

Symbolifying Amplitudes and Wilson Loops



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The View from 30,000 Feet

The past few years have seen a revolution in our understanding of *scattering amplitudes*, particularly in maximally supersymmetric theories.

Hidden Mathematical
Structure &
New Symmetries

Dramatically Improved
Computational Power

The View from 30,000 Feet

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The View from 30,000 Feet

$$\text{Amplitude} = \int d(\text{loop momenta}) \underbrace{\sum \text{Feynman diagrams}}$$

This processing of the *integrand* is the realm of recent sexy technology (generalized unitarity, on-shell recursion, Grassmannian...)

$$= \int d(\text{loop momenta}) \underbrace{\sum (\text{relatively}) \text{ simple basis of integrals}}$$

The View from 30,000 Feet

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and this is the step which still requires a lot of blood and tears, begging for some new technology

Introduction

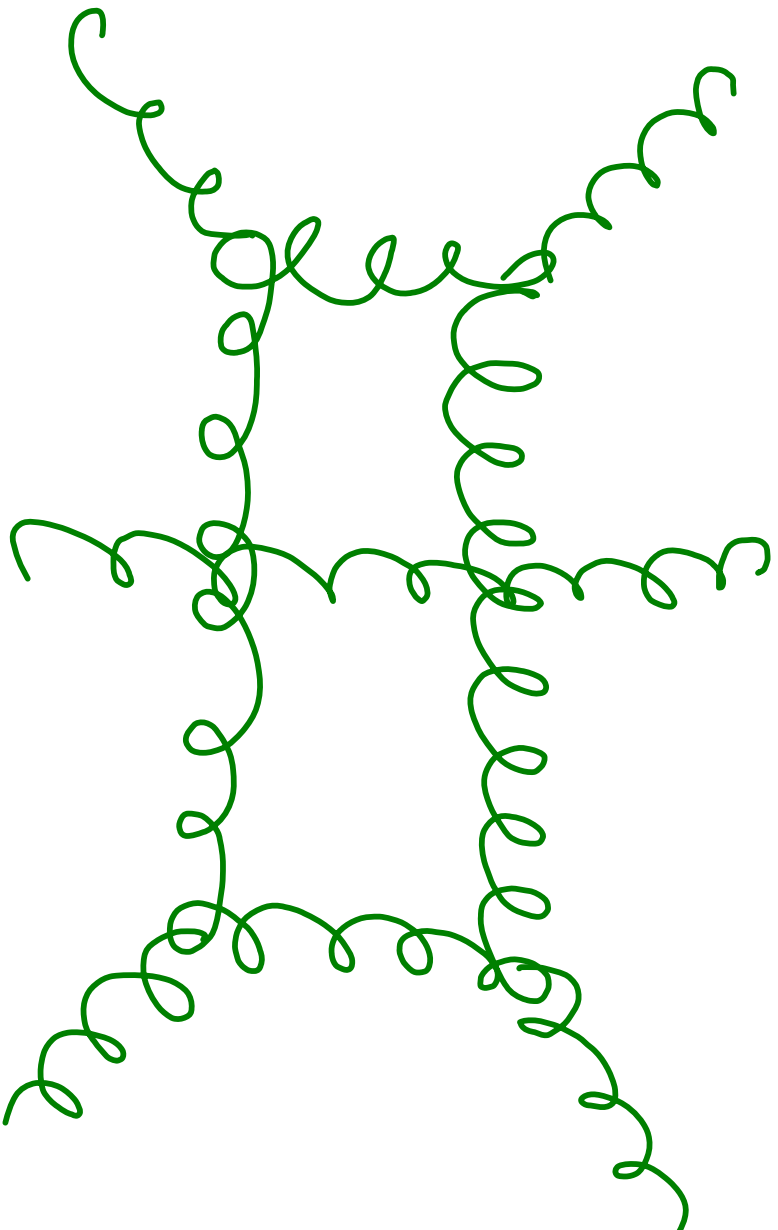
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Simplest nontrivial multi-loop scattering amplitude in $\mathcal{N}=4$ SYM:
the two-loop six-particle MHV ($++ \rightarrow ++++$) amplitude



+ many more

The ABDK/BDS Ansatz

The *infrared* and *collinear* behavior of amplitudes in massless gauge theories are tightly constrained by general field theory arguments. (Catani, Sterman, Tejada-Yeomans)

$$\log A_{MHV} \sim (\text{known IR divergent terms}) \\ + (\text{specific finite terms with prescribed collinear behavior}) \\ + (\text{finite terms with trivial collinear limits}) \\ + \mathcal{O}(\text{IR regulator}) \text{ terms}$$

Collinear Limits

A function $F_n(k_1, k_2, \dots, k_n)$ of n 4-vectors k_i has "trivial collinear limits" if

$$F_n(k_1, \dots, k_i, k_{i+1}, \dots, k_n) \rightarrow F_{n-1}(k_1, \dots, k_i + k_{i+1}, \dots, k_n)$$

when two cyclically adjacent k_i become parallel.

(Or, more generally, any number may become parallel \Rightarrow "multi-collinear" limits)

The ABDK/BDS Ansatz

Bern, Dixon and Smirnov (BDS) made a specific proposal for how to write the terms in terms of one-loop amplitudes

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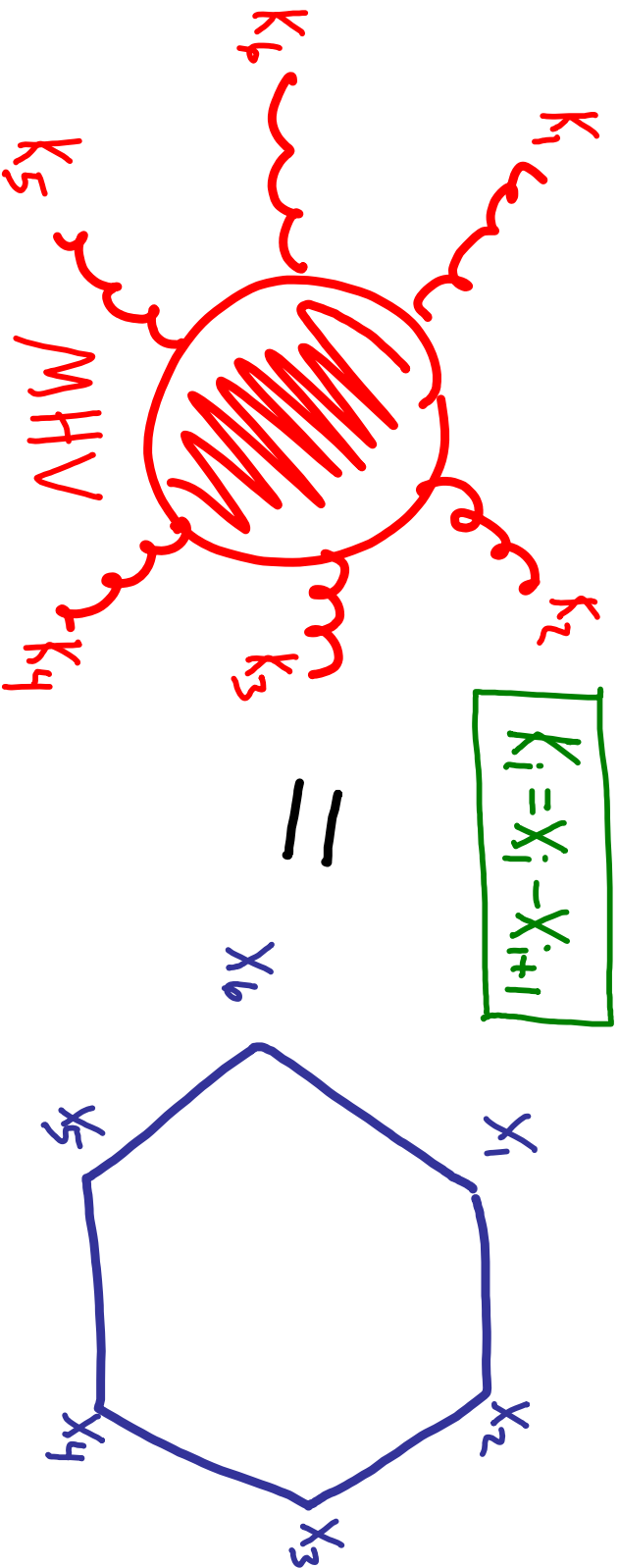
The BDS ansatz was supported by laborious calculations at $n=4,5$ through 3 loops.

But explicit calculation revealed that the BDS ansatz fails beginning at two-loops for $n=6$ particles.

\Rightarrow The BDS "remainder function" is nonzero (though, as emphasized, it must have trivial collinear limits.)

Amplitudes = Wilson Loops

In parallel developments, inspired by the work of **Alday and Maldacena** at strong coupling, it was experimentally observed that, apparently order by order in perturbation theory,



Dual Conformal Invariance

Once you accept $A=W$, then the conformal Ward identity which the Wilson loop satisfies (in X space) implies a highly non-obvious relation on A which we call the dual conformal Ward identity.

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Once you accept $A=W$, then the conformal Ward identity which the Wilson loop satisfies (in X space) implies a highly non-obvious relation on A which we call the **dual conformal Ward identity**.

The most general solution to the DCWI is

$$\log A = (\text{BDS ansatz}) + (\text{any finite dual conformally invariant})$$

ie any function of cross-ratios $\frac{x_{ij}^2 x_{kl}^2}{x_{ik}^2 x_{jl}^2}$

Modern Understanding

Given the presumption of dual conformal invariance, the BDS ansatz $\log A \sim A^{1-\text{loop}}$ is **trivially** true for $n=4,5$ to all loops, because there are no possible cross-ratios!

$$\frac{X_{12}^2 X_{34}^2}{X_{23}^2 X_{14}^2} = 0 \quad \text{if } X_1 - X_2 \text{ are null!}$$
$$X_3 - X_4$$

The Simplest Non-Trivial Amplitude/Wilson Loop in SYM

The rest of this talk is about the 2-loop 6-particle remainder function $\mathcal{R}_6^{(2)}$. What do we know?

- A function of three dual conformal cross-ratios

$$U_1 = \frac{S_{12} S_{45}}{S_{123} S_{345}}$$

$$U_2 = \frac{S_{23} S_{56}}{S_{234} S_{123}}$$

$$U_3 = \frac{S_{35} S_{61}}{S_{345} S_{234}}$$

$$S_{ij\dots} = (k_i + k_j + \dots)^2$$

- Symmetric under any permutation of the U 's.
- Vanishes in any collinear limit

$$\mathcal{R}(0, u, 1-u) = \begin{array}{c} \text{Hexagon with dashed line and red arrow} \\ = \\ \text{Pentagon} \\ = \\ 0 \end{array}$$

Motivation

In a heroic effort, Del Duca, Duhr and Smirnov found a manageable way to evaluate the appropriate Wilson loop diagrams, and obtained

an analytic formula for $R(u_1, u_2, u_3)$

$$\begin{aligned}
& \frac{3}{4}\mathcal{G}\left(\frac{1}{v_{231}}, 1, \frac{1}{1-v_2}\right) H(0; u_3) + \frac{3}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-v_2}, 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{v_{312}}{u_1}, 1, \frac{1}{1-u_3}\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{v_{312}}{1-u_3}, 1, 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{231}, 1, \frac{1}{1-u_3}\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_3}, 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}\right) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{5}{24}\pi^2 H(0; u_1) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}\right) H(0; u_2) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}\right) H(0; u_2) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; u_2) H(0; u_3) + \frac{5}{24}\pi^2 H(0; u_2) H(0; u_3) + \\
& \frac{3}{4}H(0; u_2) H(0; u_1) H(0; u_3) + 3H(0; u_1) H(0; u_2) H(0; u_3) + \\
& \frac{1}{4}H(0; u_2) H\left(0, 1; \frac{u_2+u_3-1}{u_1+u_2-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_1) H(0, 1; (u_1+u_3)) H(0; u_3) + \\
& \frac{1}{4}H(0; u_1) H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_2) H(0, 1; (u_2+u_3)) H(0; u_3) + \\
& \frac{3}{4}H(0; u_2) H(1, 0; u_1) H(0; u_3) + \frac{3}{4}H(0; u_1) H(1, 0; u_2) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0; u_1) - \frac{23}{24}\pi^2 H(0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; u_2) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{312}; 1\right) H(0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{321}; 1\right) H(0; u_2) - \\
& \frac{25}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0; u_2) - \\
& \frac{25}{4}H(0; u_1) H(0; u_2) - \frac{23}{24}\pi^2 H(0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{231}; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_3) + 3H(0; u_1) H(0; u_2) H(0; u_3) - \\
& \frac{25}{4}H(0; u_1) H(0; u_3) - \frac{25}{4}H(0; u_2) H(0; u_3) - \frac{23}{24}\pi^2 H(0; u_3) + \frac{1}{12}\pi^2 H(0, 1; u_1) + \\
& \frac{1}{12}\pi^2 H(0, 1; u_2) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_2-1}{u_2-1}\right) + \frac{1}{2}H(0; u_1) H(0; u_2) H(0, 1; (u_1+u_2)) +
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{12}\pi^2 H(0, 1; (u_1+u_2)) + \frac{1}{12}\pi^2 H(0, 1; u_3) + \frac{1}{4}H(0; u_1) H(0; u_2) H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) - \\
& \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) + \frac{1}{12}\pi^2 H(0, 1; (u_1+u_3)) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) + \\
& \frac{1}{12}\pi^2 H(0, 1; (u_2+u_3)) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{u_2}{u_3}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, \frac{u_1-1}{u_1+u_3-1}\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(1, 0; u_1) - \frac{3}{4}H(0; u_2) H(1, 0; u_1) - \frac{3}{4}H(0; u_3) H(1, 0; u_1) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) H(1, 0; u_1) - \frac{1}{3}\pi^2 H(1, 0; u_1) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{u_2}{u_3}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_1}, u_{123}; 1\right) H(1, 0; u_2) - \frac{3}{4}H(0; u_1) H(1, 0; u_2) - \frac{3}{4}H(0; u_3) H(1, 0; u_2) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) H(1, 0; u_2) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_2) - \frac{1}{3}\pi^2 H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{3}\pi^2 H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{u_2+u_3}{u_3}; 1\right) H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{u_2+u_3}{u_3}; 1\right) H(1, 0; u_3) - \frac{1}{4}G\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(1, 0; u_3) + \\
& \frac{3}{4}H(0; u_1) H(0; u_2) H(1, 0; u_3) - \frac{3}{4}H(0; u_1) H(1, 0; u_3) - \frac{3}{4}H(0; u_2) H(1, 0; u_3) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_2) H(1, 0; u_3) + \\
& \frac{1}{24}\pi^2 H(1, 1; u_1) + \frac{1}{24}\pi^2 H(1, 1; u_2) + \frac{1}{24}\pi^2 H(1, 1; u_3) + \frac{1}{2}H(0; u_2) H(0, 0; u_1) + \\
& \frac{1}{2}H(0; u_3) H(0, 0; u_2) + \frac{1}{2}H(0; u_1) H(0, 0; u_3) - \frac{1}{2}H(0; u_2) H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - \\
& \frac{1}{2}H(0; u_3) H\left(0, 0, 1; \frac{u_1+u_3-1}{u_3-1}\right) - H(0; u_1) H(0, 0, 1; (u_1+u_2)) - \\
& \frac{1}{2}H(0; u_2) H(0, 0, 1; (u_1+u_2)) - \frac{1}{2}H(0; u_1) H\left(0, 0, 1; \frac{u_1+u_3-1}{u_3-1}\right) -
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{2}H(0;u_2)H\left(0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)-H(0;u_1)H(0,0,1;(u_1+u_3))- \\
& H(0;u_3)H(0,0,1;(u_1+u_3))-\frac{1}{2}H(0;u_1)H\left(0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)- \\
& \frac{1}{2}H(0;u_3)H\left(0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)-H(0;u_2)H(0,0,1;(u_2+u_3))- \\
& H(0;u_3)H(0,0,1;(u_2+u_3))-\frac{1}{2}H(0;u_2)H(0,1,0;u_1)-\frac{1}{2}H(0;u_3)H(0,1,0;u_2)- \\
& \frac{1}{2}H(0;u_1)H(0,1,0;u_3)+\frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{u_1+u_2-1}{u_2-1}\right)- \\
& \frac{1}{4}H(0;u_3)H\left(0,1,1;\frac{u_1+u_2-1}{u_2-1}\right)+\frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{u_1+u_3-1}{u_1-1}\right)- \\
& \frac{1}{4}H(0;u_2)H\left(0,1,1;\frac{u_1+u_3-1}{u_1-1}\right)-\frac{1}{4}H(0;u_1)H\left(0,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+ \\
& \frac{1}{4}H(0;u_3)H\left(0,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+\frac{1}{2}H(0;u_2)H(1,0,0;u_1)-\frac{1}{2}H(0;u_3)H(1,0,0;u_1)- \\
& \frac{1}{2}H(0;u_1)H(1,0,0;u_2)+\frac{1}{2}H(0;u_3)H(1,0,0;u_2)+\frac{1}{2}H(0;u_1)H(1,0,0;u_3)- \\
& \frac{1}{2}H(0;u_2)H(1,0,0;u_3)-\frac{1}{4}H(0;u_3)H\left(1,0,1;\frac{u_1+u_2-1}{u_2-1}\right)- \\
& \frac{1}{4}H(0;u_2)H\left(1,0,1;\frac{u_1+u_3-1}{u_1-1}\right)-\frac{1}{4}H(0;u_1)H\left(1,0,1;\frac{u_2+u_3-1}{u_3-1}\right)- \\
& \frac{1}{4}H(0,0,0;0;u_1)-7H(0,0,0,0;0;u_2)-7H(0,0,0,0;0;u_3)+\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_2-1}{u_2-1}\right)+ \\
& 7H(0,0,0,0;0;u_1)-\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+3H(0,0,0,1;(u_1+u_3))+ \\
& 3H(0,0,0,1;(u_1+u_2))+\frac{3}{2}H\left(0,0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+3H(0,0,0,1;(u_1+u_3))+ \\
& \frac{3}{2}H\left(0,0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)+3H(0,0,0,1;(u_2+u_3))+\frac{9}{4}H(0,0,1,0;u_1)+ \\
& \frac{9}{4}H(0,0,1,0;u_2)+\frac{9}{4}H(0,0,1,0;u_3)-\frac{1}{2}H(0,1,0,0;u_1)-\frac{1}{2}H(0,1,0,0;u_2)- \\
& \frac{1}{2}H(0,1,0,0;u_3)+\frac{1}{2}H\left(0,1,0,1;\frac{u_1+u_2-1}{u_2-1}\right)+\frac{1}{2}H\left(0,1,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+ \\
& \frac{1}{2}H\left(0,1,0,1;\frac{u_2+u_3-1}{u_3-1}\right)+H(0,1,1,0;u_1)+H(0,1,1,0;u_2)+H(0,1,1,0;u_3)- \\
& \frac{1}{4}H\left(0,1,1,1;\frac{u_1+u_2-1}{u_2-1}\right)-\frac{1}{4}H\left(0,1,1,1;\frac{u_1+u_3-1}{u_1-1}\right)- \\
& \frac{1}{4}H\left(0,1,1,1;\frac{u_2+u_3-1}{u_3-1}\right)+H(1,0,0,1;\frac{u_1+u_2-1}{u_2-1})+H\left(1,0,0,1;\frac{u_1+u_3-1}{u_1-1}\right)+ \\
& H\left(1,0,0,1;\frac{u_2+u_3-1}{u_3-1}\right)+2H(1,0,1,0;u_1)+2H(1,0,1,0;u_2)+2H(1,0,1,0;u_3)+ \\
& \frac{1}{4}H(1,1,0,1;\frac{u_1+u_2-1}{u_2-1})+\frac{1}{4}H(1,1,0,1;\frac{u_1+u_3-1}{u_1-1})+ \\
& \frac{1}{4}H(1,1,0,1;\frac{u_2+u_3-1}{u_3-1})+\frac{1}{2}H(1,1,1,0;u_1)+\frac{1}{2}H(1,1,1,0;u_2)+\frac{1}{2}H(1,1,1,0;u_3)- \\
& \frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{123}}\right)-\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{231}}\right)-\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{8}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{123}}\right)-\frac{1}{8}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{123}}\right)+\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{132}}\right)-
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{132}}\right)-\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{213}}\right)+\frac{1}{24}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{213}}\right)- \\
& \frac{1}{8}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{231}}\right)+\frac{1}{8}\pi^2H(0;u_3)\mathcal{H}\left(1;\frac{1}{u_{231}}\right)+\frac{1}{8}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{312}}\right)- \\
& \frac{1}{8}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{312}}\right)+\frac{1}{24}\pi^2H(0;u_1)\mathcal{H}\left(1;\frac{1}{u_{321}}\right)-\frac{1}{24}\pi^2H(0;u_2)\mathcal{H}\left(1;\frac{1}{u_{321}}\right)- \\
& \frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)-\frac{1}{4}H(1,0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)+\frac{1}{24}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{24}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)-\frac{1}{4}H(0;u_1)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)-\frac{1}{4}H(1,0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)-\frac{1}{4}H(1,0;u_1)\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)+\frac{1}{24}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)- \\
& \frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)+\frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{123}}\right)-\frac{1}{4}H(0;u_2)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{132}}\right)+ \\
& \frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{132}}\right)+\frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{132}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{132}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{213}}\right)+\frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1;\frac{1}{u_{213}}\right)+ \\
& \frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{213}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{213}}\right)-\frac{1}{4}H(0;u_1)H(0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)+ \\
& \frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)+\frac{1}{4}H(0,0;u_3)\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{231}}\right)- \\
& \frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)+\frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{312}}\right)-\frac{1}{4}H(0;u_1)H(0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{321}}\right)+ \\
& \frac{1}{4}H(0,0;u_1)\mathcal{H}\left(0,1;\frac{1}{u_{321}}\right)+\frac{1}{4}H(0,0;u_2)\mathcal{H}\left(0,1;\frac{1}{u_{321}}\right)+\frac{1}{6}\pi^2\mathcal{H}\left(0,1;\frac{1}{u_{321}}\right)- \\
& \frac{1}{2}H(0,0;u_3)\mathcal{H}\left(1,1;\frac{1}{u_{123}}\right)+\frac{1}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{231}}\right)-\frac{1}{2}H(0;u_1)H(0;u_3)\mathcal{H}\left(1,1;\frac{1}{u_{231}}\right)+\frac{1}{2}H(0,0;u_1)\mathcal{H}\left(1,1;\frac{1}{u_{231}}\right)+ \\
& \frac{1}{2}H(0,0;u_3)\mathcal{H}\left(1,1;\frac{1}{u_{231}}\right)+\frac{11}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{231}}\right)-\frac{1}{2}H(0;u_1)H(0;u_2)\mathcal{H}\left(1,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{2}H(0,0;u_1)\mathcal{H}\left(1,1;\frac{1}{u_{312}}\right)+\frac{1}{2}H(0,0;u_2)\mathcal{H}\left(1,1;\frac{1}{u_{312}}\right)+\frac{11}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{312}}\right)- \\
& \frac{1}{2}H(0;u_1)\mathcal{H}\left(1,1;\frac{1}{u_{321}}\right)+\frac{1}{2}H(0;u_2)\mathcal{H}\left(1,1;\frac{1}{u_{321}}\right)+\frac{1}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{321}}\right)- \\
& \frac{1}{24}\pi^2\mathcal{H}\left(1,1;\frac{1}{u_{321}}\right)+\frac{1}{2}H(0;u_2)\mathcal{H}\left(0,0,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{2}H(0;u_1)\mathcal{H}\left(0,0,1;\frac{1}{u_{231}}\right)+\frac{1}{2}H(0;u_3)\mathcal{H}\left(0,0,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{2}H(0;u_2)\mathcal{H}\left(0,0,1;\frac{1}{u_{321}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{u_{123}}\right)+ \\
& \frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{u_{231}}\right)+\frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{u_{312}}\right)+
\end{aligned}$$

Motivation

There has to be a better way...

(otherwise we should all abandon $N=4$ SYM!)

But how can we possibly proceed?

Cast of Characters - Functions

"Classical Polylogarithms"

$$Li_n(z) = \sum_{k=1}^{\infty} \frac{z^k}{k^n} = \int_0^z Li_{n-1}(t) dt$$

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$$Li_n(z) = \sum_{k=1}^{\infty} \frac{z^k}{k^n} = \int_0^z Li_{n-1}(t) d \log t$$

"Harmonic Polylogarithms"

$$H_{a_k, \dots, a_1}(z) = \int_0^z H_{a_{k-1}, \dots, a_1}(t) \begin{cases} d \log t & a_k = 0 \\ d(-\log(1-t)) & a_k = 1 \end{cases}$$

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"Goncharov Polylogarithms"

$$G(a_k, a_{k-1}, \dots, a_1; z) = \int_0^z G(a_{k-1}, \dots, a_1; t) \frac{dt}{t - a_k}$$

Polylogarithms satisfy numerous functional identities,

$$- \operatorname{Li}_2(1-x) = \operatorname{Li}_2(x) + \frac{1}{2} \log(x)^2, \text{ etc, etc, etc}$$

Polylogarithms satisfy numerous functional identities,

$$- \operatorname{Li}_2(1-1/x) = \operatorname{Li}_2(1-x) + \frac{1}{2} \log(x)^2, \text{ etc, etc, etc}$$

For harmonic polylogarithms, there is a beautiful
Mathematica package by Daniel Maitre

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But to do battle against the Goncharov polylogarithms we need a serious weapon —

In particular we want to assign some kind of "invariant object" to any linear combination of such special functions — the "symbol".

Transcendentality

The two-loop remainder function has uniform

transcendentality degree 4.

A function of degree k is one which can be written as

$$T_k = \int_0^{t_1} \int_0^{t_2} \dots \int_0^{t_{k-1}} g^*(d \log R_1(t_1)) \dots g^*(d \log R_k(t_k))$$

where R_i are rational functions with rational coefficients

- $g(0)=a$, $g(1)=b$, a, b are rational points in \mathbb{C}^n
- $\int d \log R_1 \circ \dots \circ d \log R_k = \int_a^b \left(\int_a^t d \log R_1 \circ \dots \circ d \log R_{k-1} \right) d \log R_k(t)$
- The integral is taken along a path in \mathbb{C}^n and one needs to check local homotopy invariance.

The Symbol of a Transcendental Function

A very useful quantity associated to a function of uniform degree k is its **symbol**.

The symbol is an element of the

k -fold tensor product of the multiplicative group of rational functions (modulo constants)

$$\text{symbol}(T_k) = R_1 \otimes R_2 \otimes \dots \otimes R_k$$

The Symbol made Simple

A function of degree k is one which can be written as a k -fold iterated integral of a rational integrand

$$T_k(x_1, \dots, x_m) = \int^1 dt_1 \int^{t_1} dt_2 \dots \int^{t_{k-1}} dt_k R(x_1, \dots, x_m; t_1, \dots, t_k)$$

The symbol is a way to express the integrand in a way which preserves information about the order (but not the path) of integration.

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The symbol is a way to express the integrand in a way which preserves information about the order (but not the path) of integration.

The symbol can be computed recursively. If

$$dT_k = \sum_i T_{k-1}^i d \log R_i$$

Then $\text{symbol}(T_k) = \sum_i \text{symbol}(T_{k-1}^i) \otimes R_i$

Note that the symbol satisfies

$$\begin{aligned} R_1 \otimes \dots \otimes (R_a R_b) \otimes \dots \otimes R_k \\ = R_1 \otimes \dots \otimes R_a \otimes \dots \otimes R_k + R_1 \otimes \dots \otimes R_b \otimes \dots \otimes R_k \end{aligned}$$

$$\begin{aligned} R_1 \otimes \dots \otimes (c R_i) \otimes \dots \otimes R_k \\ = R_1 \otimes \dots \otimes R_i \otimes \dots \otimes R_k \end{aligned} \quad \text{for any constant } c$$

... properties it inherits from $\text{diag} R$.

Examples

• degree 0 $T_0 = R$ $\text{symbol}(T_0) = 0$

• degree 1 $T_1 = \log R$ $dT_1 = d \log R \Rightarrow \text{symbol}(T_1) = R$

• degree 2 $T_2 = \text{Li}_2(R) = \int_0^R -\log(1-t) d \log t$

$$dT_2 = -\log(1-R) d \log R$$

$$\Rightarrow \text{symbol}(T_2) = -(1-R) \otimes R$$

or consider $T_2 = \log R_1 \log R_2$

$$dT_2 = \log R_1 d \log R_2 + \log R_2 d \log R_1$$

$$\Rightarrow \text{symbol}(T_2) = R_1 \otimes R_2 + R_2 \otimes R_1$$

Classical Polylogarithms

The functions Li_k are defined recursively by

$$Li_k(z) = \int_0^z Li_{k-1}(t) dt \quad Li_1(z) = -\log(1-z)$$

$$\Rightarrow \text{symbol}(Li_k(z)) = - (1-z) \otimes \underbrace{z \otimes \dots \otimes z}_{k-1 \text{ times}}$$

Using these definitions it is straightforward to calculate the symbol of all harmonic & Goncharov polylogs, and hence the symbol S of the DBS formula.

But what is it good for?

Uses of the Symbol

The symbol converts polylog functional equations into rational function identities

$$\text{symbol}(Li_2(z)) = -(1-z) \otimes z$$


$$\begin{aligned}\Rightarrow \text{symbol}(Li_2(1-1/x)) &= -(1/x) \otimes (1-1/x) = x \otimes (x-1/x) \\ &= x \otimes (1/x) = x \otimes (1-x) - x \otimes x \\ &= \text{symbol}\left(-Li_2(1-x) - \frac{1}{2} \log(x)^2\right)\end{aligned}$$

$$\Rightarrow Li_2(1-1/x) = -Li_2(1-x) - \frac{1}{2} \log(x)^2$$

Ambiguities

Here we got lucky, but consider

$$\operatorname{Li}_2(1-x) + \log(x) \log(1-x) = -\operatorname{Li}_2(x) + \frac{\pi^2}{6}$$

The symbol doesn't see this term 

\Rightarrow The symbol only fixes the "leading transcendentalty" piece, i.e. modulo constants times functions of lower degree

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A related ambiguity is that the symbol obviously has no knowledge of where to put branch cuts

$\log(x)$, $\log(-x)$, $-\log(\frac{1}{x})$, $\frac{1}{2} \log(x^2)$, $\frac{3}{2} \log(-x^{2/3})$ have the same symbol

Cast of Characters - Arguments

The symbol S is a mess! It involves arguments

$$u_i,$$

$$1-u_i,$$

$$u_i+u_j,$$

$$1-u_i-u_j$$

$$v_{jkr}^{\pm} = \frac{u_k - u_r \pm \sqrt{(u_k+u_r)^2 - 4u_j u_k u_r}}{2(1-u_j)u_k}$$

$$u_{jkr}^{\pm} = \frac{1-u_j - u_k + u_r \pm \sqrt{\Delta}}{2(1-u_j)u_k}$$

$$\Delta = (1-u_1-u_2-u_3)^2 - 4u_1u_2u_3$$

The Symbol of the DDS Function

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$$u_{jkr}^{\pm} = \frac{1-u_j-u_k \pm \sqrt{\Delta}}{2(1-u_j)u_r}$$

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This is a redundant set; they satisfy many algebraic identities.

Eliminate these in terms of the rest (not obvious that this had to be possible!)

Momentum Twistors

We still have a $\sqrt{\Delta}$ we don't want; we need to find variables on a covering space.

The "right" variables are

$$u_1 = \frac{z_{23} z_{56}}{z_{25} z_{36}}$$

$$u_2 = \frac{z_{16} z_{34}}{z_{14} z_{36}}$$

$$u_3 = \frac{z_{12} z_{45}}{z_{14} z_{25}}$$

$$z_{ij} = z_i - z_j$$

$$z_i \in \mathbb{P}^1$$

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$$z_{ij} = z_i - z_j \\ z_i \in \mathbb{P}^1$$

or, if you prefer momentum twistors,

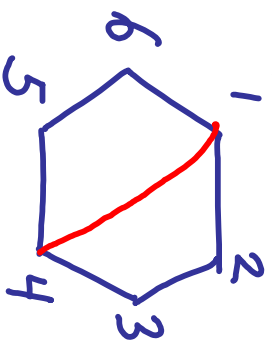
$$u_1 = \frac{\langle 1234 \rangle \langle 4561 \rangle}{\langle 1245 \rangle \langle 3461 \rangle}$$

etc (cyclically)

Then Δ becomes a perfect square, so the symbol can be expressed in terms of cross-ratios of z 's.

Cross-Ratios

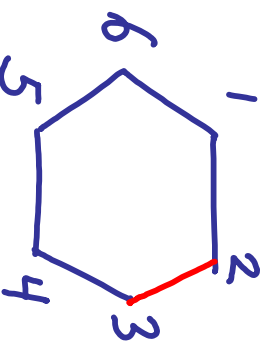
There are 15 different kinds of cross-ratios (of course, only 3 are independent, they satisfy algebraic relations)



$$\frac{z_{23}z_{56}}{z_{25}z_{36}} = u_1$$

"diagonal"

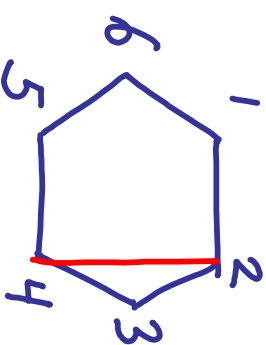
u_1, u_2, u_3



$$\frac{z_{14}z_{56}}{z_{16}z_{45}} = -X_1^{\pm}$$

"edge"

$X_1^{\pm}, X_2^{\pm}, X_3^{\pm}$



$$\frac{z_{13}z_{56}}{z_{16}z_{35}}$$

"chord"

6 of these

$$X_1^{\pm} = u_1 X_1^{\pm}$$

$$X_1^{\pm} =$$

$$\frac{u_1 + u_2 + u_3 - 1 \pm \sqrt{\Delta}}{2u_1 u_2 u_3}$$

$$X_1^{\pm} X_1^{\mp} = \frac{u_1^2}{u_1 u_2 u_3}$$

Miracle #1

In Z coordinates, the symbol could have been a linear combination of terms like

$$p_1(\tau_i) \otimes p_2(\tau_i) \otimes p_3(\tau_i) \otimes p_4(\tau_i)$$

for arbitrary polynomials of the cross-ratios τ_i

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for arbitrary polynomials of the cross-ratios τ_i

We found that in fact it takes the form

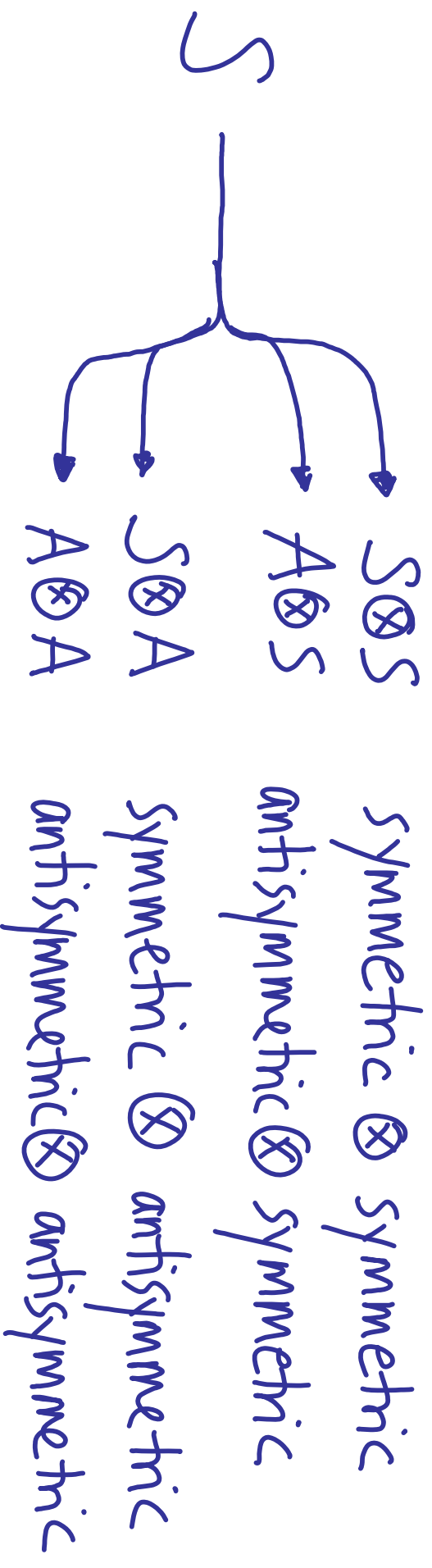
$$S = \sum_{i,j,k,r=1}^{15} c_{ijkr} \tau_i \otimes \tau_j \otimes \tau_k \otimes \tau_r \quad \text{for rational } c_{ijkr}$$

Miracle #2

We find that S can be expressed in terms of the diagonal and edge cross-ratios only; chords drop out.

(This should provide a possible hint about the structure of the n -point remainder function for $n > 6$)

Let us decompose the symbol S into four pieces depending on the symmetry properties under exchange of the first two entries, and last two entries



i.e. $(A \otimes A)_{ijkl} = \frac{1}{4}(C_{ijkr} - C_{jikr} - C_{ijrk} + C_{jirk})$ etc.

Miracle #3

We find that $AA = 0$

(This means $(A \otimes A)_{jk} - (A \otimes A)_{ki} = 0$)

Miracle #3

We find that $A \wedge A = 0$

(This means $(A \otimes A)_{ijk} - (A \otimes A)_{kij} = 0$)

Motivic hi-tech: any function of transcendental degree 4 whose symbol satisfies $A \wedge A = 0$ can be expressed in terms of the classical (poly)logs $Li_4(x)$, $Li_3(x)$, $Li_2(x)$, $Li_1(x)$ & $\log(x)$ only!

A comment: there is no constructive algorithm for determining what arguments can appear inside the polylogs, they could be in principle arbitrary rational functions of your original variables.

$$\begin{aligned} \text{e.g. } (1+x+x^2) \otimes x &= \left(\frac{1-x^3}{1-x} \right) \otimes x \\ &= (1-x^3) \otimes x - (1-x) \otimes x \\ &= \frac{1}{3} (1-x^3) \otimes x^3 - (1-x) \otimes x \\ &\Rightarrow -\frac{1}{3} \text{Li}_2(x^3) + \text{Li}_2(x) \end{aligned}$$

In practice, guessing, experimentation & luck required!

A Divide & Conquer Algorithm

We want to find a function with given symbol S in terms of

	AAA	SBA	AAS	SAS	
$L_{i_4}(x)$	x	x	✓	✓	} → do third
$L_{i_3}(x) \log(y)$	x	x	✓	✓	
$L_{i_2}(x) L_{i_2}(y)$	✓	✓	✓	✓	→ do first
$L_{i_2}(x) \log(y) \log(z)$	x	✓	✓	✓	→ do second
$\log(x) \log(y) \log(z) \log(w)$	x	x	x	✓	→ do last

(Note that $L_{i_2}(x) L_{i_2}(y)$ satisfies $A \wedge A = 0!$)

Analyticity

This algorithm gives a prototype, \tilde{R} , which has the same symbol as that of the DOS formula for $R_{\text{G}}^{(2)}$.

The most annoying part was finding an expression which put all branch cuts in the right place.

Physical input specifies that we expect the remainder function to be smooth (and real-valued) everywhere in the Euclidean regime where $u_i > 0$.

This frequently involved "unsimplifying"

$$\text{ie. } \frac{1}{2} \log(x^+/x^-) \rightarrow \sum_{i=1}^3 \rho_i(x_i^+) - \rho_i(x_i^-) \equiv J$$

where

$$\rho_n(x) = \frac{1}{2} \ln(x) - \frac{1}{2} (-1)^n \ln(1/x)$$

$$X^\pm = \frac{n_1 + n_2 + n_3 - 1 \pm \sqrt{(v_1 + v_2 + v_3 - 1)^2 - 4v_1v_2v_3}}{v_1v_2v_3}$$

$$X_i^\pm = u_i X^\pm$$

This frequently involved "unsimplifying"

ie. $\frac{1}{2} \log(x^+ / x^-) \rightarrow \sum_{i=1}^3 \rho_i(x_i^+) - \rho_i(x_i^-)$

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In terms of the function

$$\begin{aligned} L_4(x^+, x^-) = & \rho_4(x^+) - \frac{1}{2} \log(x^+ x^-) \rho_3(x^+) \\ & + \frac{1}{8} \log(x^+ x^-)^2 \rho_2(x^+) - \frac{1}{48} \log(x^+ x^-)^3 \rho_1(x) \\ & + \frac{1}{384} \log(x^+ x^-)^4 + (x^+ \leftrightarrow x^-) \end{aligned}$$

We find...

Final Result

$$\mathcal{R} = \sum_{i=1}^3 \left(L_4(x_i^+, x_i^-) - \frac{1}{2} L_{i4}(1-y_{i1}) \right) - \frac{1}{8} \left(\sum_{i=1}^3 L_{i2}(1-y_{i1}) \right)^2 \\ + \frac{5^4}{12} + \frac{\pi^2}{12} 5^2 + \frac{\pi^4}{72}$$

The full function is completely smooth,
manifestly symmetric in n_1, n_2, n_3 , and valid
for all $n_i > 0$.

Comments

Even though this is just the simplest non-trivial scattering amplitude in SYM, one data point is *infinitely* better than none!

Our formula provides hope to the idea that we might be able to really unlock the secrets of multi-loop SYM amplitudes, and hopefully connect to strong coupling

Prospects

Symbiology should be applicable to all amplitudes/wilson loops

- any number of sides
- any number of loops
- any helicity configuration (not only MHV)

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But there is very little data in the literature on any other cases (except the Wilson loop OPE construction of *Gaiotto, Maldacena, Sever, Vieira*)

[We know that by five loops, if not earlier, classical polylogarithms do not suffice even for $n=6$]

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What we urgently need is technology for computing symbols without actually evaluating integrals first. \Rightarrow Hard

An Example

$$F(x, y) = \int_1^{\infty} \frac{da}{a} \frac{db}{b} \frac{1}{1+ax+by}$$

It's usually not even obvious whether or not an integral even has well-defined transcendentality!

Empirical rule: unit leading singularities.

$$(a, b) = \left. \begin{array}{l} (0, 0) \\ (0, -1/y) \\ (-1/x, 0) \end{array} \right\}$$

all three poles
have residue ± 1 .

$$dF = -\frac{dx}{x} \int_1^{\infty} \frac{db}{b} \frac{1}{1+ax+by} + (x \leftrightarrow y)$$

$F'(x, y)$ ←

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$$= -\frac{dx}{x} \frac{1}{1+x} \int_1^{\infty} \frac{db}{b} \frac{1+x}{1+x+by} + (x \leftrightarrow y)$$

poles @ $b=0$, $b=-\frac{(1+x)}{y}$ with residue ± 1 !

$$= \log\left(\frac{1+x+y}{y}\right) d\log\left(\frac{1+x}{x}\right) + (x \leftrightarrow y)$$

$$\text{Symbol}(F) = \frac{1+x+y}{y} \otimes \frac{1+x}{x} + (x \leftrightarrow y)$$