

New string vacua from simple topologies

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Introduction

Supersymmetry must be spontaneously broken.

- Superpartners?
- $\Lambda > 0$ incompatible with unbroken supersymmetry.

$$\left[\begin{array}{l} \text{for susy vacua:} \\ V \sim (|D\cancel{W}|^2 - |W|^2) \end{array} \right]$$

Diagram description: The equation is enclosed in large square brackets. The text "for susy vacua:" is at the top. Below it is the potential $V \sim (|D\cancel{W}|^2 - |W|^2)$. A dashed arrow points from the word "superpotential" (written in teal) to the $|W|^2$ term. Another dashed arrow points from the word "superpotential" to the $|D\cancel{W}|^2$ term, which has a diagonal line through it.

Strategy:

Step 1. Start from vacuum with $\Lambda < 0$

Step 2. $\left\{ \begin{array}{l} \text{break susy} \\ \text{lift to } \Lambda > 0 \end{array} \right.$

Examples:

[Kachru, Kallosh, Linde, Trivedi'03]

Step 1. $\text{AdS}_4 \times \text{CY}_6$ in IIB

using

- quantum corrections (brane instantons)
- O_3, D_3, D_7

Step 2. using $D_3\text{-}\overline{D_3}$ pairs

[de Wolfe, Giryavets, Kachru, Taylor'05]

Step 1. $\text{AdS}_4 \times \text{CY}_6$ in IIA

using

- classical ingredients
- O6

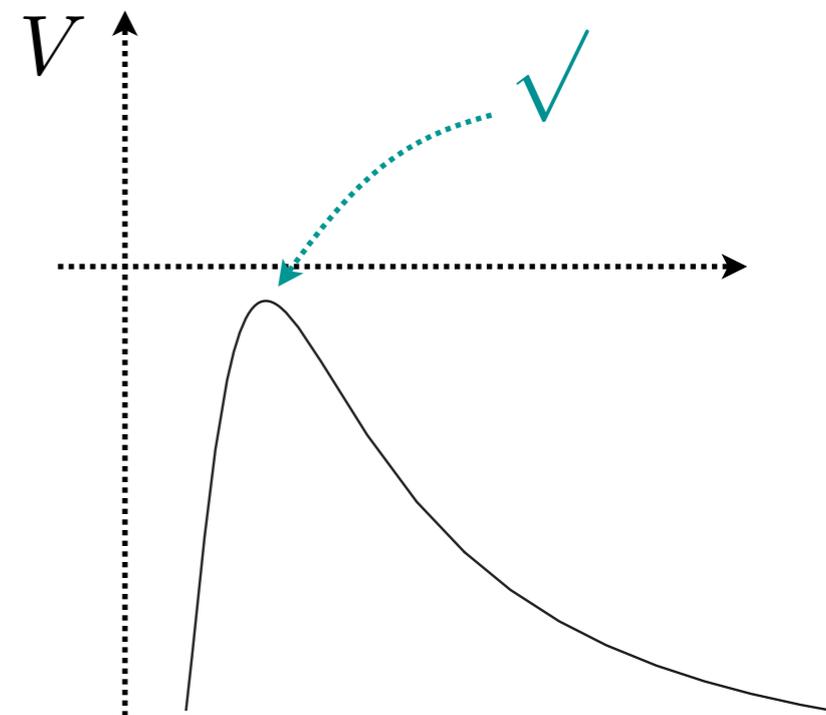
Step 2. Not easy (no-go, in some sense)

[Hertzberg, Kachru,
Taylor, Tegmark'07]

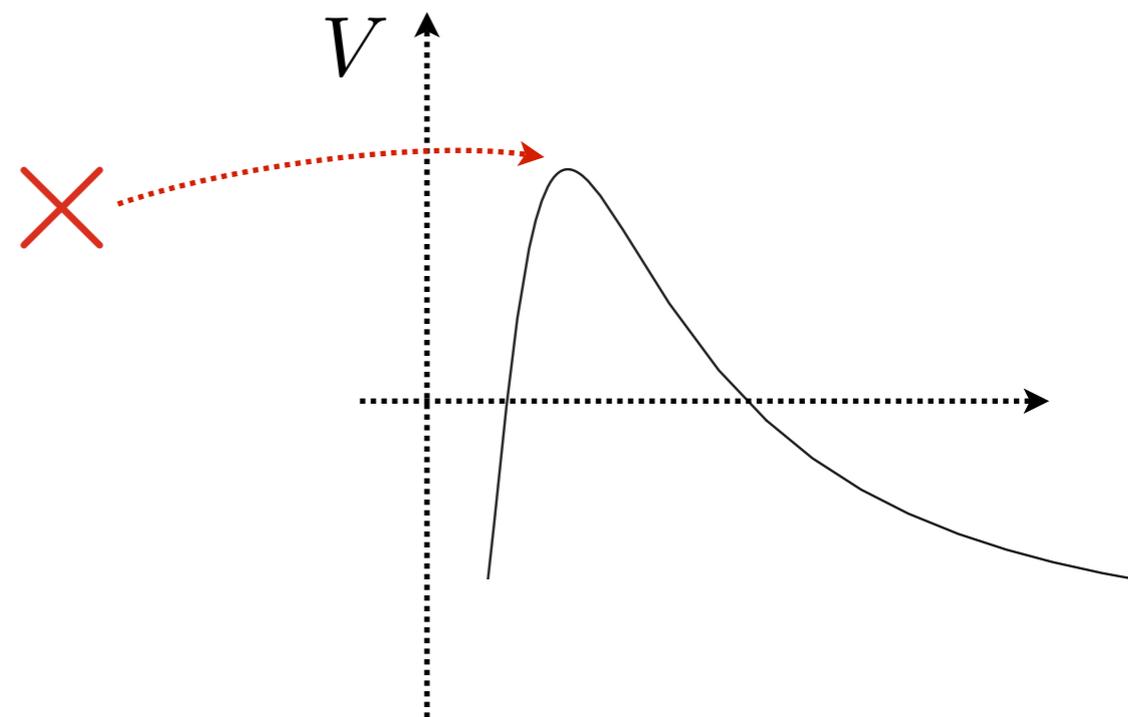
General caveat about step 2:

from the point of view of the 4d effective action

At Step 1. a tachyon might be acceptable
(if above Breitenlohner-Freedman bound)



But it will not be so
after lifting (at Step 2.)

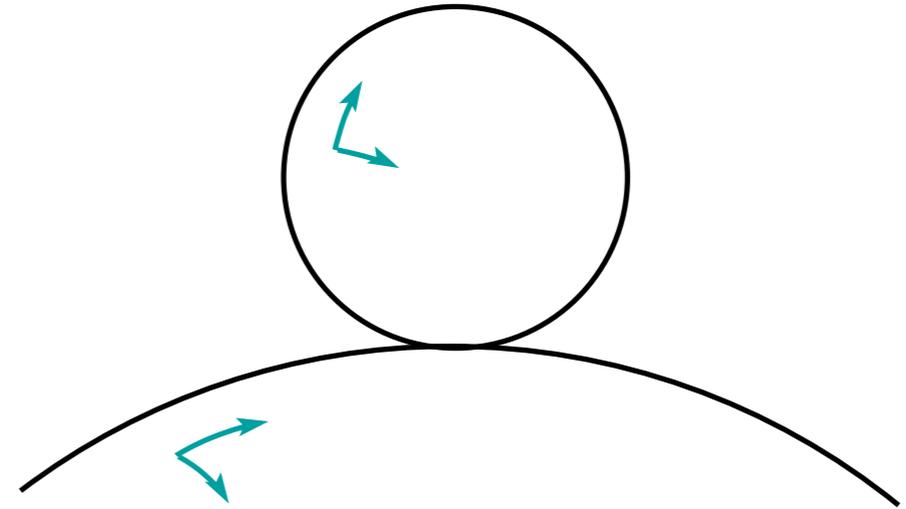


In this talk:

{only **Step I.** will be considered}

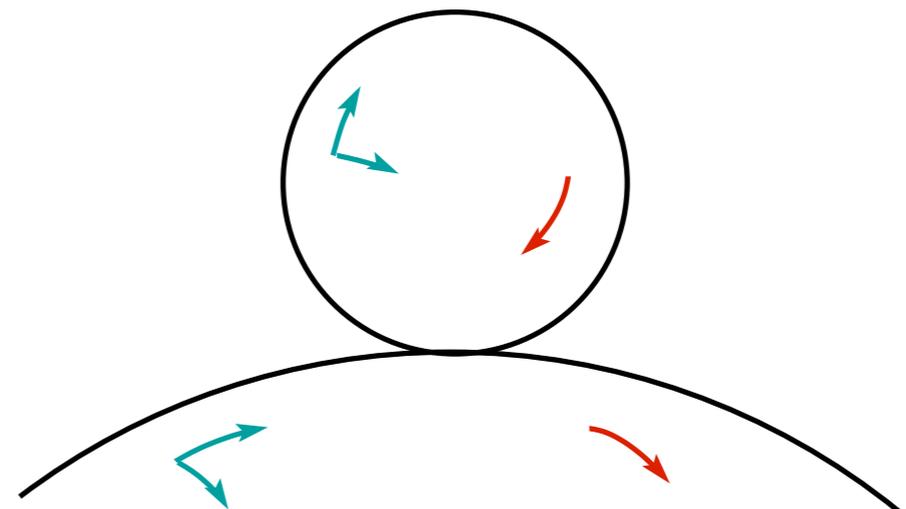
- $\text{AdS}_4 \times \mathbb{CP}^3$ vacua in IIA
- no orientifolds, no brane instantons
- infinitely many
- all moduli stabilized: **few to begin with**

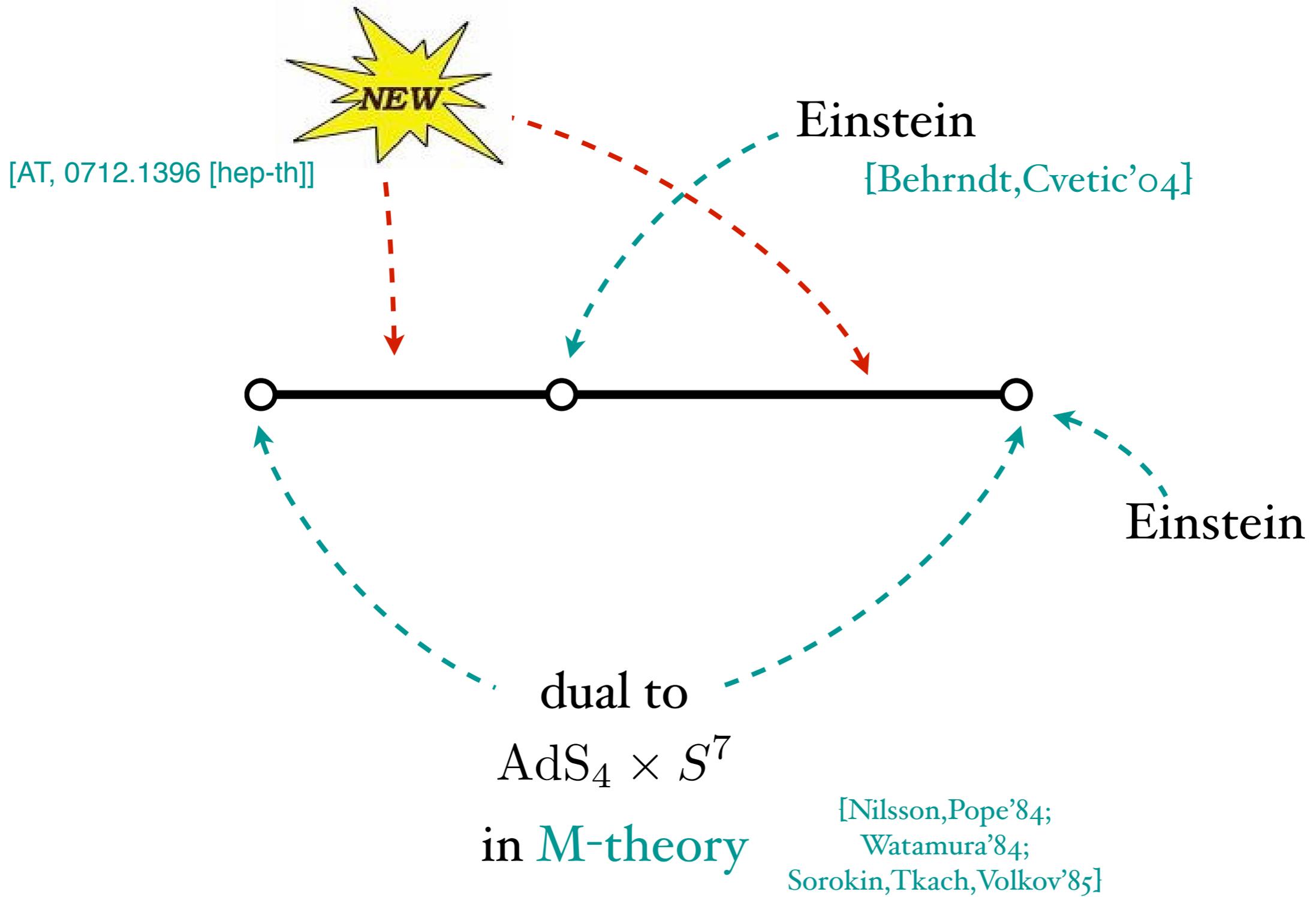
unlike usual
Freund-Rubin construction



(fluxes proportional to **volume forms**)

all fluxes will be on





Plan

- General considerations about **supersymmetry**
 - Some geometry of $\mathbb{C}P^3$
 - Finding the **new vacua**

Supersymmetry

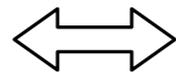
Minkowski₄ × M_6

would be more interesting geometrically:

has to be ‘generalized Calabi-Yau’

[Graña, Minasian,
Petrini, AT’05]

susy



‘calibrations’

$$d\Phi_1 = 0$$

$$d\Phi_2 = i * F$$

RR fluxes

(includes all the particular cases you know already:

[Klebanov, Strassler’00; Maldacena, Nuñez’00; ...])

$\text{AdS}_4 \times M_6$

slightly less so

has to be ‘generalized half-flat’

[Graña, Minasian,
Petrini, AT’06]

However, AdS vacua are easier to find.

Important subclass: “SU(3) vacua”

fluxes* are determined by two real numbers...

$$g_s F_0 = 5m \qquad g_s F_2 = \frac{1}{3} \tilde{m} J - W_2$$

... and by a two-form

$$g_s F_4 = \frac{3}{2} m J^2 \qquad g_s F_6 = -\frac{1}{2} \tilde{m} J^3$$

$$H = 2m \operatorname{Re} \Omega$$

metric determined by
2-form J and 3-form Ω

Bianchi
becomes

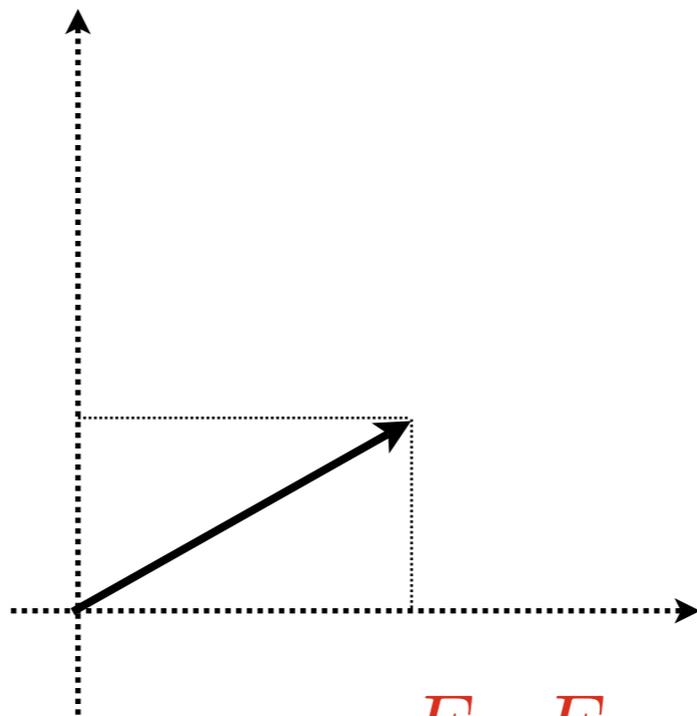
such that

$$dJ = 2\tilde{m} \operatorname{Re} \Omega \qquad dW_2^- = \frac{2}{3} (\tilde{m}^2 - 15m^2) \operatorname{Re} \Omega$$

$$d\operatorname{Im} \Omega = \frac{8}{3} \tilde{m} J^2 + W_2 \wedge J$$

*these are internal; external determined by duality

F_0, F_4, H



F_2, F_6

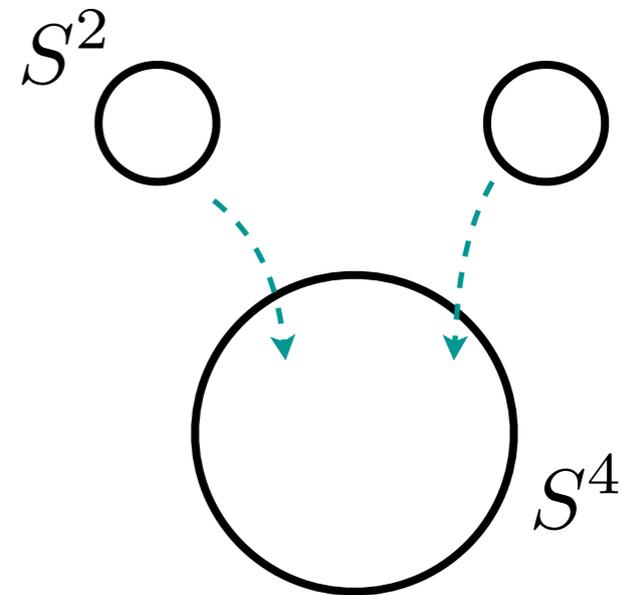
$$\Lambda = -3(m^2 + \tilde{m}^2)$$

Some geometry of $\mathbb{C}P^3$

- topology

$$S^2 \hookrightarrow \mathbb{C}P^3 \downarrow S^4$$

‘twistor fibration’



fibre = $\frac{\text{spinors}}{\text{complex rescalings}} \equiv \text{twistors}$

$$\parallel$$

$$\mathbb{C}P^1 = S^2$$

cohomology:

h^0	h^1	h^2	h^3	h^4	h^5	h^6
1	0	1	0	1	0	1

- complex structure and metric

Usually on \mathbb{CP}^3 : Fubini-Study (Kähler, Einstein) $I_{\text{FS}}, g_{\text{FS}}$

It can't be useful for us:

- susy $\Rightarrow dJ = 2\tilde{m} \text{Re}\Omega \Rightarrow I$ not integrable
 $(1,1)$ $(3,0) + (0,3)$

- I_{FS} has no globally defined (3,0)-form: $\nexists \Omega_{\text{FS}}$

We need a different (almost) complex structure

 not integrable

not

$$I_{\text{FS}} = \begin{pmatrix} I_2 & \\ & I_4 \end{pmatrix}$$

but

$$I_{\text{susy}} = \begin{pmatrix} -I_2 & \\ & I_4 \end{pmatrix}$$

- integrable

- $c_1 = 4$



$$\nexists \Omega_{\text{FS}}$$

not integrable

$$c_1 = 0$$



$$\exists \Omega_{\text{susy}}$$

- We will also change the metric

rather than $g_{\text{FS}} = g_2 + 2g_4$

$$g_{\text{susy}} = R^2 (g_2 + \sigma g_4)$$

overall scale

‘squashing parameter’

notice that we have few parameters to begin with:

‘stabilizing moduli’ will be easy

Finding the new vacua

$$\begin{array}{l}
 g_{\text{susy}} \\
 I_{\text{susy}}
 \end{array}
 \begin{array}{l}
 \text{---} \\
 \text{---}
 \end{array}
 \begin{array}{l}
 \text{---} \\
 \text{---}
 \end{array}
 \longrightarrow
 J_{\text{susy}} \equiv g_{\text{susy}} I_{\text{susy}}$$

can solve
(susy)

$$\left\{ \begin{array}{l}
 dJ = 2\tilde{m} \text{Re}\Omega \\
 d\text{Im}\Omega = \frac{8}{3}\tilde{m}J^2 + W_2^- \\
 dW_2^- = \frac{2}{3}(\tilde{m}^2 - 15m^2) \text{Re}\Omega
 \end{array} \right.$$

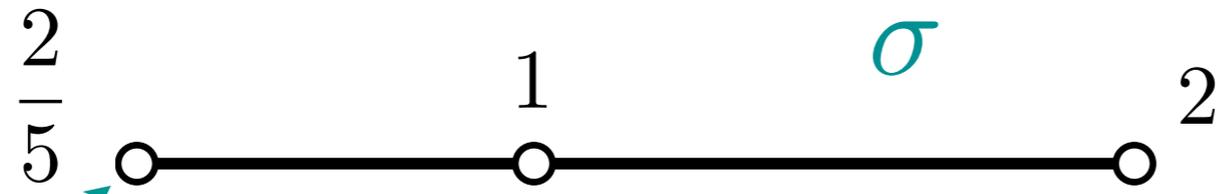
with

$$m = \frac{1}{2R} \sqrt{\left(\sigma - \frac{2}{5}\right)(2 - \sigma)} \qquad \tilde{m} = -\frac{1}{2R}(\sigma + 2)$$

$$(\sigma, R) \longrightarrow (m, \tilde{m}) \longrightarrow \text{fluxes}$$

$$m = \frac{1}{2R} \sqrt{\left(\sigma - \frac{2}{5}\right)(2 - \sigma)} \quad \Rightarrow$$

$$\frac{2}{5} \leq \sigma \leq 2$$



not Einstein

can be lifted to another Einstein metric on S^7

$$W_2 = 0$$

also Einstein;
can be generalized to 'nearly Kähler'

[Behrndt, Cvetic'04]

$g = g_{\text{FS}}$ Einstein;

$$I \neq I_{\text{FS}}$$

(would have been Kähler)

can be lifted to S^7

[Nilsson, Pope'84;
Watamura'84;
Sorokin, Tkach, Volkov'85]

Flux quantization

- H is exact
$$\begin{aligned} H &= 2m\text{Re}\Omega \\ dJ &= 2\tilde{m}\text{Re}\Omega \end{aligned} \Rightarrow H = d\left(\frac{m}{\tilde{m}}J\right)$$
- F_k are not closed
$$\begin{aligned} dF_k &= H \wedge F_{k-2} \Rightarrow d(e^{-B \wedge} F)_k = 0 \\ &\quad \parallel \\ &\quad d\tilde{F}_k = 0 \end{aligned}$$

using

$$m = \frac{1}{2R} \sqrt{\left(\sigma - \frac{2}{5}\right) (2 - \sigma)}$$

$$\tilde{m} = -\frac{1}{2R} (\sigma + 2) \quad \Rightarrow$$

$$r \equiv \frac{R}{2\pi l_s}$$

$$\int \tilde{F}_k = f_k(\sigma) r^{k-1}$$

||
 n_k

flux quantization

four equations for g_s, R, σ

- not all n_k are allowed
- all three are **fixed**
- for each vacuum, \exists infinitely many others
- it can be arranged: $r \gg 1, g_s \ll 1$

Some lessons

- complexity of vacua not unique to Calabi-Yau
- “vacua first, then effective field theories”

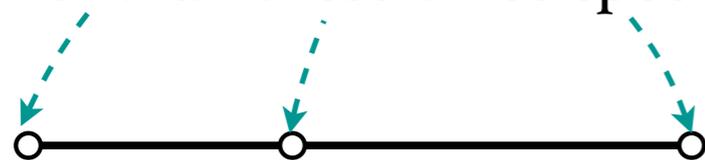
KK reduction on general M_6 is hard!

it can be done on M_6

- Calabi-Yau
- “twisted tori” [Scherk-Schwarz]

For the vacua we found, 4d theory **not known**

other than these three special cases



[Kashani-Poor'07]

$$\mathcal{N} = 2;$$

one vector multiplet,
one hypermultiplet

Conclusions

- “Landscape” of vacua not exclusive to Calabi-Yau’s
 - Even just classical physics goes a long way
 - KK more difficult than finding vacua