

# IMPROVING IDENTIFICATION OF DIJET RESONANCES AT THE LHC

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arXiv:1407.7037 & work in progress

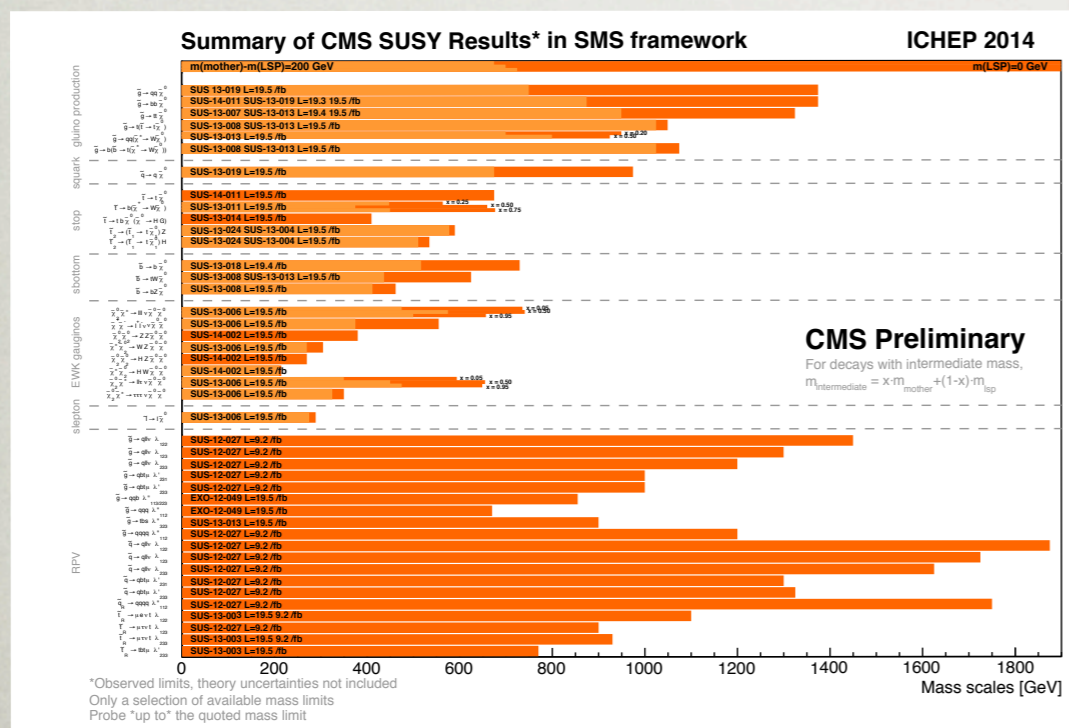
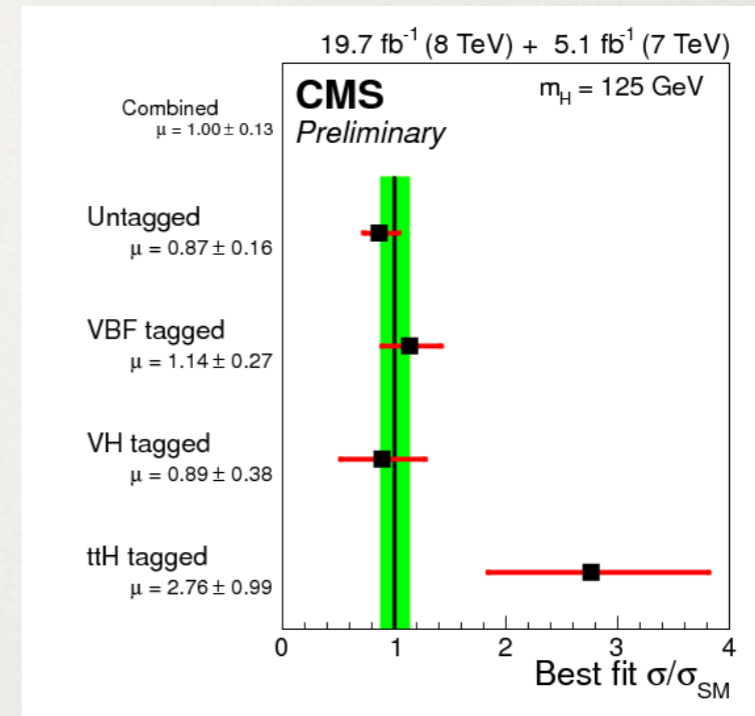
Rutgers University

14 October 2014



# Summary of LHC Run I

- Entered the era of precision Higgs physics

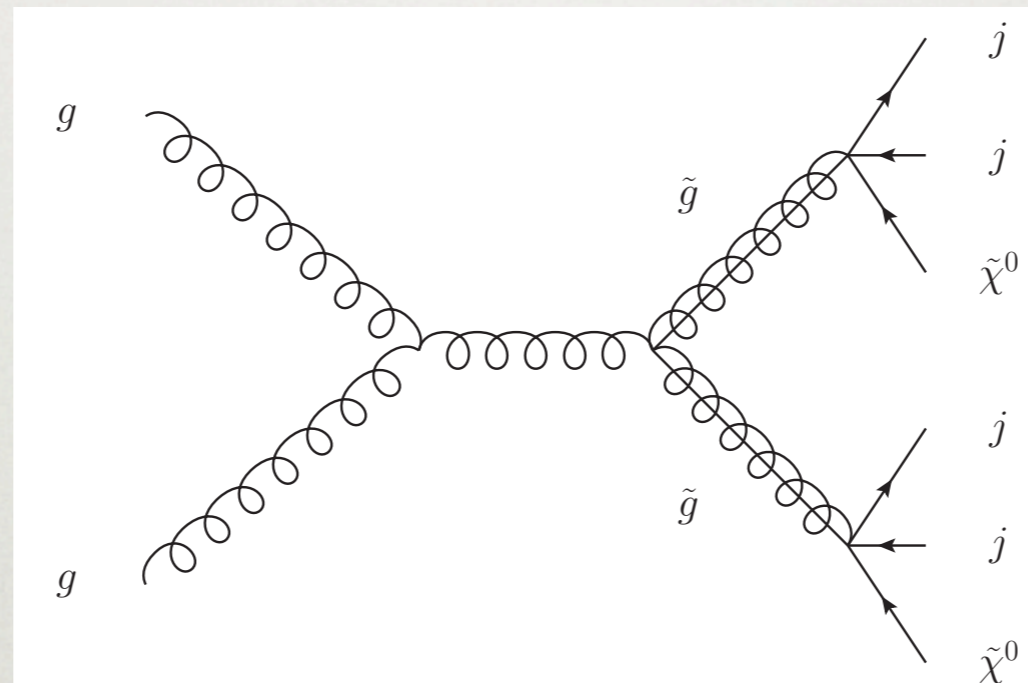


- Still no signs of
  - SUSY
  - Exotic EWSB mechanisms
  - Dark matter
  - ...



# Looking forward to Run II+

- Reach improved due to **higher energies**
  - Not much to do here: scale up cuts and look for obvious signs of new physics





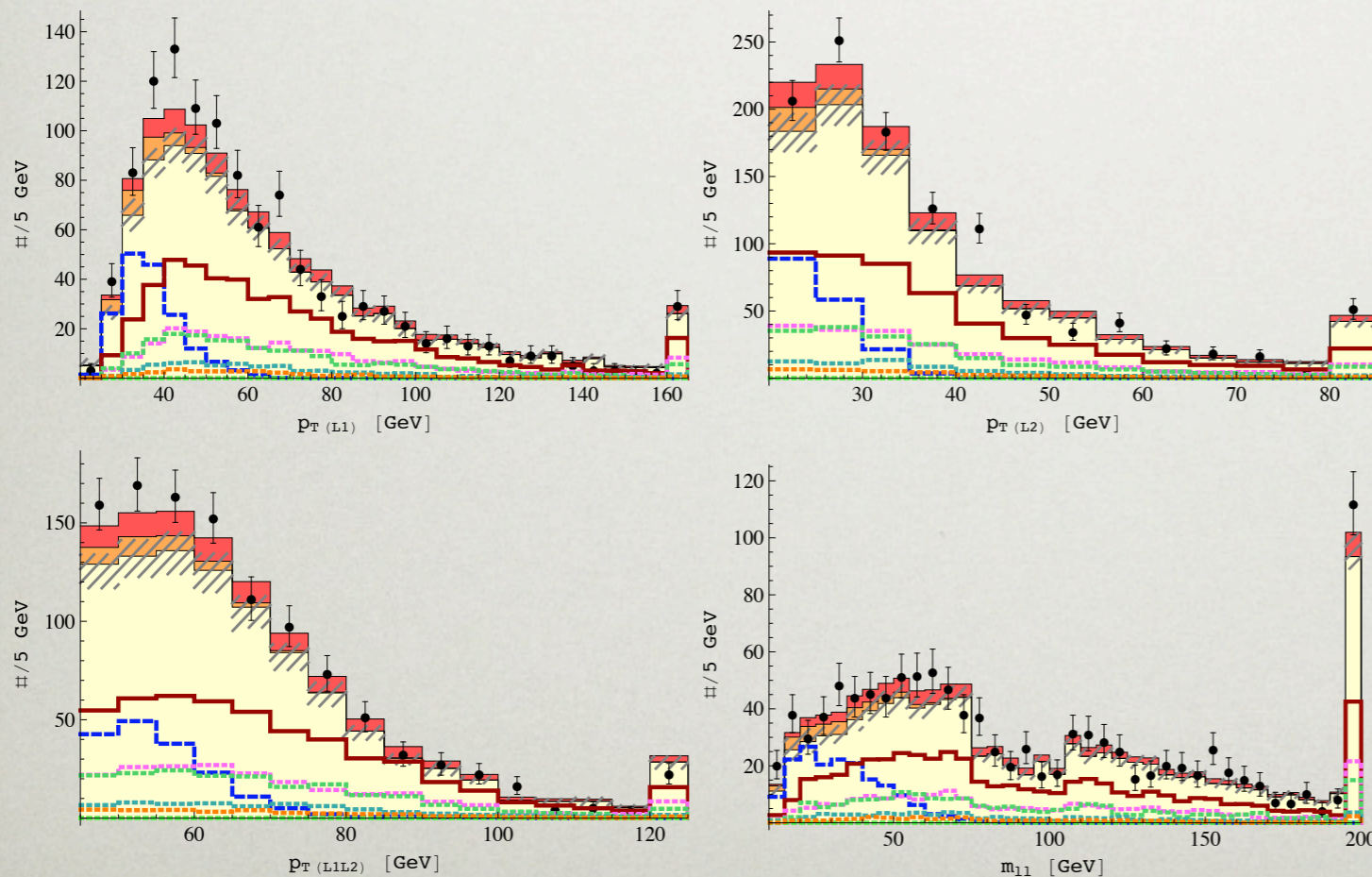
# Looking forward to Run II+

- What if new physics lies around the electroweak scale?
  - Can look very similar to SM physics
- **High integrated luminosity** gives a lot of signal events but....
  - **Trigger thresholds** are higher
  - **SM backgrounds** are higher
  - **Pile-up** is higher
- Cutting out the SM background cuts out most of the signal, too
  - Study BSM physics via precision SM physics
- Requires dedicated strategies to ensure we don't miss anything



# Hadronic resonances

- Ubiquitous in the Standard Model
  - W/Z/H/t
  - SM resonances can also decay leptonically, but suffer from smaller branching fractions
  - Want as many handles as possible on SM rates
  - (Possible) discrepancy in fully leptonic WW cross section



But see:

Meade, Ramani, Zeng  
arXiv:1407.4481

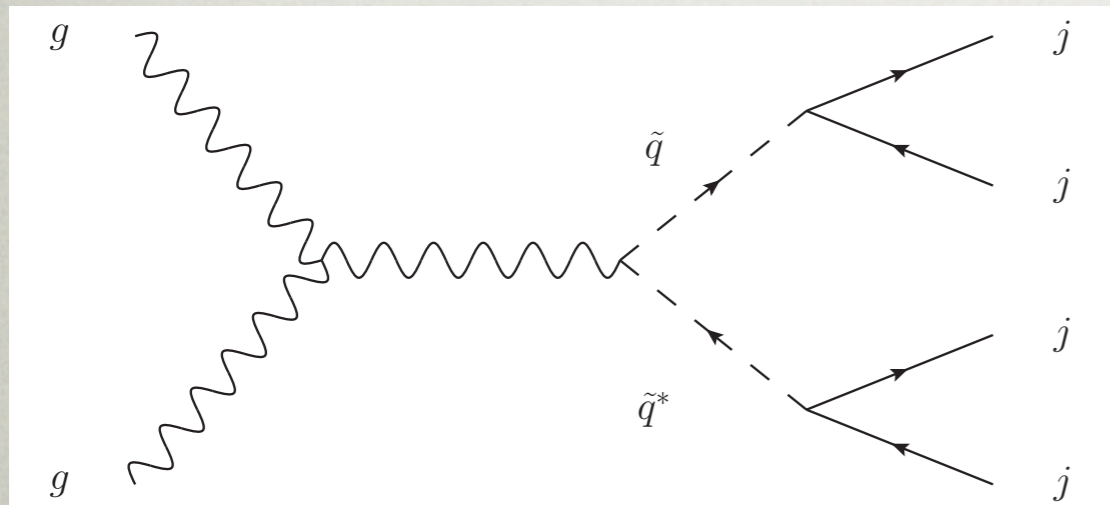
Jaiswal, Okui  
arXiv:1407.4537

Taken from Curtin, Jaiswal, Meade, arXiv:1206.6888

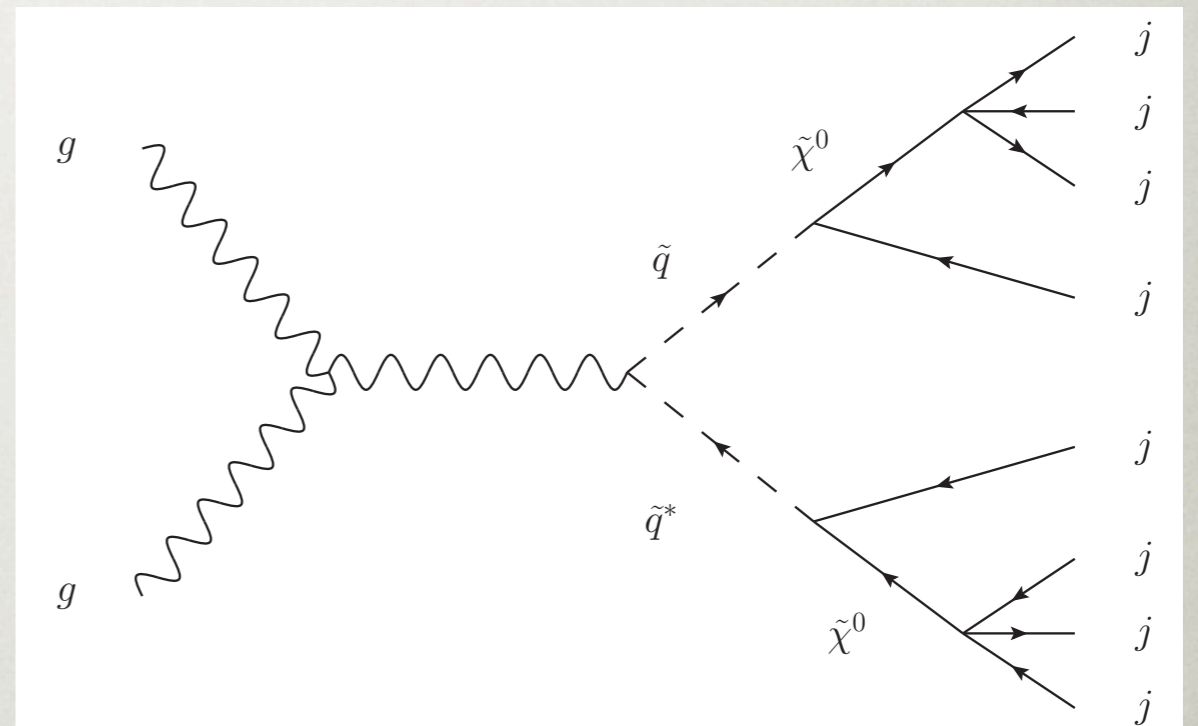


# Hadronic resonances

- Also ubiquitous **beyond** the Standard Model
  - Extended Higgs sectors
  - R-parity-violating supersymmetry
  - Supersymmetric cascade decays
  - Extra dimensions
  - New gauge interactions



(threshold)



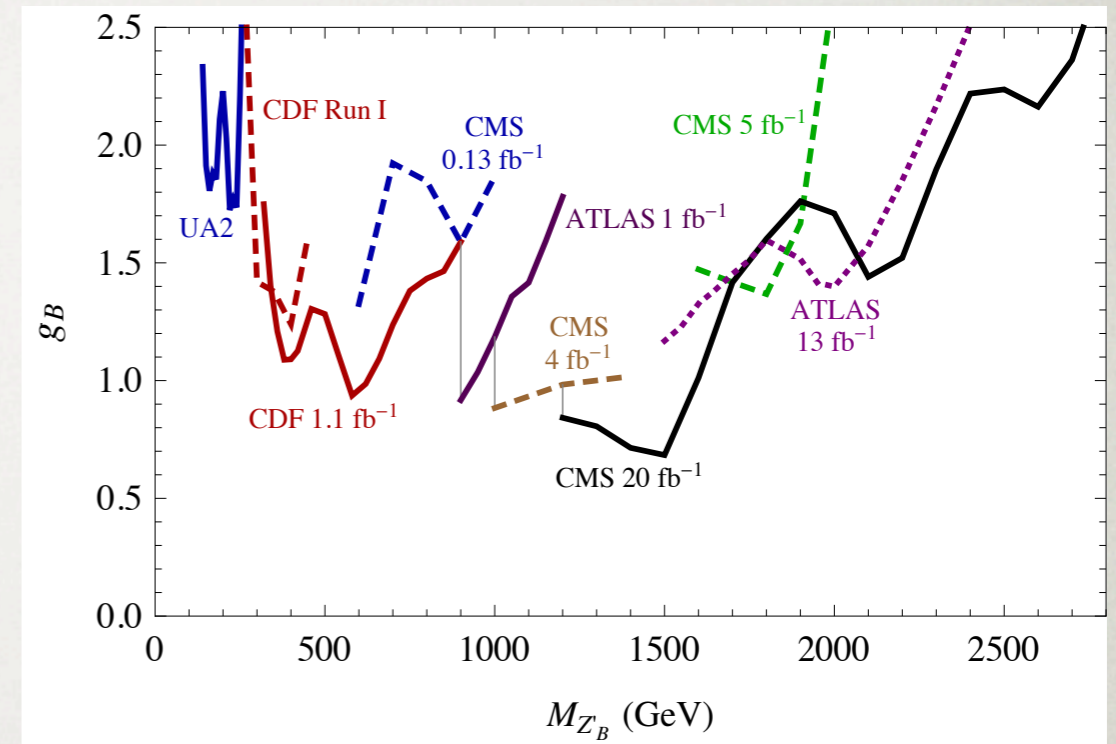
(boosted)



# Hadronic resonances

- We are not guaranteed to do better at the LHC
- Extreme example: baryonic  $Z'$

(taken from Dobrescu, Yu  
arXiv:1306.2629)

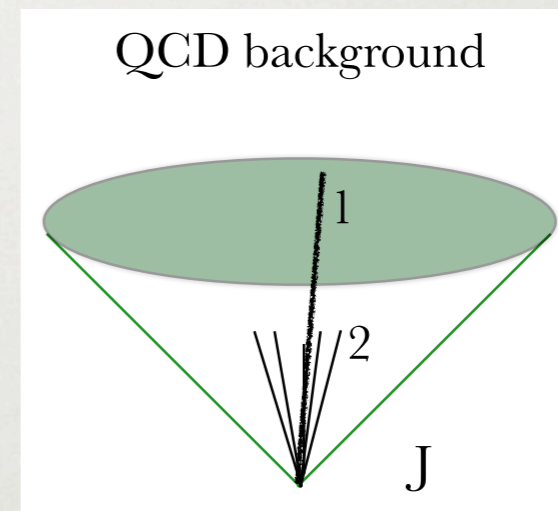
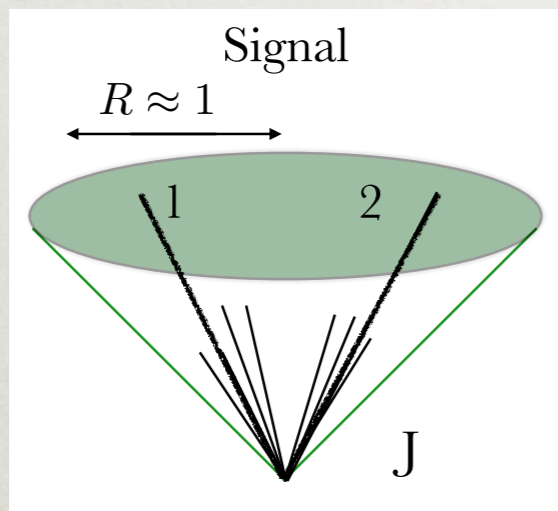


- Current approaches:
  - Some searches highly optimized, using sophisticated multivariable techniques (H to bb searches)
  - Others place simple cuts on jet kinematics and do a bump hunt ( $Z' \rightarrow WW$  semileptonic, SM  $WW+WZ$  semileptonic,...)
  - Can we do better?



# Hadronic resonances

- One method to improve sensitivity: move to new kinematic regime
  - Large boost means resonance decay products merge into “fat jet”
  - QCD jet looks nothing like  $t/W/Z/H$  jet - much recent progress on this!
  - Good for highly boosted regime, but suffer large kinematic penalties
  - Extra sensitivity to pile-up, underlying event ( $\sim R^3$ )

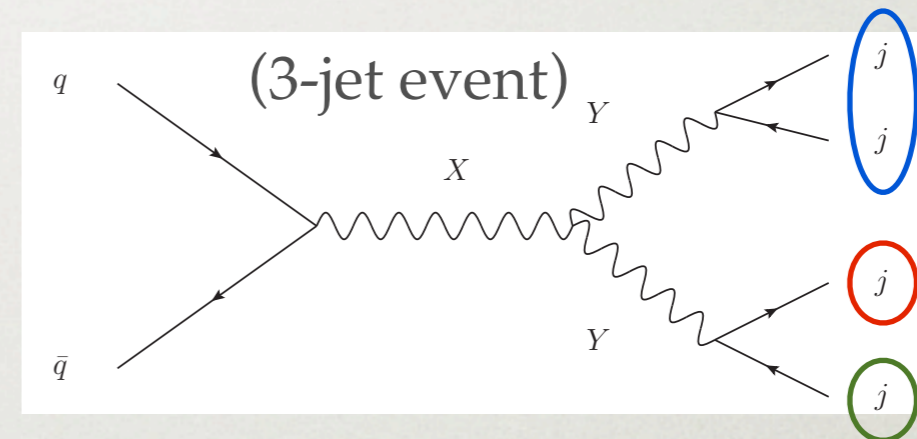
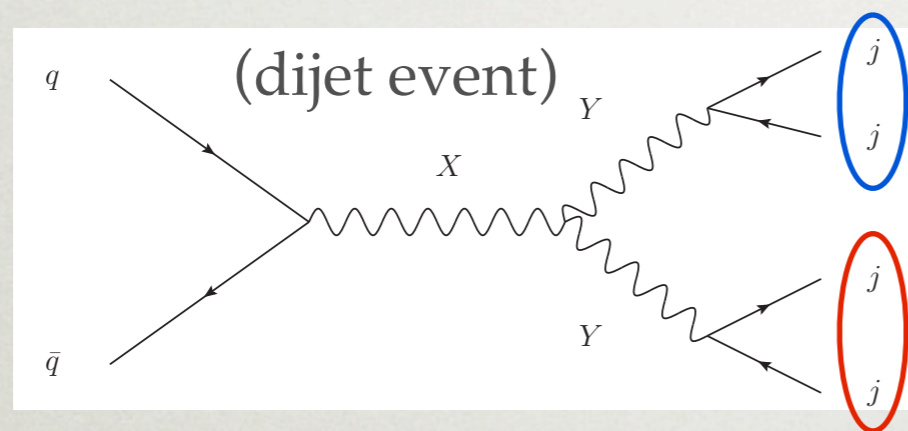


- Can we still use some of the same strategies to improve identification of **moderately boosted** resonances?
  - In this case, the decay products of the resonance are separately resolved



# Hadronic resonances

- Past work attempted to combine resolved and substructure searches
  - Gouzevitch *et al.*, arXiv:1303.6636
  - Look at pair production of dijet resonances in 2j, 3j, 4j final states
  - Use substructure when decay products merged, apply similar cuts to resolved jets when not
  - Find that jet substructure observables don't give much discriminating power away from highly boosted regime



- Are there still ways to separate resonance decay from “hard” QCD splitting?
  - We define a new observable that can be added to existing searches with a factor 2-6 gain in  $S/B$
  - Outperforms other possible cuts; includes resolved jet masses



# Outline

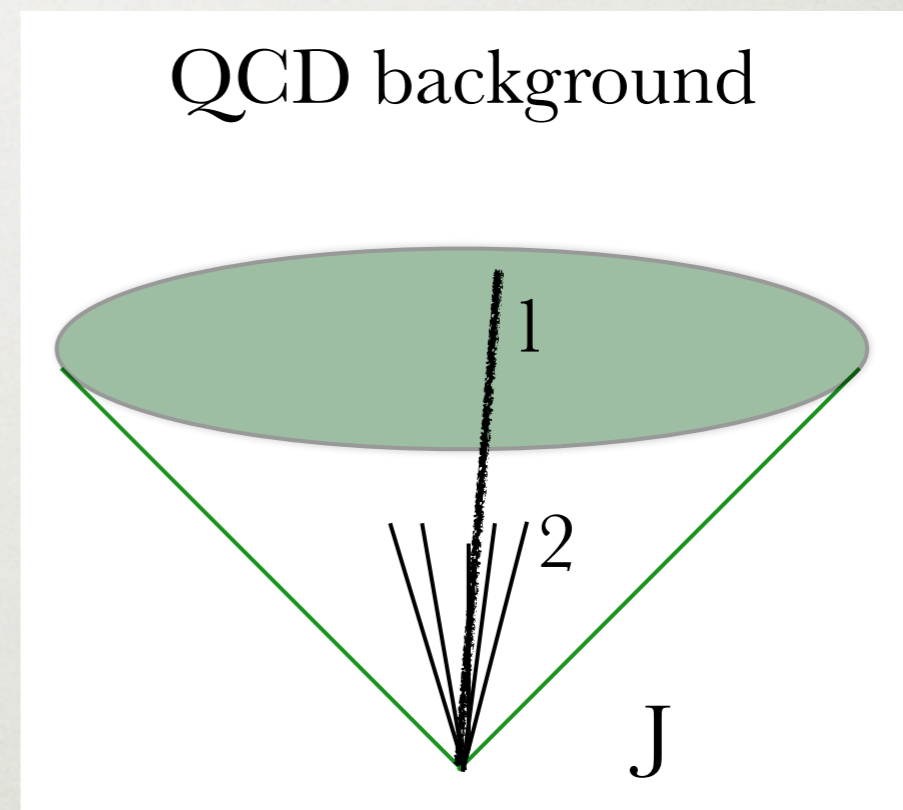
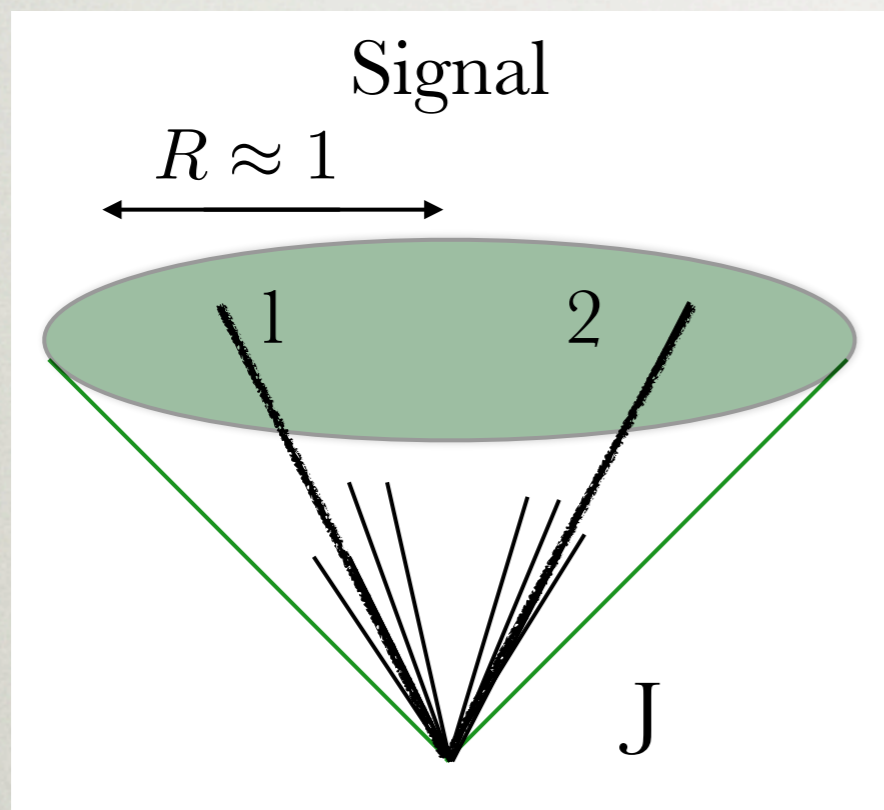
1. Jet substructure and the highly boosted regime
2. Resonance tagging in the mildly boosted regime
3. Examples
  - SM:  $WW+WZ$
  - SM:  $V(H\rightarrow bb)$
  - BSM:  $Z' \rightarrow WW$
4. Future directions



# Jet substructure at high boost

- When an object is highly boosted, its decay products are collimated
  - Can be clustered together into a single, “fat” jet

$$p_T \gtrsim (\text{few}) \times m_{\text{resonance}}$$

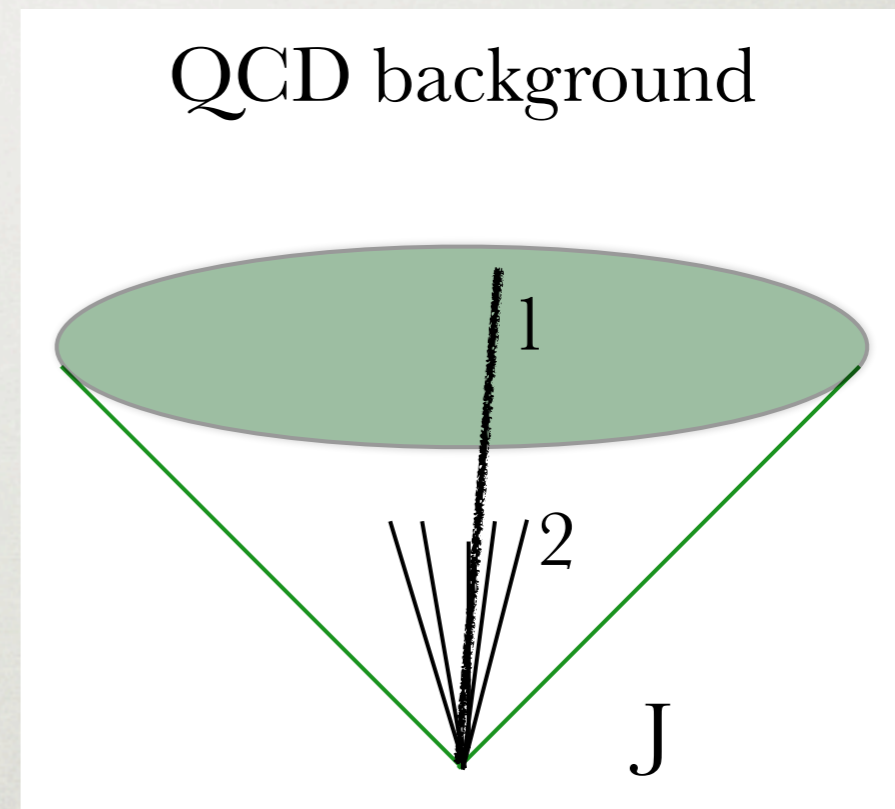
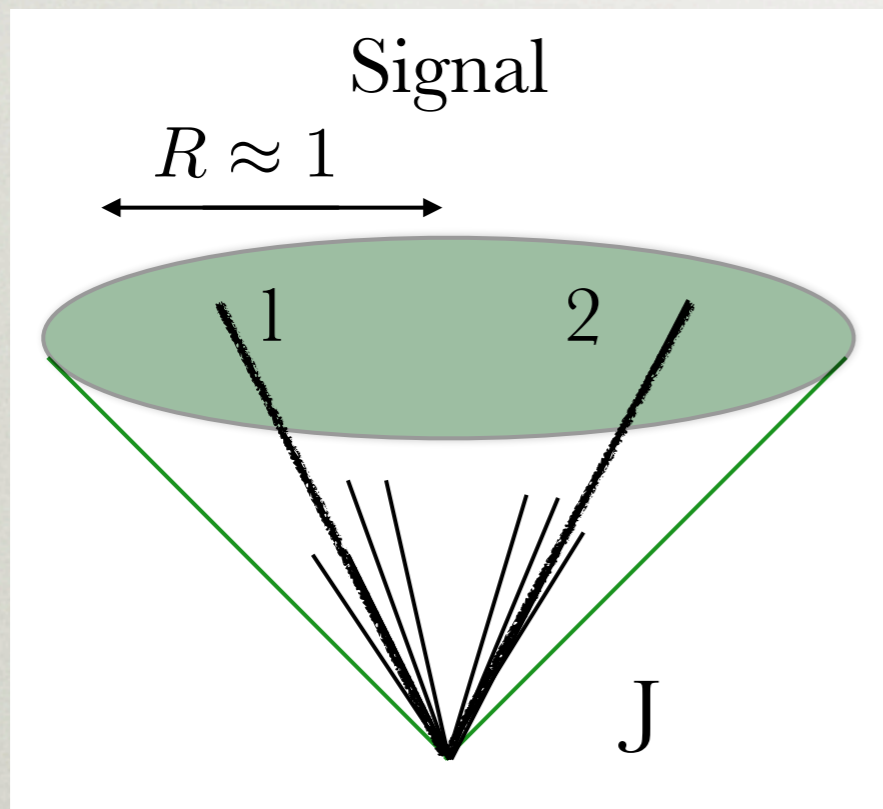


- Dominant background originates from a single QCD parton



# Jet substructure at high boost

- The signal typically gives two hard **subjets** from the decay of a resonance, while the QCD subjects typically come from parton shower
  - Can take either a **decomposition** approach or **energy-flow** approach





# Decomposition Approach

- Canonical example: BDRS mass-drop tagger (arXiv:0802.2470)  
(similar mass-drop procedure in HEPTopTagger)
  1. Cluster a jet  $j$  with the Cambridge / Aachen algorithm ( $R = 1.2$ )
  2. Undo the last step, splitting  $j$  into subjets  $j_1, j_2$  with  $m_{j_1} > m_{j_2}$
  3. Discard  $j_2$ , set  $j = j_1$ , and continue de-clustering until **both**:

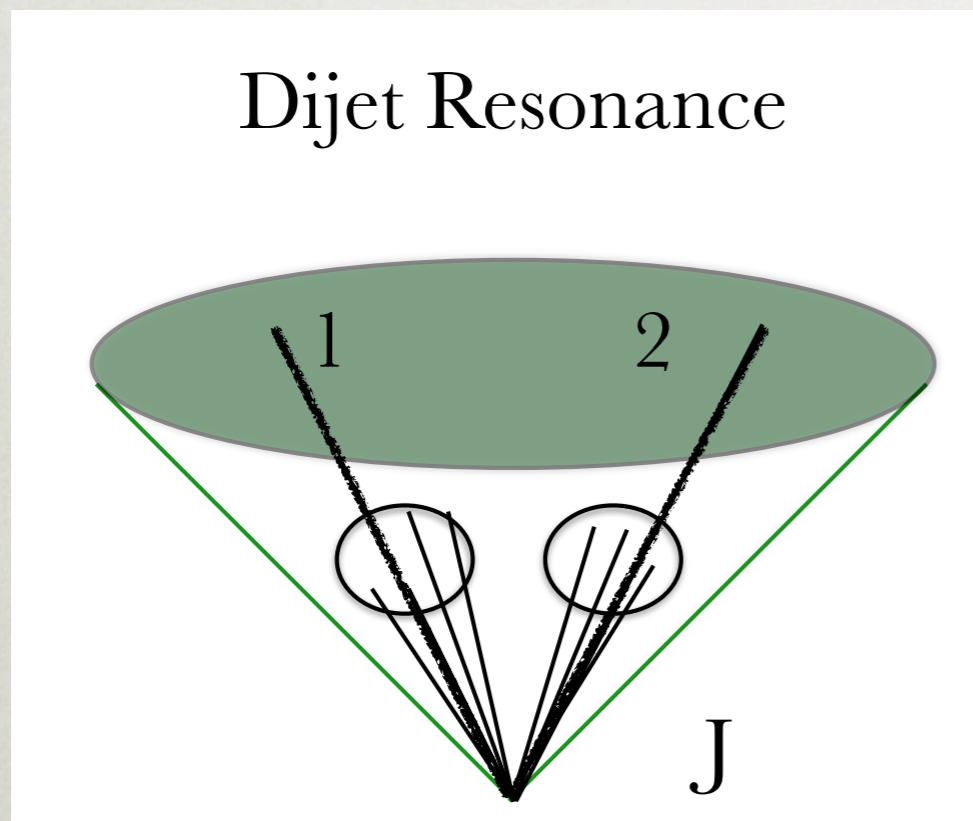
$$\frac{m_1}{m_j} < 0.67 \quad (\text{mass drop})$$

$$\frac{\min(p_{T1}^2, p_{T2}^2)}{m_j^2} > 0.09 \quad (\text{symmetric splitting})$$

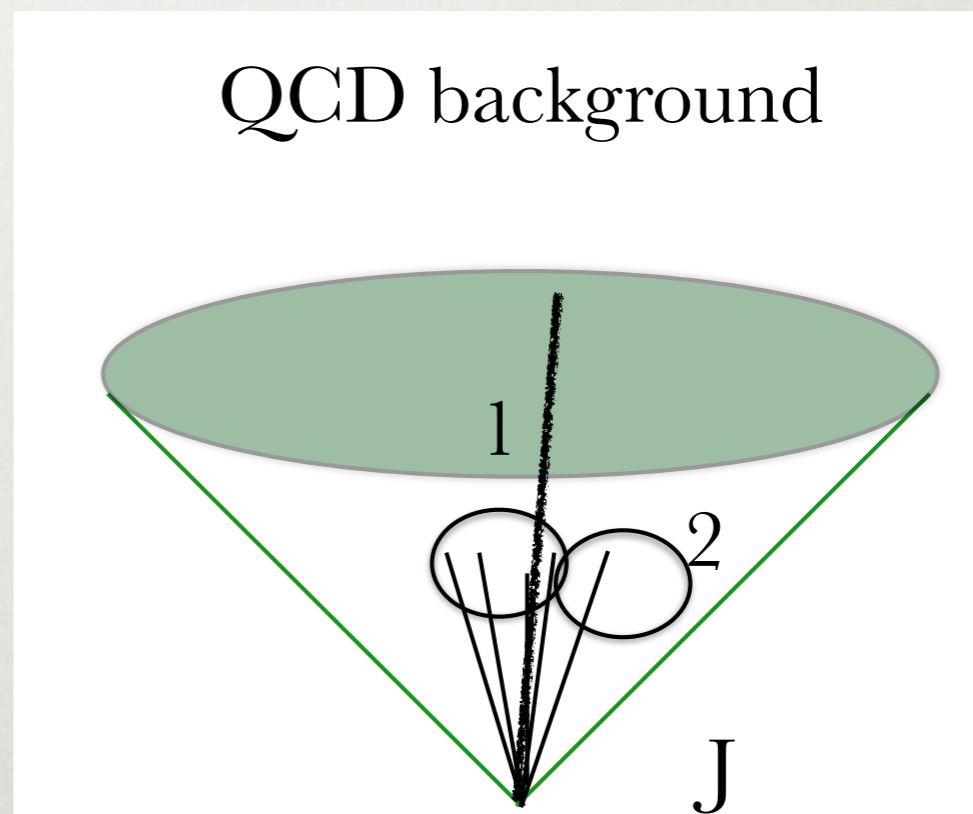


# Decomposition Approach

- This works because of their respective origins of jet masses
  - QCD masses come largely from sequences of asymmetric, wide-angle splittings
  - Signal masses come from the hard resonance decay



$$\frac{m_1}{m_j} \ll \mathcal{O}(1)$$



$$\frac{m_1}{m_j} \sim \mathcal{O}(1)$$



# Energy-flow approach

- Uses inclusive, energy-flow information (algorithm-independent)
  - ex.  $N$ -subjettiness (Thaler, van Tilburg, arXiv:1011.2268)
  - To compute  $N$ -subjettiness, define  $N$  axes (denoted by Greek letters) and associate each particle  $i$  in the jet to the closest axis

$$\tau_N^\beta = \frac{1}{\mathcal{N}} \sum_{i \in j} p_{Ti} \Delta R_{i\alpha}^\beta$$

- This gives a measure of how well the radiation is aligned along  $N$  axes
  - For dijet resonances,  $\tau_2/\tau_1$  performs better than  $\tau_2$
  - Combines information from hard and soft radiation
- Generally, both approaches work well and are complementary
  - See upcoming BOOST 2013 working group report



# **Moderately boosted resonances**



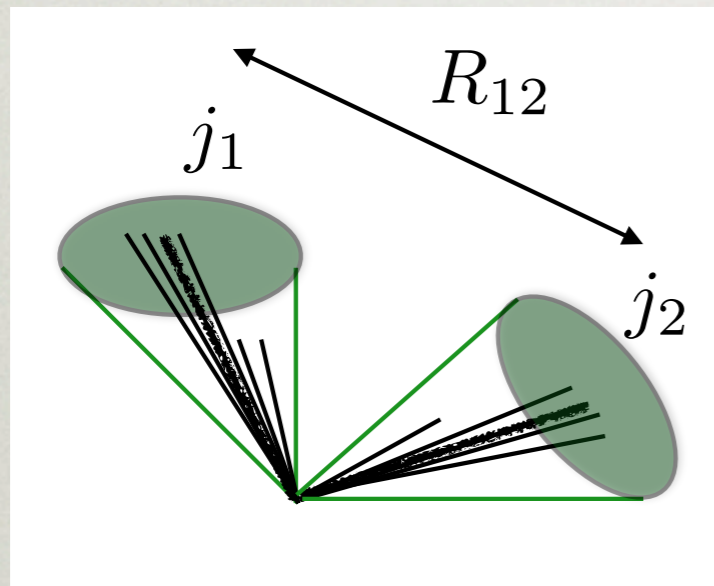
# Tagging at moderate boosts

- Often, resonances are produced near threshold, paying a high penalty in signal acceptance for going to the boosted regime
  - Direct  $t\bar{t}$ , diboson, ...
- We can try to recover some of the signal discrimination in the **moderately** boosted regime
  - Still dominated by QCD splittings instead of uncorrelated emission
  - Requiring two separately resolved jets already eliminates most of the backgrounds from soft splittings
  - For further improvement, must now separate relatively hard QCD splittings from the signal
- Use decomposition approach because energy-flow observables are sensitive to all radiation between the jets
  - No longer need large- $R$  jets



# Tagging at moderate boosts

- Analogy of mass drop
  - The lax cut on mass drop from the boosted regime ( $<0.67$ ) does not veto a hard QCD splitting
  - As jets become more widely separated, the mass drop becomes **smaller**
  - For background:



$$\langle m_1^2 \rangle \approx C \frac{\alpha_s}{\pi} R^2 p_{T1}^2$$

$$m_{12}^2 \sim p_{T1} p_{T2} \Delta R_{12}^2$$

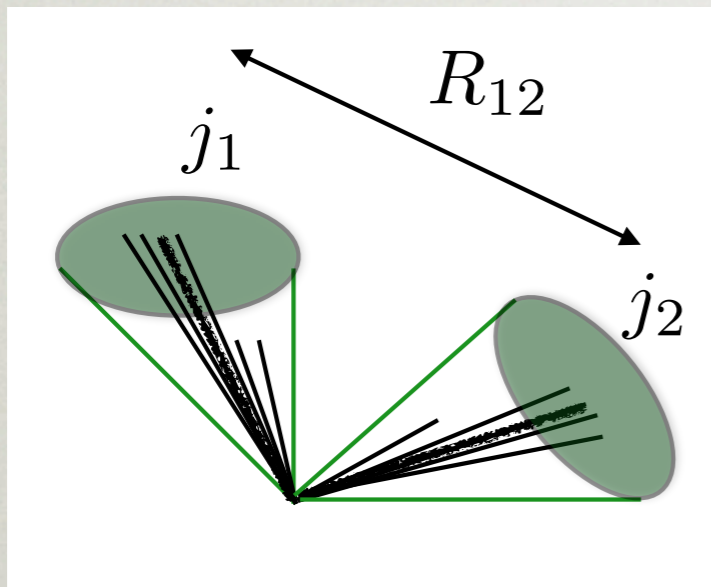
$$\frac{m_1}{m_{12}} \propto \frac{1}{\Delta R_{12}}$$

- Interpolates between boosted and unboosted regimes



# Tagging at moderate boosts

- This suggests that we **scale** the mass drop cut with  $\Delta R$
- Exploit differences between QCD splitting and resonance decay
  - Signal has a mass drop that is more constant in  $\Delta R$
  - QCD prefers asymmetric splittings, giving rise to **larger**  $m_1$
  - In many examples with QCD backgrounds, one of the radiated partons is a gluon ( $C_A > C_F$ ), giving rise to **larger**  $m_1$  (on average)



- This motivates a new observable:

$$\zeta \equiv \frac{m_1}{m_{12}} \Delta R_{12}$$

$$\zeta < \zeta_c$$



# Tagging at moderate boosts

- Other functional forms could accomplish a similar scaling
  - For example:

$$\zeta(R_c) = \frac{m_1}{m_{12}} (\Delta R_{12} - R_c)$$

- These types of observables can be very effective at enhancing  $S/B$  when added on top of existing searches
  - Outperform other observables we studied
  - Robust performance under simple smearing and with different MC
  - Uses simple, small- $R$  jet properties



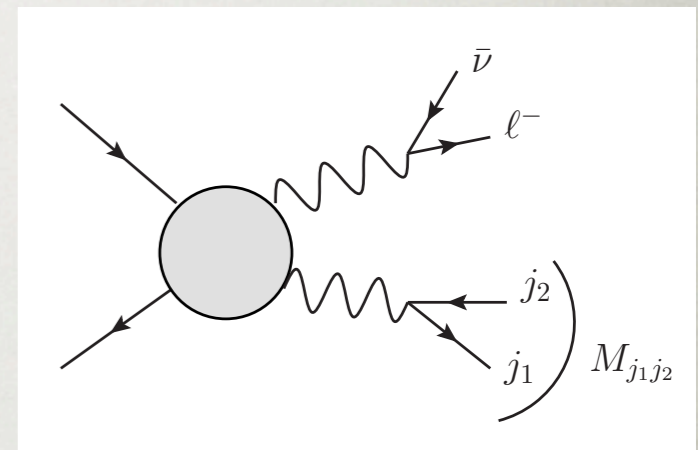
# Examples

- SM:  $WW+WZ$
- SM:  $V(H\rightarrow bb)$
- BSM:  $Z' \rightarrow WW$



# WW+WZ Analysis

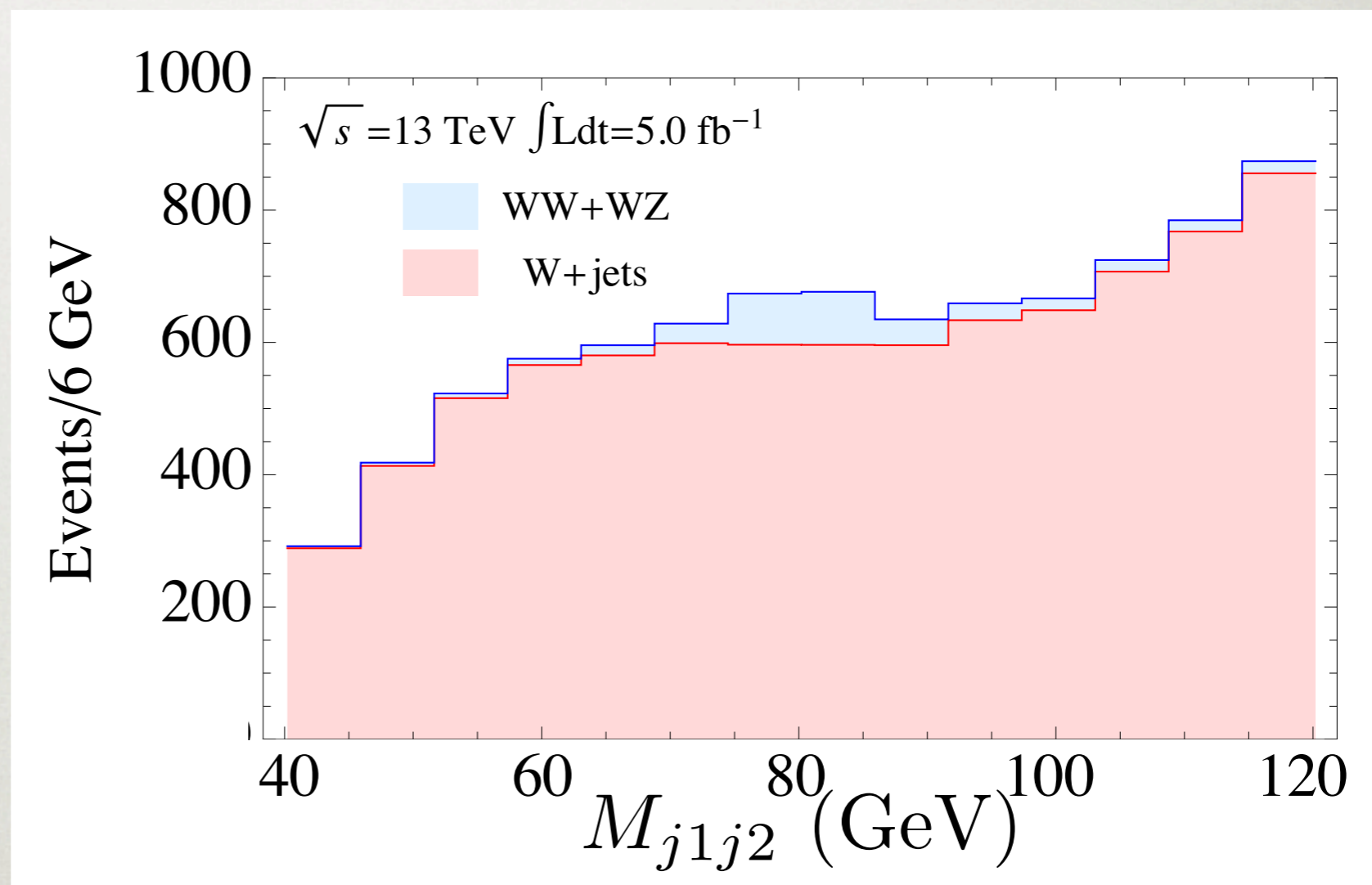
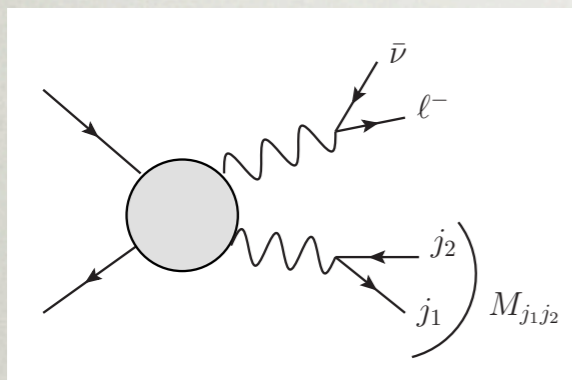
- Semileptonic channel is an independent check of the (possible) excess in the fully leptonic channel and an important SM measurement
- Simulate WW+WZ, W+jets events with Madgraph 5
  - Match matrix element to Pythia 6 parton shower using shower- $k_{\perp}$  scheme
  - Cluster and analyze events with Fastjet 3
  - Validated MC with CMS analysis
  - Include UE but no pile-up (more on this later)
- Use similar cuts as CMS 7 TeV (arXiv: 1210.7544), re-scaled to 13 TeV
  - Two jets with  $p_T > 50$  GeV
  - One lepton with  $p_T > 25$  GeV
  - MET > 50 GeV
  - $M_T > 50$  GeV





# WW+WZ Analysis

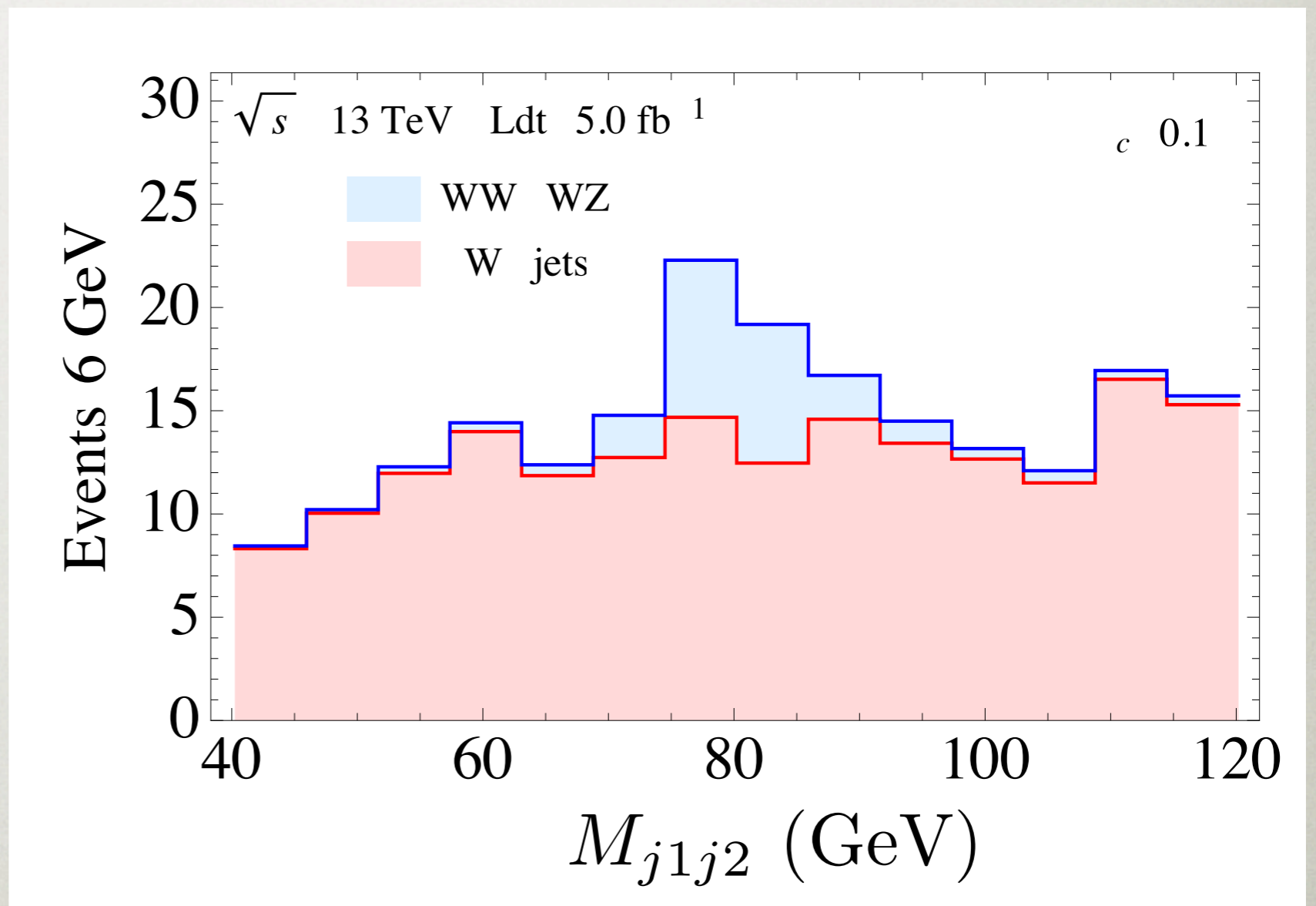
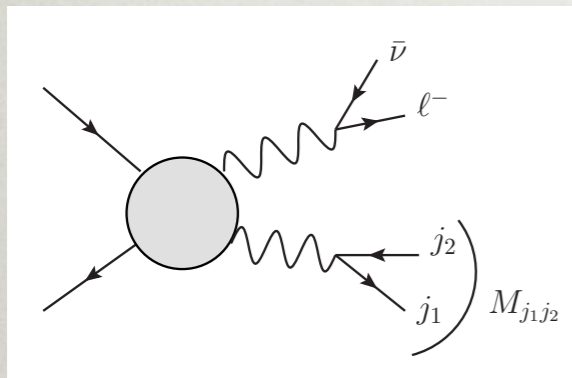
- After CMS selection cuts:





# WW+WZ Analysis

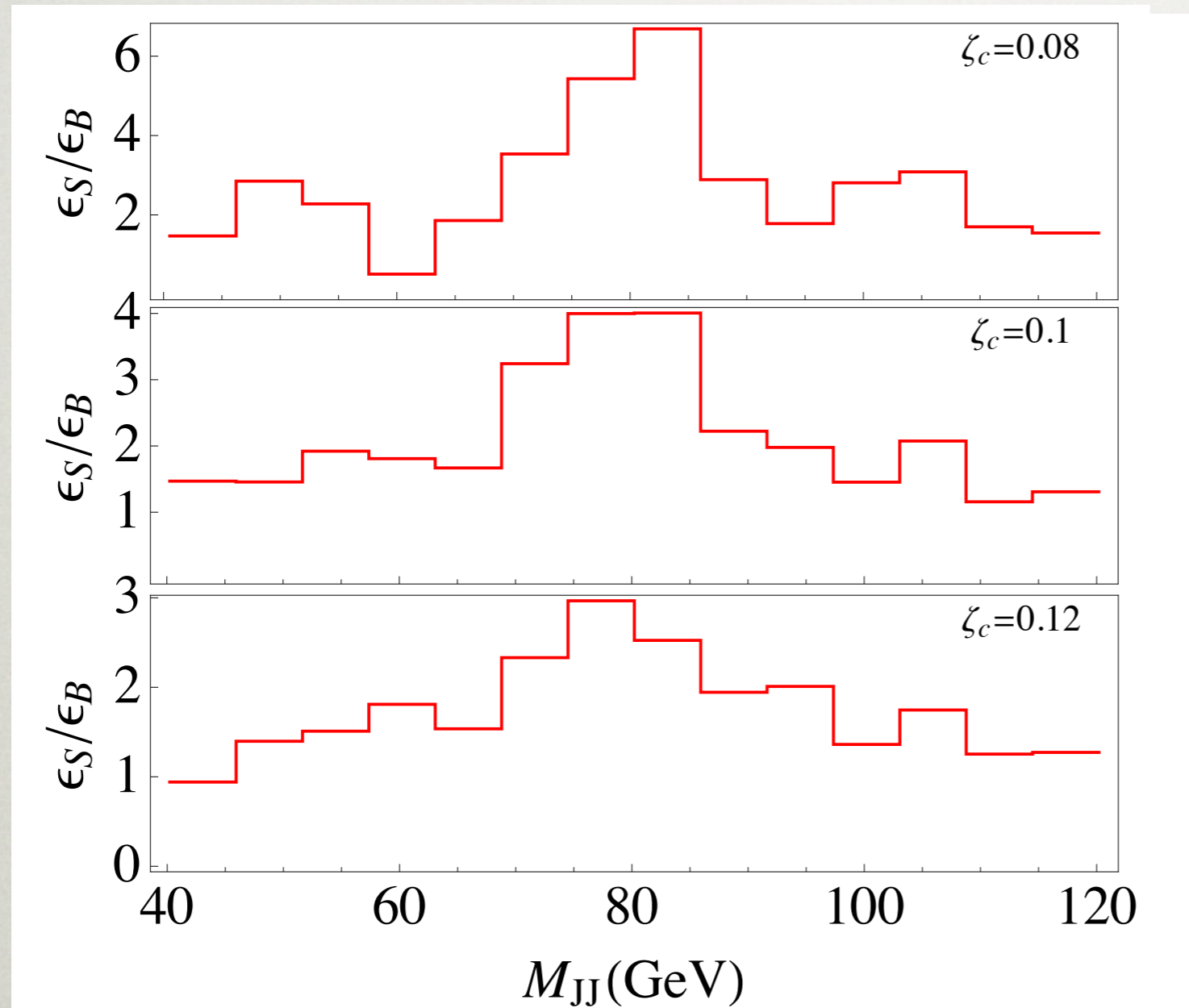
- After CMS selection AND cut on  $\zeta < \zeta_c$





# WW+WZ Analysis

- Gains for different choices of the cut



$$\epsilon_S(\zeta < \zeta_c) \approx 3\%$$

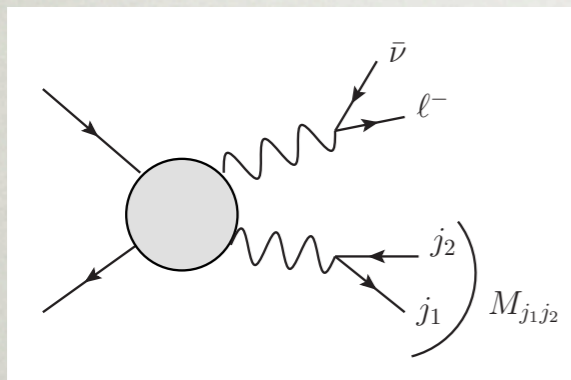
$$\epsilon_S(\zeta < \zeta_c) \approx 10\%$$

$$\epsilon_S(\zeta < \zeta_c) \approx 25\%$$

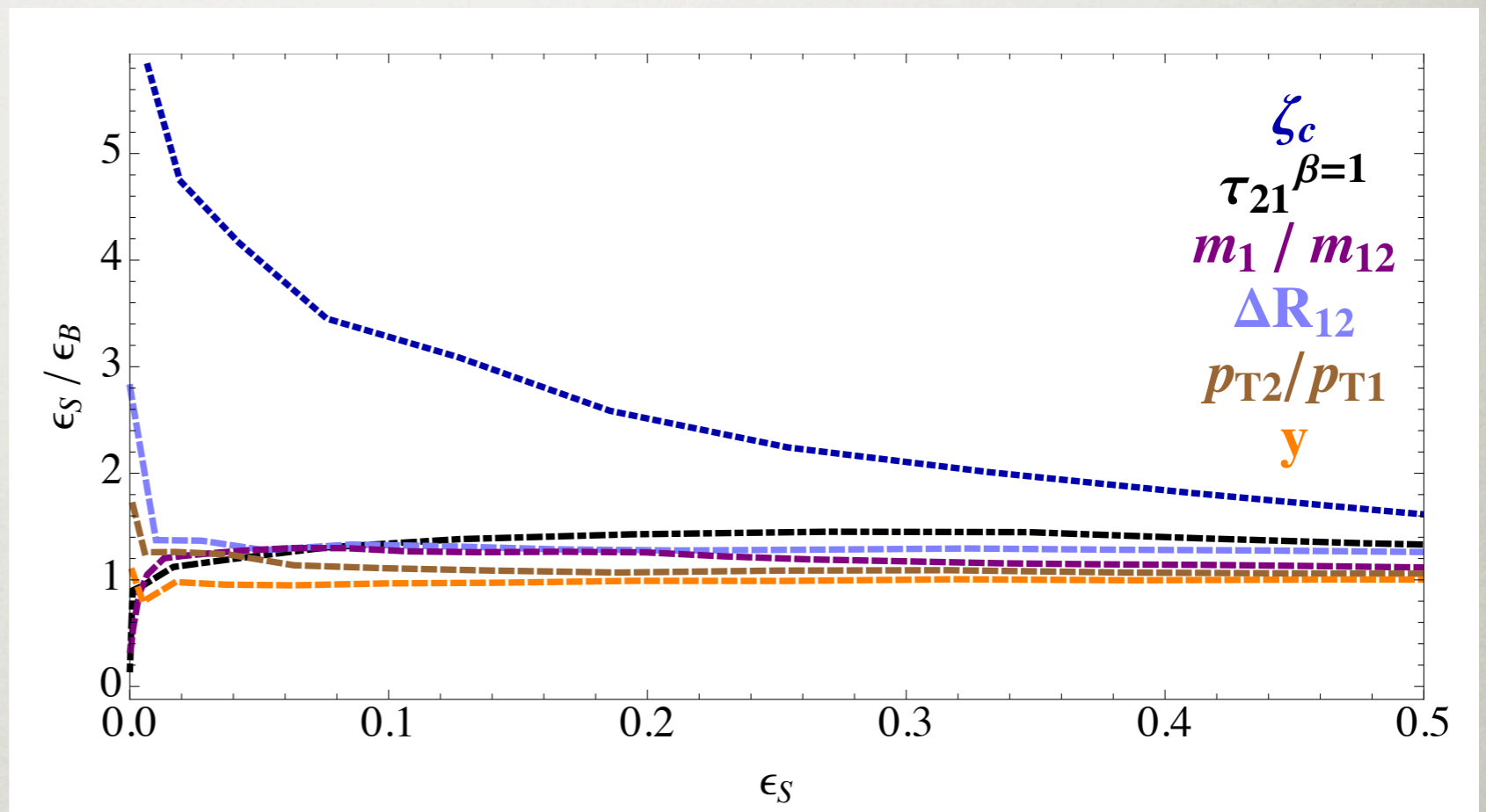


# WW+WZ Analysis

- How does this compare to other possible cuts we could have used?
  - Look in  $M_{j_1 j_2}$  window between 70-100 GeV
  - Use **filtered** version of shape / energy-flow observables



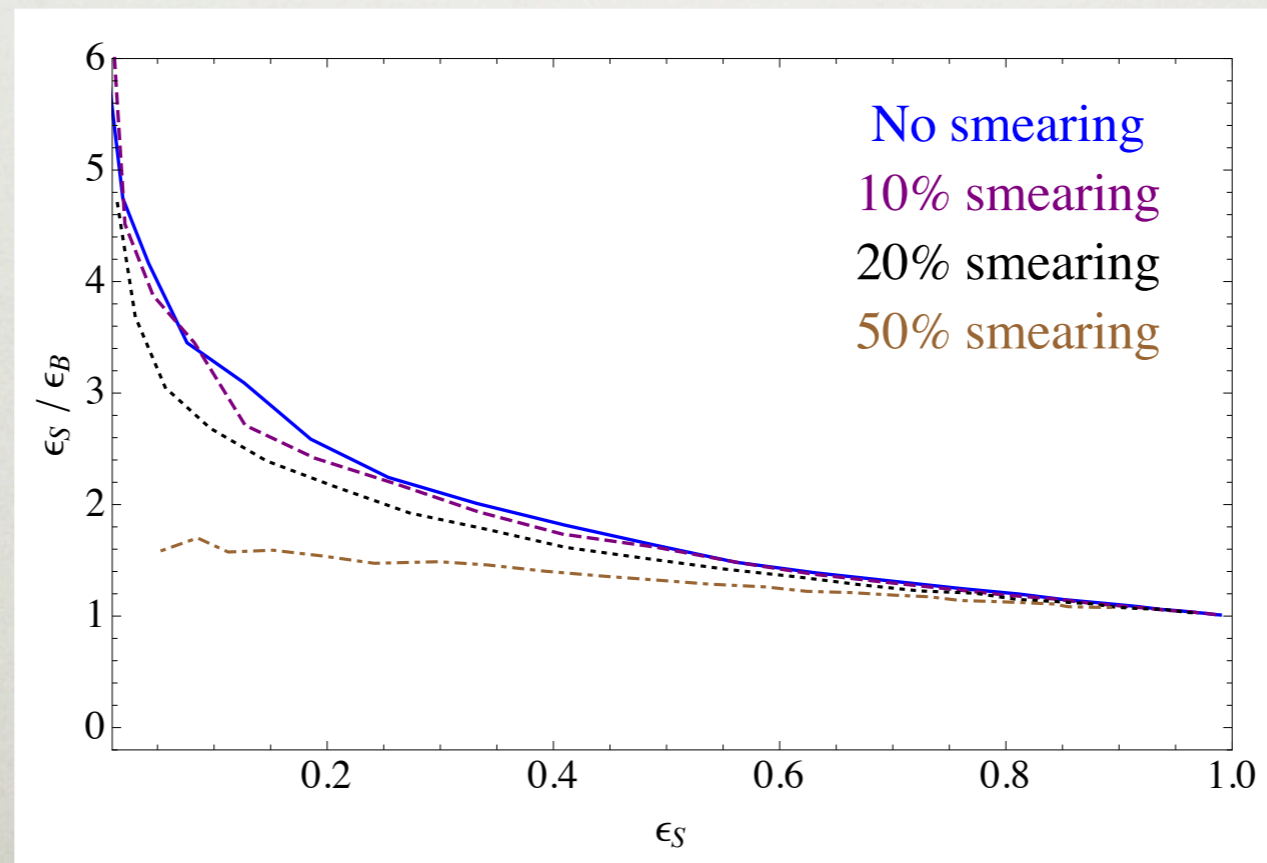
$$y = \frac{p_{T2}^2}{M_{12}^2} \Delta R_{12}^2$$





# WW+WZ Analysis

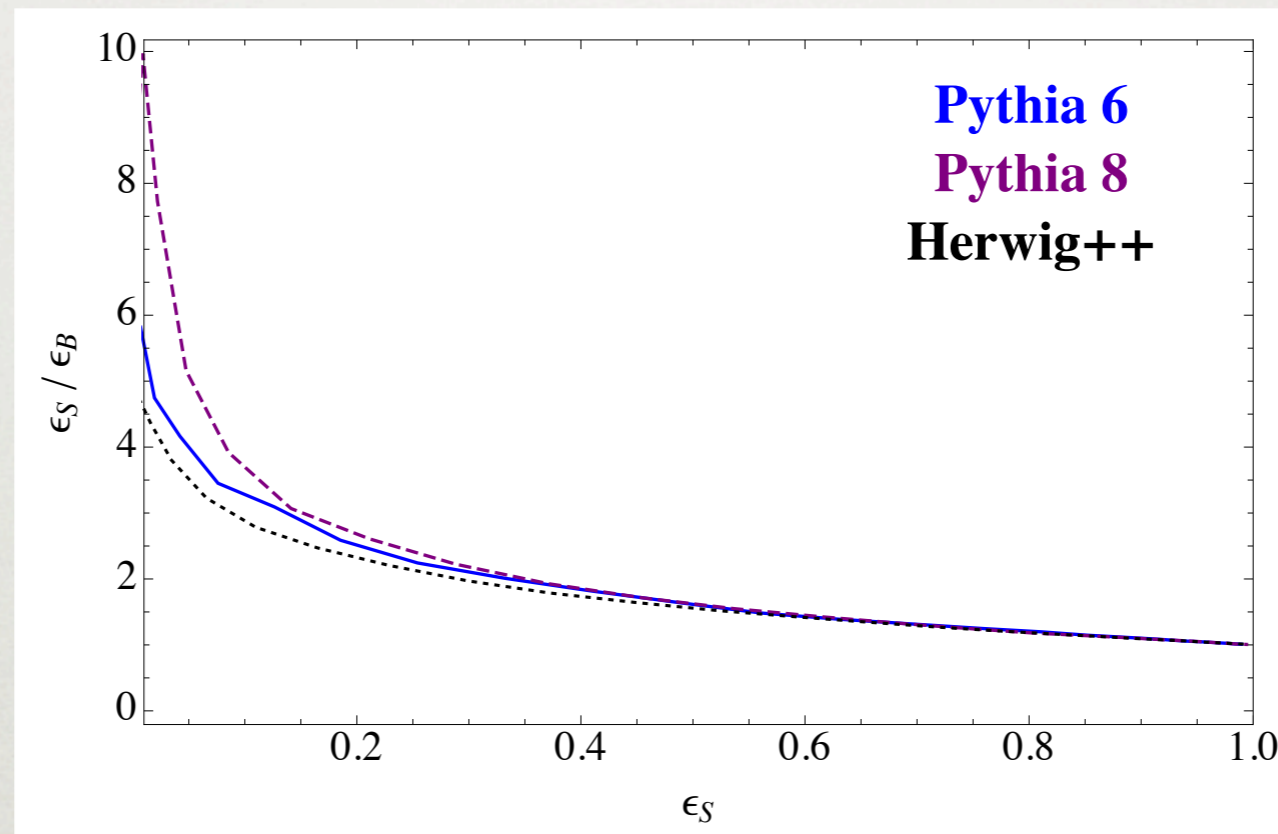
- Would this be included in a BDT analysis?
  - Not currently used for SM WW+WZ
  - Seems there is substantial gain that comes from using resolved jet *masses*, which are not included in most BDT analyses
- Possible worry: jet masses are subject to uncertainties in shower mechanism & reconstruction





# WW+WZ Analysis

- Possible worry: jet masses are subject to uncertainties in shower mechanism & reconstruction
  - Show Pythia 6 vs. Pythia 8 vs. Herwig++



- Zeta performs well and is robust against various uncertainties



# Limitations and Caveats

- Our observable gives a significant enhancement in  $S/B$  at the cost of a mild reduction in statistical significance
  - Most applicable to searches dominated by systematic uncertainties
  - Will become more relevant for later LHC running
- What about pile-up?
  - Serious challenge facing high-luminosity running
  - We simulated WW+WZ search with  $\langle N_{PV} \rangle = 50$ , found that a more aggressive form of **jet trimming** recovered  $S/B$  gains to within 10-20%
  - Ongoing work needed for pile-up mitigation of small- $R$  jet masses
  - Our observable only involves small- $R$  jets



# Examples

- SM:  $WW+WZ$
- SM:  $V(H\rightarrow bb)$
- BSM:  $Z' \rightarrow WW$



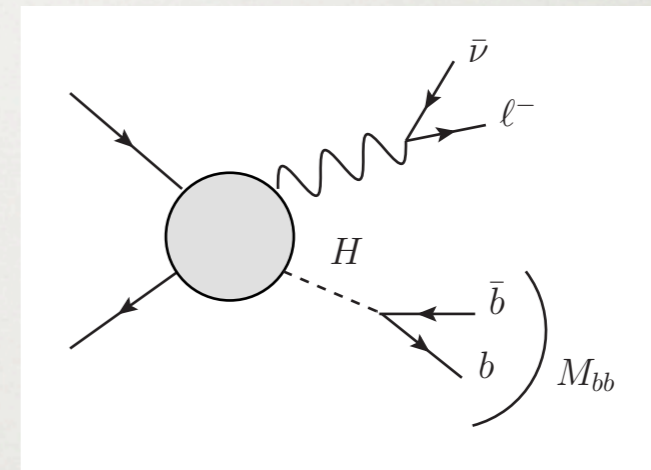
# W(H→bb) Analysis

- ATLAS and CMS have both dijet-mass and multivariate analyses
  - We follow the ATLAS 7+8 TeV analysis (now arXiv:1409.6212)

- Focus on dijet search, associated leptonic W
  - Dominant backgrounds are W+b+jets, tt

- Use same selection cuts as ATLAS

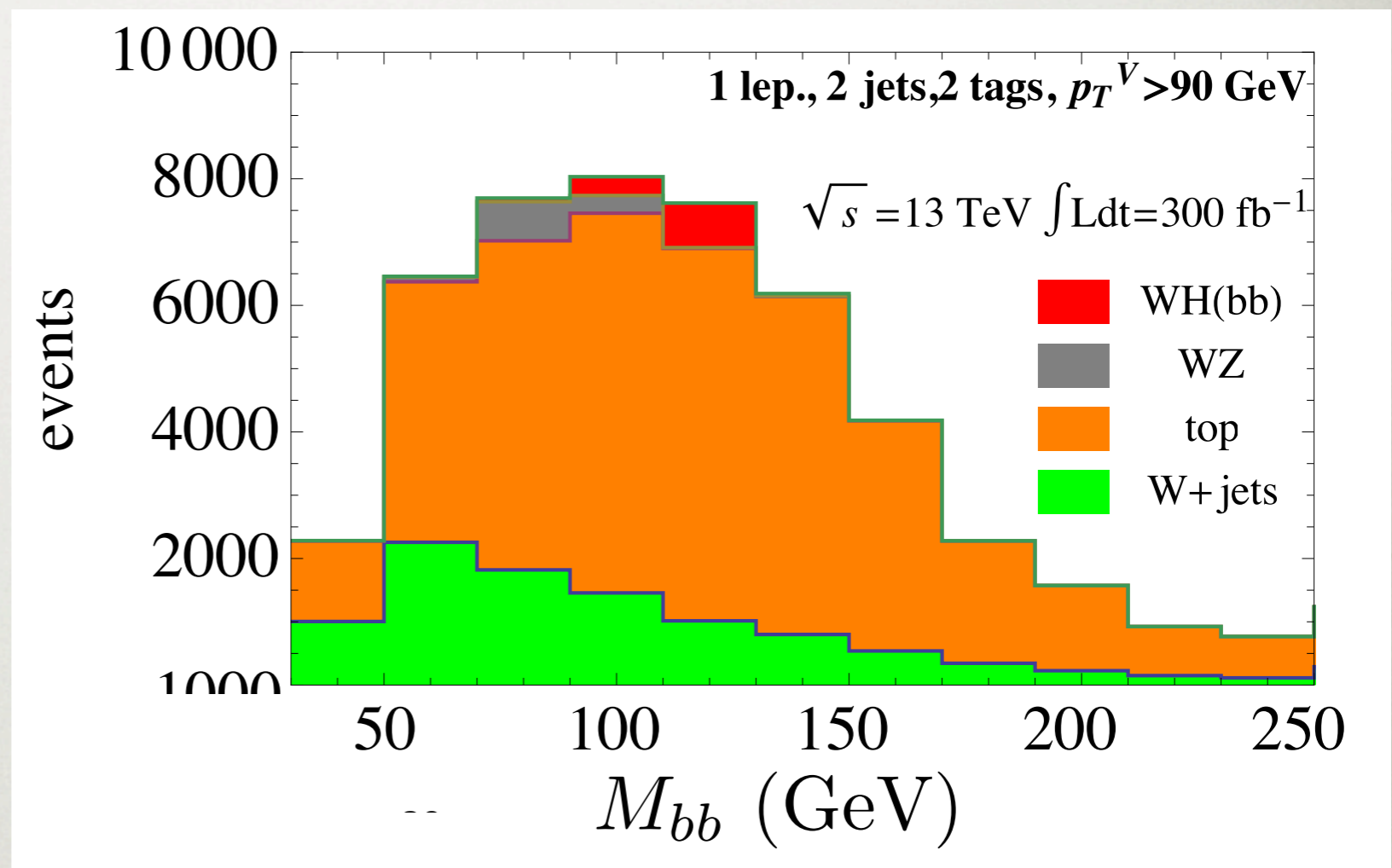
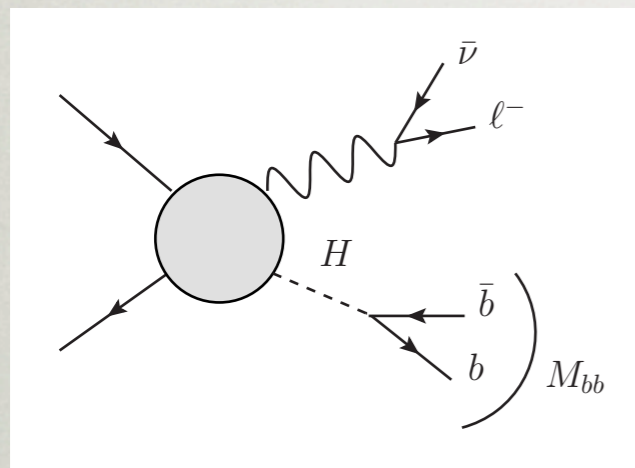
- One tight lepton,  $p_T > 25$  GeV
- Exactly 2 b-tagged jets,  $p_T > 20$  GeV (leading jet  $p_T > 45$  GeV)
- $MET > 25$  GeV
- $120 \text{ GeV} > M_T > 40 \text{ GeV}$
- Loose selections on  $\Delta R_{bb}$  as a function of  $p_T$
- Associate muons with adjacent b-jets to improve mass reconstruction





# W(H → bb) Analysis

- After ATLAS selection cuts:

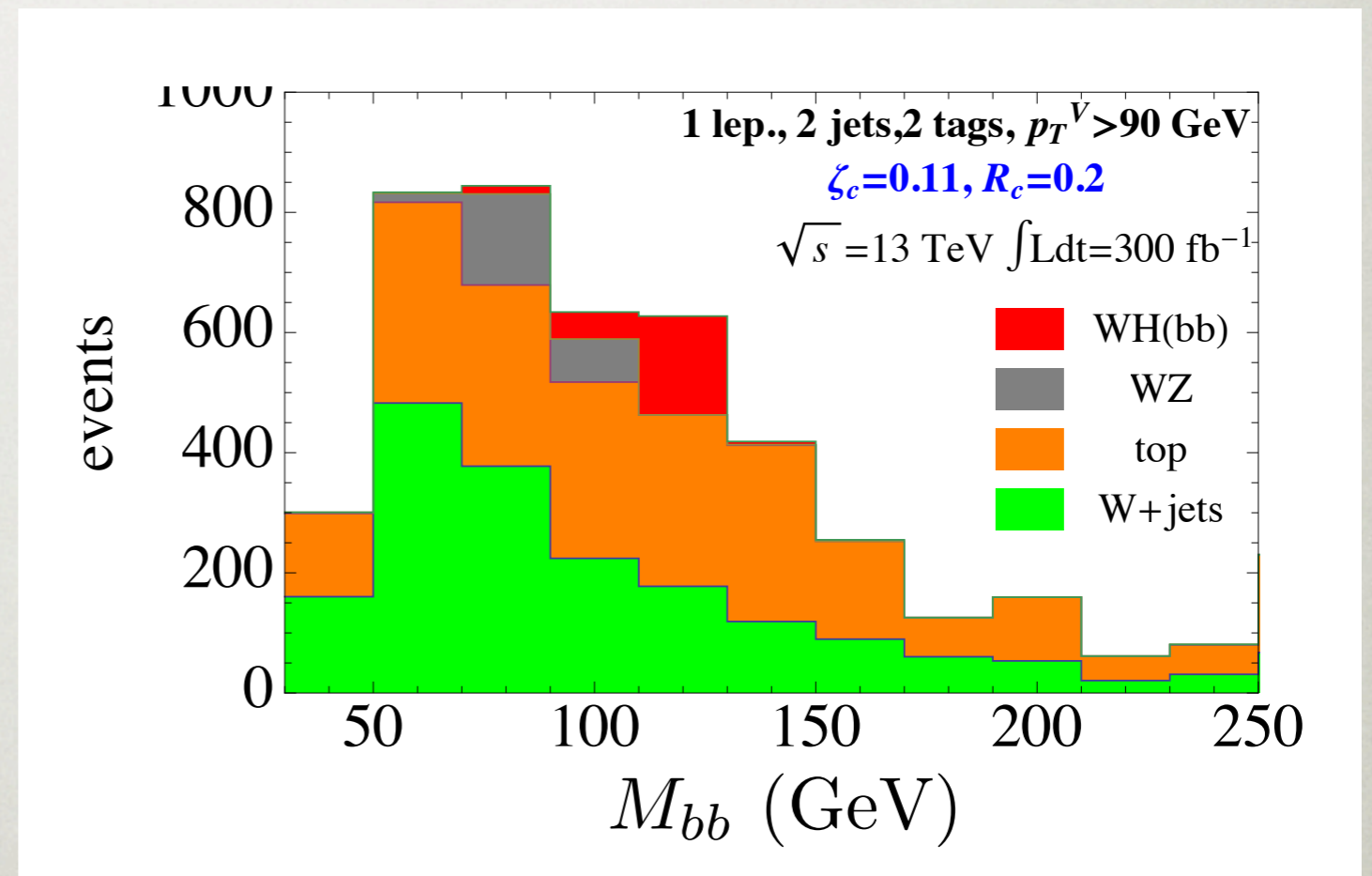
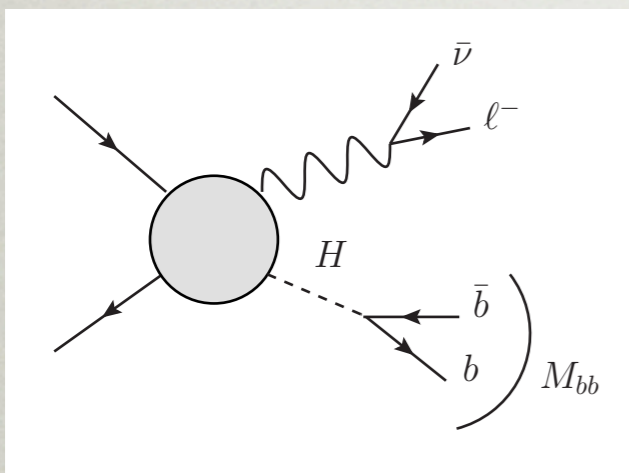




# W(H→bb) Analysis

- After ATLAS selection and a cut on the shifted version of  $\zeta$ :
  - Better at balancing preserving statistics and  $S/B$  gain

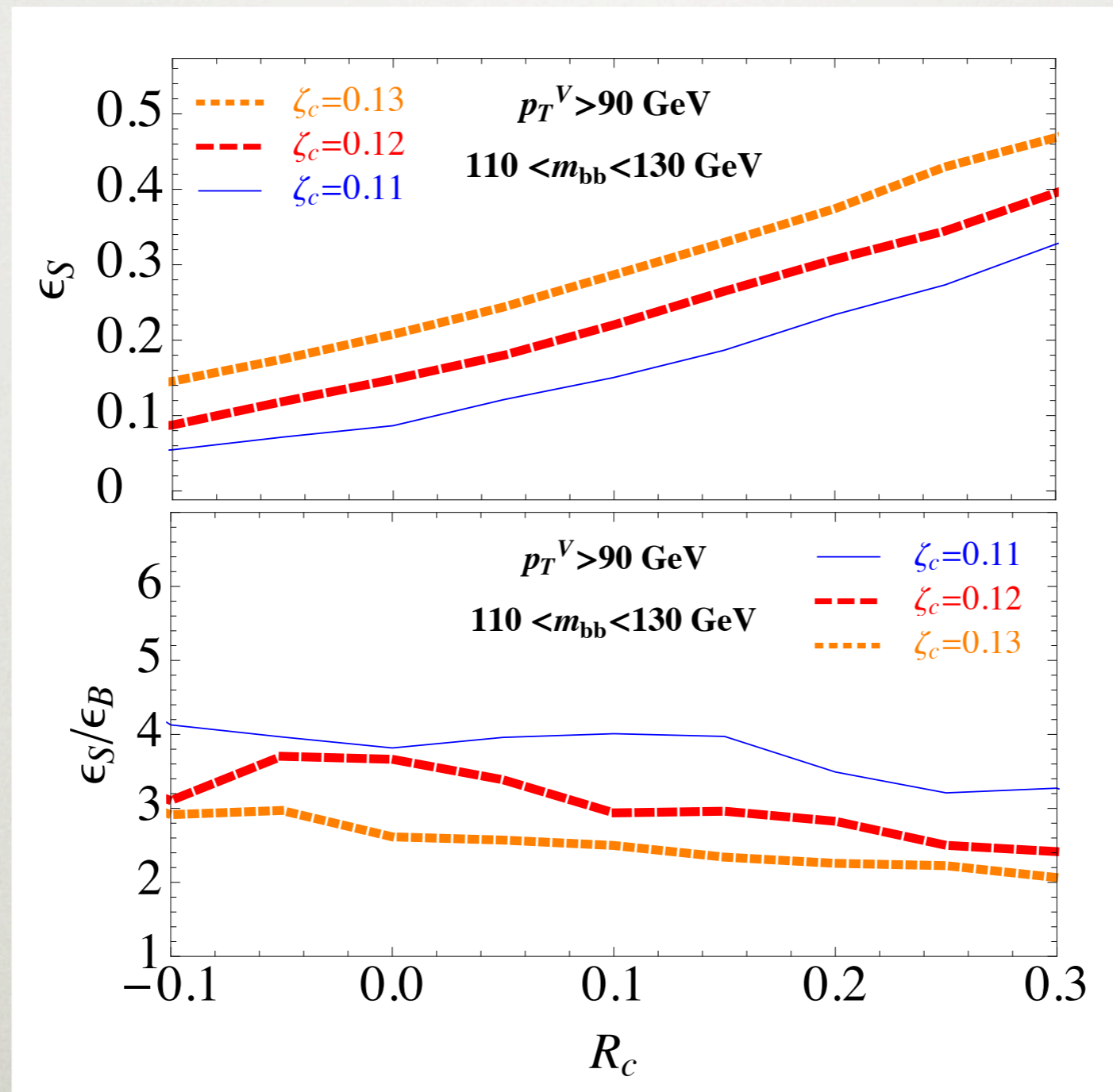
$$\zeta(R_c) = \frac{m_{j_1}}{m_{j_1 j_2}} (\Delta R_{12} - R_c) < \zeta_c$$





# W(H → bb) Analysis

- Gains for different choices of the cut:





# **W(H→bb) Analysis**

- Is our gain just coming from the highly boosted region?
  - BDRS requires  $p_{TV} > 200$  GeV
- If we restrict ourselves to the **moderately boosted** regime,  $90 \text{ GeV} < p_{TV} < 200 \text{ GeV}$ :
  - We still find an  $S/B$  gain of  $\sim 2-3$  (reduction of  $\sim 25\%$ )
- Our observable is effective in a boost range complementary to BDRS and other substructure methods
- Consider inclusion of jet masses in more sophisticated BDT as well



# Examples

- SM:  $WW+WZ$
- SM:  $V(H\rightarrow bb)$
- BSM:  $Z' \rightarrow WW$



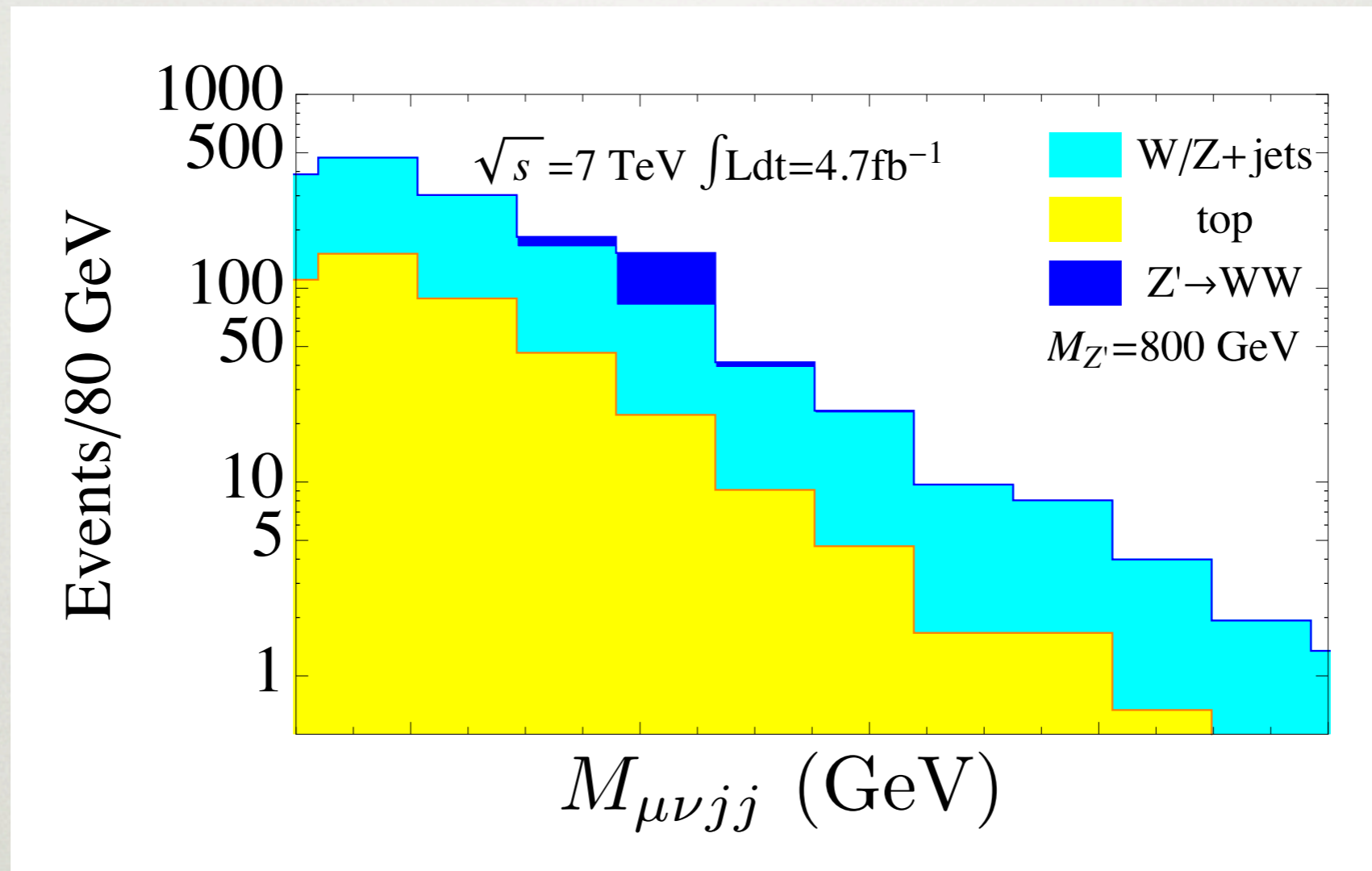
# $Z' \rightarrow WW$ Analysis

- ATLAS has a search for resonant semileptonic  $WW/WZ$  production for masses up to 1 TeV (arXiv:1305.0125)
  - At higher masses, use jet substructure techniques
  - We consider a sequential SM  $Z'$  decaying to  $WW$
  - Dominant background is  $W$ +jets
- Use same selection cuts as ATLAS
  - Two jets, at least one with  $p_T > 100$  GeV
  - One tight lepton,  $p_T > 35$  GeV
  - $MET > 40$  GeV
  - $p_{TV} > 200$  GeV for each candidate gauge boson
  - $65 \text{ GeV} < m_{jj} < 115 \text{ GeV}$
  - Various cuts on  $\Delta\phi_{\ell\nu}$



# $Z' \rightarrow WW$ Analysis

- After ATLAS selection cuts:

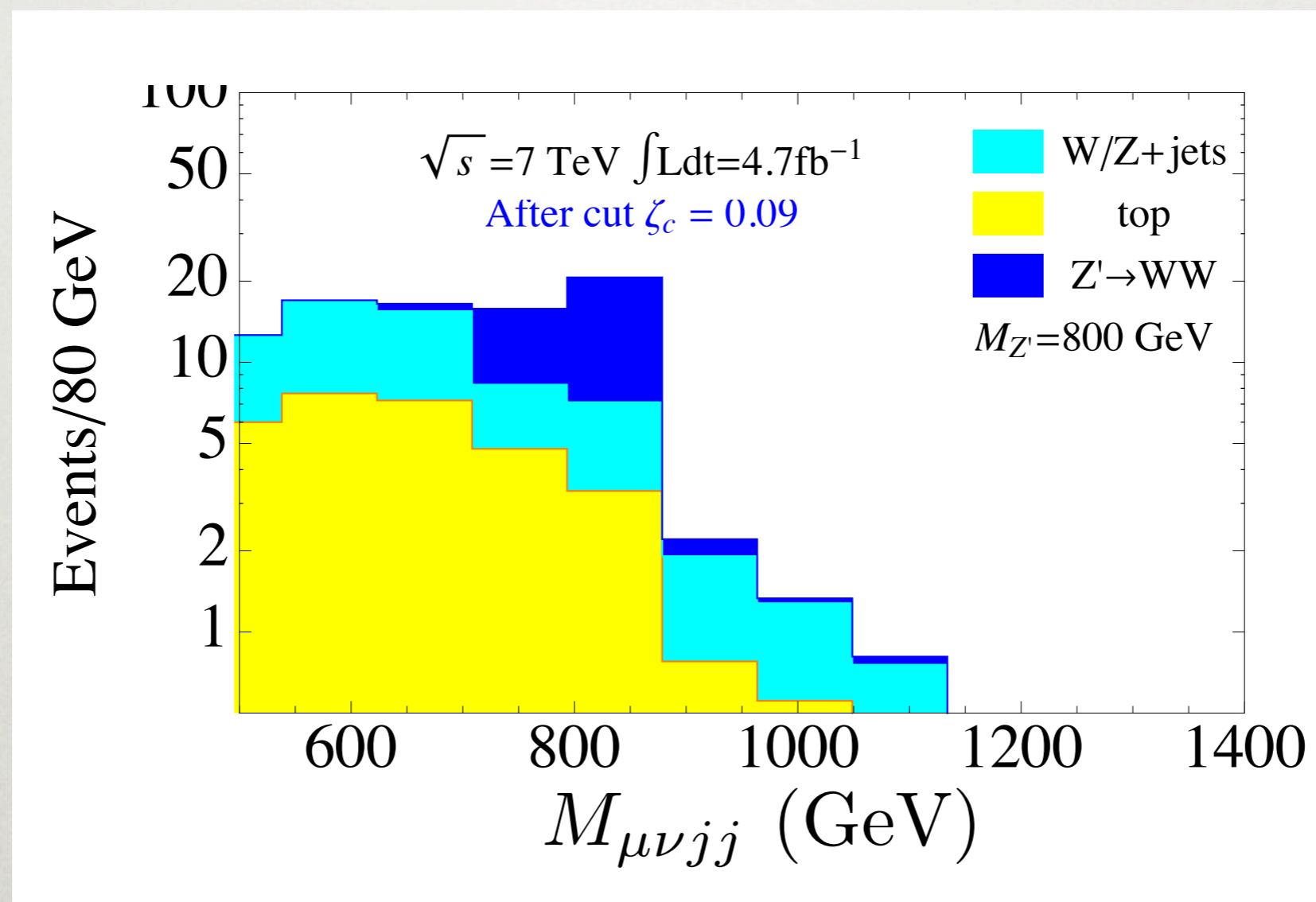


- Note: **large** systematic uncertainties ( $\sim 30\%$ )



# $Z' \rightarrow WW$ Analysis

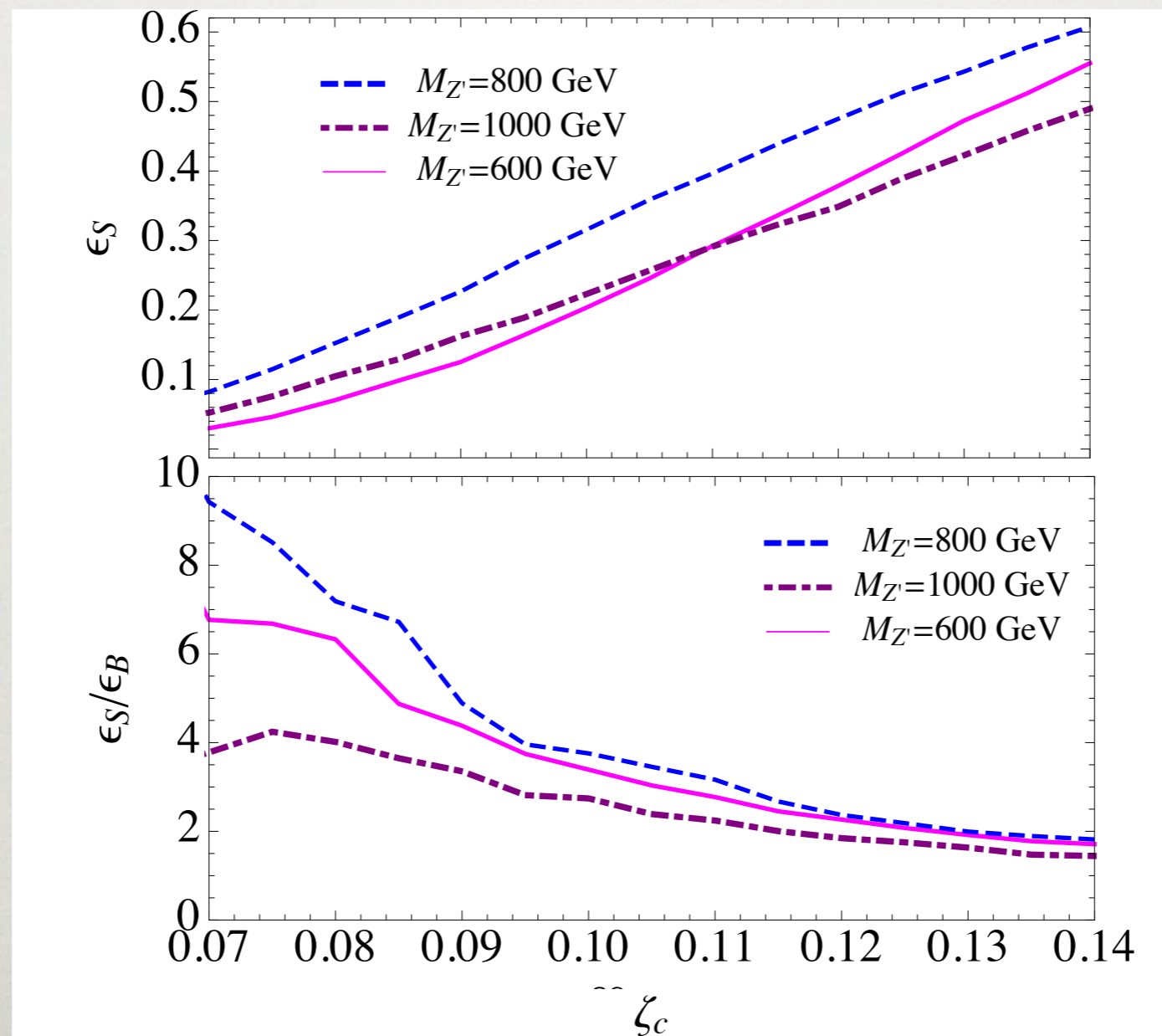
- After ATLAS selection AND cut on  $\zeta < \zeta_c$ :





# $Z' \rightarrow WW$ Analysis

- $S/B$  gains and efficiency change:



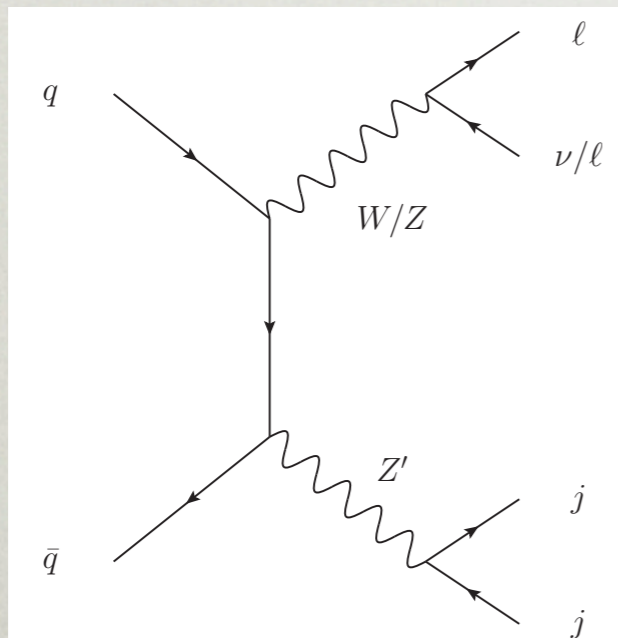
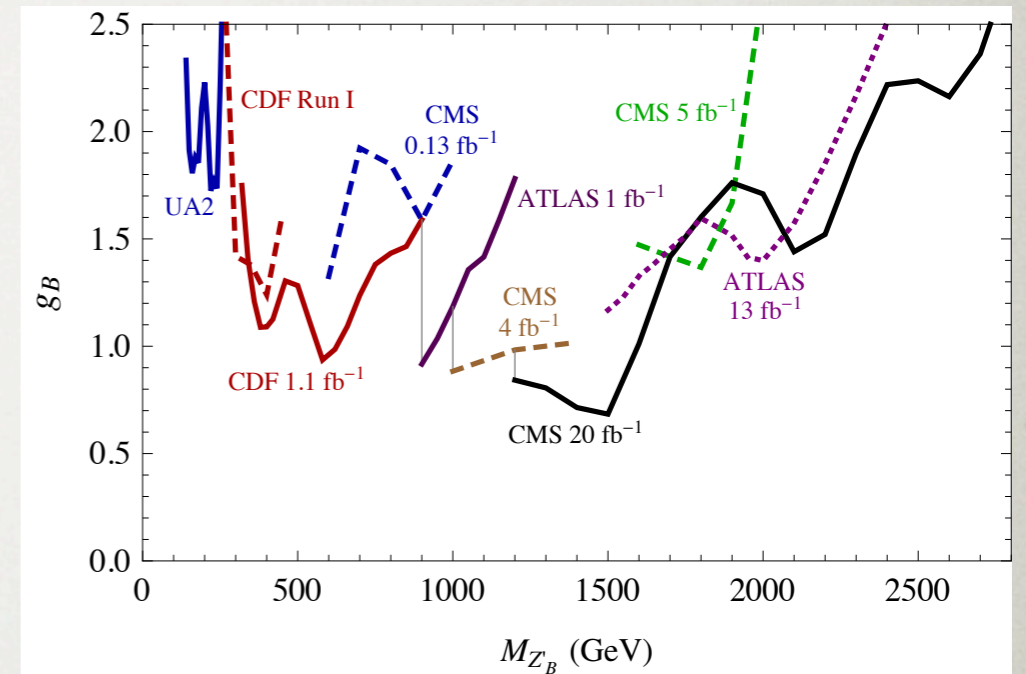


# **Future directions**



# Direct resonance production

- Best bounds come from UA2/Tevatron
- At LHC, hard to pass triggers and discriminate from backgrounds
- Consider associated production
  - Provides handle for trigger
  - Gives resonance a (mild) boost

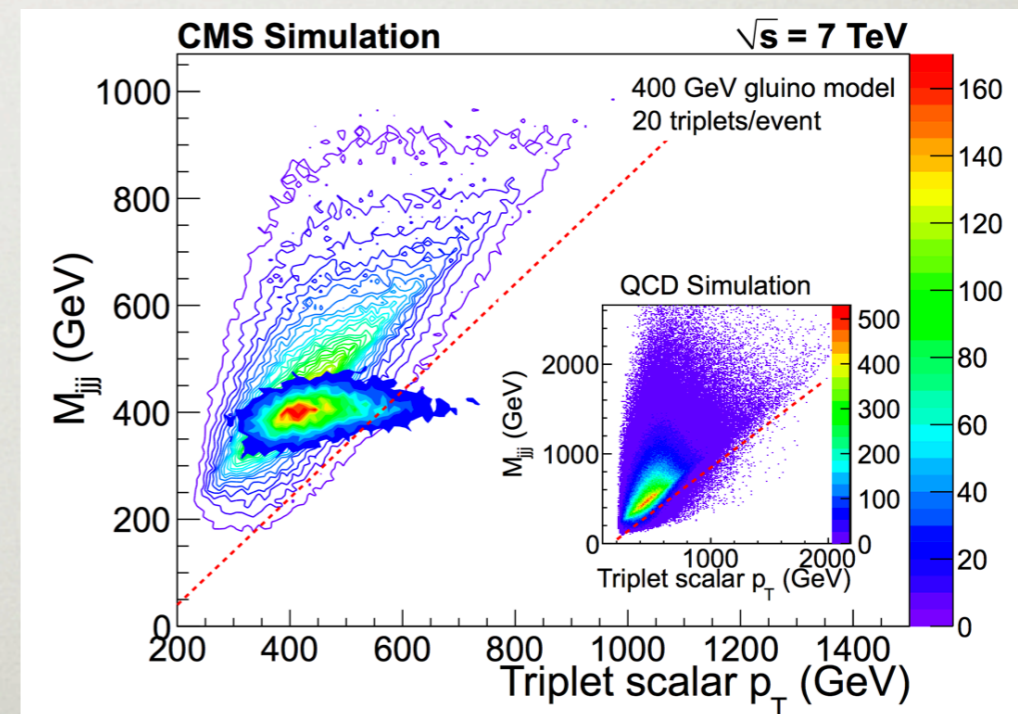
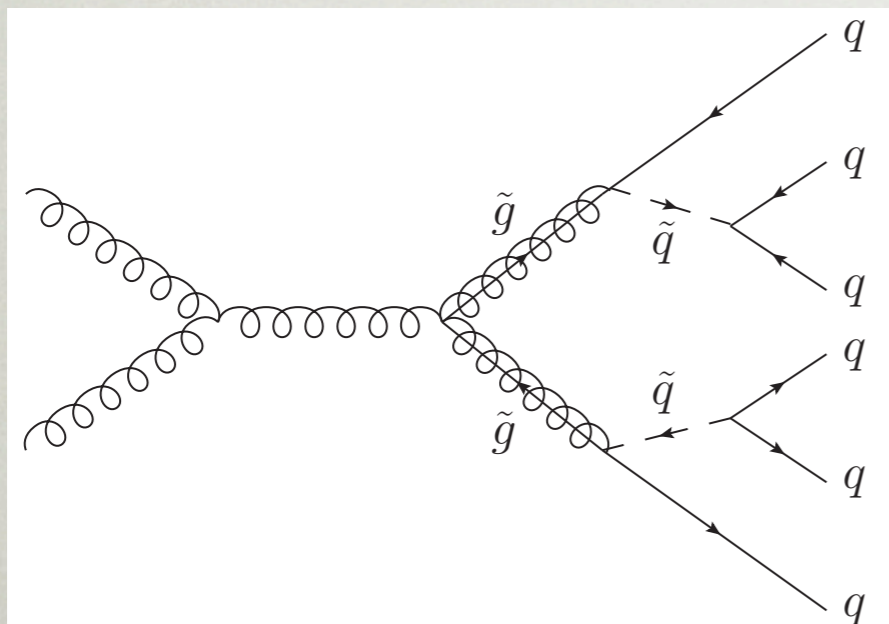


- Recast of ATLAS techni-rho W+dijet search can beat Tevatron by a factor of a few in cross section
- Can we do better with an optimized search?
- What about  $\zeta$ /some similar observable?
- Work in progress



# Multijet resonances

- Jet substructure can also be useful for three-jet resonances, but come at a cost of producing them well above threshold (ex. RPV gluinos in Curtin, Essig, BS arXiv:1210.5523)
- There are already good resolved 3-jet resonance searches (ex. Rutgers gp., CMS analysis arXiv:1311.1799)
  - Already in somewhat boosted regime





# Conclusions

- Jet-substructure-inspired observables can improve identification of dijet resonances, even in the moderate boost regime / resolved limit
  - Interpolate between different kinematic regimes

$$\zeta \equiv \frac{m_1}{m_{12}} \Delta R_{12} \quad (\text{and variations})$$

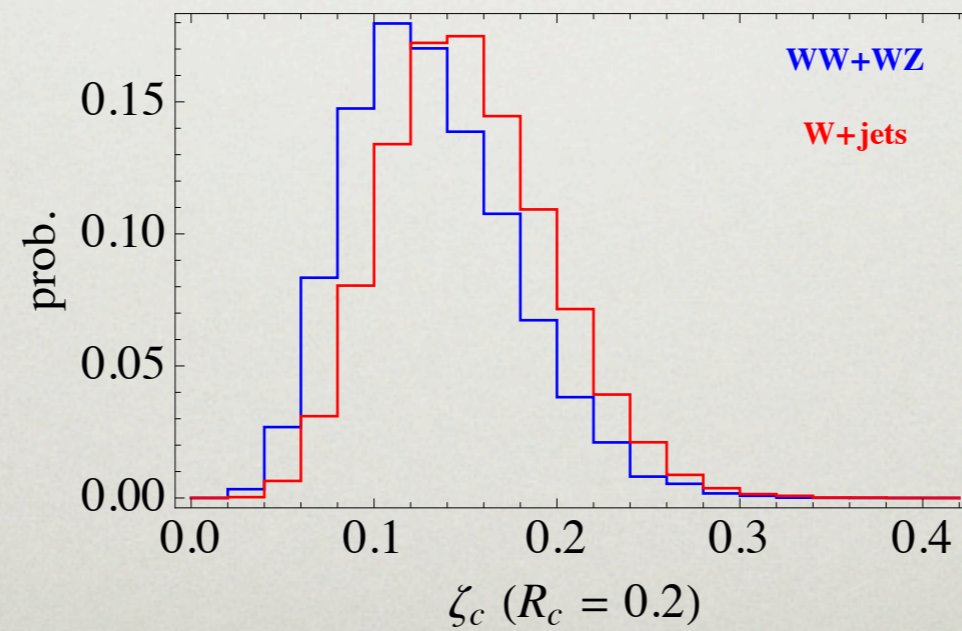
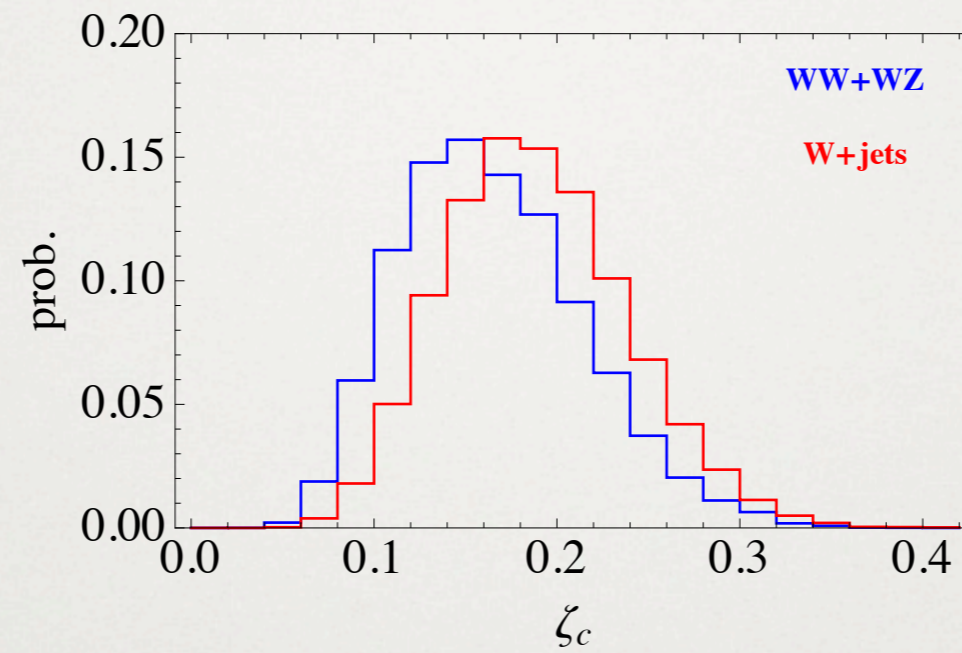
- Works well for two important examples of SM hadronic resonances
  - $WW+WZ$
  - $V + (H \rightarrow bb)$
- Also useful in beyond-SM physics searches
  - $Z' \rightarrow WW$
  - $Z' \rightarrow jj$
- Uses standard-radius jets, no optimization for different  $R$
- Let's find out what LHC13 has in store!



# **Back-up slides**

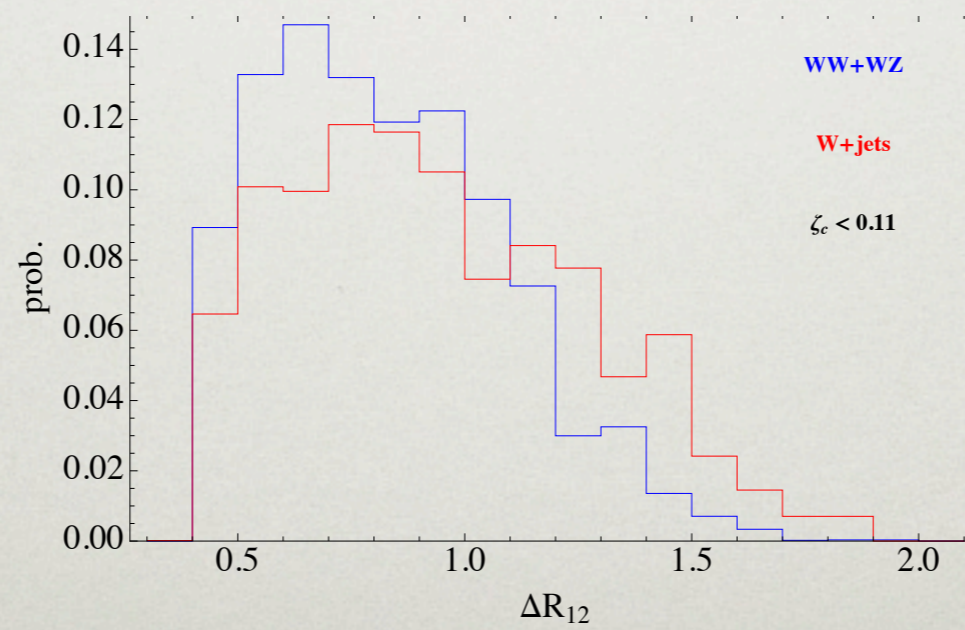
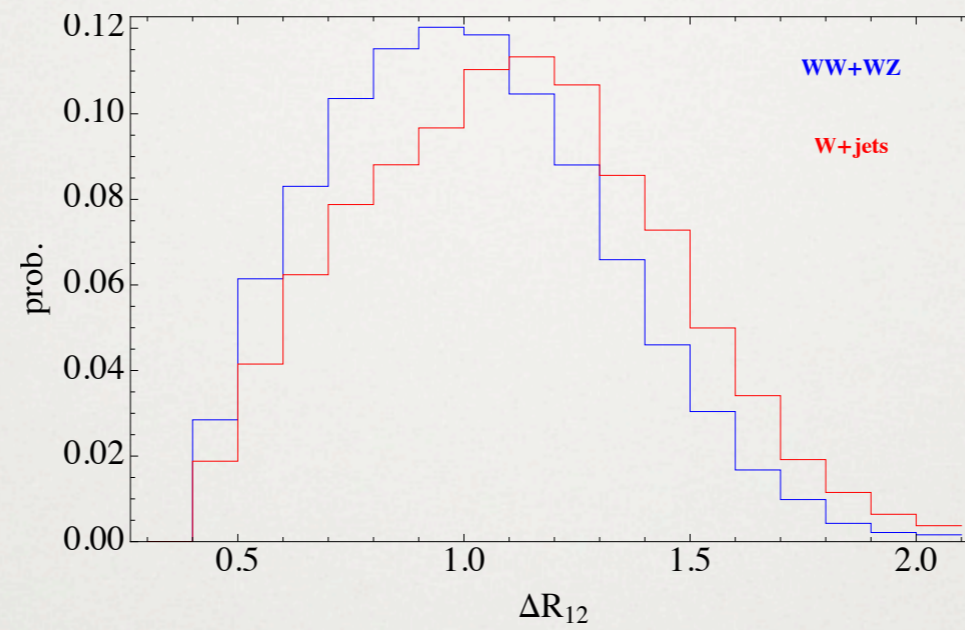


# zeta distribution



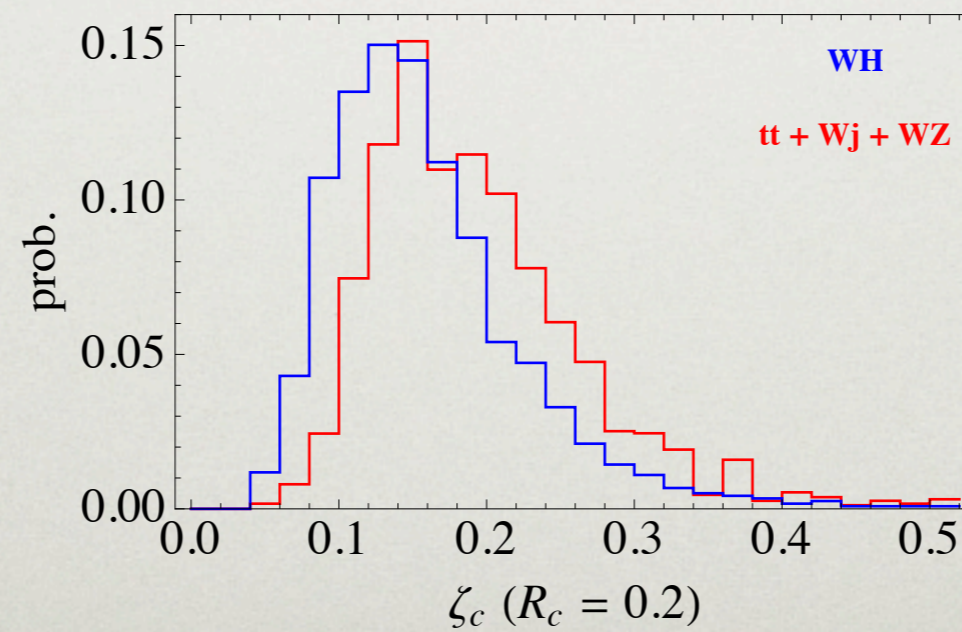
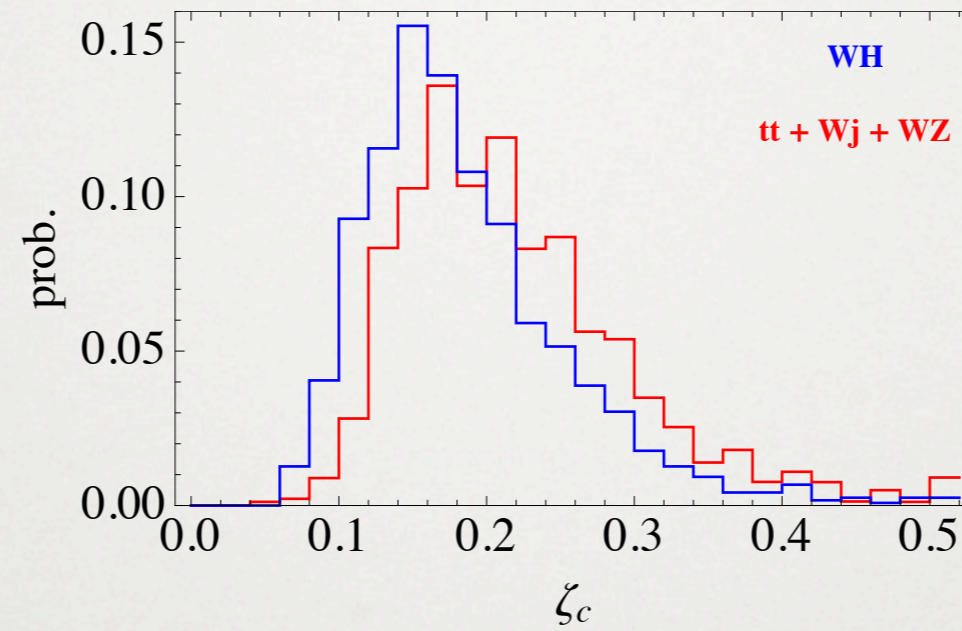


# $\Delta R$ distribution



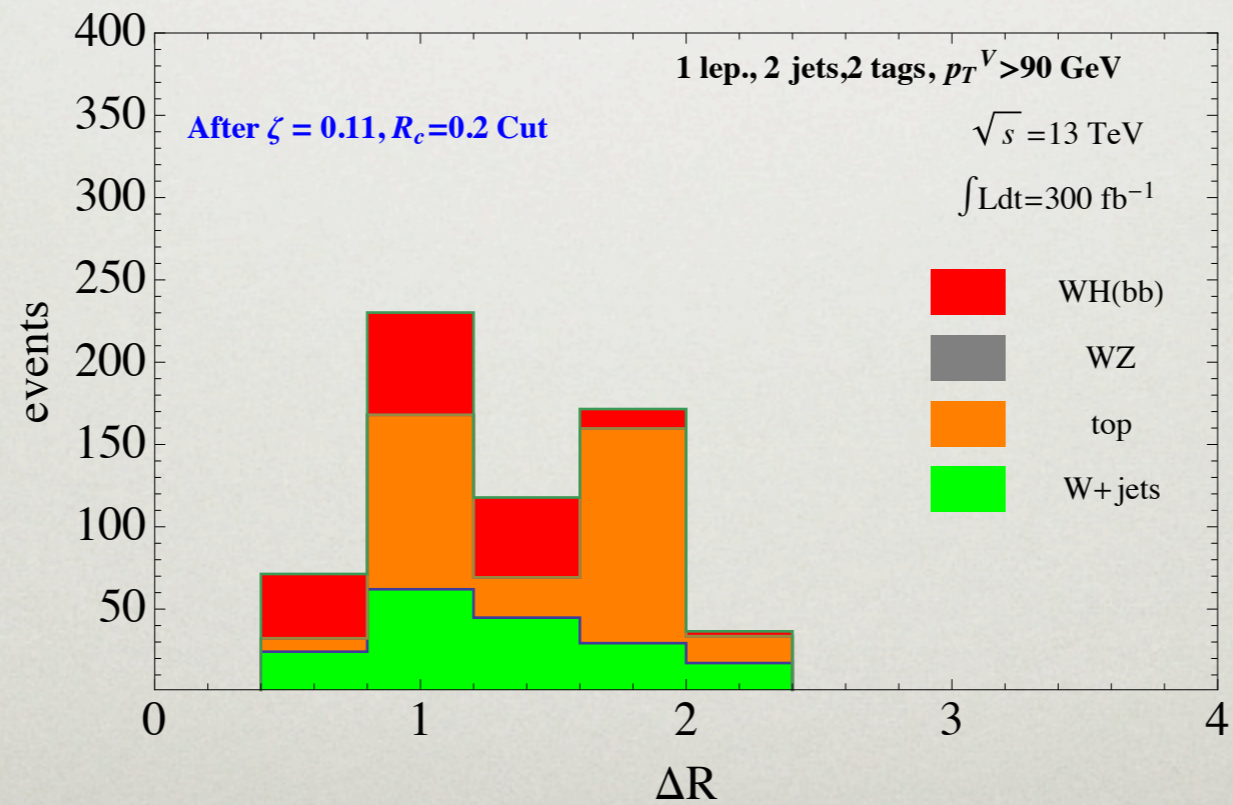
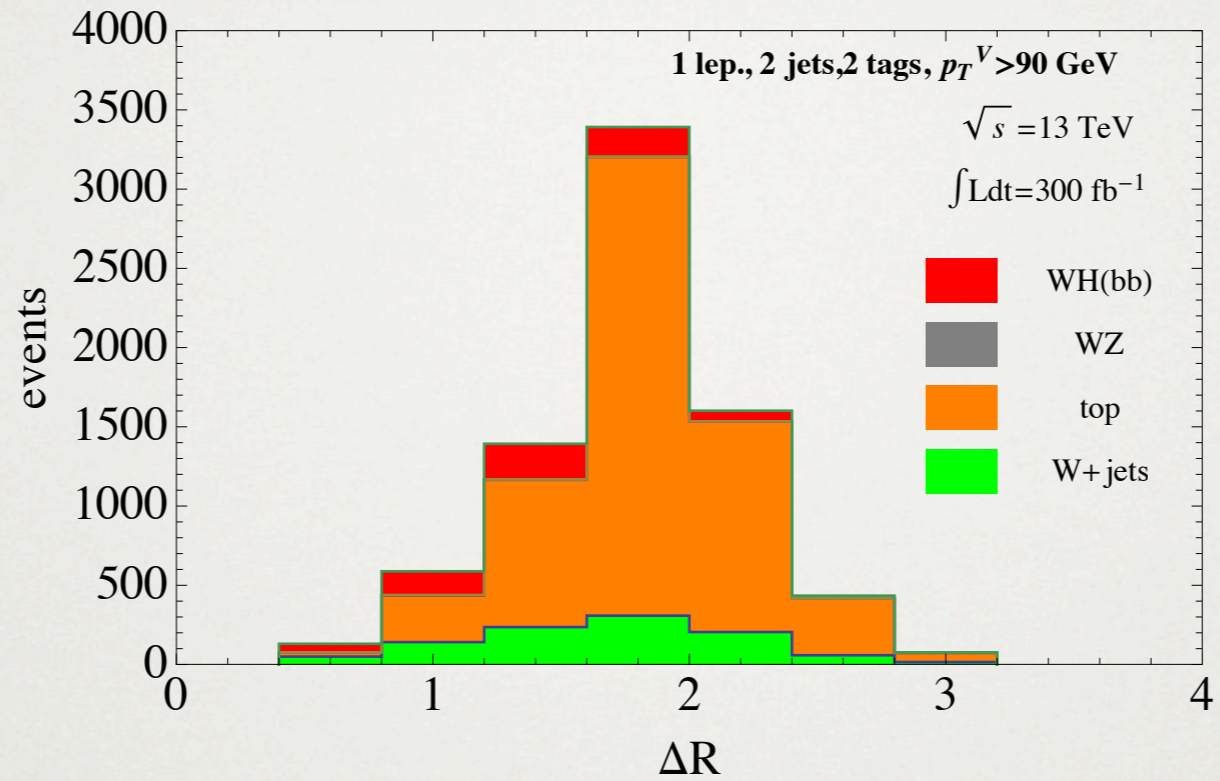


# zeta distribution



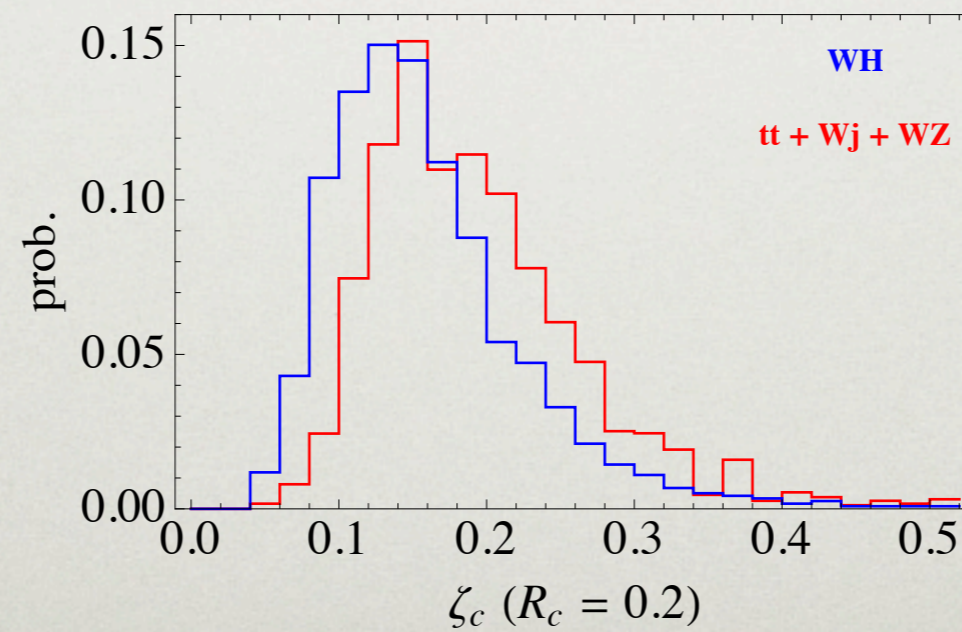
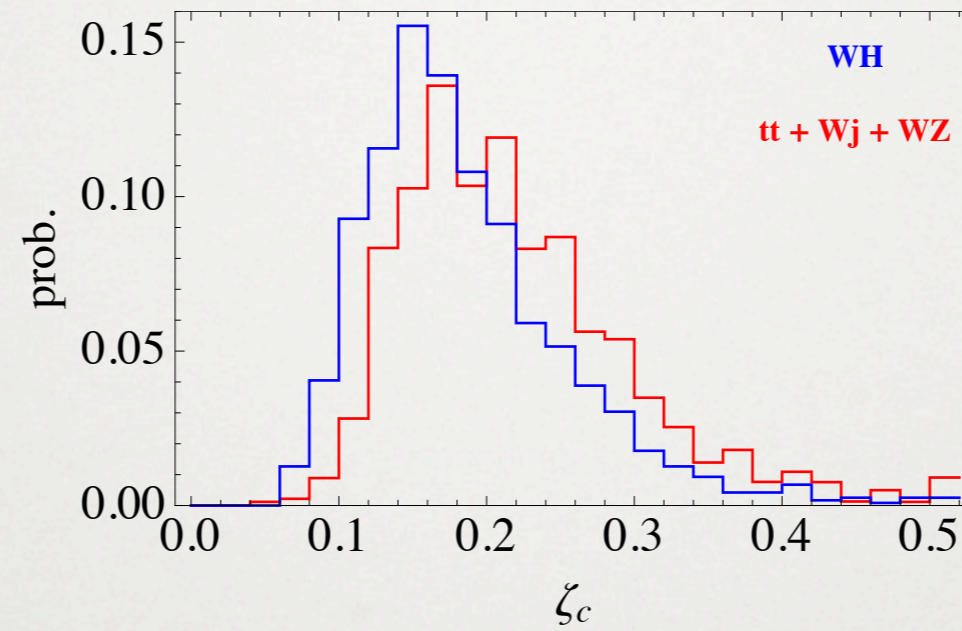


# $\Delta R$ distribution



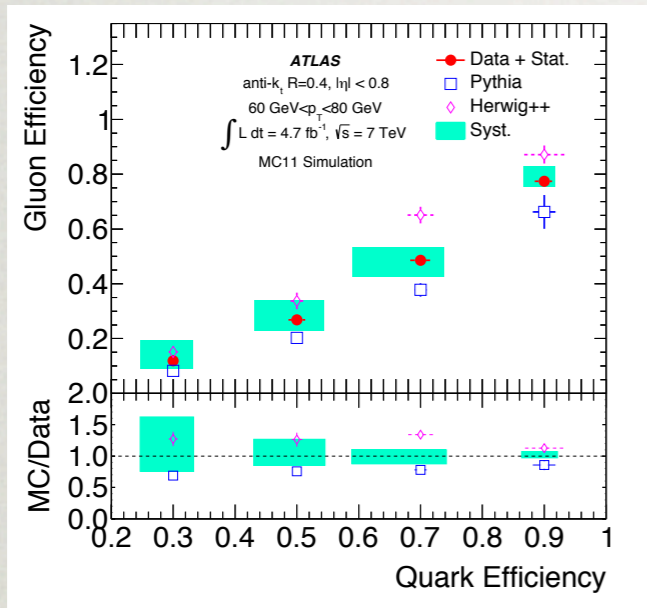


# zeta distribution





# q/g tagging



(a)

Taken from ATLAS q/g tagging study  
arXiv:1405.6583

