

Lecture 1

→ XIX century: deterministic diff'l eq's.

XX century: QM, chaos, fluctuating phenomena

("curve + noise")

Brownian motion

Robert Brown (1827) \Rightarrow particles suspended in water undergo random motion

Einstein's explanation: introduce $\tau \ll t$
 & independently, (1905)
 Smoluchowski

Subseq. collisions are indep.; part's also move independent

In a time τ , x changes by Δ :

$\Psi(\Delta)$ is the prob. distr'n: $\int_{-\infty}^{\infty} \Psi(\Delta) d\Delta = 1$,
 total # part. suspended in liquid

$$dn = n \Psi(\Delta) d\Delta$$

part. whose coordinates change in the $(\Delta, \Delta+d\Delta)$ interval

Assume: 1. $\Psi(\Delta) \neq 0$ only for "small" Δ ;

2. $\Psi(\Delta) = \Psi(-\Delta)$ [symm.]

Consider $f(x,t)$ - # part. per unit volume @ x , @ t .

Then $f(x,t+\tau) dx = dx \int_{-\infty}^{\infty} d\Delta \Psi(\Delta) f(x+\Delta, t)$

τ small: $f(x,t) + \tau \frac{\partial f}{\partial t} + \dots$

Δ small: $f(x+\Delta, t) \approx f(x, t) + \Delta \frac{\partial f}{\partial x} + \frac{\Delta^2}{2} \frac{\partial^2 f}{\partial x^2} + \dots$

Then $f + \tau \frac{\partial f}{\partial t} = f \int_{-\infty}^{\infty} g(\Delta) d\Delta + \frac{\partial f}{\partial x} \int_{-\infty}^{\infty} d\Delta / \Delta g(\Delta) +$

 $+ \frac{\partial^2 f}{\partial x^2} \int_{-\infty}^{\infty} d\Delta \frac{\Delta^2}{2} g(\Delta)$

$\underbrace{\qquad\qquad\qquad}_{\sim D^2}$

$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$

diff'n eq'n

$$f(x, t) = \frac{n}{\sqrt{4\pi D}} \frac{e^{-x^2/4Dt}}{\sqrt{t}}$$

$$\sigma_x = \sqrt{2Dt}$$

Tangential's eq'n

& diam. a

Consider part. of mass $m \checkmark$ suspended in liquid. Two forces:

1. Viscous drag : from hydrodynamics

$$-6\pi\eta a v$$

↑
viscosity

2. Fluctuating force X from random collisions

stochastic diff'l eq'n

So, $m \frac{d^2 x}{dt^2} = -6\pi\eta a \frac{dx}{dt} + X$ $| \ast x$

$$\underbrace{\frac{m}{2} \frac{d^2}{dt^2} (x^2) - m \left(\frac{dx}{dt} \right)^2}_{\text{stochastic diff'l eq'n}} = -3\pi\eta a \frac{d(x^2)}{dt} + X x$$

$$\frac{m}{2} \frac{d}{dt} \left(2x \frac{dx}{dt} \right) - m \left(\frac{dx}{dt} \right)^2 =$$

$$= m x \frac{d^2 x}{dt^2}$$

$\langle \dots \rangle$ - average over all particles:

use $\langle \frac{mv^2}{2} \rangle = \frac{k_B T}{2}$ equip. theorem

$$\frac{m}{2} \frac{d^2}{dt^2} \langle x^2 \rangle + 3\pi \eta a \frac{d\langle x^2 \rangle}{dt} = k_B T + \underbrace{\langle Xx \rangle}_{=0}$$

Then $\frac{d\langle x^2 \rangle}{dt} = \frac{k_B T}{3\pi \eta a} + C \stackrel{\text{const}}{\underset{\text{const}}{\ell}} \frac{6\pi \eta a}{m} \langle x \rangle \langle x \rangle = 0$

$$\frac{6\pi \eta a}{m} \gg 1 \Rightarrow \text{neglect } 2^{\text{nd}} \text{ term}$$

Then $\langle x^2 \rangle - \langle x_0^2 \rangle = \frac{k_B T}{3\pi \eta a} t$

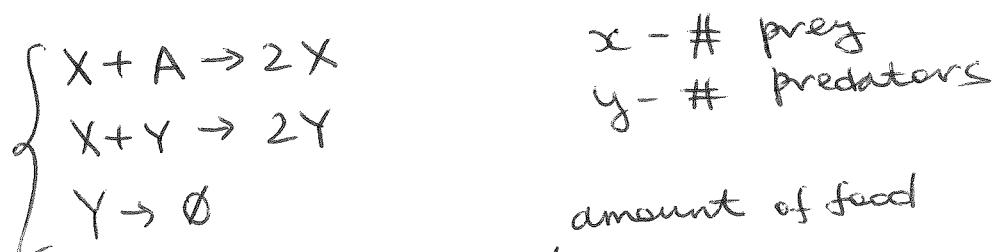
"2D, same as Einstein's result
(see 6x above)

Birth-death processes

Consider X : ~~possible~~ prey

Y : predator

A : food for prey



Then $\begin{cases} \frac{dx}{dt} = k_1 \overset{\leftarrow}{dx} - k_2 xy, \\ \frac{dy}{dt} = k_2 xy - k_3 y \end{cases} \quad \{k_i\} \text{ are rates}$

Lotka
Volterra
eq'n's



oscillatory
sol's

But: real data
has noise...

How to include fluct's?

Birth-death master eq'n (ME):

Consider $P(x, y, t)$ ← prob. distr'n

Consider small Δt :

$$\left\{ \begin{array}{l} P(x \rightarrow x+1; y \rightarrow y) = k_1 a x \Delta t, \\ P(x \rightarrow x-1; y \rightarrow y+1) = k_2 x y \Delta t, \\ P(x \rightarrow x; y \rightarrow y-1) = k_3 y \Delta t, \\ P(x \rightarrow x; y \rightarrow y) = 1 - (k_1 a x + k_2 x y + k_3 y) \Delta t \end{array} \right.$$

no memory
(Markov assumption):
probs can
be constructed
simply from
(x, y)

Then
$$\frac{P(x, y, t + \Delta t) - P(x, y, t)}{\Delta t} = \frac{\partial P}{\partial t}$$

$$= k_1 a (x-1) P(x-1, y, t) + k_2 (x+1)(y-1) \times \\ \times P(x+1, y-1, t) + k_3 (y+1) P(x, y+1, t) - \\ - (k_1 a x + k_2 x y + k_3 y) P(x, y, t)$$

↗ Is Markov postulate justified?
(not clear...)

ME ; determines both fluct's and average behavior