

HW #5 solutions (2026)

1. (a) Uniform $m \Rightarrow \left(\frac{dm}{dx}\right)^2 = 0$, we just need to minimize $V(m) = \frac{a}{2}m^2 + \frac{b}{4}m^4$.

For $T < T_c$, $\frac{dV}{dm} = 0$ yields

$$m=0, \quad m^2 = -\frac{a}{b} = \frac{|a|}{b}.$$

The minima are given by $m = \pm m_0$,

where $m_0 = \sqrt{\frac{|a|}{b}}$.

Indeed, $V(m_0) = V(-m_0) = \frac{a}{2} \left(-\frac{a}{b}\right) + \frac{b}{4} \frac{a^2}{b^2} = -\frac{a^2}{4b} < V(0) = 0$.

(b) Use Euler-Lagrange equation:

$$\frac{d}{dx} \frac{\partial \mathcal{I}}{\partial (m')} = \frac{\partial \mathcal{I}}{\partial m}, \quad \text{where the "Lagrangian" is the free energy density:}$$

$$F[m] = \int_{-\infty}^{\infty} dx \mathcal{I}(m, m', x).$$

We obtain: $c \frac{d^2 m}{dx^2} = a m + b m^3$ (*)

"m" for brevity

Finally, $\frac{1}{m_0} \tanh^{-1} \left(\frac{m}{m_0} \right) = \sqrt{\frac{b}{2c}} x$, or

$$m(x) = m_0 \tanh \left(\sqrt{\frac{b}{2c}} m_0 x \right).$$

$$\underbrace{\sqrt{\frac{b}{2c}} \sqrt{\frac{|a|}{b}}}_{m_0 \text{ from (a)}} = \frac{1}{\sqrt{2}} \sqrt{\frac{|a|}{c}} \Rightarrow \underline{\underline{\xi = \sqrt{\frac{c}{|a|}}}}$$

$$\text{So, } \begin{cases} m_0 = \sqrt{\frac{|a|}{b}}, \\ \xi = \sqrt{\frac{c}{|a|}} \end{cases} \quad \text{and} \quad m(x) = m_0 \tanh \left(\frac{x}{\sqrt{2} \xi} \right),$$

as desired.

$$\text{Since } a = \tilde{a}(T - T_c), \quad \xi \sim \frac{1}{\sqrt{T_c - T}}.$$

ξ is the correlation length;

$\xi \rightarrow \infty$ as $T \rightarrow T_c^-$, signifying that the domain walls become very broad.

② (a) Represent each philosopher by $\sigma_i = \{-1, +1\}$; $i=1, \dots, N$
 \uparrow spin \uparrow # spins

Circular table \rightarrow periodic BCs

Note that if $\sigma_i = +1$, $\sigma_{i-1} = \sigma_{i+1} = -1$ are not allowed.

(b) $N=2$: $\begin{matrix} - & - \\ - & + \\ + & - \end{matrix} \left. \vphantom{\begin{matrix} - & - \\ - & + \\ + & - \end{matrix}} \right\} M_N = 3$

$N=3$: $\begin{matrix} - & - & - \\ + & - & - \\ - & + & - \\ - & \bullet & + \end{matrix} \left. \vphantom{\begin{matrix} - & - & - \\ + & - & - \\ - & + & - \\ - & \bullet & + \end{matrix}} \right\} M_N = 4$

$N=4$: $\begin{matrix} - & - & - & - \\ 4 \text{ states with one } + \\ - & + & - & + \\ + & - & + & - \end{matrix} \left. \vphantom{\begin{matrix} - & - & - & - \\ - & + & - & + \\ + & - & + & - \end{matrix}} \right\} M_N = 7$

$N=5$: $\begin{matrix} - & - & - & - & - \\ 5 \text{ states with one } + \\ + & - & + & - & - \\ + & - & - & + & - \\ - & + & - & + & - \\ - & + & - & - & + \\ - & - & + & - & + \end{matrix} \left. \vphantom{\begin{matrix} + & - & + & - & - \\ + & - & - & + & - \\ - & + & - & + & - \\ - & + & - & - & + \\ - & - & + & - & + \end{matrix}} \right\} \begin{matrix} M_N = 11 \\ \text{two } +\text{'s} \end{matrix}$

(c) Note that

$$M_4 = M_3 + M_2 = 4 + 3,$$

" 7

$$M_5 = M_4 + M_3 = 7 + 4.$$

" 11

In general, $M_N = M_{N(-1)} + M_{N(+1)}.$

Now, $M_{N+1}(-1) = M_N + \underbrace{M_{N-1}(+1)}_{\substack{\text{just add} \\ -1 \text{ @ pos. } N+1 \\ \text{to any } N\text{-spin} \\ \text{configuration}}}$

$\underbrace{\dots \uparrow}_{N-1}$
 becomes
 $\underbrace{\dots \uparrow \downarrow \uparrow}_{N+1}$

$$M_{N+1}(+1) = \underbrace{M_{N-1}(-1)}_{\substack{\dots \downarrow \\ N-1 \\ \text{becomes} \\ \dots \downarrow \uparrow \downarrow \\ N+1}}$$

Therefore, $M_{N+1} = \underline{\underline{M_N + M_{N-1}}}. \quad (*)$

(d) The contribution of 2 neighboring spins to \mathcal{H} is

$$\Delta E(\sigma, \sigma') = -J\sigma\sigma' - \frac{h}{2}(\sigma + \sigma')$$

with $J = -1$, $h = -2$ we obtain:

σ/σ'	+1	-1
+1	3	-1
-1	-1	-1

← table of $\Delta E(\sigma, \sigma')$ values

all states that are allowed in the philosopher problem have the same ground-state energy $E_0(N) = -N$.
nnb pairs in 1D chain

The forbidden states have higher energies. Thus, it suffices to count the # ground states to find M_N .

(e) In the $\beta \rightarrow \infty$ limit,

$$Z_N \rightarrow \sum_{\text{ground states}} e^{-\beta E_0(N)} = e^{-\beta E_0(N)} M_N.$$

↑ excited states do not contribute

Thus, $M_N = \lim_{\beta \rightarrow \infty} e^{\beta E_0(N)} Z_N.$

Now, recall that

$$Z_N = \lambda_0^N + \lambda_1^N, \text{ where}$$

$$\lambda_{0,1} = e^{\beta J} \cosh(\beta H) \pm \sqrt{e^{2\beta J} \sinh^2(\beta H) + e^{-2\beta J}}$$

with $J=-1$, $h=-2$ we have:

$$\lambda_{0,1} = e^{-\beta} \cosh(2\beta) \pm \sqrt{e^{-2\beta} \sinh^2(2\beta) + e^{2\beta}}$$

In the $\beta \rightarrow \infty$ limit,

$$\lambda_{0,1} \rightarrow e^{-\beta} \frac{e^{2\beta}}{2} \pm \sqrt{e^{-2\beta} \frac{e^{4\beta}}{4} + e^{2\beta}} =$$

$$= e^{\beta} \frac{1}{2} \pm e^{\beta} \frac{\sqrt{5}}{2} = e^{\beta} \left(\frac{1 \pm \sqrt{5}}{2} \right)$$

Finally,

$$M_N \xrightarrow{\beta \rightarrow \infty} e^{-\beta N} \left[e^{\beta N} \left(\frac{1+\sqrt{5}}{2} \right)^N + e^{\beta N} \left(\frac{1-\sqrt{5}}{2} \right)^N \right] =$$

$$= \left(\frac{1+\sqrt{5}}{2} \right)^N + \left(\frac{1-\sqrt{5}}{2} \right)^N$$

"non-physical states"

Note that $M_0 = 2$, $M_1 = 1$,

$$M_2 = \frac{1+2\sqrt{5}+5}{4} + \frac{1-2\sqrt{5}+5}{4} = 3, \text{ etc.}$$

Finally, note that

$$M_N + M_{N-1} = \left(\frac{1+\sqrt{5}}{2}\right)^N + \left(\frac{1-\sqrt{5}}{2}\right)^N +$$

$$+ \left(\frac{1+\sqrt{5}}{2}\right)^{N-1} + \left(\frac{1-\sqrt{5}}{2}\right)^{N-1} =$$

$$= \left(\frac{1+\sqrt{5}}{2}\right)^{N-1} \underbrace{\left(\frac{3+\sqrt{5}}{2}\right)}_{\left(\frac{1+\sqrt{5}}{2}\right)^2} + \left(\frac{1-\sqrt{5}}{2}\right)^{N-1} \underbrace{\left(\frac{3-\sqrt{5}}{2}\right)}_{\left(\frac{1-\sqrt{5}}{2}\right)^2} =$$

$$= M_{N+1}, \quad \text{consistent with } (*).$$