

Variational Principle

Hamilton's Principle: Action integral is stationary on the actual path.



Paths
in configuration
space.

$$I = \int_{t_1}^{t_2} L dt$$

Action integral

[Monogenic: All applied force comes from a generalized scalar potential]

In other words:

$$\delta I = \delta \int_{t_1}^{t_2} L(q_1, \dots, q_n, \dot{q}_1, \dots, \dot{q}_n, t) dt = 0$$

Comments on least vs. Stationary

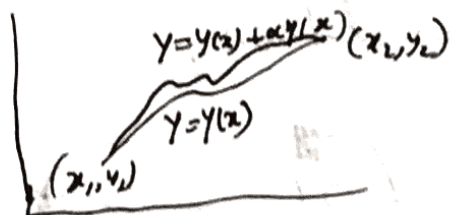
Calculus of Variations

One dim example

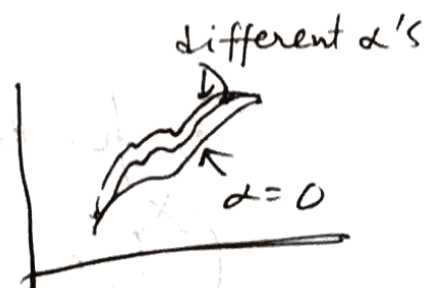
$f(y, \dot{y}, x)$ defined on a path $y = y(x)$

$$\dot{y} = \frac{dy}{dx}$$

$$J = \int_{x_1}^{x_2} f(y, \dot{y}, x) dx$$



Deformed path



$$y(x, \alpha) = y(x, 0) + \alpha \eta(x)$$

$\uparrow$
 $y(x)$

$$\eta(x_1) = \eta(x_2) = 0$$

$$J(\alpha) = \int_{x_1}^{x_2} f(y(x, \alpha), \dot{y}(x, \alpha), x) dx$$

Stationary path condition

$$\left(\frac{dJ(\alpha)}{d\alpha} \right)_{\alpha=0} = 0$$

$$\frac{dJ}{d\alpha} = \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} \frac{\partial y}{\partial \alpha} + \frac{\partial f}{\partial \dot{y}} \frac{\partial \dot{y}}{\partial \alpha} \right) dx$$

Notice that $\frac{\partial \dot{y}}{\partial \alpha} = \frac{\partial^2 y(x)}{\partial x \partial \alpha} = \frac{\partial^2 y(x)}{\partial x \partial \alpha}$

So, for the ~~second~~ second term,

$$\int_{x_1}^{x_2} \frac{\partial f}{\partial \dot{y}} \frac{\partial \dot{y}}{\partial \alpha} dx = \int_{x_1}^{x_2} \frac{\partial f}{\partial \dot{y}} \frac{\partial}{\partial x} \left(\frac{\partial y}{\partial \alpha} \right) dx$$

$$= \left. \frac{\partial f}{\partial \dot{y}} \frac{\partial y}{\partial \alpha} \right|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial \dot{y}} \right) \frac{\partial y}{\partial \alpha} dx \left. \vphantom{\int_{x_1}^{x_2}} \right\} \begin{array}{l} \text{integration} \\ \text{by} \\ \text{parts} \end{array}$$

If the varied ~~curve~~ curve passes through the boundary points $\Rightarrow \frac{\partial y}{\partial \alpha} = \eta = 0$ at x_1, x_2 .

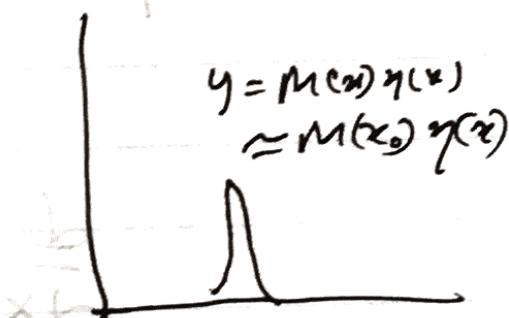
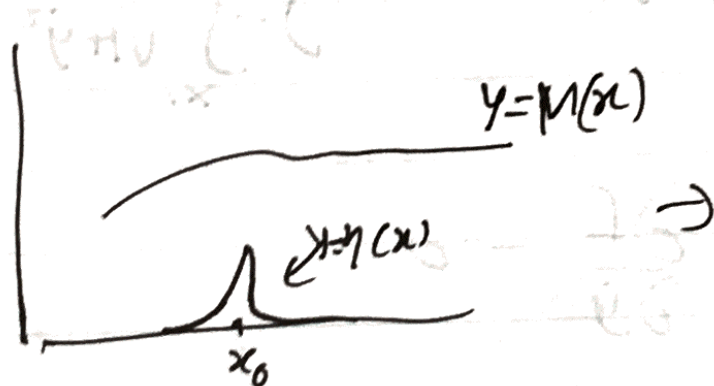
Thus
$$\frac{\partial J}{\partial \alpha} = \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial \dot{y}} \right) \right] \frac{\partial y}{\partial \alpha} dx$$

going to $\alpha = 0$

$$\int_{x_1}^{x_2} M(x) \eta(x) dx = 0$$

$$M(x) = \frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial \dot{y}}$$

for arbitrary "smooth" η .



$$\int_{x_1}^{x_2} M \eta dx \approx M(x_0) \int_{x_1}^{x_2} \eta(x) dx = 0$$

$$M(x_0) = 0 \text{ for every } x_0$$

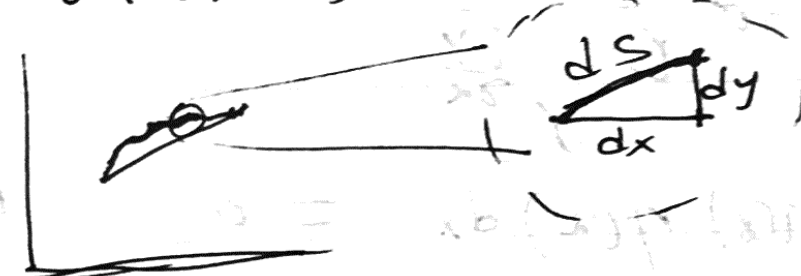
$$\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial \dot{y}} = 0$$

Usual notation
$$\frac{\partial y}{\partial \alpha} \delta \alpha = \delta y$$

$$\delta J = \int_{x_1}^{x_2} \left(\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial \dot{y}} \right) \delta y dx$$

Examples:

Shortest distance between two point


$$ds^2 = dx^2 + dy^2$$
$$ds = \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$
$$= \sqrt{1 + \dot{y}^2} dx$$

$$f = \sqrt{1 + \dot{y}^2}$$

$$J = \int_{x_1}^{x_2} \sqrt{1 + \dot{y}^2} dx$$

$$\frac{d}{dx} \frac{\partial f}{\partial \dot{y}} = 0$$

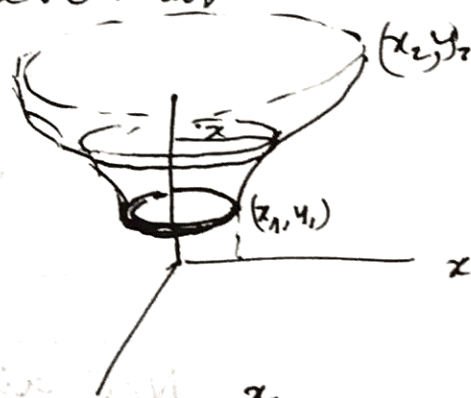
$$\frac{\dot{y}}{\sqrt{1 + \dot{y}^2}} = c \Rightarrow \dot{y} = a$$

$$a = \frac{c}{\sqrt{1 - c^2}}$$

$$y = ax + b$$

Straight line

Minimum Surface of revolution



$$\text{Area} = \int_{x_1}^{x_2} 2\pi x \, ds = 2\pi \int_{x_1}^{x_2} x \sqrt{1+y'^2} \, dx$$

$$f = x \sqrt{1+y'^2}$$

$$\frac{xy'}{\sqrt{1+y'^2}} = a$$

$$(x^2 - a^2)y' = a^2$$

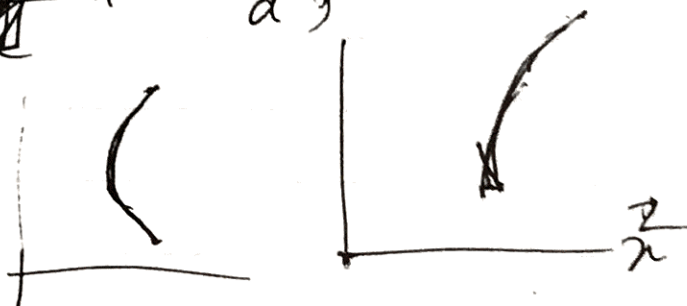
$$y = \int \frac{dy}{dx} dx = \int \frac{a \, dx}{\sqrt{x^2 - a^2}}$$

$$x = a \cosh t$$

$$\int \frac{a \sinh t \, dt}{a \cosh t} = a \tanh t + b$$

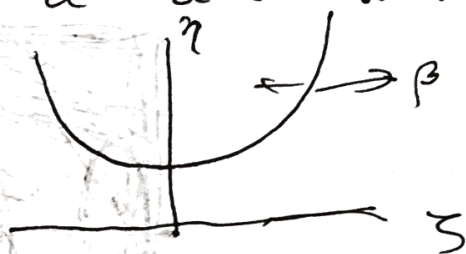
$$x = a \cosh t = a \cosh \frac{y-b}{a}$$

Even b like $\rightarrow$



Catenary: Define $\frac{x}{a} = \xi$ $\frac{y}{a} = \eta$ $\frac{b}{a} = \beta$

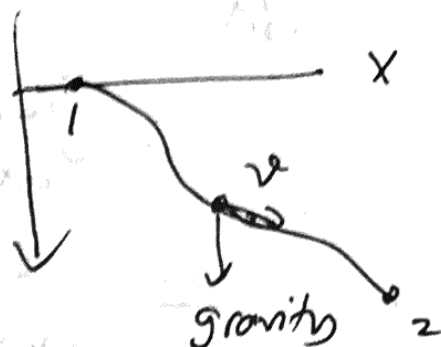
$$\xi = \text{ch}(\eta - \beta)$$



Not all ~~initial~~ initial and final conditions lead to unique ~~solutions~~ solutions. Some lead to none.

The brachistochrone problem

t_{12}



$$t_{12} = \int \frac{ds}{v}$$

$$\frac{1}{2}mv^2 = mgy$$

$$v = \sqrt{2gy}$$

$$t_{12} = \int \frac{\sqrt{1+y'^2}}{\sqrt{2gy}} dx$$

$$f = \sqrt{\frac{1+y'^2}{2gy}}$$

$$\frac{\partial f}{\partial y} = -\frac{1}{2y} \sqrt{\frac{1+y'^2}{2gy}}$$

$$\frac{\partial f}{\partial y'} = \frac{1}{\sqrt{2gy}} \frac{y'}{\sqrt{1+y'^2}}$$

$$\frac{\partial f}{\partial y} = \frac{d}{dx} \frac{\partial f}{\partial y'}$$

looks a bit complicated.

If $f = f(y, y')$, no x dependence,

$$\frac{df}{dx} = y' \frac{\partial f}{\partial y} + y'' \frac{\partial f}{\partial y'} + \frac{\partial f}{\partial x} = 0$$

$$\begin{aligned}
 \frac{df}{dx} &= \dot{y} \frac{\partial f}{\partial y} + \dot{y} \frac{\partial f}{\partial \dot{y}} \\
 &= \dot{y} \frac{d}{dx} \frac{\partial f}{\partial \dot{y}} + \dot{y} \frac{\partial f}{\partial y} \\
 &= \frac{d}{dx} \left(\dot{y} \frac{\partial f}{\partial \dot{y}} \right)
 \end{aligned}$$

That means $\frac{d}{dx} \left(\dot{y} \frac{\partial f}{\partial \dot{y}} - f \right) = 0$

or $\dot{y} \frac{\partial f}{\partial \dot{y}} - f = \text{constant}.$

[Aside: $L = \frac{m}{2} \dot{x}^2 - V(x)$

$$\begin{aligned}
 \dot{x} \frac{\partial L}{\partial \dot{x}} - L &= m \dot{x}^2 - \left\{ \frac{m}{2} \dot{x}^2 - V(x) \right\} \\
 &= \frac{1}{2} m \dot{x}^2 + V(x)
 \end{aligned}$$

Energy = constant

Using it for the ~~brachistochrone~~ brachistochrone problem:

$$\begin{aligned}
 & \frac{y^2}{\sqrt{2gy} \sqrt{1+y^2}} - \frac{\sqrt{1+y^2}}{\sqrt{2gy}} \\
 &= \frac{y^2}{\sqrt{2gy} \sqrt{1+y^2}} - \frac{1+y^2}{\sqrt{2gy} \sqrt{1+y^2}} \\
 &= - \frac{1}{\sqrt{2gy} \sqrt{1+y^2}} \rightarrow \text{constant}
 \end{aligned}$$

$$\sqrt{y} \sqrt{1+y^2} = k$$

$$\frac{dy}{dx} = \sqrt{\frac{k^2 - y}{y}}$$

$$\int dx = \int \sqrt{\frac{y}{k^2 - y}} dy$$

Natural transformation

$$y = k^2 \sin^2 \theta$$

$$\int \sqrt{\frac{y}{k^2 - y}} dy = 2k^2 \int \sin^2 \theta d\theta$$

$$= k^2 \int (1 - \cos 2\theta) d\theta$$

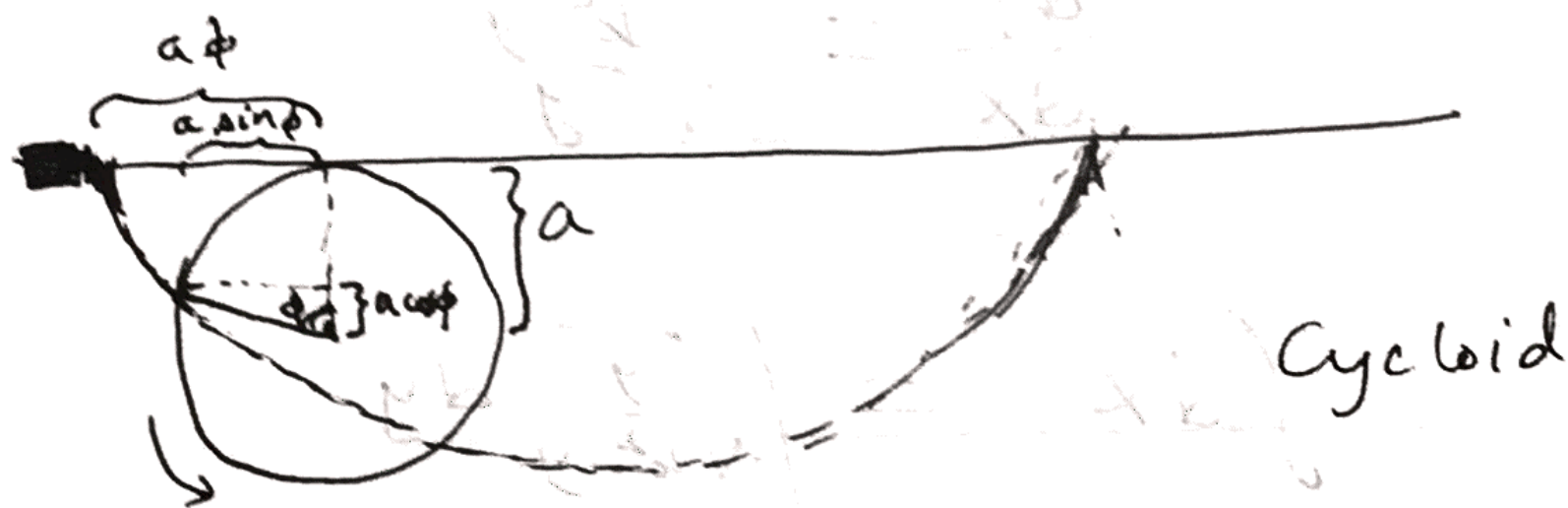
$$= k^2 \left(\theta - \frac{\sin 2\theta}{2} \right) + C$$

Use $2\theta = \phi$ and $\frac{k^2}{2} = a$
to get

$$x = a(\phi - \sin \phi)$$

$$y = a(1 - \cos \phi)$$

We used the initial condition $x=0, y=0$
to set $C=0$.

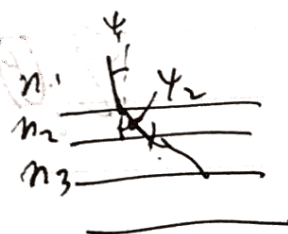
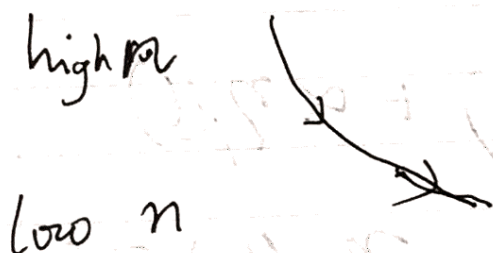


What is this 'conserved' quantity

$$\frac{1}{\sqrt{2gy} \sqrt{1+\dot{y}^2}}$$

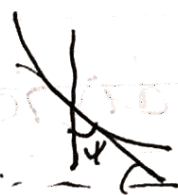
Analogy to Fermat's Principle and Snell's law.

$$\int_1^2 n ds \leftrightarrow \int_1^2 \frac{1}{v} ds$$



$$n_1 \sin \psi_1 = n_2 \sin \psi_2 = n_3 \sin \psi_3 = \dots$$

$n \sin \psi$ is 'conserved'.



$$\tan\left(\frac{\pi}{2} - \psi\right) = y$$

$$\begin{aligned} \sqrt{1+y^2} &= \sqrt{\sec^2\left(\frac{\pi}{2} - \psi\right)} = \sec\left(\frac{\pi}{2} - \psi\right) \\ &= \frac{1}{\sin \psi} \end{aligned}$$

$$n \sin \psi \rightarrow \frac{\sin \psi}{v} = \frac{1}{\sqrt{2g} \sqrt{1+y^2}}$$