

Classification of Singular Point

$$L[u] = 0 \quad \text{with} \quad L = \frac{d^2}{dz^2} + p(z) \frac{d}{dz} + q(z)$$

Sturm-Liouville equation.

If p, q are $\blacksquare$ analytic $\blacksquare$ around $z=z_0$, it is regular (ordinary) point. One can Taylor expand u around $z=z_0$ and find a solution.

The singular point $z=z_0$ is called ^aregular ^{singular} point if $(z-z_0)p(z)$ and $(z-z_0)^2 q(z)$ are $\blacksquare$ restrictions of analytic functions ^{p, q} in the neighborhood of z_0 .

Otherwise, the singular pt is an irregular singular pt. *

$$\text{Let } P(z) = \sum_{n=0}^{\infty} p_n (z-z_0)^n = (z-z_0) p(z)$$

$$Q(z) = \sum_{n=0}^{\infty} q_n (z-z_0)^n = (z-z_0)^2 q(z)$$

Try a solution of the form $u(z) = (z-z_0)^\alpha \left[1 + \sum_{n=1}^{\infty} a_n (z-z_0)^n \right]$

$$a_0 = 1$$

$\blacksquare$ This is the Frobenius method of series solution.

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* $u'' + \frac{B}{z^2} u' + \frac{2B}{z^3} u = 0$ has a soln understood By

Like $u(z) = A e^{B/z}$
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Essential singularity!

$$(u')' = \left(\frac{AB}{z^2} e^{B/z} \right)' = \frac{AB^2}{z^4} e^{B/z} + \frac{2AB}{z^3} e^{B/z} = \frac{(B^2 + 2B)}{z^4} e^{B/z}$$

$$u' = -\frac{AB}{z^2} e^{B/z} = -\frac{B}{z} u$$

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$$u''(z) = \sum_{n=0}^{\infty} (\alpha+n)(\alpha+n-1) a_n (z-z_0)^{n+\alpha-2}$$

$$p(z)u(z) = \sum_{m=0}^{\infty} p_m (z-z_0)^{m-1} \sum_{k=0}^{\infty} (\alpha+k) a_k (z-z_0)^{k+\alpha-1}$$

$$q(z)u(z) = \sum_{m=0}^{\infty} q_m (z-z_0)^{m-2} \sum_{k=0}^{\infty} a_k (z-z_0)^{k+\alpha}$$

$$\Rightarrow (\alpha+n)(\alpha+n-1)a_n + \sum_{k=0}^n (\alpha+k)p_{n-k} a_k + \sum_{k=0}^n q_{n-k} a_k = 0$$

$n=0$ gives

$$\alpha(\alpha-1)a_0 + \alpha p_0 a_0 + q_0 a_0 = 0$$

So $\alpha(\alpha-1) + p_0 \alpha + q_0 = 0$ Indicial Eqn

This is a quadratic equation for α with two roots $\alpha_{\pm} = \frac{1-p_0 \pm \sqrt{(1-p_0)^2 - 4q_0}}{2}$

Starting with an α root one can obtain the other coefficients

$$[\alpha(\alpha+1) + (\alpha+1)p_0 + q_0]a_1 + \alpha p_1 + q_1 = 0$$

Determines a_1

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$$[(\alpha+n)(\alpha+n-1) + (\alpha+n)p_0 + q_0]a_n + [(\alpha+n-1)p_1 + q_1]a_{n-1} + \dots$$

Determines an
interm of $a_1, a_2, \dots$

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$$+ [(\alpha+1)p_{n-1} + q_{n-1}]a_1 + \alpha p_n + q_n = 0$$

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unless $(\alpha+n)(\alpha+n-1) + (\alpha+n)p_0 + q_0 = 0$ for some n . When we can determine these α we get convergent series (not showing that).

This observation leads to the different possible cases we have to deal with

$$\text{Call } \alpha_{\pm} = \frac{1}{2}(1 - p_0 \pm \sqrt{(1-p_0)^2 - 4q_0})$$

$\alpha_+ - \alpha_- = n$, a non negative integer,

We have a problem. In that case, the "series"

for α_- cannot be continued, since $(\alpha_+ + n)(\alpha_+ + n - 1) + (\alpha_+ + n)p_0 + q_0 = \alpha_+ (\alpha_+ - 1) + \alpha_+ p_0 + q_0 = 0$

So there are three cases

a) $\alpha_+ - \alpha_-$ not an integer!

Then we get two independent solutions

$$u_{\pm}(z) = (z - z_0)^{\alpha_{\pm}} \left[1 + \sum_{n=1}^{\infty} a_n^{(\pm)} (z - z_0)^n \right] = (z - z_0)^{\alpha_{\pm}} f_{\pm}(z)$$

b) $\alpha_+ = \alpha_- = \alpha$!

$$\text{We get only one series } u(z) = (z - z_0)^{\alpha} \left[1 + \sum_{n=1}^{\infty} a_n (z - z_0)^n \right] = (z - z_0)^{\alpha} f(z)$$

c) $\alpha_+ = \alpha_- = n$, a positive integer!

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We only get one solution from the method, namely

$$u_+(z) = (z - z_0)^{\alpha_+} \left[1 + \sum_{n=1}^{\infty} a_n^{(+)} (z - z_0)^n \right] = (z - z_0)^{\alpha_+} f_+(z)$$

$f_+(z)$, $f_-(z)$ have value 1 at $z = z_0$ and are analytic in the neighborhood of z_0 .

With second order differential equations, we expect two independent solutions.

What to do in case b) and c) ?

Here, the Wronskian comes to rescue.

If we have only one solution, $u_+(z)$, and we are try to find another independent solution, $u_-(z)$

$$W(z) = \begin{vmatrix} u_+(z) & u_-(z) \\ u'_+(z) & u'_-(z) \end{vmatrix} \quad \text{must be non zero somewhere.}$$

$$\text{We also know that } W(z) = C \exp \left\{ - \int^z p(\xi) d\xi \right\}$$

$$\begin{aligned} W(z) &= u_+(z) u'_-(z) - u'_+(z) u_-(z) = \frac{u_+(z) u'_-(z) - u'_+(z) u_-(z)}{u_+(z)^2} u_+(z)^2 \\ &= (u_+(z))^2 \frac{d}{dz} \left(\frac{u_-(z)}{u_+(z)} \right) \end{aligned}$$

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$$\text{So } \frac{d}{dz} \left(\frac{u(z)}{u_+(z)} \right) = \frac{C}{[u_+(z)]^2} e^{-\int^z p(\xi) d\xi}$$

$$\text{If } p(z) = \frac{p_0}{z-z_0} + \underset{\substack{\uparrow \\ \text{analytic}}}{r(z)} = \frac{1-\alpha_+-\alpha_-}{z-z_0} + r(z) = \frac{n+1-2\alpha_+}{z-z_0} + r(z)$$

$$\begin{aligned} e^{-\int^z p(\xi) d\xi} &= \exp \left[(n+1-2\alpha_+) \ln(z-z_0) - \int^z r(\xi) d\xi \right] \\ &= \frac{(z-z_0)^{2\alpha_+}}{(z-z_0)^{n+1}} e^{-\int^z r(\xi) d\xi} \end{aligned}$$

$$\text{Since } u_+(z) = (z-z_0)^{\alpha_+} f_+(z)$$

$$\frac{d}{dz} \left(\frac{u(z)}{u_+(z)} \right) = \frac{C}{(z-z_0)^{n+1}} \frac{e^{-\int^z r(\xi) d\xi}}{[f_+(z)]^2} = \frac{K}{(z-z_0)^{n+1}} g(z)$$

$\uparrow$
analytic
around $z=z_0$

$$\text{Let } g(z) = \sum_{m=0}^{\infty} \frac{g^{(m)}(z_0)}{m!} (z-z_0)^m$$

$$\frac{u(z)}{u_+(z)} = K_1 + K \int^z \frac{g(\xi)}{(\xi-z_0)^{n+1}} d\xi$$

$$= K_1 + K \left[\sum_{\substack{m=0 \\ m \neq n}}^{\infty} \frac{g^{(m)}(z_0)}{m-n} \frac{(z-z_0)^{m-n}}{n!} + \frac{g^{(n)}(z_0)}{n!} \ln(z-z_0) \right]$$

$$\text{So the other soln is } u_-(z) = u_+(z) \left[\sum_{\substack{m=0 \\ m \neq n}}^{\infty} \frac{g^{(m)}(z_0)}{m-n} \frac{(z-z_0)^{m-n}}{n!} + \frac{g^{(n)}(z_0)}{n!} \ln(z-z_0) \right]$$

$$\text{Using } u_+(z) = (z-z_0)^{\alpha_+} f_+(z) = (z-z_0)^{\alpha_+ + n} f_+(z)$$

$$u_-(z) = (z-z_0)^{\alpha_-} f_-(z) + \frac{g^{(n)}(z_0)}{n!} u_+(z) \ln(z-z_0)$$

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where
$$f_-(z) = f_+(z) \sum_{\substack{m=0 \\ m \neq n}}^{\infty} \frac{g^{(m)}(z_0)}{n-n} \frac{(z-z_0)^m}{m!}$$

$f_-(z)$ is analytic around $z=z_0$

The key point is that $u_-(z)$ has a logarithmic singularity from the term
~~the~~ $\frac{g^{(n)}(z_0)}{n!} \ln(z-z_0) u_+(z)$, that behaves

like $(z-z_0)^{\alpha} + \ln(z-z_0)$ near the singularity,

(in addition to possible singularity of the form $(z-z_0)^{\alpha} \sim (z-z_0)^{\alpha-n}$).

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Legendre's Equation

Legendre Polynomial

$$(z^2 - 1) u''(z) + 2z u'(z) - \lambda(\lambda + 1) u(z) = 0$$

[Came from $\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d}{d\theta} u(\theta) \right) = -\lambda(\lambda + 1) u(\theta)$

with $z = \cos \theta$ transformation Originally $\frac{1}{\sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d}{d\theta} f \right) = -\lambda(\lambda + 1) f$

$$u''(z) + \underbrace{\frac{2z}{z^2 - 1}}_{\uparrow \frac{1}{z}} u'(z) - \underbrace{\frac{\lambda(\lambda + 1)}{z^2 - 1}}_{\uparrow \frac{1}{z^2}} u(z) = 0$$

Reg. Singularities at $z = \pm 1$

Let us try to expand around $z = 0$, a regular point,

$$u(z) = \sum_{k=0}^{\infty} a_k z^k$$

$$(z^2 - 1) \sum_k \frac{k(k-1)}{z^2} a_k z^{k-2} + 2z \sum_k a_k z^{k-1} - \lambda(\lambda + 1) \sum_k a_k z^k = 0$$

$$\sum_k k(k-1) a_k z^k - \sum_k k(k-1) a_k z^{k-2} + \sum_k 2k a_k z^k - \sum_k \lambda(\lambda + 1) a_k z^k = 0$$

$$(k+1)(k+2) a_{k+2} = [k(k+1) - \lambda(\lambda + 1)] a_k$$

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$$So \quad \frac{a_{k+2}}{a_k} = \frac{k(k+1) - \lambda(\lambda+1)}{(k+1)(k+2)} = \frac{(k-\lambda)(k+\lambda+1)}{(k+1)(k+2)}$$

As $k \rightarrow \infty$, this ratio ~~■~~ tends to 1, telling us that the radius of convergence of the series is 1. This, we expected from the singularities ~~■~~ of $P(z)$, $Q(z)$.

However, if λ is a nonnegative integer n , this series would terminate and we get a polynomial, an entire function.

$P_n(z)$ are the Legendre polynomials

$$P_0(z) = 1, \quad P_1(z) = z, \quad P_2(z) = \frac{1}{2}(3z^2 - 1), \dots$$

with the normalization $P_n(1) = 1$

What about general λ ? Let us try to expand around $z=1$ and use variable t s.t.

$$z = 1 - 2t$$

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$$\begin{aligned}
 & (z^2-1) \frac{d^2}{dz^2} + 2z \frac{d}{dz} - \lambda(\lambda+1) \\
 &= \left[\frac{(1-z)^2}{4} - 1 \right] \frac{1}{4} \frac{d^2}{dt^2} + 2(1-2t) \left(-\frac{1}{2} \right) \frac{d}{dt} - \lambda(\lambda+1) \\
 &= t(t-1) \frac{d^2}{dt^2} + (2t-1) \frac{d}{dt} - \lambda(\lambda+1) \quad \Rightarrow
 \end{aligned}$$

$$t(t-1)u''(t) + (2t-1)u'(t) - \lambda(\lambda+1)u(t) = 0$$

$$p(t) = \frac{2t-1}{t(t-1)}, \quad q(t) = -\frac{\lambda(\lambda+1)}{t(t-1)} \quad \checkmark$$

$$\begin{aligned}
 p(t) &\sim +\frac{1}{t}, & q(t) &\sim \frac{\lambda(\lambda+1)}{t} \\
 \Rightarrow p_0 &= +1, & \Rightarrow q_0 &= 0
 \end{aligned}$$

Indicial eqn

$$\begin{aligned}
 \alpha(\alpha-1) + \alpha p_0 + q_0 &= 0 \quad \Rightarrow \alpha(\alpha-1) + \alpha = 0 \\
 \text{So } \alpha^2 &= 0
 \end{aligned}$$

So we have $p_\lambda(z) = u\left(\frac{1-z}{2}\right)$

$$U(t) = 1 + \sum_{k=1}^{\infty} c_k t^k \quad c_0 = 1$$

$$k(k-1)c_k - (k+1)kc_{k+1} + 2kc_k - (k+1)c_{k+1} - \lambda(\lambda+1)c_k = 0$$

$$(k+1)^2 c_{k+1} = [k(k+1) - \lambda(\lambda+1)] c_k = (k-\lambda)(k+\lambda+1) c_k$$

$$c_k = \frac{\Gamma(k-\lambda)\Gamma(k+\lambda+1)}{\Gamma(-\lambda)\Gamma(\lambda+1)} \frac{1}{(k!)^2}$$

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$P_\lambda(z)$ is the Legendre function of the first kind.

There ought to be a second solution $Q_\lambda(z)$ with logarithmic ~~singularity~~ singularity around $z = \pm 1$.

This is the Legendre function of the second kind.

Alternatively one can look at behavior for $|z| > 1$.
One way to do that is use variable $\xi = \frac{1}{z}$.

The Legendre eqn now becomes

$$\left[\xi^2 \frac{d}{d\xi} (\xi^2 - 1) \frac{d}{d\xi} + \lambda(\lambda+1) \right] u(\xi) = 0$$

$$\Rightarrow \left(\frac{d^2}{d\xi^2} + \frac{2\xi}{(\xi^2-1)} \frac{d}{d\xi} + \frac{\lambda(\lambda+1)}{\xi^2(\xi^2-1)} \right) u(\xi) = 0$$

The differential operator in ξ , in the final version has coefficients $p(\xi) = \frac{2\xi}{\xi^2-1}$, $q(\xi) = \frac{\lambda(\lambda+1)}{\xi^2(\xi^2-1)}$

So it has regular singular points at $\xi = 0, \pm 1$

Q_λ is the series solution around $\xi = 0$.

$$Q_\lambda\left(\frac{1}{\xi}\right) \sim \xi^\alpha (1 + \xi^2 + \dots)$$

α satisfies $\alpha(\alpha-1) = \lambda(\lambda+1) \Rightarrow \alpha = \lambda+1, -\lambda$

$\alpha = \lambda+1$ solution gives Q_λ .

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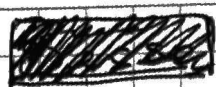
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Schl\"afli Integral representation

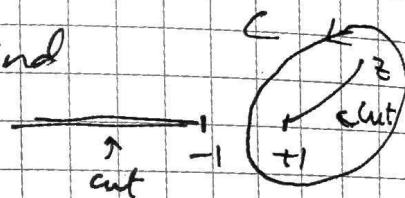
$$I_{\lambda}(z) = K_{\lambda} \int_C \frac{(1-\xi^2)^{\lambda}}{(z-\xi)^{\lambda+1}} d\xi$$

~~By showing that~~ By showing that $\left[\frac{d}{dz} \left(\frac{(1-\xi^2)^{\lambda}}{(z-\xi)^{\lambda+1}} \right) - \lambda(\lambda+1) \frac{(1-\xi^2)^{\lambda}}{(z-\xi)^{\lambda+1}} \right]$

$$= -\frac{d}{dz} \frac{(1-\xi^2)^{\lambda+1}}{(z-\xi)^{\lambda+2}}$$

namely a total derivative in ξ , we guarantee that $I_{\lambda}(z)$ would satisfy the Legendre equation, because the closed contour integral of the total derivative would vanish. Different choices of C correspond to different solutions.

Legendre function of the first kind



See Vaughan page 233 for $\nu = 0, 1/2, \dots$

For comment on Schl\"afli representation of $Q_{\nu}(z)$, see Vaughan, page 236.

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Bessel's Equation

Bessel Function

$$\frac{d^2 u}{dz^2} + \frac{1}{z} \frac{du}{dz} + \left(1 - \frac{\lambda^2}{z^2}\right) u = 0$$

[Came from $\frac{1}{z} \frac{d}{dz} z \frac{d}{dz} u - \frac{\lambda^2}{z^2} u = -u$. In cylindrical coordinates $\left(\frac{1}{\rho} \frac{d}{d\rho} \rho \frac{d}{d\rho} - \frac{m^2}{\rho^2}\right) f = -k^2 f$]

We have a regular singular point at $z=0$

$$p(z) = \frac{1}{z}, \quad q(z) = 1 - \frac{\lambda^2}{z^2} \Rightarrow p_0 = 1, \quad q_0 = -\lambda^2$$

Indicial eq: $\lambda(\lambda-1) + \lambda - \lambda^2 = 0$

$$\Rightarrow \lambda^2 = \lambda^2 \Rightarrow \lambda = \pm \lambda$$

Chose $\alpha = 1$, $u(z) = z^\lambda \left(1 + \sum_{k=1}^{\infty} a_k z^k\right)$, $a_0 = 1$

$$[(k+\lambda+2)^2 - \lambda^2] a_{k+2} + a_k = 0$$

If we start with $a_0 = 1$ we only have even powers

$$a_{2m} = - \frac{1}{(2m+\lambda)^2 - \lambda^2} a_{2m-2} = \frac{1}{4m(m+\lambda)} a_{2m-2}$$

$$J_\lambda(z) = \left(\frac{z}{2}\right)^\lambda \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+\lambda+1)} \left(\frac{z}{2}\right)^{2m}$$

This is the Bessel function of order λ .

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To get the other solution ^{for noninteger λ ,} one considers $\alpha = -\lambda$.
Conventionally, one defines

$$N_\lambda(z) = \frac{J_\lambda(z) \cos \pi \lambda - J_{-\lambda}(z)}{\sin \pi \lambda}$$

which is called Bessel func. of 2nd kind, also called Neumann func. This ~~com~~ combination has a well defined limit when λ goes to an integer value. The limit, as expected, has logarithms showing up in expansion around $z=0$.

Bessel functions also have integral representation. For integer order,

$$J_n(z) = \frac{1}{2\pi i} \oint_C e^{\frac{z}{t} - \frac{1}{2}t} \frac{dt}{t^{n+1}}$$



These relations are useful for asymptotics

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Legendre equation and Bessel equation are related to a broader class of differential equations. Equations with only three regular singular points (like $\pm 1, \infty$ for Legendre eqn) in the Riemann sphere could be transformed into the hypergeometric equation. Legendre equation, therefore, is related to a special case of hypergeometric function.

A confluence of two singular points lead to the confluent hypergeometric equation. Bessel equation, with singularities only at $z=0$ and $z=\infty$ is related to a special case of this category.

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