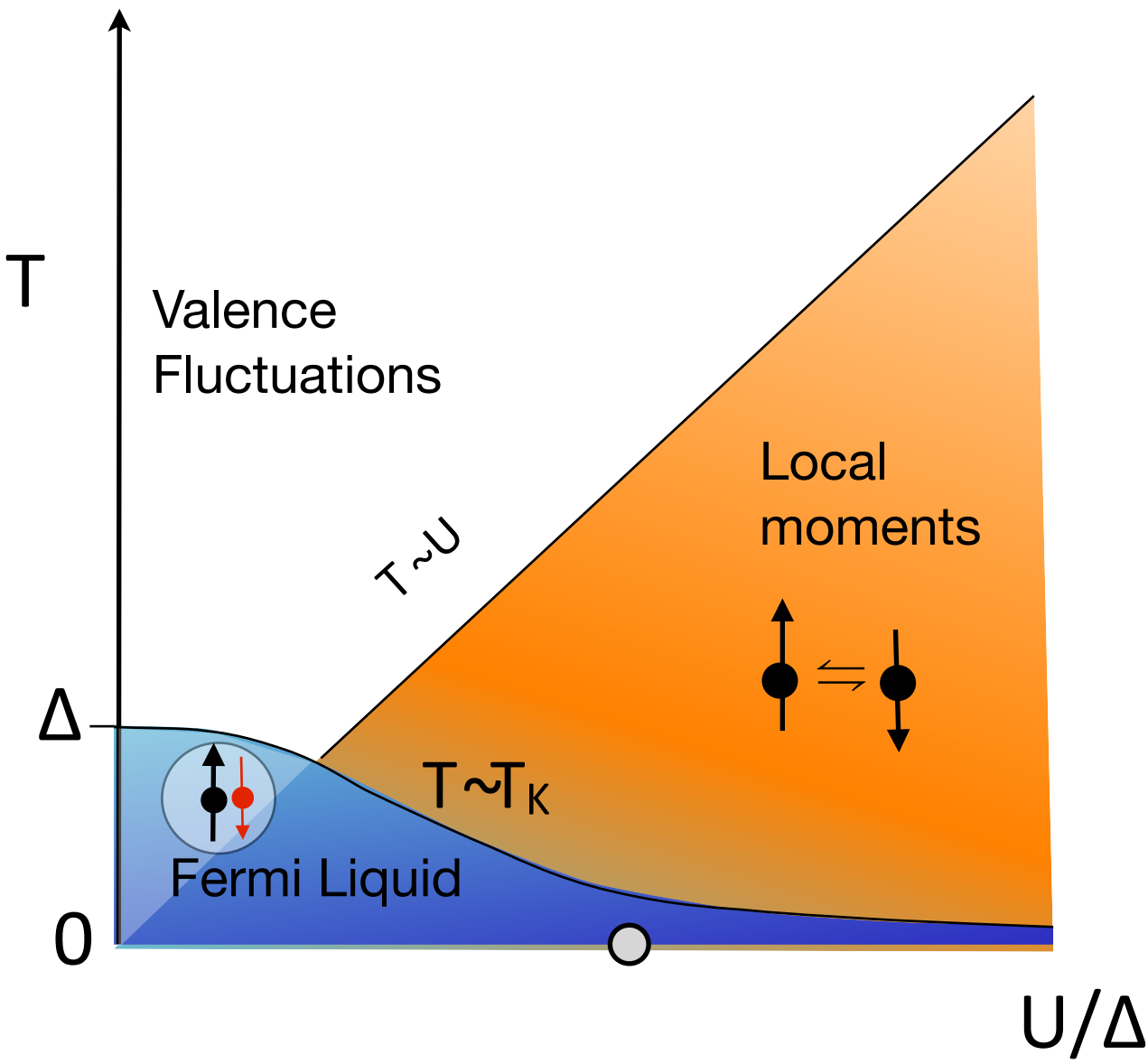
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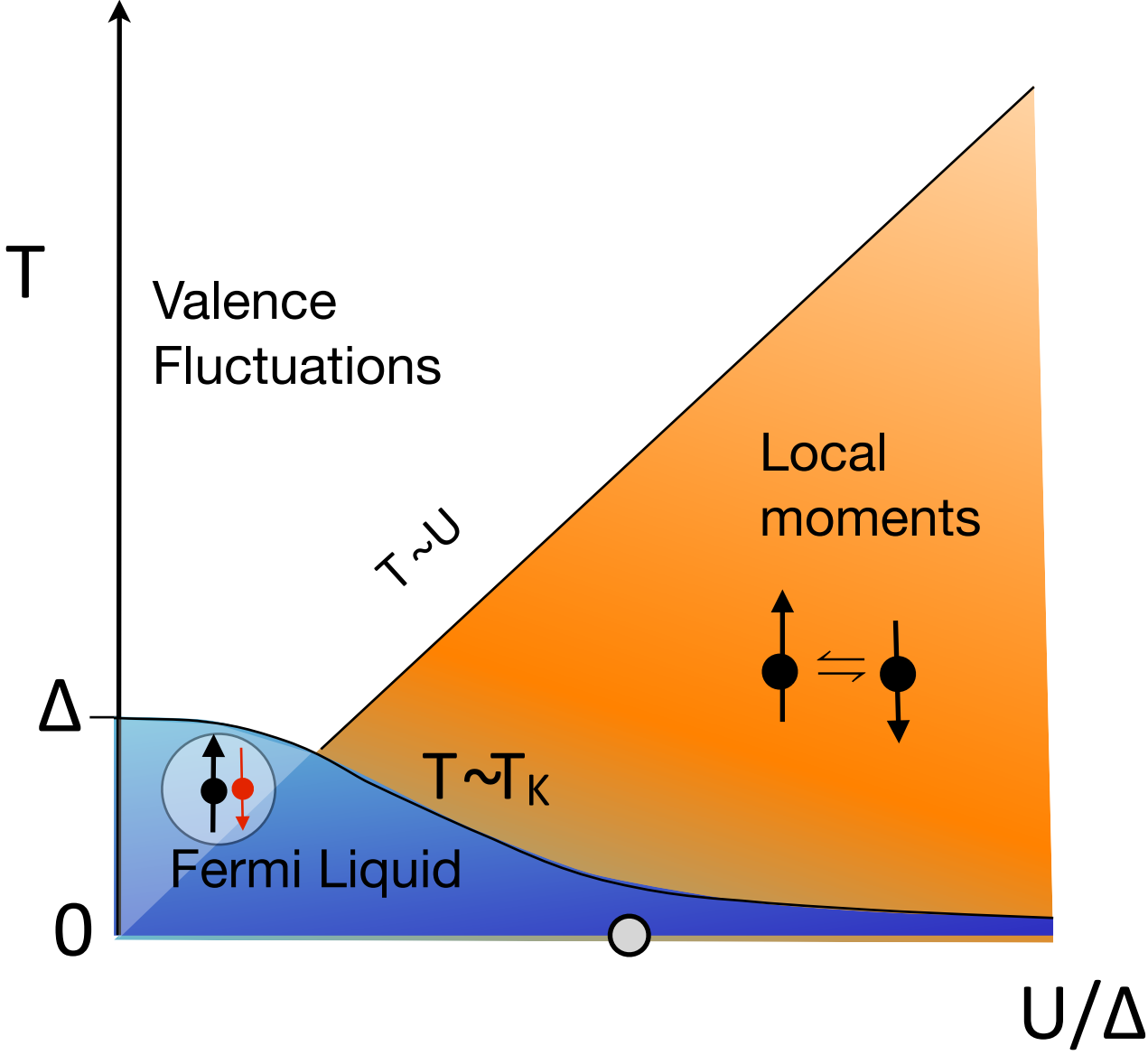
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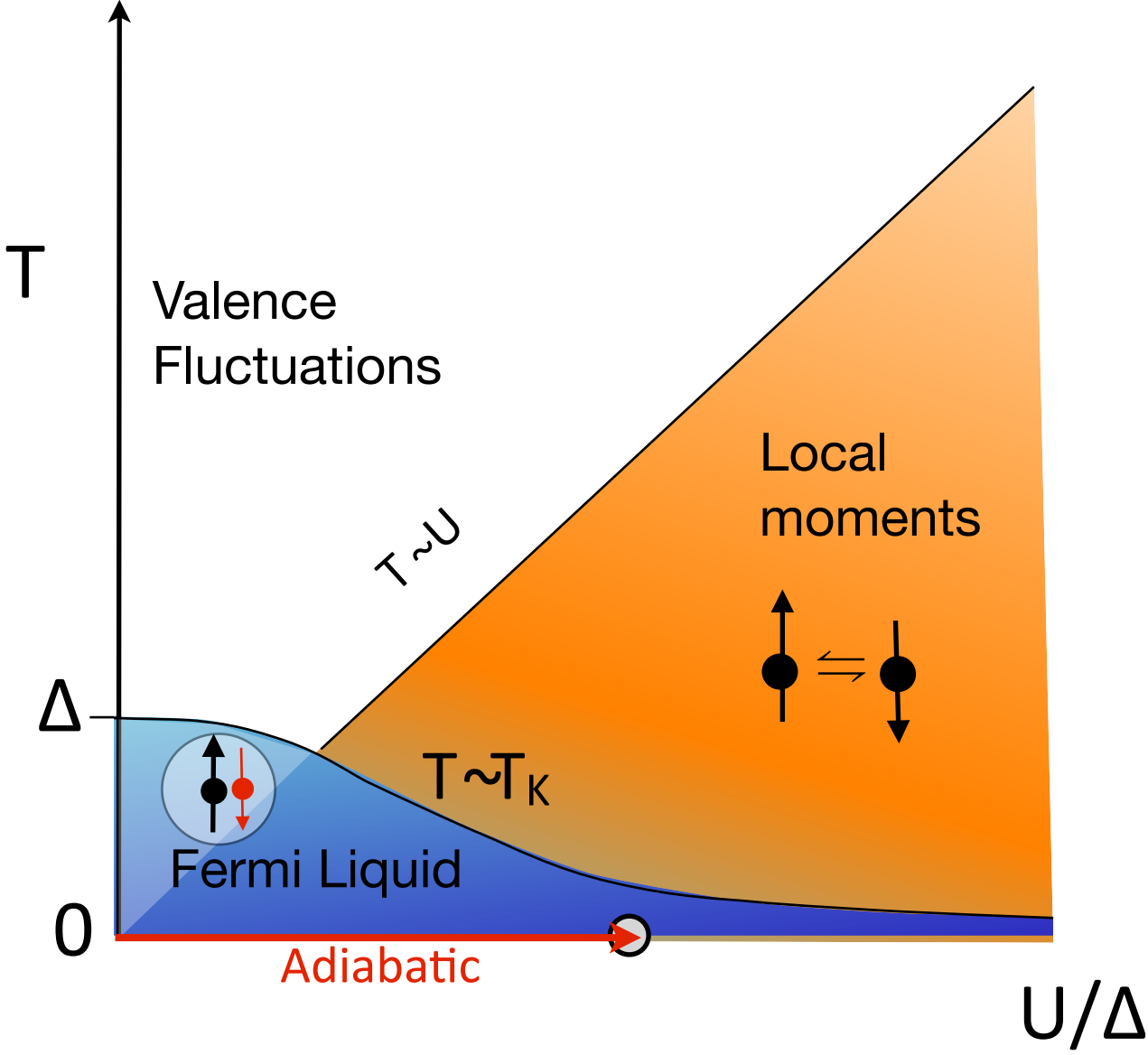
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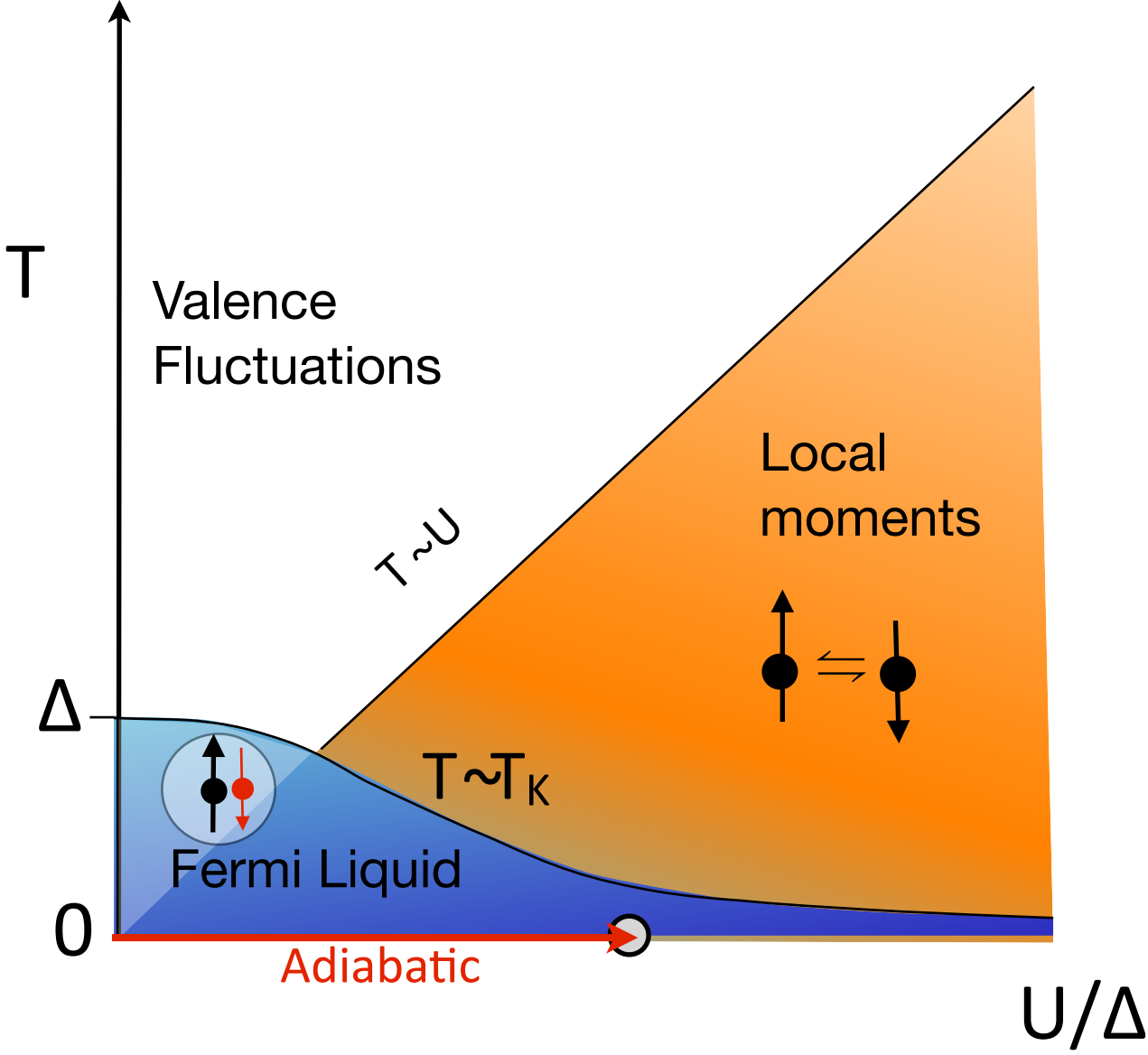
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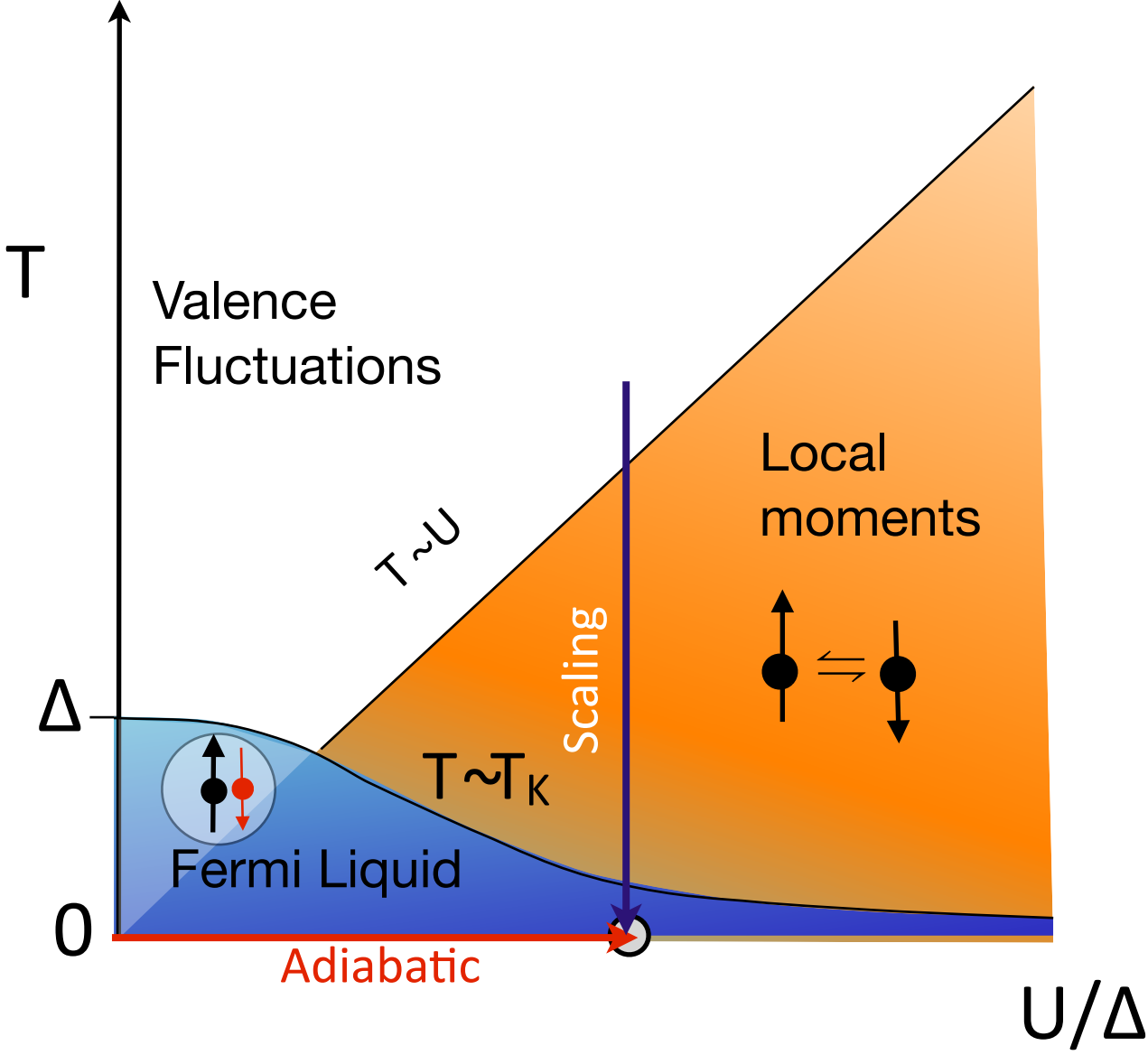
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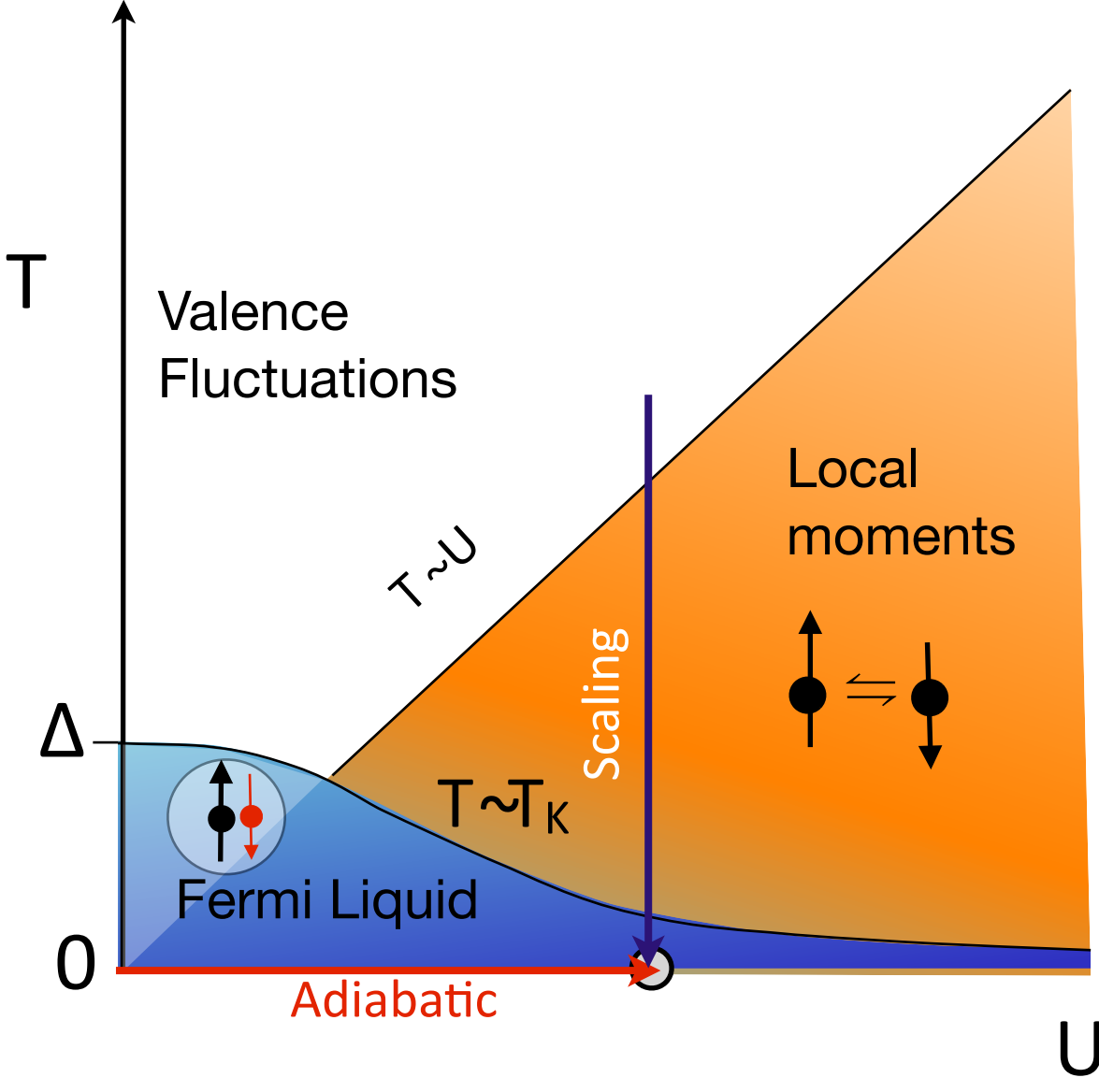
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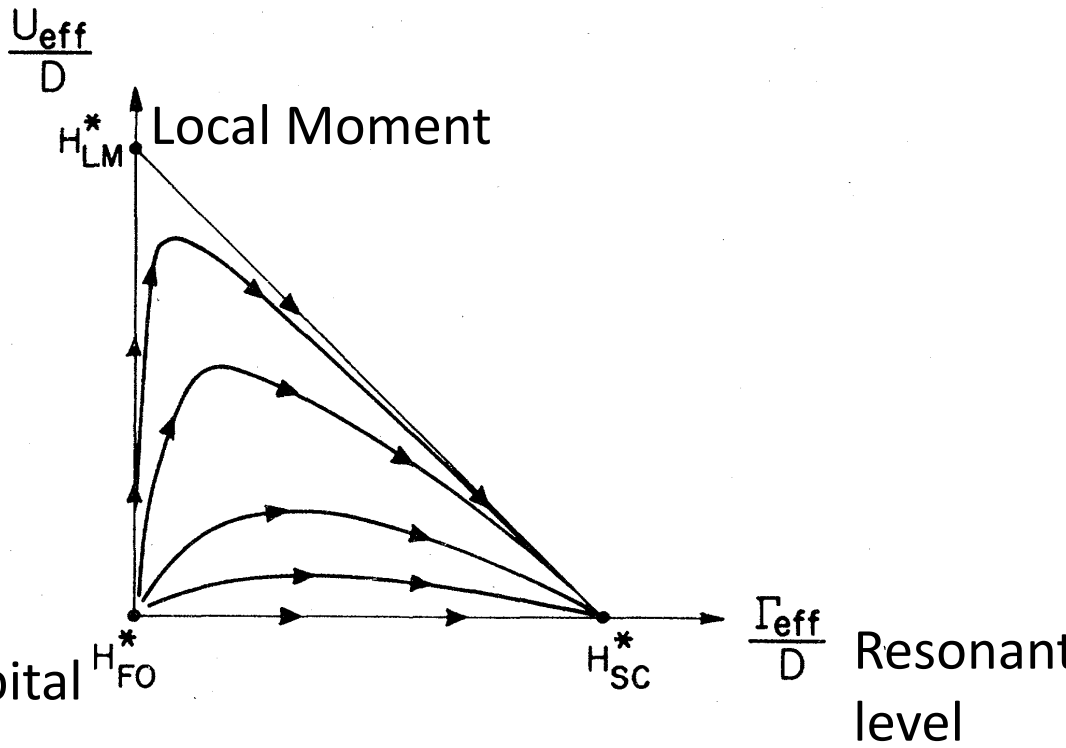
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PHYSICAL REVIEW B VOLUME 21, NUMBER 3 1 FEBRUARY 1980

Renormalization-group approach to the Anderson model of dilute magnetic alloys.
I. Static properties for the symmetric case

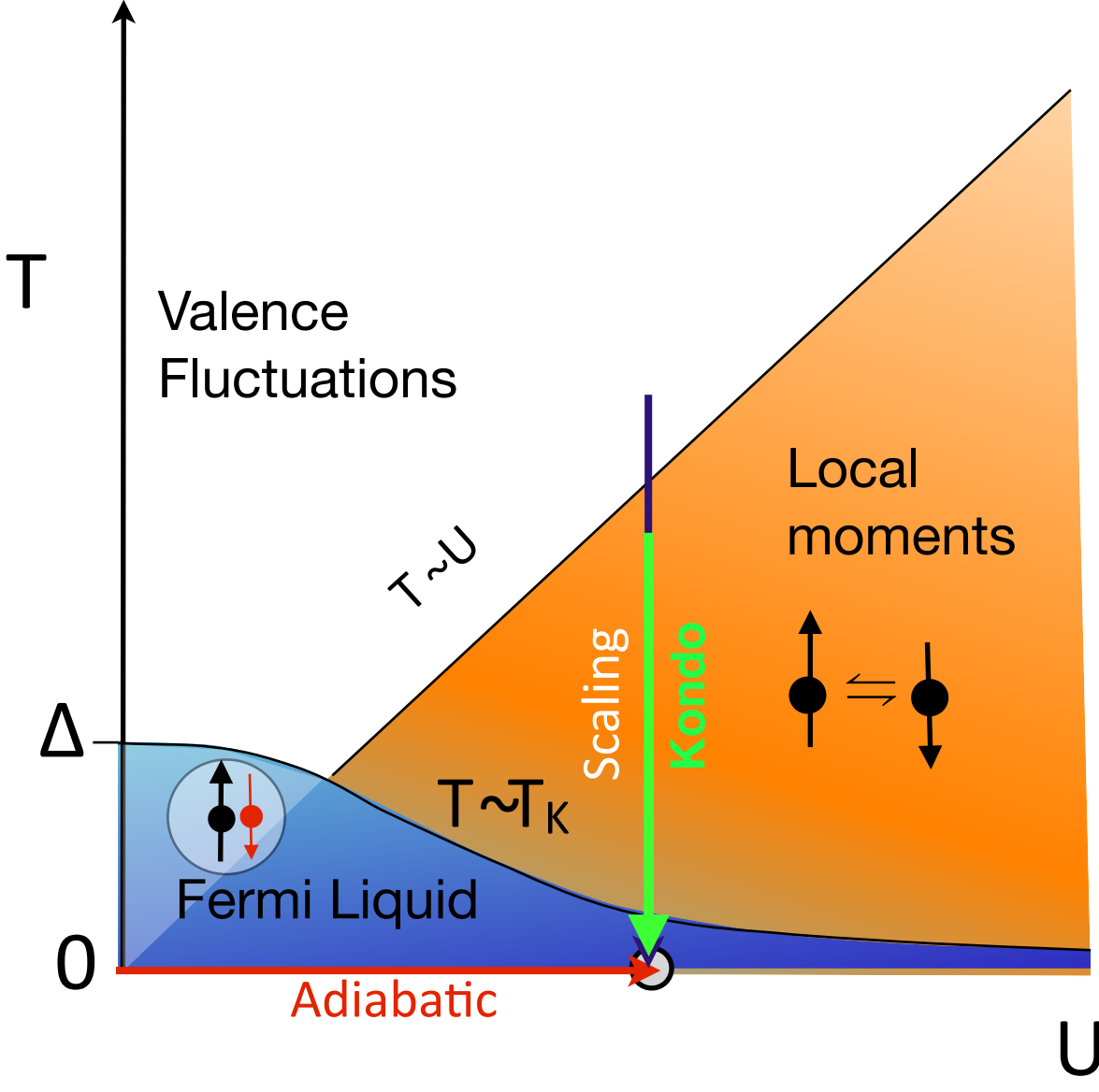
H. R. Krishna-murthy,* J. W. Wilkins, and K. G. Wilson
Physics Department, Cornell University, Ithaca, New York 14853
(Received 10 September 1979)



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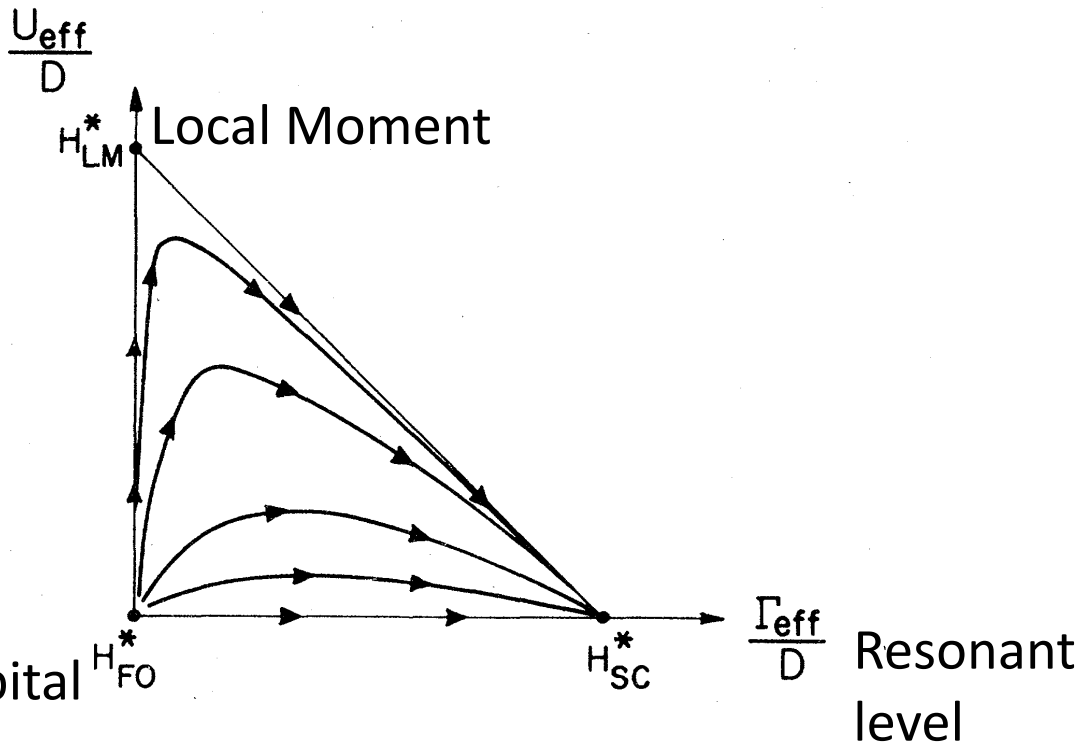
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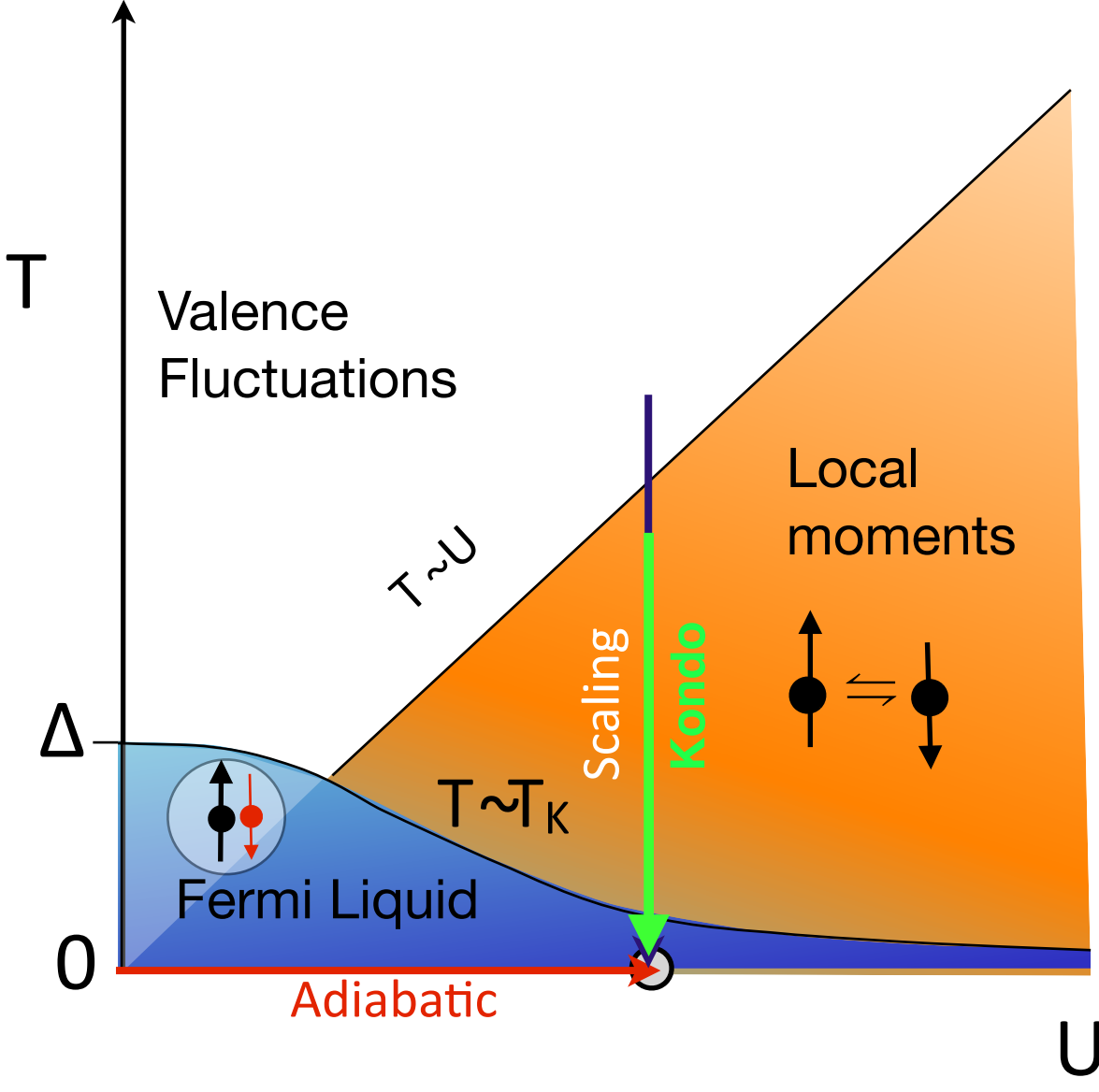
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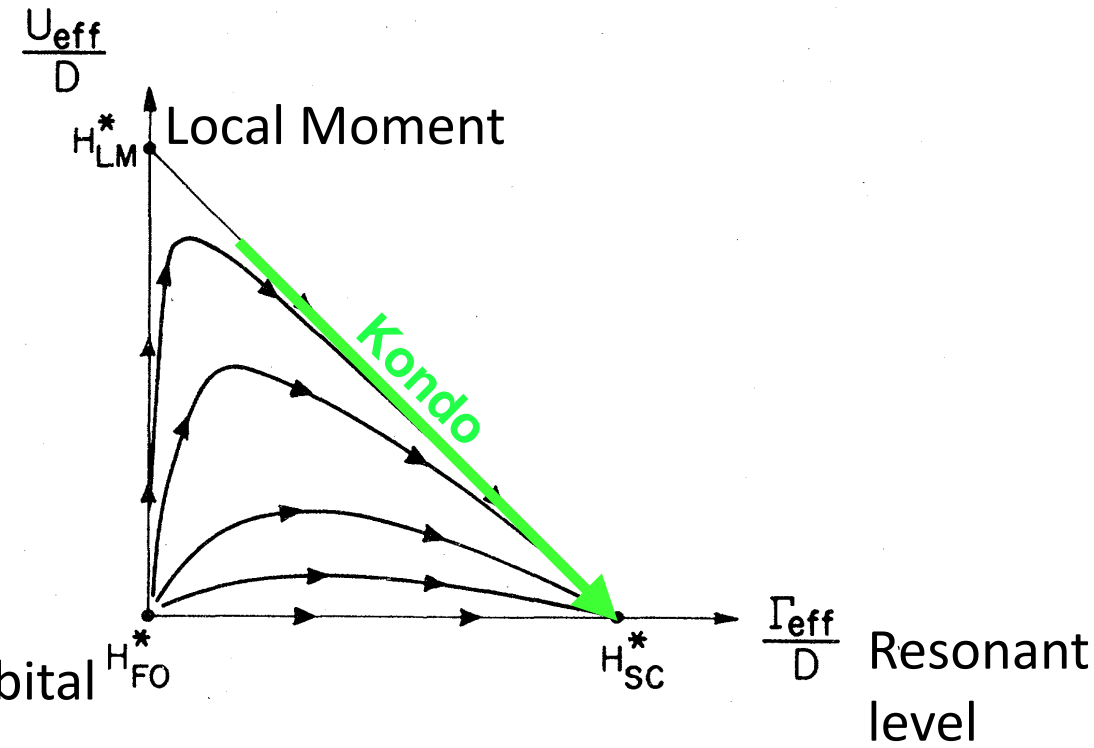
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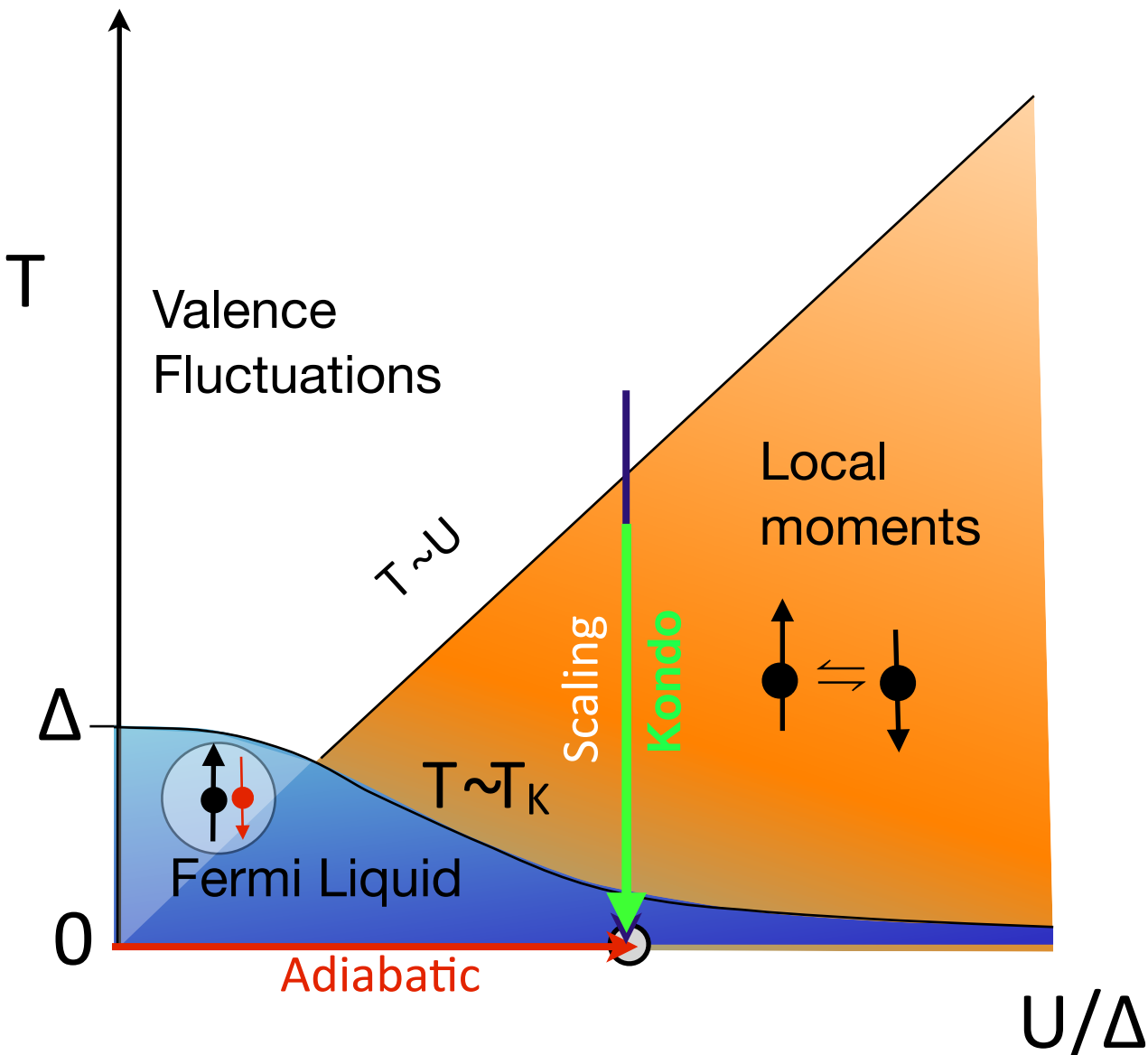
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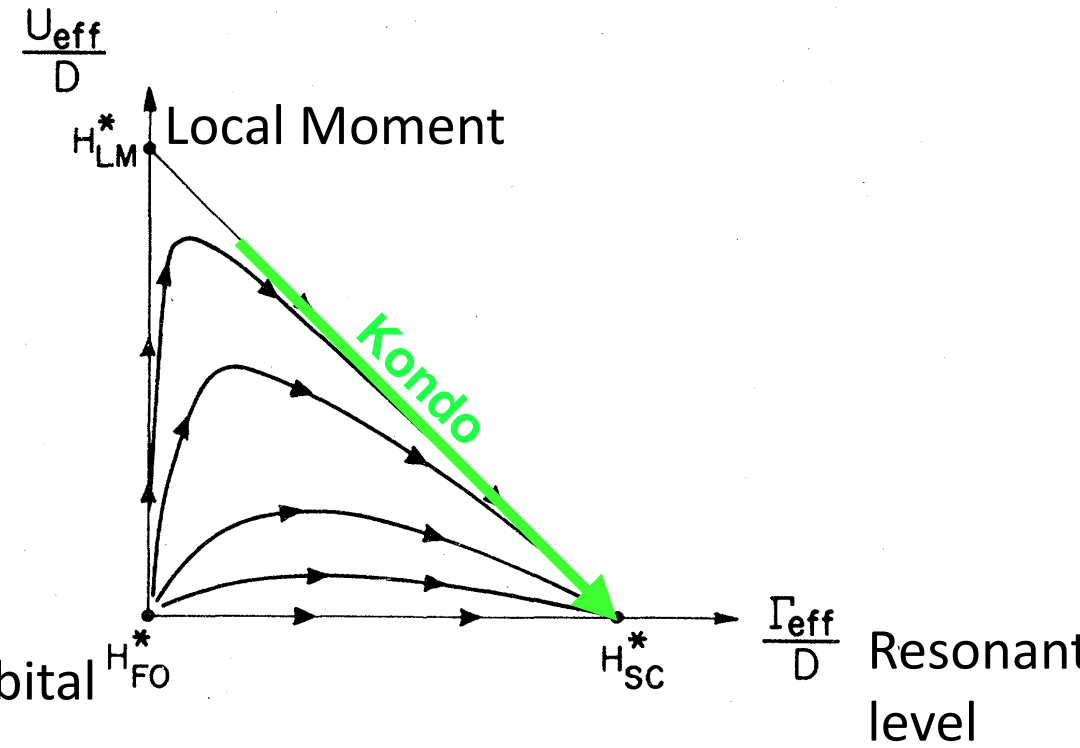


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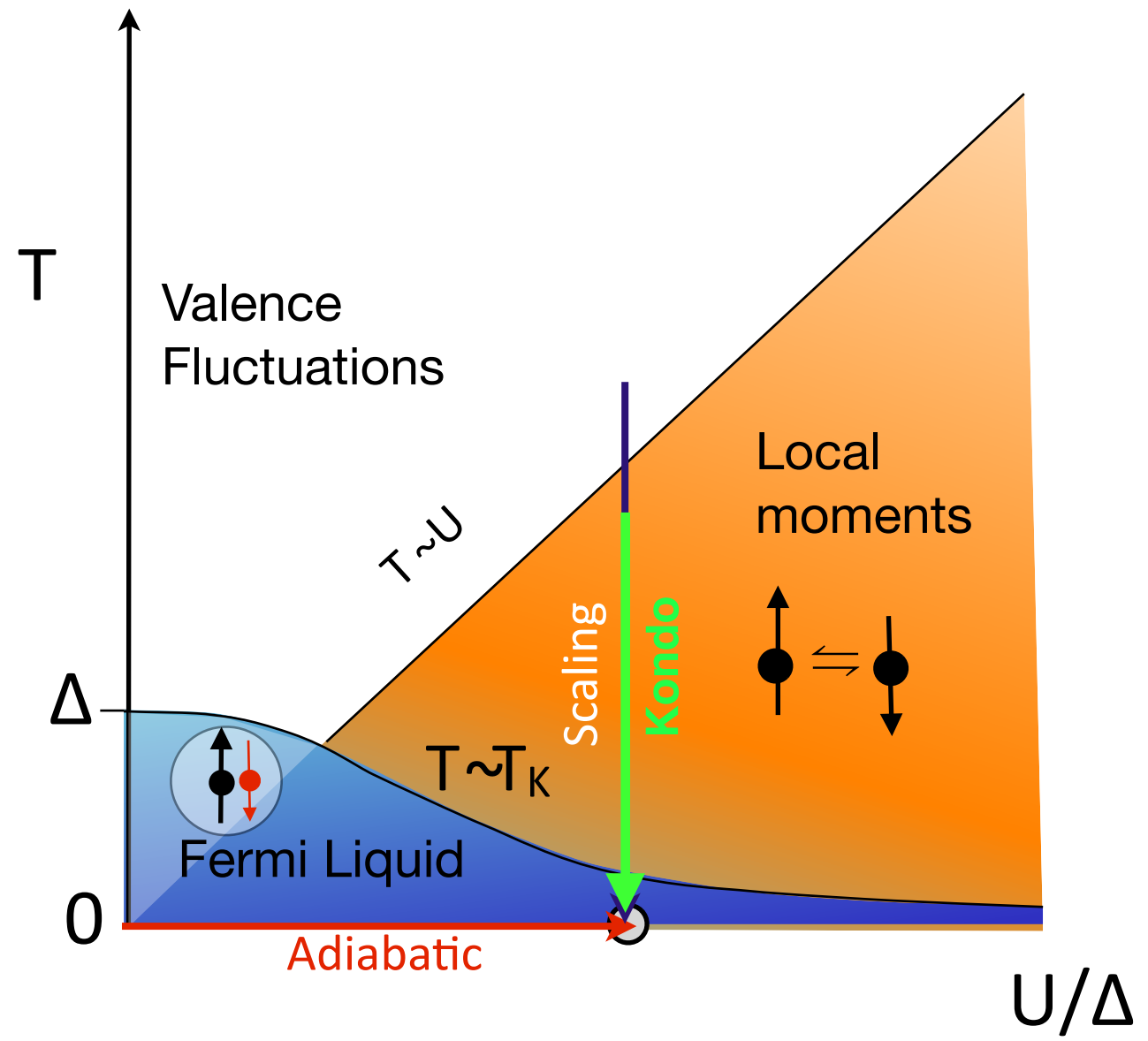


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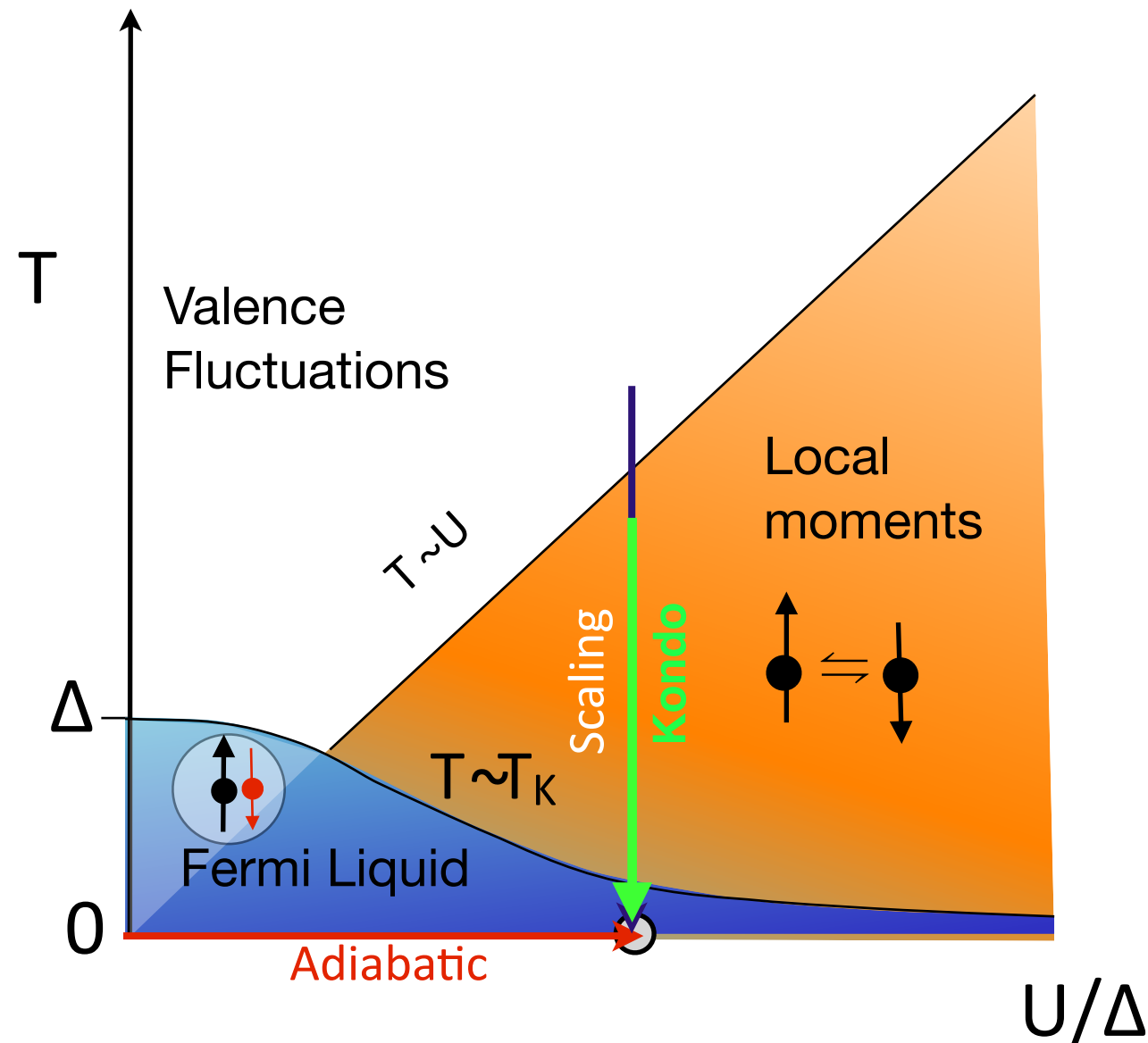
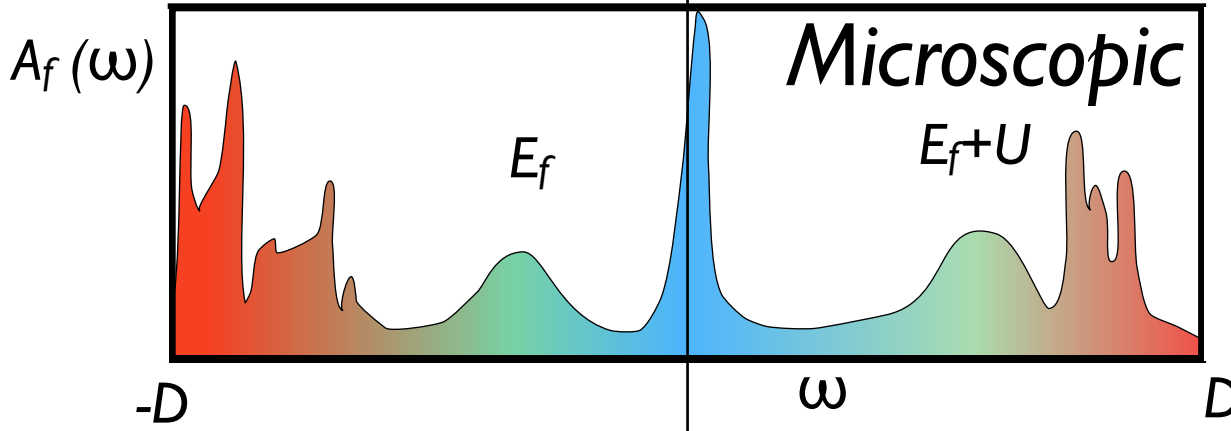
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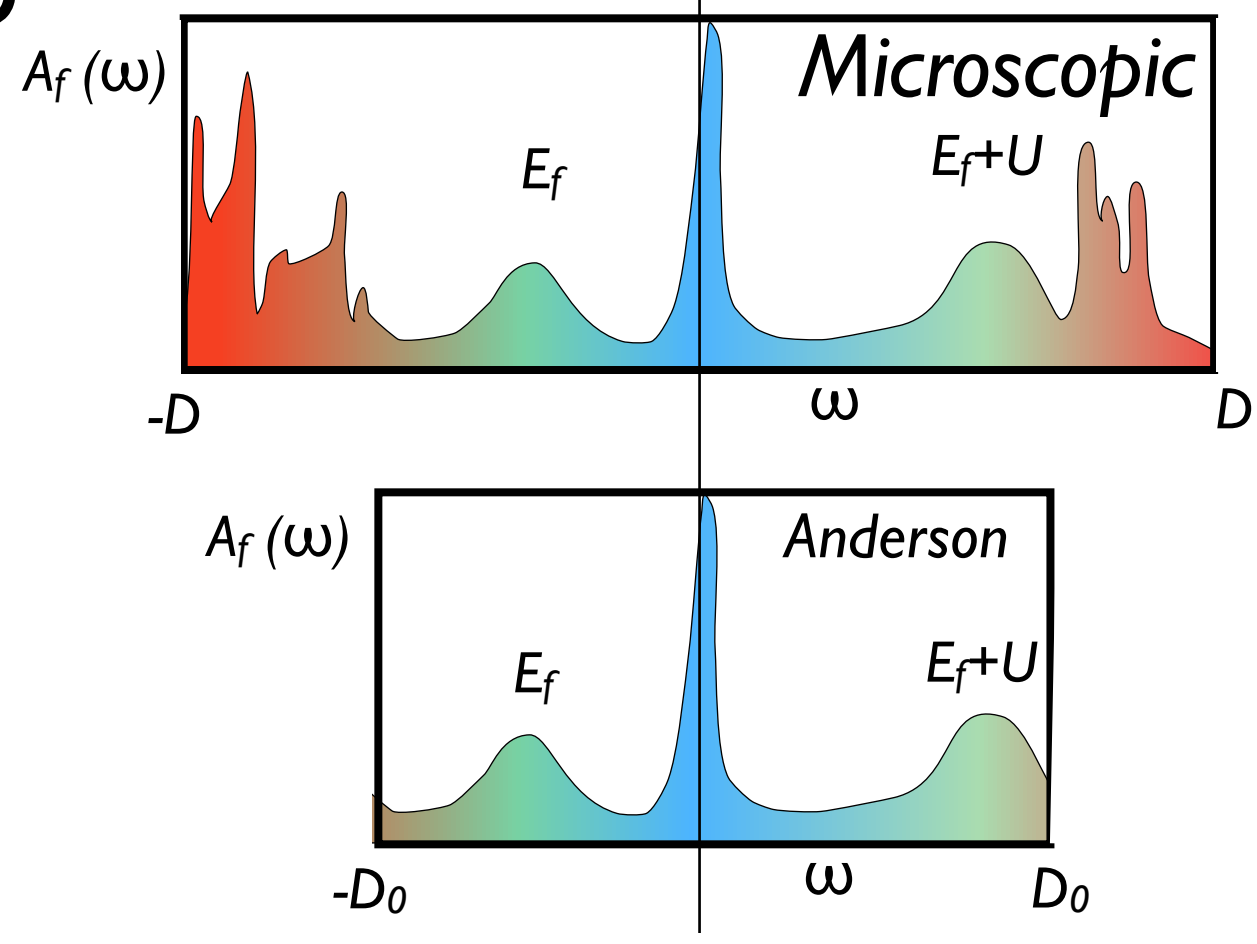
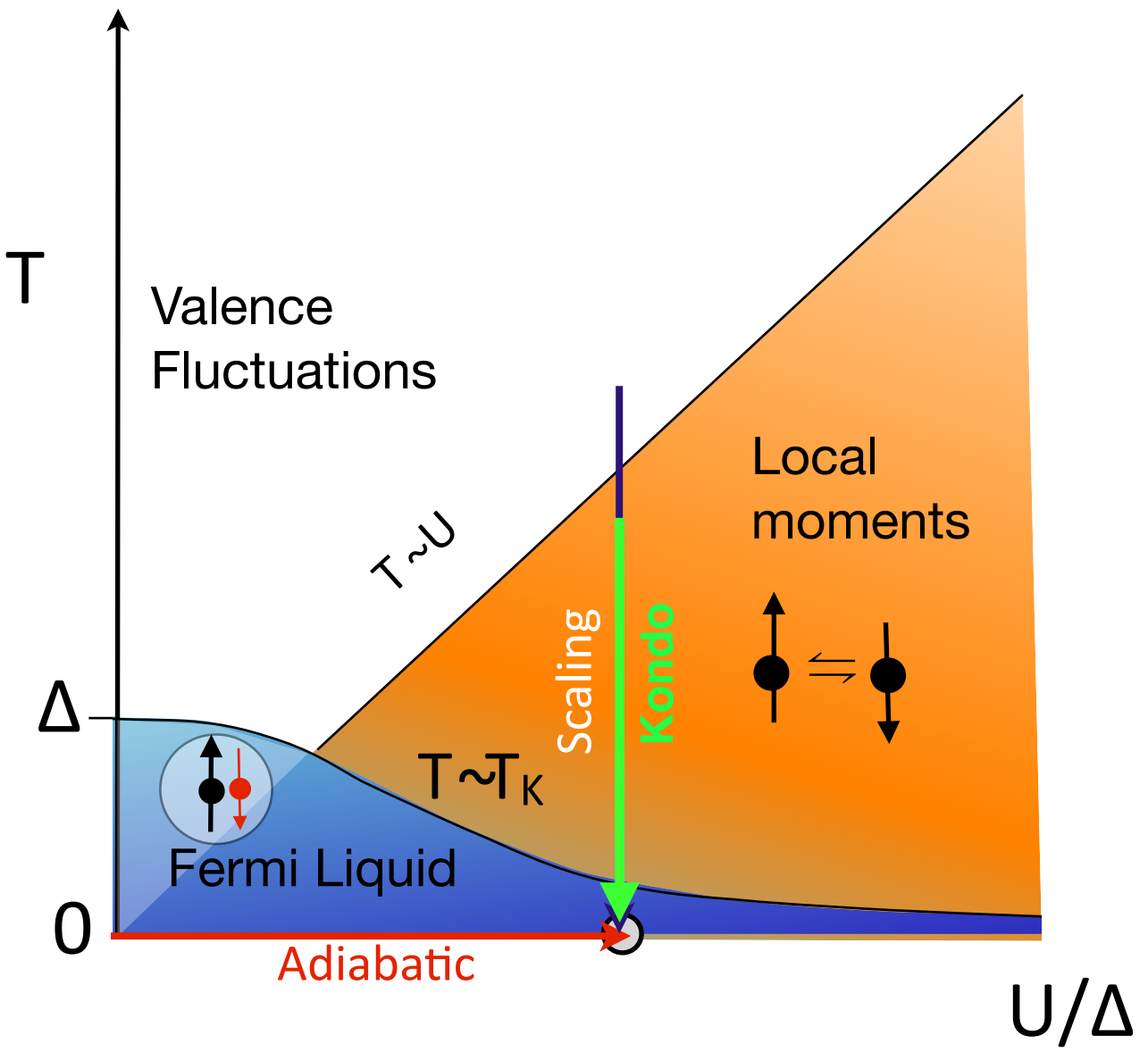
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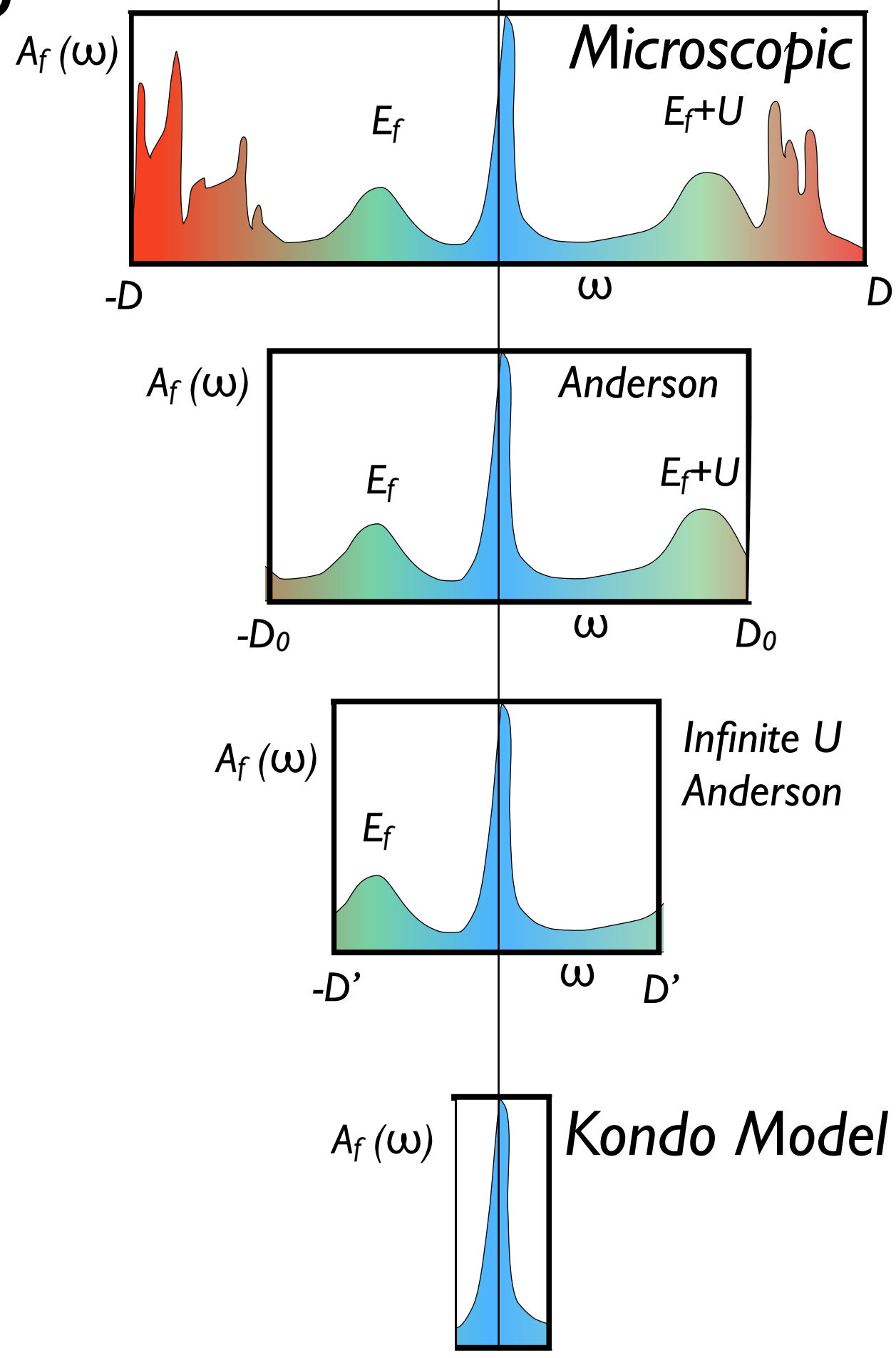
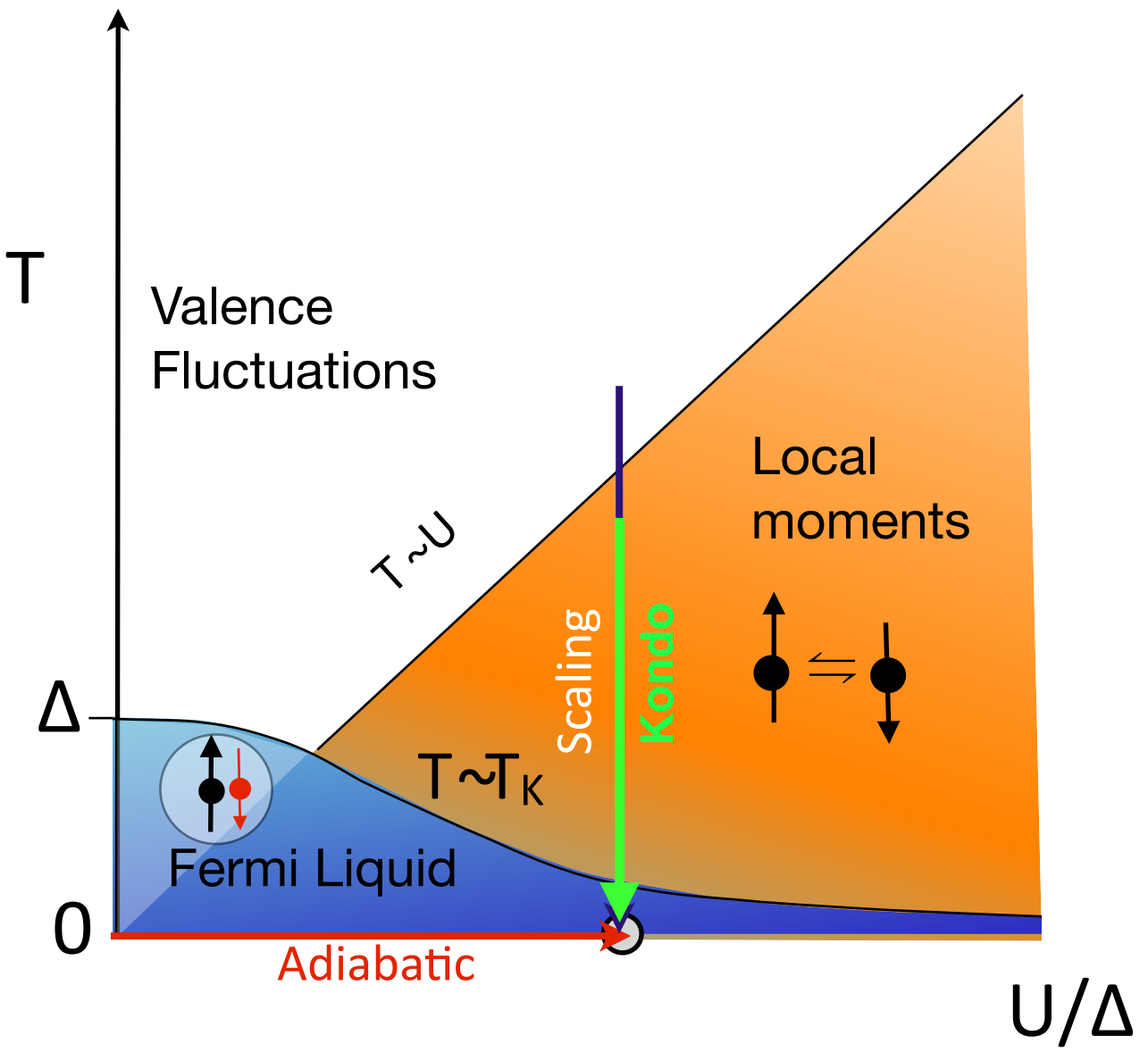
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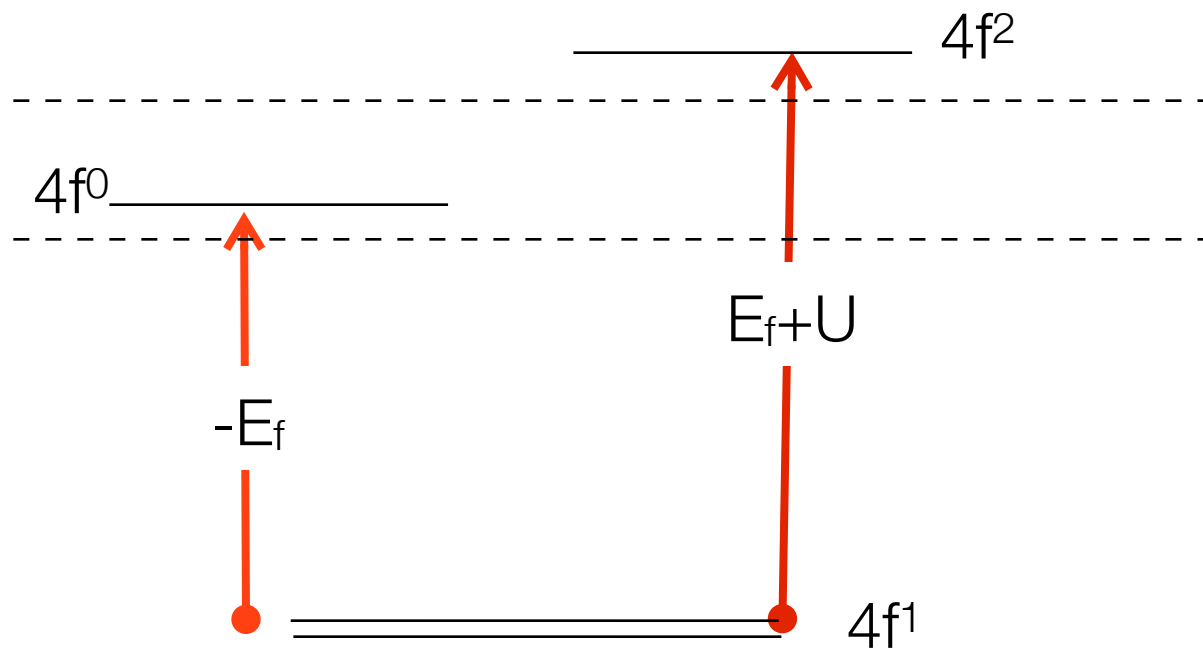


From Anderson to Kondo



Kondo effect

1. Schrieffer Wolff Transformation

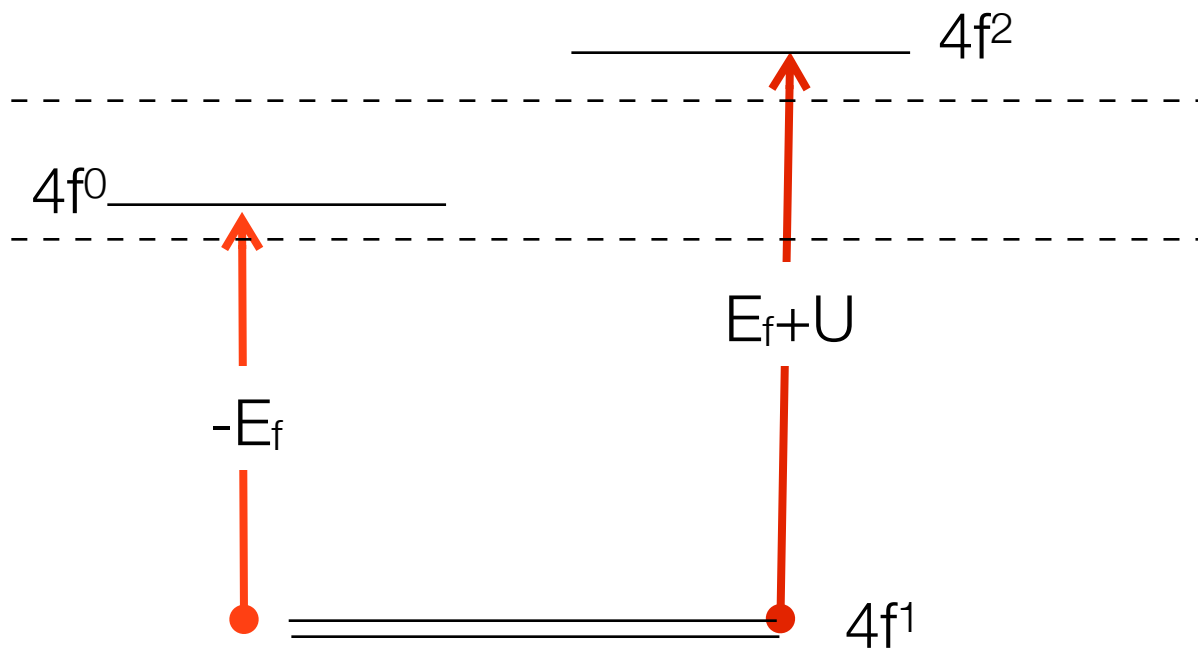


Virtual Valence fluctuations in the singlet channel, induced by hybridization

$$\begin{array}{ll}
 e_{\uparrow}^{-} + f_{\downarrow}^1 \leftrightarrow f^2 \leftrightarrow e_{\downarrow}^{-} + f_{\uparrow}^1 & \Delta E_I \sim U + E_f \\
 h_{\uparrow}^{+} + f_{\downarrow}^1 \leftrightarrow f^0 \leftrightarrow h_{\downarrow}^{+} + f_{\uparrow}^1 & \Delta E_{II} \sim -E_f
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Kondo effect

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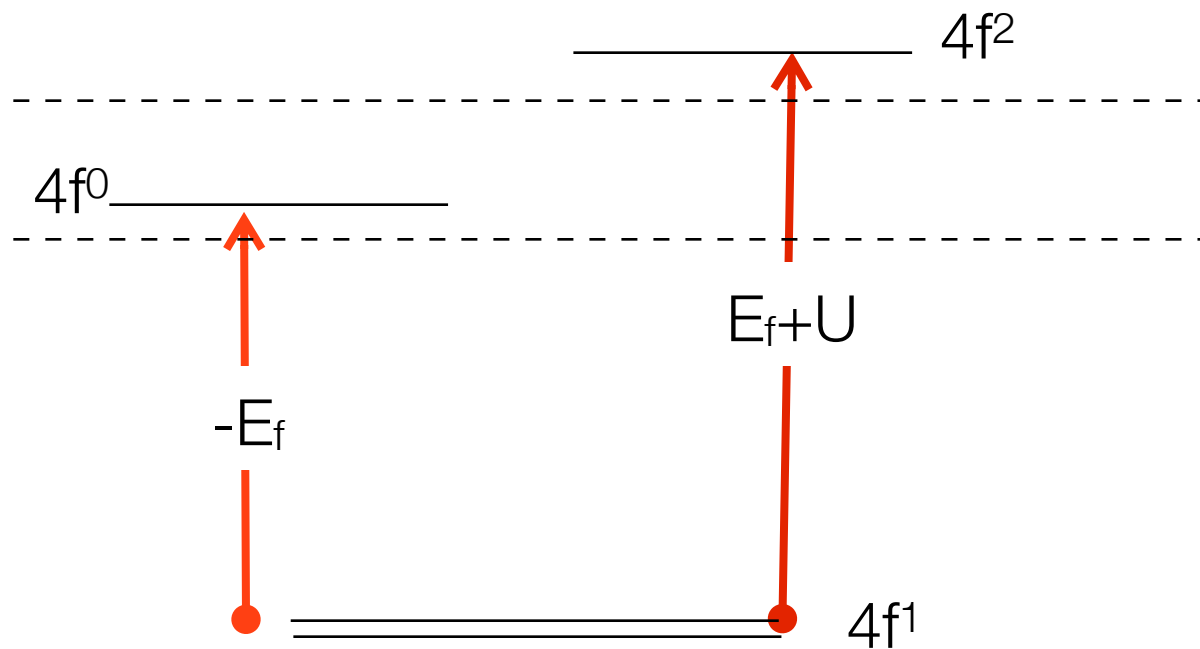
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From second order perturbation theory, the energy of c-f singlets **reduces** by an amount $2J$, where

$$J = V^2 \left[\frac{1}{\Delta E_1} + \frac{1}{\Delta E_2} \right]$$

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$$H_K = -2JP_{S=0} = -2J \left[\frac{1}{4} - \frac{1}{2} \vec{\sigma}_c(0) \cdot \vec{S}_f \right] \rightarrow J \vec{\sigma}_c(0) \cdot \vec{S}_f$$

Antiferromagnetic interaction