

Due: Thursday, Oct. 7, 2010

**5.1** Consider a particle constrained to move on the surface described in cylindrical coordinates by  $z = \alpha r^3$ , subject to a constant gravitational force  $\vec{F} = -mg\hat{e}_z$ . Find the Lagrangian, two conserved quantities, and reduce the problem to a one dimensional problem. What is the condition for circular motion at constant  $r$ ?

**5.2** Suppose a particle of mass  $m$  moves under the influence of a power-law central force,  $\vec{F} = -cr^p\hat{e}_r$ , and is observed to have an orbit which is a circle of radius  $R$  passing through the point of attraction.

Find what values the power  $p$  could be, what is the angular momentum about the center of force, and what is the energy relative to  $U(\infty)$ .

How do  $\dot{\theta}$ ,  $\dot{y}$ , and  $\dot{x}$  behave as the particle approaches the origin, as a function of  $r$  as  $r \rightarrow 0$ ?. Is this consistent with  $x$  taking its minimum value at that point?

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Hamilton's Principle tells us that the motion of a particle is determined by the action functional being stationary under small variations of the path  $\Gamma$  in extended configuration space  $(t, \vec{x})$ . The unsymmetrical treatment of  $t$  and  $\vec{x}(t)$  is not suitable for relativity, but we may still associate an action with each path, which we can parameterize with  $\lambda$ , so  $\Gamma$  is the trajectory  $\lambda \rightarrow (t(\lambda), \vec{x}(\lambda))$ .

In the general relativistic treatment of a particle's motion in a gravitational field, the action is given by  $mc^2\Delta\tau$ , where  $\Delta\tau$  is the elapsed proper time,  $\Delta\tau = \int d\tau$ . But distances and time intervals are measured with a spatial varying metric  $g_{\mu\nu}$ , with  $\mu$  and  $\nu$  ranging from 0 to 3, with the zeroth component referring to time. The four components of extended configuration space are written  $x^\mu$ , with a superscript rather than a subscript, and  $x^0 = ct$ . The other three  $x^\mu$  can be generalized coordinates, as long as  $g_{\mu\nu}$  is appropriate. In the next problem they are similar to spherical coordinates. The gravitational field is described by the space-time dependence of the metric  $g_{\mu\nu}(x^\rho)$ . In this language, an infinitesimal element of the path of a particle corresponds to a proper time  $d\tau = (1/c)\sqrt{\sum_{\mu\nu} g_{\mu\nu}dx^\mu dx^\nu}$ , so

$$S = mc^2\Delta\tau = mc \int d\lambda \sqrt{\sum_{\mu\nu} g_{\mu\nu}(x^\rho) \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}.$$

This is in preparation for the next problem which is worth 20 points, twice normal.

**5.3** In problem 2.12 we learned that the general-relativistic motion of a particle in a gravitational field is given by Hamilton's variational principle on the path  $x^\mu(\lambda)$  with the action

$$S = \int d\lambda \mathcal{L} \quad \text{with} \quad \mathcal{L} = mc \sqrt{\sum_{\mu\nu} g_{\mu\nu}(x^\rho) \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}},$$

where we may freely choose the path parameter  $\lambda$  to be the proper time (after doing the variation), so that the  $\sqrt{\quad}$  is  $c$ , the speed of light.

The gravitational field of a static point mass  $M$  is given by the Schwarzschild metric

$$g_{00} = 1 - \frac{2GM}{rc^2}, \quad g_{rr} = -1 / \left(1 - \frac{2GM}{rc^2}\right), \quad g_{\theta\theta} = -r^2, \quad g_{\phi\phi} = -r^2 \sin^2 \theta,$$

where all other components of  $g_{\mu\nu}$  are zero. Treating the four  $x^\mu(\lambda)$  as the coordinates, with  $\lambda$  playing the role of time, find the four conjugate momenta  $p_\mu$ , show that  $p_0$  and  $p_\phi = L$  are constants, and use the freedom to choose

$$\lambda = \tau = \frac{1}{c} \int \sqrt{\sum_{\mu\nu} g_{\mu\nu}(x^\rho) \frac{dx^\mu}{d\lambda} \frac{dx^\nu}{d\lambda}}$$

to show  $m^2c^2 = \sum_{\mu\nu} g^{\mu\nu} p_\mu p_\nu$ , where  $g^{\mu\nu}$  is the inverse matrix to  $g_{\alpha\beta}$ . Use this to show that

$$\frac{dr}{d\tau} = \sqrt{\kappa - \left(-\frac{2GM}{r} + \frac{L^2}{m^2r^2} - \frac{2GML^2}{m^2r^3c^2}\right)},$$

where  $\kappa$  is a constant. For an almost circular orbit at the minimum  $r = a$  of the effective potential this implies, show that the precession of the perihelion is  $6\pi GM/ac^2$ .

Find the rate of precession for Mercury, with  $G = 6.67 \times 10^{-11} \text{ Nm}^2/\text{kg}^2$ ,  $M = 1.99 \times 10^{30} \text{ kg}$  and  $a = 5.79 \times 10^{10} \text{ m}$ , per revolution, and also per century, using the period of the orbit as 0.241 years.