

Back to Lagrange's equation

$$\delta \int_{t_1}^{t_2} L(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n, t) dt = 0$$

$$q_i(t) \rightarrow q_i(t) + \alpha \eta_i(t)$$

$$\eta_i(t_1, t_2) = 0$$

$$\frac{d}{d\alpha} \int_{t_1}^{t_2} L dt = \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q_i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} \right) \eta_i dt$$

Choosing individual η 's
to be like

$$\frac{\partial L}{\partial q_i} = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i}$$

Ex $L = \frac{m\dot{x}^2}{2} - V(x)$

Dealing with additional constraints.

$$\sum_a f_a(q_1, \dots, q_n, t) = 0$$

$$\int_1 \left[L(q_1, \dot{q}_1, \dots, t) + \sum_a \lambda_a(t) f_a(q_1, \dots, t) \right] dt$$

$$\frac{\partial L}{\partial q_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} + \sum_a \lambda_a \frac{\partial f_a}{\partial q_j} = 0$$

$\underbrace{\hspace{10em}}_{Q_j^c}$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} = \frac{\partial L}{\partial q_j} + Q_j^c$$

forces due to constraint



$$L = \frac{1}{2} M (\dot{x}^2 + \dot{z}^2) - Mgz$$

$$x = r \sin \theta, z = r \cos \theta$$

$$\sqrt{x^2 + z^2} = a$$

$$L = \frac{1}{2} M (\dot{r}^2 + r^2 \dot{\theta}^2) - Mgr \cos \theta$$

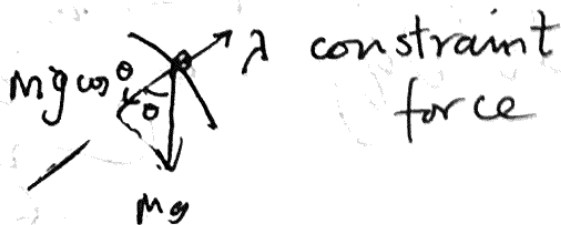
Constraint $r - a = 0$

$$\int (L + \lambda(t)(r - a)) dt = 0$$

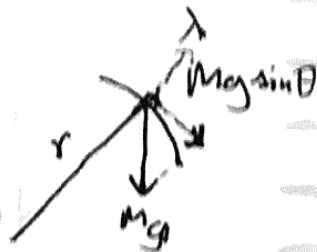
~~$$\int (L + \lambda(t)(r - a)) dt = 0$$~~

Equation of motion

$$M \ddot{r} = M r \dot{\theta}^2 - M g \cos \theta + \lambda(t)$$



$$\frac{d}{dt} (M r^2 \dot{\theta}) = M g r \sin \theta$$



Three equations, including the constraint equation, determining $r(t)$, $\theta(t)$, $\lambda(t)$

Now we use $r(t) = a$

$$0 = M a \dot{\theta}^2 - M g \cos \theta + \lambda$$

$$M a^2 \ddot{\theta} = M g a \sin \theta$$

Second equation $\Rightarrow \frac{M a^2}{2} \dot{\theta}^2 + M g a \cos \theta = E \text{ constant}$

Initially $\dot{\theta} = 0$ $\cos \theta = 1$

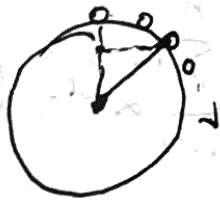
$$\Rightarrow \dot{\theta}^2 = \frac{2g}{a} (1 - \cos \theta)$$

Use that in the first equation

$$\lambda = M g \cos \theta - M a \dot{\theta}^2$$

$$= M g \cos \theta - M a \frac{2g}{a} (1 - \cos \theta)$$

$$= M g (3 \cos \theta - 2)$$



$$\lambda = 0 \quad \text{when} \quad \cos \theta = \frac{2}{3}$$

$$\text{when } \cos \theta < \frac{2}{3} \quad \lambda < 0.$$

$\Rightarrow$ That is not possible in this case
[non-holonomic nature of constraint]

$\Rightarrow$ Ball leaves constraint surface.



Semi-holonomic Constraints

$$f_\alpha(q_1, \dots, q_n; \dot{q}_1, \dots, \dot{q}_n; t) = 0$$

$$\alpha = 1, \dots, m$$

Let us take a special case:

$$f_\alpha = \sum_{k=1}^n a_{\alpha k}(t) \dot{q}_k + a_{\alpha 0}(t) = 0$$

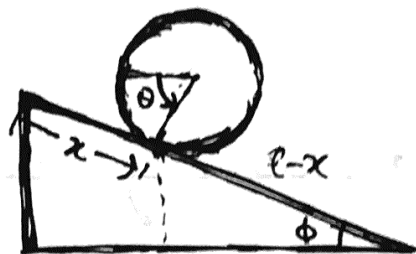
$$\delta \int_{t_1}^{t_2} \left[L + \sum_{\alpha=1}^m \mu_\alpha f_\alpha \right] dt$$

$$\frac{\partial L}{\partial q_k} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) - \frac{d}{dt} \sum_{\alpha=1}^m \mu_\alpha(t) a_{\alpha k}(t) \stackrel{!}{=} 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_k} \right) - \frac{\partial L}{\partial q_k} = \underbrace{\frac{d}{dt} \sum_{\alpha=1}^m \mu_\alpha a_{\alpha k}}_{Q_k^c}$$

$$\left(= \sum_{\alpha=1}^m \frac{d}{dt} \left(\mu_\alpha \frac{\partial f_\alpha}{\partial \dot{q}_k} \right) \right)_{\text{in this case}}$$

Generalization for $f_\alpha(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n; t)$ possible.



$$dx = r d\theta$$

$$\dot{x} = r \dot{\theta}$$

r is constant.

$$T = \frac{1}{2} M \dot{x}^2 + \frac{1}{2} M r^2 \dot{\theta}^2$$

$$V = Mg(l-x) \sin \phi$$

$$f = r \ddot{\theta} - \ddot{x} = 0$$

$$a_{\theta} = r \quad a_x = -1$$

$$\int_0^2 dt \left[\frac{1}{2} M \dot{x}^2 + \frac{1}{2} M r^2 \dot{\theta}^2 - Mg(l-x) \sin \phi + \mu(r\dot{\theta} - \dot{x}) \right] = 0$$

$$M \ddot{x} - Mg \sin \phi - \mu \dot{x} = 0$$

$$M r^2 \ddot{\theta} + \mu \dot{x} r = 0$$

Let us call $(-\mu \dot{x}) = \lambda$.

$$M \ddot{x} = Mg \sin \phi - \lambda$$

$$M r^2 \ddot{\theta} = \lambda r$$

λ is the force applied by the plane.
Force & torque equations.



These two, along with

$$\boxed{r\dot{\theta} = \dot{x}}$$

are the three equations that determine $x(t)$, $\theta(t)$, $\lambda(t)$

$$r\dot{\theta} = \dot{x} \Rightarrow r\ddot{\theta} = \ddot{x}$$

$$Mr^2\ddot{\theta} = \lambda r \quad \text{then says } M\ddot{x} = \lambda$$

$$\text{Using it in } M\ddot{x} = mg \sin \phi - \lambda$$

$$\ddot{x} = \frac{g \sin \phi}{2}$$

$$\lambda = \frac{mg \sin \phi}{2}$$

$$\ddot{\theta} = \frac{g \sin \phi}{2r}$$

←
← Constants
←

Of course this is a very fancy way of solving a simpler problem. In this particular case,

We could have integrated $\dot{x} = r\dot{\theta}$ and eliminate θ .
(that is not possible, in general, for semi holonomic constraints).

$$T = \frac{1}{2} M \dot{x}^2 + \frac{1}{2} M r^2 \dot{\theta}^2 = M \dot{x}^2$$

$$\delta \int [M \dot{x}^2 - Mg(l-x) \sin \phi] dt$$

$$2M \ddot{x} = Mg \sin \phi$$

$$\ddot{x} = \frac{g \sin \phi}{2}$$