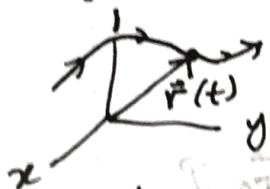


Notations and Preliminary Discussions

Particle motion



$$\vec{v} = \frac{d\vec{r}}{dt} = \dot{\vec{r}}, \quad \vec{a} = \frac{d\vec{v}}{dt} = \frac{d^2\vec{r}}{dt^2}$$

Linear momentum:

$$\vec{p} = m\vec{v}$$

Newton 2: $\vec{F} = \frac{d\vec{p}}{dt} = \dot{\vec{p}}$

[Newton 1: $\vec{F} = 0$, m const, $\Rightarrow \dot{\vec{v}} = 0$]

Momentum turns out to be something more fundamental than just being a concept derived from velocity and mass: think of particles which could change into other particles.

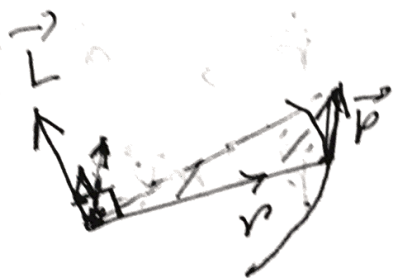
Conservation of Linear momentum

Total force zero $\Rightarrow$ Linear momentum conserved $\Rightarrow \vec{p}$ constant

$$\vec{F} = 0 \Rightarrow \vec{p} \text{ constant}$$

Angular Momentum

$$\vec{L} = \vec{r} \times \vec{p}$$



$$\begin{aligned}\vec{L} &= \vec{r} \times \vec{p} \\ &= m(\vec{r} \times \vec{v})\end{aligned}$$

Relation to Kepler 2: $\frac{1}{2} \vec{r} \times \vec{v}$ = area swept at origin.

Note that $\frac{d\vec{L}}{dt} = \frac{d}{dt}(\vec{r} \times \vec{p})$

$$= \frac{d\vec{r}}{dt} \times \vec{p} + \vec{r} \times \frac{d\vec{p}}{dt}$$

$$= \underbrace{\vec{v} \times (m\vec{v})}_0 + \vec{r} \times \vec{F}$$

$\vec{N} = \vec{r} \times \vec{F}$ is the  torque.

$$\frac{d\vec{L}}{dt} = \vec{N}$$

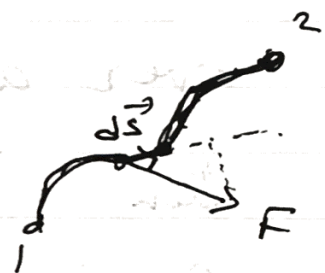
Conservation of Angular Momentum

Total torque zero $\Rightarrow$ Angular Mom. conserved
 $\vec{N} = 0 \Rightarrow \vec{L}$ constant

Ex. Central

Work done by external force $\vec{F}$

$$W_{12} = \int_1^2 \vec{F} \cdot d\vec{s}$$



$$\int_1^2 \vec{F} \cdot d\vec{s} = \int_1^2 m \frac{d\vec{v}}{dt} \cdot \vec{v} dt$$

$$= \int_1^2 \frac{d}{dt} \left(\frac{1}{2} m \vec{v}^2 \right) dt$$

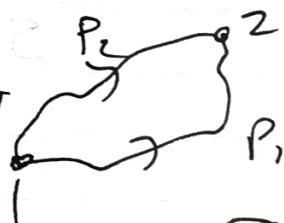
$$= \frac{m}{2} (v_2^2 - v_1^2)$$

$$W_{12} = T_2 - T_1$$

Consider $\vec{F} = \vec{F}(\vec{r})$

If W_{12} is path-independent

$$W_{12}(P_1) = W_{12}(P_2)$$



$$\oint \vec{F} \cdot d\vec{r} = W_{12}(P_2) - W_{12}(P_1) = 0$$

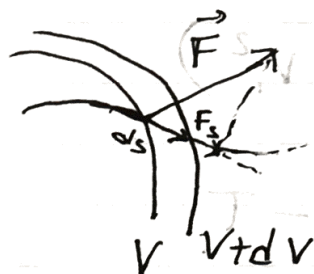
Conservative force/system

$$\Rightarrow \oint \vec{F} \cdot d\vec{s} = 0$$

If $\vec{F} = \vec{F}(\vec{r})$ satisfied $\oint \vec{F} \cdot d\vec{s} = 0$ over any closed path, then $\vec{F}(\vec{r}) = -\nabla V(\vec{r})$ for some function $V(\vec{r})$.

$$\int_P^r \vec{F} \cdot d\vec{s} = V(0) - V(r)$$

Defines $V(r)$ upto a constant.



$$\vec{F} \cdot d\vec{s} = F_s ds = -\frac{\partial V}{\partial s} ds$$

Since

$$W_{12} = T_2 - T_1$$

and

$$W_{12} = V_1 - V_2$$

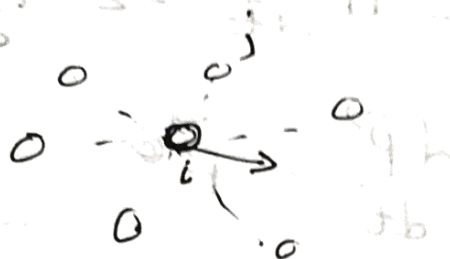
for conservative systems

$$T_1 + V_1 = T_2 + V_2$$

Conservation of Energy

Forces conservative $\Rightarrow T + V$ conserved -

System of Particles



Force on i th particle due to j th particle

$$\vec{F}_i = \sum_j \vec{F}_{ji} + \vec{F}_i^{(e)}$$

↑
External force

Newton 3: $\vec{F}_{ij} = -\vec{F}_{ji}$
(Weak) Law of action & reaction

Note that then $\vec{F}_{ii} = 0$

$$\frac{d^2}{dt^2} \sum_i m_i \vec{r}_i = \sum_i \vec{F}_i^{(e)} + \underbrace{\sum_{ij} \vec{F}_{ji}}_{\text{vanishes}}$$

Think of the two particle example
 $\sum_i \vec{F}_i = \vec{F}_{12} + \vec{F}_{21}$

Define center of mass

$$\vec{R} = \frac{\sum_i m_i \vec{r}_i}{\sum_i m_i} = \frac{\sum_i m_i \vec{r}_i}{M}$$

M total mass

$$M \frac{d^2 \vec{R}}{dt^2} = \sum_i \vec{F}_i^{(e)} = \vec{F}^{(e)}$$

If we define total momentum to be

$$\vec{P} = M \frac{d\vec{R}}{dt} = \sum m_i \frac{d\vec{r}_i}{dt}$$

then

$$\frac{d\vec{P}}{dt} = \vec{F}^{\text{ext}}$$

Conservation theorem for total linear momentum

~~Total~~ external force is zero $\Rightarrow$ Total linear momentum is conserved

Interactions between the particles only lead to exchange of momentum among ~~themselves~~ themselves. The total momentum remains unaffected by those interactions.

What about total angular momentum?

$$\frac{dL}{dt} = \frac{d}{dt} \sum_i \vec{r}_i \times \vec{p}_i = \sum_i \vec{r}_i \times \dot{\vec{p}}_i$$

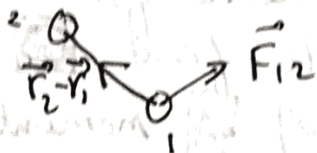
~~Remember~~ $\frac{d}{dt} \vec{r}_i \times \vec{p}_i = \vec{v}_i \times (m \vec{v}_i) = 0$, remember?

So

$$\frac{d\vec{L}}{dt} = \sum_i \vec{r}_i \times \vec{F}_i = \sum_i \vec{r}_i \times \vec{F}_i^{(e)} + \sum_i \vec{r}_i \times \vec{F}_{ji}$$

How do we deal with this second term?
Look at the two particle case again.

$$\vec{r}_1 \times \vec{F}_{12} + \vec{r}_2 \times \vec{F}_{21} = (\vec{r}_2 - \vec{r}_1) \times \vec{F}_{12}$$



If $\vec{F}_{12}$ is in the direction of $\vec{r}_2 - \vec{r}_1$ (strong law of action and reaction) then $(\vec{r}_2 - \vec{r}_1) \times \vec{F}_{12} = 0$



In that case

$$\sum_i \vec{r}_i \times \vec{F}_{ji} = \sum_{i,j} (\vec{r}_i - \vec{r}_j) \times \vec{F}_{ji} = 0$$

$$\Rightarrow \frac{d\vec{L}}{dt} = \sum_i \vec{r}_i \times \vec{F}_i^{(e)} = \vec{N}^{(e)} \leftarrow \text{External torque}$$

In that case, we have

Conservation theorem for total angular
momentum! $N_{\text{external torque}}^{\text{total}}$

⇒ Angular momentum
is conserved.

Before we push on, we should note that
 $\vec{F}_{ij} = -\vec{F}_{ji}$ and $\vec{F}_{ij} \times (\vec{r}_i - \vec{r}_j) = 0$ are not always
true, e.g. magnetic forces/Biot-Savart law.

Physicists tend to find nontrivial
generalizations of momentum/angular
momentum/energy when things do
not work out. For Biot-Savart law, one
needs to invoke momentum/angular
momentum of electromagnetic fields.

The more sophisticated formulation of
dynamics allows us to find such
~~generalization~~ generalization by
relating conservation laws to symmetries.

$$\vec{B}_{12} = 0 \quad \vec{B}_{21} \neq 0$$

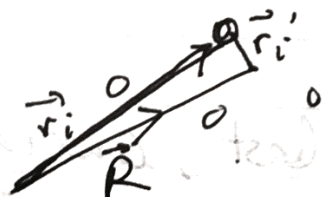
Note that the total angular momentum could be broken up into the angular momentum of all particles concentrated at the center of mass, plus the angular momentum of the particles around the center of mass.

$$\vec{r}_i = \vec{r}_i' + \vec{R}$$

$$\vec{R} = \frac{\sum m_i \vec{r}_i}{\sum m_i}$$

$$\left(\sum_i m_i \vec{r}_i' = 0 \right)$$

$$\vec{v}_i = \vec{v}_i' + \vec{V}$$



$$\vec{v}_i' = \frac{d\vec{r}_i'}{dt}, \quad \vec{V} = \frac{d\vec{R}}{dt}$$

$$\left(\sum_i m_i \vec{v}_i' = 0 \right)$$

$$\sum_i m_i \vec{v}_i' = \vec{p}'$$

$$\vec{L} = \sum_i \vec{r}_i \times m_i \vec{v}_i = \sum_i \vec{R} \times m_i \vec{V}$$

$$+ \sum_i \vec{r}_i' \times m_i \vec{v}_i'$$

$$+ \sum_i \vec{r}_i' \times (m_i \vec{V}) + \sum_i \vec{R} \times m_i \vec{v}_i'$$

$\swarrow$ $\uparrow$ Cross terms
 $(\sum_i m_i \vec{r}_i') \times \vec{V}$

$\searrow$
 $\vec{R} \times \sum_i m_i \vec{v}_i'$



$$\vec{L} = \vec{R} \times M \vec{V} + \sum_i \vec{r}_i' \times m_i \vec{v}_i'$$

$$\vec{L} = \vec{R} \times \vec{P} + \sum \vec{r}_i' \times \vec{p}_i'$$

↑
'C.M.'
angular mom.

↑
Angular mom.
around C.M



Under some conditions
they could each be
approximately conserved.

Last but not the least, energy.

$$W_{12} = \sum_i \int_1^2 \vec{F}_i \cdot d\vec{s}_i = \sum_i \int_1^2 d\left(\frac{1}{2} m_i v_i^2\right) = T_2 - T_1$$

$$\text{Where } T = \frac{1}{2} \sum_i m_i v_i^2$$

$$\begin{aligned} \text{But } W_{12} &= \sum_i \int_1^2 \vec{F}_i \cdot d\vec{s}_i \\ &= \sum_i \int_1^2 \vec{F}_i^{(e)} \cdot d\vec{s}_i + \sum_{i,j} \int_1^2 \vec{F}_{ji} \cdot d\vec{s}_i \end{aligned}$$



If $\vec{F}_i^{(e)} = -\nabla_i V_i(\vec{r}_i)$
 and $\vec{F}_{ji} = -\nabla_i V_{ij}(\vec{r}_i - \vec{r}_j),$

~~$\vec{F}_{ji} = -\nabla_j V_{ji}(\vec{r}_j - \vec{r}_i)$~~ with
 $V_{ij}(\vec{r}_i - \vec{r}_j) = V_{ji}(\vec{r}_j - \vec{r}_i)$

[special case $V_{ij}(\vec{r}_i - \vec{r}_j) = V_{ij}(|\vec{r}_i - \vec{r}_j|)$]

then $\sum_i \int \vec{F}_i^{(e)} \cdot d\vec{S}_i = -\sum_i \int \nabla_i V_i \cdot d\vec{S}_i$

~~$\vec{F}_{ji} = -\nabla_j V_{ji}(\vec{r}_j - \vec{r}_i)$~~
 $\frac{1}{2} \sum_{ij} (\nabla_i V_{ji} \cdot d\vec{S}_i + \nabla_j V_{ij} \cdot d\vec{S}_j)$

The first term is just
 $-\sum_i V_i$

The second one could be simplified using

$$\nabla_i V_{ji} \cdot d\vec{S}_i + \nabla_j V_{ij} \cdot d\vec{S}_j$$

$$\nabla_i V_{ji}(\vec{r}_i - \vec{r}_j) \cdot d\vec{S}_i + \nabla_j V_{ji}(\vec{r}_j - \vec{r}_i) \cdot d\vec{S}_j$$

$$\nabla_i V_{ji} \cdot d\vec{S}_i = -\nabla_j V_{ji} \cdot d\vec{S}_i = \nabla_j V_{ji} \cdot d\vec{S}_j$$

Then the second term is just

$$\frac{1}{2} \sum_{ij} V_{ij} (\vec{r}_i - \vec{r}_j)^2$$

So if $V = \sum_i V_i(\vec{r}_i) + \frac{1}{2} \sum_{ij} V_{ij}(\vec{r}_i - \vec{r}_j)$

Then once more $\left. \frac{d}{dt} (T+V) \right|_1 = 0$

or $T+V$ is conserved

→ Theorem of Conservation of Energy for a system of particles.

Note that
$$\begin{aligned} T &= \frac{1}{2} \sum m_i v_i^2 \\ &= \frac{1}{2} \sum m_i (V + v_i')^2 \\ &= \frac{1}{2} M V^2 + \frac{1}{2} \sum m_i v_i'^2 \end{aligned}$$

$\begin{matrix} \text{C.M.} \\ \text{K.E.} \\ \downarrow \end{matrix}$ $\begin{matrix} \text{'relative'} \\ \text{K.E.} \end{matrix}$ $+ \sum m_i v_i'^2$ $\begin{matrix} \text{'0'} \\ \text{V} \end{matrix}$

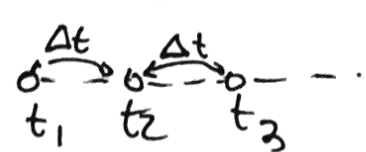
So $T = \frac{1}{2} M V^2 + \frac{1}{2} \sum m_i v_i'^2$

A brief word on numerically solving Newton's Equation.

Quedim: $\dot{x}(t) = v(t)$
 $\dot{v}(t) = f(x(t))$ $f(x) = \frac{F(x)}{m}$

'Obvious' discretization

$$x(t_{n+1}) = x(t_n) + v(t_n) \Delta t$$

$$v(t_{n+1}) = v(t_n) + f(x(t_n)) \Delta t$$


The diagram shows a horizontal timeline with three points labeled t_1 , t_2 , and t_3 . Above t_1 and t_2 are small circles representing states. A solid arrow points from the circle at t_1 to the circle at t_2 , with Δt written above it. A dashed arrow points from the circle at t_2 to the circle at t_3 , also with Δt written above it. The timeline continues with a dashed line after t_3 .

Problem; error $O(\Delta t)$, lack of reversibility, stability, ...

Cleverer discretization: Leapfrog method

