

# Recent Results in the Conformal Bootstrap

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# Outline

- 1 Bootstrap Review
- 2 Mixed Correlator Bootstrap
- 3 Spinning Bootstrap
- 4 Future

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# Conformal Bootstrap

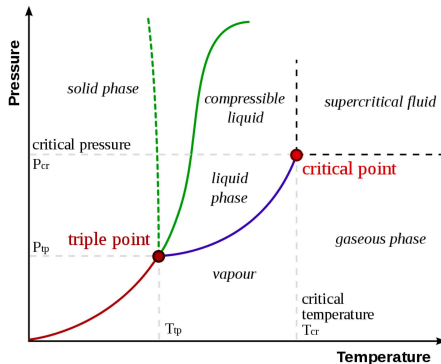
- ▶ The conformal bootstrap aims to use basic consistency conditions to map out and solve the space of CFTs
  - ▶ Conformal symmetry
  - ▶ Associativity of the OPE (crossing symmetry)
  - ▶ Unitarity
- ▶ Beautiful success story in 2D  
[Ferrara, Gatto, Grillo '73; Polyakov '74; Belavin, Polyakov, Zamolodchikov '84]
- ▶ Great progress in  $D > 2$  starting in 2008  
[Rattazzi, Rychkov, Tonni, Vichi '08; ...]

# Motivations

Many motivations to learn about CFTs in  $D > 2$ :

- ▶  $3D$ : Condensed Matter and Statistical Systems at Phase Transitions
- ▶  $4D$ : Scenarios for Physics Beyond the Standard Model
- ▶ Structure of QFT and space of CFTs
- ▶ AdS/CFT Correspondence (precise way to study quantum gravity)

# Example: 3D Ising Model



- ▶ CFTs describe condensed matter systems at phase transitions
  - ▶ Liquid-vapor critical point in fluids ( $^3\text{He}$ ,  $\text{O}_2$ ,  $\text{Ar}$ ,  $\text{Xe}$ , ...)
  - ▶ Continuum limit of 3D spin lattice:  $H = \sum_{\langle ij \rangle} \sigma_i \cdot \sigma_j$  at  $T_c$   
 → Both described by 3D Ising CFT: scalar with  $\phi^4$  interaction

# Single Correlator Bootstrap

Simplest bootstrap involves evaluating scalar 4-point functions with OPE:

$$\begin{aligned}
 & \langle \overbrace{\phi(x_1)\phi(x_2)} \overbrace{\phi(x_3)\phi(x_4)} \rangle \\
 &= \sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\mathcal{O}}^2 C_I(x_{12}, \partial_2) C_J(x_{34}, \partial_4) \langle \mathcal{O}^I(x_2) \mathcal{O}^J(x_4) \rangle \\
 &\equiv \frac{1}{x_{12}^{2\Delta_\phi} x_{34}^{2\Delta_\phi}} \sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\mathcal{O}}^2 g_{\Delta, \ell}(u, v)
 \end{aligned}$$

- ▶  $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2}$ ,  $v = \frac{x_{14}^2 x_{23}^2}{x_{13}^2 x_{24}^2}$  conformally-invariant cross ratios
- ▶  $g_{\Delta, \ell}(u, v)$  conformal block, labeled by  $\Delta = \dim(\mathcal{O})$  and  $\ell = \text{spin}(\mathcal{O})$  (known in any  $D$  using explicit formulas or recursion relations)

# Crossing Symmetry

- ▶  $\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle$  is symmetric under permutations of  $x_i$
- ▶ Switching  $x_1 \leftrightarrow x_3$  after OPE gives the crossing relation:

$$\sum \text{Diagram 1} = \sum \text{Diagram 2}$$

The diagrammatic equation shows two sums separated by an equals sign. The first sum is over a diagram with four external legs labeled 1, 2, 3, and 4. Legs 1 and 2 meet at a vertex on the left, and legs 3 and 4 meet at a vertex on the right. A horizontal double line representing an operator  $\mathcal{O}$  connects the two vertices. The second sum is over a similar diagram, but the double line representing  $\mathcal{O}$  is vertical, connecting the top vertex (legs 1 and 4) and the bottom vertex (legs 2 and 3).

$$v^{\Delta_\phi} \sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\mathcal{O}}^2 g_{\Delta, \ell}(u, v) = u^{\Delta_\phi} \sum_{\mathcal{O} \in \phi \times \phi} \lambda_{\mathcal{O}}^2 g_{\Delta, \ell}(v, u)$$

- ▶ This is a *constraint* on the spectrum of primary  $\Delta$ 's,  $\ell$ 's, and  $\lambda_{\mathcal{O}}$ 's



# Crossing Symmetry

Convenient to write as a sum rule (separating out  $\phi \times \phi \sim \mathbb{1} + \dots$ )

$$0 = \underbrace{F_{0,0}(u, v)}_{\text{unit op.}} + \underbrace{\sum \lambda_{\mathcal{O}}^2 F_{\Delta, \ell}(u, v)}_{\text{everything else}}$$

where

$$F_{\Delta, \ell}(u, v) \equiv v^{\Delta_\phi} g_{\Delta, \ell}(u, v) - u^{\Delta_\phi} g_{\Delta, \ell}(v, u).$$

# How Does Crossing Symmetry Lead to CFT Bounds?

Crossing relation for real scalar  $\phi$ :

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- ▶ Make an assumption: all scalars have dimension  $\Delta > \Delta_{\min}$
- ▶ Search for a linear functional  $\alpha$  such that

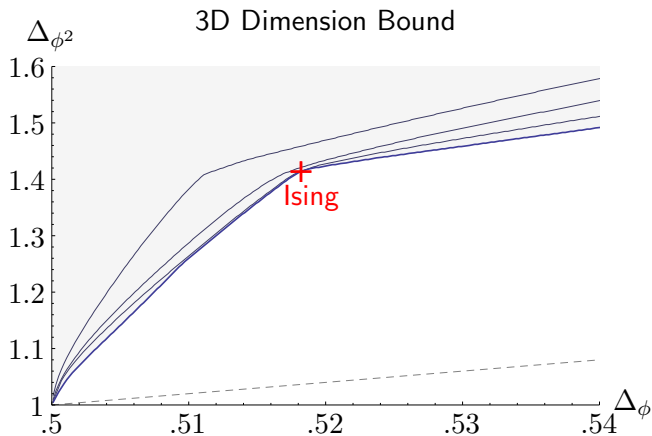
$$\begin{aligned} \alpha(F_{0,0}) &= 1, \quad \text{and} \\ \alpha(F_{\Delta, \ell}) &\geq 0, \quad \text{for all other } \mathcal{O} \in \phi \times \phi. \end{aligned}$$

- ▶ If you find one, the assumption is ruled out!

# CFT Bounds

- ▶ Can be solved with linear or semidefinite programming techniques  
[Rattazzi, Rychkov, Tonni, Vichi '08; DP, Simmons-Duffin, Vichi '11]
- ▶ Many nice results between 2008-2014 following this approach in (2-6)D, as well as generalizations to SUSY and other global symmetries
- ▶ Here I will focus on the 3D story...

# 3D Dimension Bounds

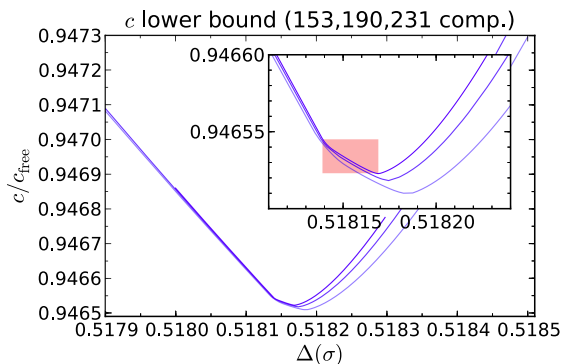


[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '12]

- ▶ 3D Ising dimensions from numerical simulations:

$$\Delta_{\sigma} \simeq 0.51813(5), \quad \Delta_{\epsilon} \simeq 1.41275(25) \quad [\text{Hasenbusch '10}]$$

# $c$ -minimization and Spectrum Extraction

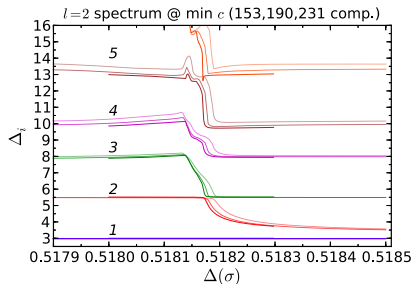
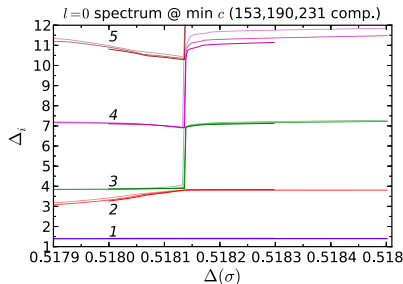


[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '14]

- Under the conjecture that the central charge  $\langle TT \rangle \propto c$  is minimized, a precise spectrum in  $\sigma \times \sigma \sim \mathbb{1} + \epsilon + \epsilon' + \dots$  can be extracted:

$$\Delta_\sigma \simeq 0.518154(15), \quad \Delta_\epsilon \simeq 1.41267(13), \quad \Delta_{\epsilon'} = 3.8303(18), \quad \dots$$

# $c$ -minimization and Spectrum Extraction

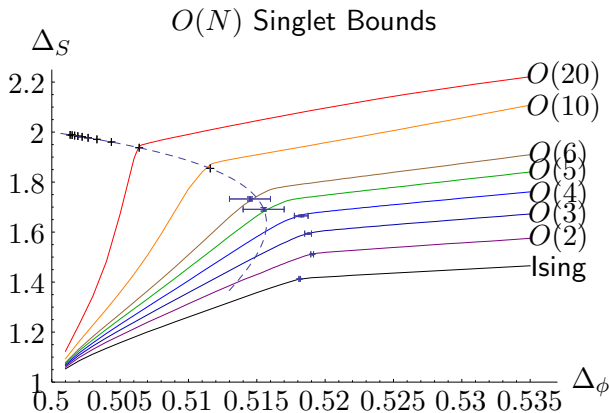


[El-Showk, Paulos, DP, Rychkov, Simmons-Duffin, Vichi, '14]

- ▶ Operators merge and disappear from the spectrum!
- ▶ Reminiscent of null states in 2D

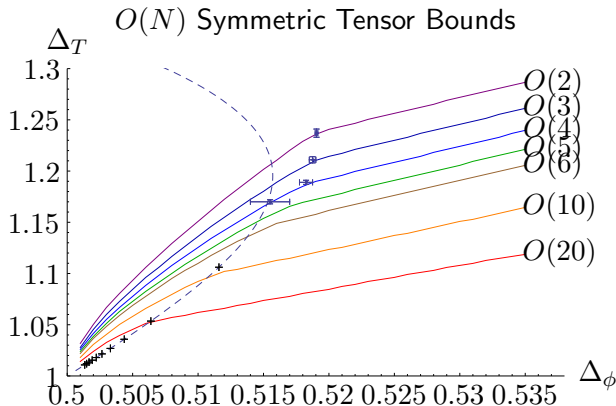


# 3D $O(N)$ Bounds



- ▶ Extension to  $\langle \phi_i \phi_j \phi_k \phi_l \rangle$ , where  $\phi_i$  is  $O(N)$  vector
- ▶ OPE  $\phi_i \times \phi_j \sim \mathbb{1} + S + T_{ij} + \dots$  contains singlets and two-index tensors

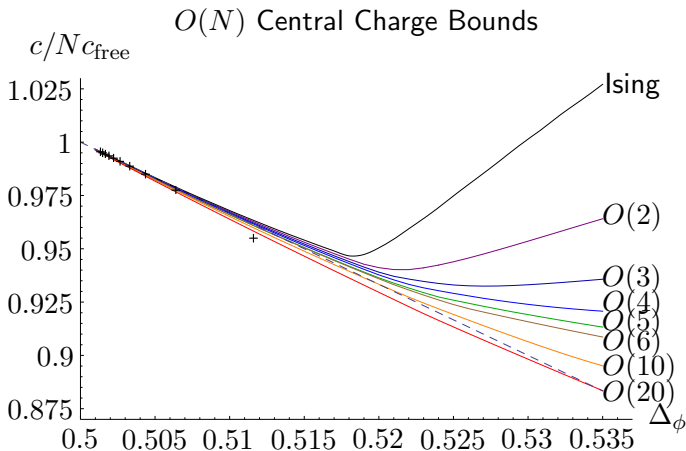
# 3D $O(N)$ Bounds



[Kos, DP, Simmons-Duffin '13]

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# 3D $O(N)$ Bounds



- In general  $c$  bounds do not show a minimum

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# Missing Operators?

Studies so far failed to access parts of the operator spectrum:

- ▶ In 3D Ising, only saw  $\mathbb{Z}_2$ -even operators in  $\sigma \times \sigma \sim \mathbb{1} + \epsilon + \dots$
- ▶ In  $O(N)$  models, only saw  $O(N)$  singlets  $S$  and two-index tensors  $T_{ij}$

In fact, there are good reasons to expect that the unaccessed operators are important...

# Non-perturbative Equations of Motion

- ▶ In  $\phi^4$  theory, one expects an equation of motion like  $\partial^2\phi \sim \phi^3 + \dots$
- ▶ This means that the  $\phi^3$  operator becomes a descendant, and is removed from the primary spectrum
- ▶ The consequence is that there is a large *gap* in the  $\mathbb{Z}_2$ -odd spectrum, along with many other relations between operators (e.g.,  $\phi\partial^2\phi \sim \phi^4$ )

It is very important to understand the role of these gaps (and operator relations) in the context of the conformal bootstrap!

# Mixed Correlators

- ▶ To probe gaps, one must consider mixed correlators like  $\langle \sigma \sigma \epsilon \epsilon \rangle$
- ▶ However, the expansion

$$\langle \overline{\sigma \sigma} \overline{\epsilon \epsilon} \rangle \sim \sum_{\mathcal{O}} \lambda_{\sigma \sigma \mathcal{O}} \lambda_{\epsilon \epsilon \mathcal{O}} g_{\Delta, \ell}(u, v)$$

does not have positive coefficients, so we cannot use the same logic

- ▶ In fact, it turns out that the positivity constraints must be phrased in terms of positive semidefinite matrices (SDP is mandatory)

# Mixed Correlators

- ▶ The positivity properties can be made manifest by considering the system  $\{\langle\sigma\sigma\sigma\sigma\rangle, \langle\sigma\sigma\epsilon\epsilon\rangle, \langle\epsilon\epsilon\epsilon\epsilon\rangle\}$ , which leads to 5 sum rules:

$$\sum_{\mathcal{O}^+} (\lambda_{\sigma\sigma\mathcal{O}} \quad \lambda_{\epsilon\epsilon\mathcal{O}}) \vec{V}_{+,\Delta,\ell}(u, v) \begin{pmatrix} \lambda_{\sigma\sigma\mathcal{O}} \\ \lambda_{\epsilon\epsilon\mathcal{O}} \end{pmatrix} + \sum_{\mathcal{O}^-} \lambda_{\sigma\epsilon\mathcal{O}}^2 \vec{V}_{-,\Delta,\ell}(u, v) = 0,$$

where  $\vec{V}_{\pm,\Delta,\ell}(u, v)$  are 5-vectors and  $\vec{V}_{+,\Delta,\ell}(u, v)$  is a  $2 \times 2$  matrix



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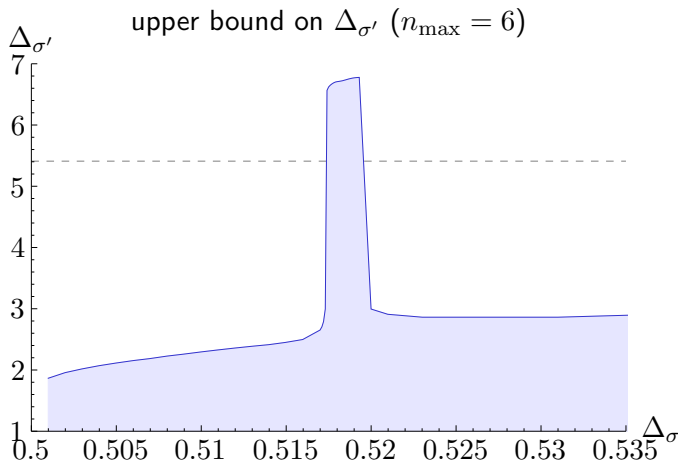
- ▶ Bounds follow from applying a 5-vector of functionals  $\vec{\alpha}$  such that

$$\begin{pmatrix} 1 & 1 \end{pmatrix} \vec{\alpha} \cdot \vec{V}_{+,0,0} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = 1,$$

$$\vec{\alpha} \cdot \vec{V}_{+,\Delta,\ell} \succeq 0, \quad \text{for all } \mathbb{Z}_2\text{-even operators } \mathcal{O}^+,$$

$$\vec{\alpha} \cdot \vec{V}_{-,\Delta,\ell} \geq 0, \quad \text{for all } \mathbb{Z}_2\text{-odd operators } \mathcal{O}^-.$$

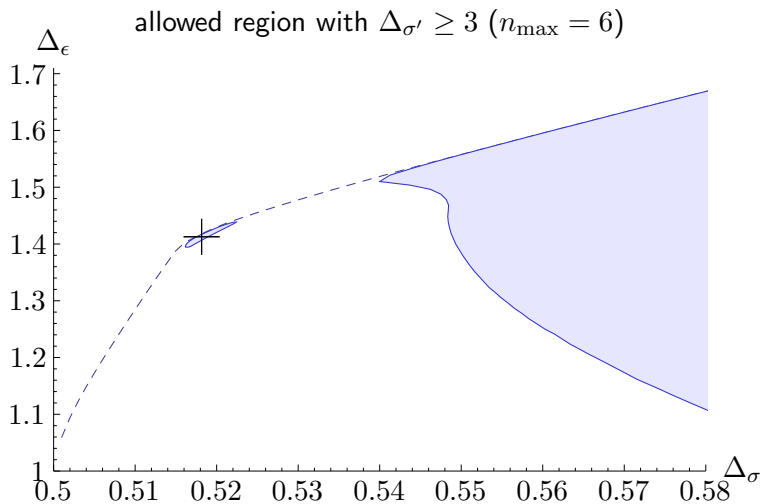
# Mixed Correlator Bounds



[Kos, DP, Simmons-Duffin '14]

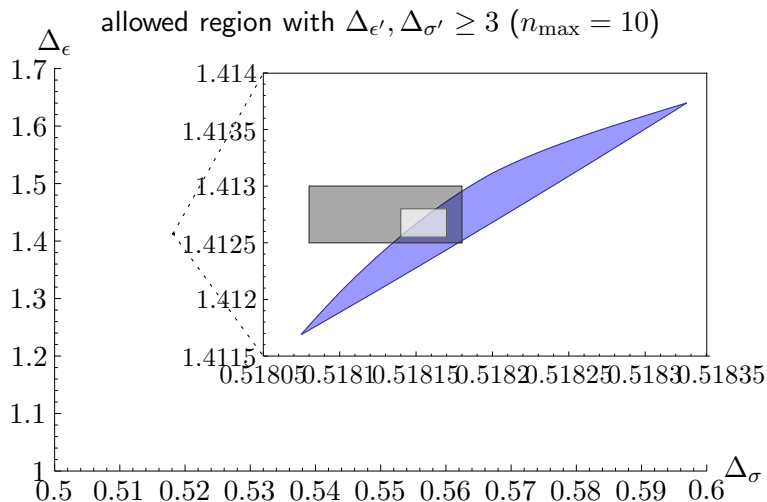
- ▶ Imposing a gap in  $\mathbb{Z}_2$ -odd spectrum restricts  $\Delta_{\sigma}$  to a small interval!

# Mixed Correlator Bounds



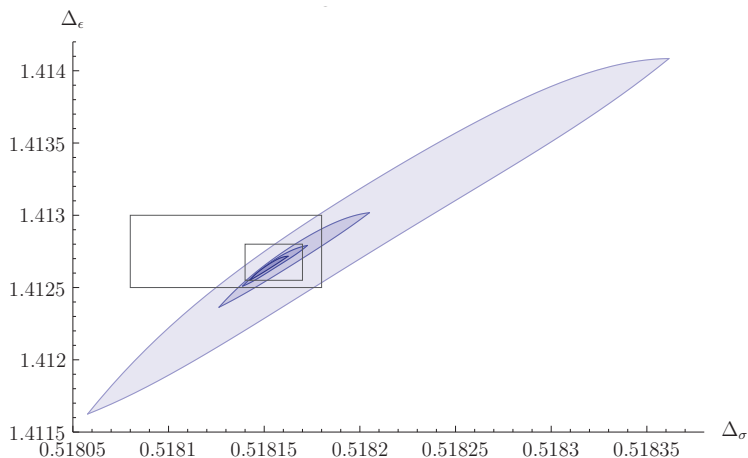
[Kos, DP, Simmons-Duffin '14]

# Mixed Correlator Bounds



[Kos, DP, Simmons-Duffin '14]

# Mixed Correlator Bounds



Pushing farther, the region keeps shrinking! [\[Simmons-Duffin '15\]](#)

$$\{\Delta_\sigma, \Delta_\epsilon\} = \{0.518151(6), 1.41264(6)\}$$

# Mixed Correlator Lessons

- ▶ The 3D Ising CFT is *isolated* in the space of 3D CFTs with  $\mathbb{Z}_2$  symmetry and 2 relevant operators
- ▶ It is a plausible conjecture that it is the only CFT with this property
- ▶ The conformal bootstrap can place the idea of *critical universality* on a rigorous footing

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- ▶ The conformal bootstrap can place the idea of *critical universality* on a rigorous footing
  
- ▶ Extension to  $O(N)$  symmetry in progress [Kos, DP, Simmons-Duffin, Vichi]
  - ▶ Challenge is many relevant operators:  
e.g.  $\phi_i, \phi_i\phi_j, \dots, \phi_i\phi_j\phi_k\phi_l\phi_m$  all relevant at large  $N$
  - ▶ However, preliminary results show isolated regions from system:  

$$\{\langle\phi_i\phi_j\phi_k\phi_l\rangle, \langle\phi_i\phi_j\phi^2\phi^2\rangle, \langle\phi^2\phi^2\phi^2\phi^2\rangle\}$$

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# Spinning Bootstrap

- ▶ Another important direction is to extend the conformal bootstrap to external operators with spin
- ▶ E.g., one would like to include fermions, global symmetry currents, the stress tensor, higher spin operators, ...
- ▶ This brings two complications:
  - ▶ Multiple tensor structures in the 3- and 4-point functions
  - ▶ Need to calculate the conformal blocks
- ▶ Current work in progress: fermion bootstrap in 3D  
[Iliesiu, Kos, DP, Pufu, Simmons-Duffin, Yacoby, in progress]

# Fermion Bootstrap

- ▶ Consider 4-point functions  $\langle \psi\psi\psi\psi \rangle$  of a Majorana fermion in 3D ( $SO(2, 1) \simeq Sp(2, \mathbb{R}) \rightarrow$  real two-component spinors)
- ▶ For now, we will also assume a parity symmetry:  $(x, y) \rightarrow (-x, y)$
- ▶ To classify 3-point and 4-point structures, we can work in an embedding space, where  $SO(3, 2) \simeq Sp(4, \mathbb{R})$  is linearly realized

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## Results:

- ▶  $\langle \psi\psi\mathcal{O}^{(\ell \text{ even})} \rangle$  has two structures of even parity and one of odd parity
- ▶  $\langle \psi\psi\mathcal{O}^{(\ell \text{ odd})} \rangle$  has one structure of odd parity
- ▶  $\langle \psi\psi\psi\psi \rangle$  has 5 independent tensor structures

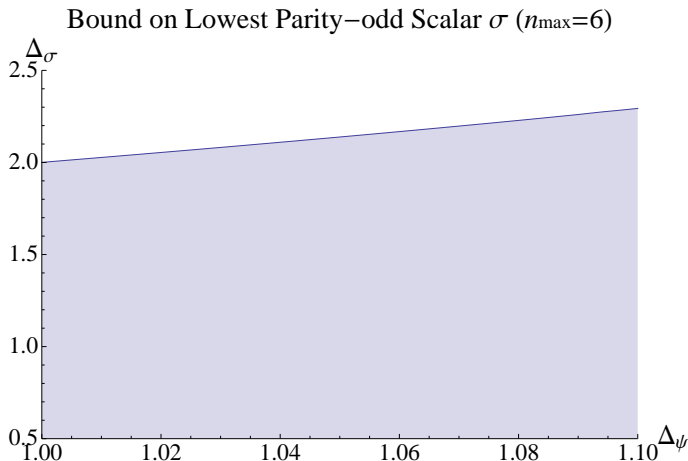
# Fermion Bootstrap

- ▶ Crossing symmetry leads to a 5-vector of sum rules:

$$\begin{aligned}
 0 = & \sum_{\mathcal{O}_+, \ell_+} \left( \lambda_{\mathcal{O}_+}^1 \quad \lambda_{\mathcal{O}_+}^2 \right) \vec{F}_{++,\Delta,\ell}(u,v) \begin{pmatrix} \lambda_{\mathcal{O}_+}^1 \\ \lambda_{\mathcal{O}_+}^2 \end{pmatrix} \\
 & + \sum_{\mathcal{O}_-, \ell_+} (\lambda_{\mathcal{O}_-}^3)^2 \vec{F}_{-+,\Delta,\ell}(u,v) + \sum_{\mathcal{O}_-, \ell_-} (\lambda_{\mathcal{O}_-}^4)^2 \vec{F}_{--,\Delta,\ell}(u,v),
 \end{aligned}$$

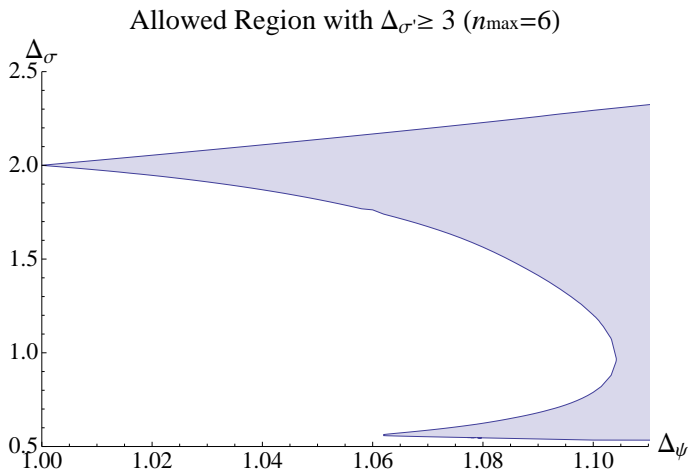
- ▶ To calculate the conformal blocks, we express  $\langle \psi\psi\mathcal{O} \rangle_a = D_a \langle \phi\phi\mathcal{O} \rangle$ , which lets us relate  $\int \langle \psi\psi\mathcal{O} \rangle_a \langle \tilde{\mathcal{O}}\psi\psi \rangle_b$  to  $\int \langle \phi\phi\mathcal{O} \rangle \langle \tilde{\mathcal{O}}\phi\phi \rangle$
- ▶ Bounds follow from applying functionals  $\alpha_I$  (again SDP is mandatory!)

# Preliminary Fermion Bounds



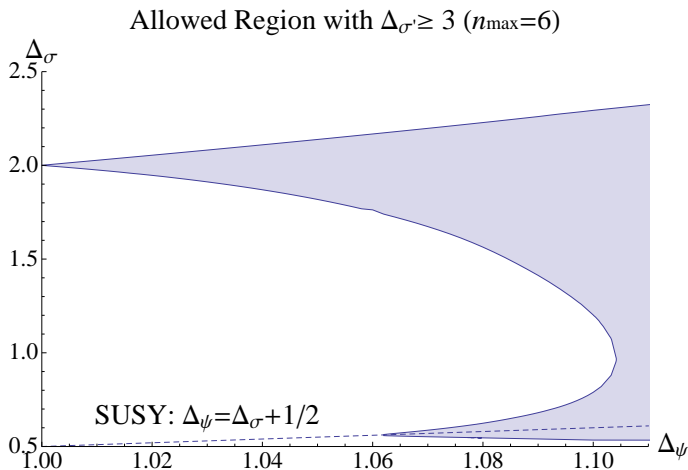
- ▶ Bound converges to free fermion values:  $\psi \times \psi \sim \mathbb{1} + \psi^2 + \dots$

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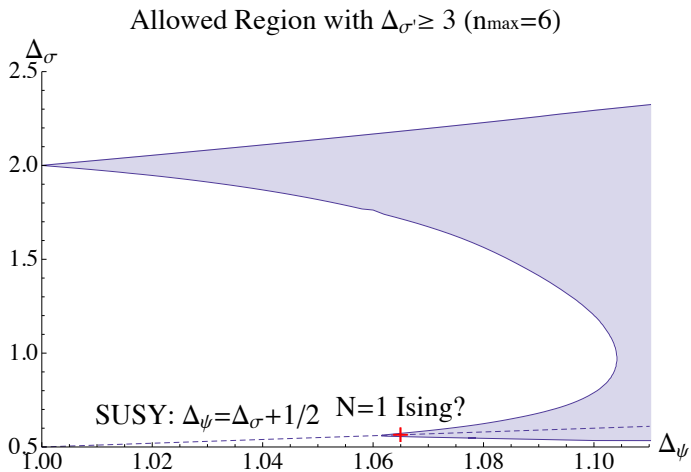
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# Preliminary Fermion Bounds



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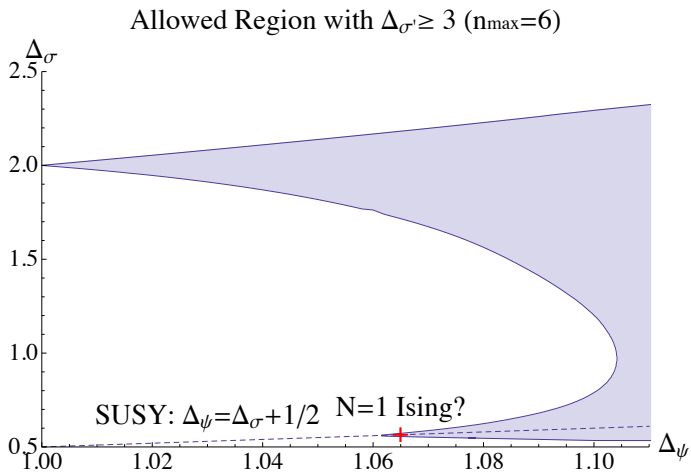
# Preliminary Fermion Bounds



- ▶ Conjecture: Lower tip will converge to the 3D  $\mathcal{N} = 1$  Ising model (perhaps with a slightly smaller  $\sigma'$  gap)
- ▶ SCFT with superpotential  $W = \Sigma^3$  and  $\Sigma = \sigma + \theta\psi + \theta^2\epsilon$



# Preliminary Fermion Bounds



- Here it is at  $\{\Delta_{\sigma}, \Delta_{\psi}\} \sim \{.562, 1.062\}$ , not far from previous estimates ( $\Delta_{\sigma} \sim .571$  [Grover, Sheng, Vishwanath '13],  $\Delta_{\sigma} \gtrsim .565$  [Bashkirov '13])

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- ▶ Extend spinning bootstrap to currents, stress tensor, higher spins, ...
- ▶ Larger systems of mixed correlators (all relevant operators!)
- ▶ Develop efficient algorithms for high-precision SDP
- ▶ Improve analytic arguments: large  $N$ , large  $\ell$ , SUSY chiral algebras, ...

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- ▶ Classify and map out space of CFTs in all dimensions!

# Targets for the Bootstrap

- ▶ 2D CFTs with  $c > 1$
- ▶ 3D Gross-Neveu Models (landscape of theories with fermions+scalars)
- ▶ 3D QED or QCD + matter (monopole ops?)
- ▶ 3D Chern-Simons + vector matter (connect to higher spin theory!)
- ▶ 4D QCD/SQCD in conformal window (archipelago for each  $N_f$ ?)
- ▶ Classify space of 5D and 6D CFTs
- ▶ Existence of CFTs in  $D > 6$ ?
- ▶ Conformal manifolds