

HW 4 Solutions

1. For crystal diffraction we need $\lambda \sim 1 \text{ \AA} \equiv \lambda_0$.
Since we need $\lambda \sim$ interatomic spacing.

Using the de Broglie relation

$$\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}} \Rightarrow \text{Needed: } E = \frac{h^2}{2m\lambda_0^2}$$

$$m_{\text{neutron}} > m_{\text{electron}}$$

\Downarrow

$$E_{\text{neutron}} < E_{\text{electron}} \quad \left(E \sim \frac{1}{m} \right)$$

2. For crystal diffraction to occur we would like

$$\lambda \sim d$$

where we note that from Bragg's law

$$\lambda = 2d \sin \theta$$

there is no diffraction for $\lambda > 2d$ (since $\sin \theta \leq 1$)

Thus visible green ($\sim 5100 \text{ \AA}$) light could provide diffraction spots when scattered from a crystal with lattice constant $d \sim 5000 \text{ \AA}$; this is three orders of magnitude larger than that of crystals we see in Nature.

3. Powder Diffraction on CaF_2 .

$\lambda = 0.1542$

We follow the procedure in Avrami p. 154.

a)

Peak	θ	$d = \frac{\lambda}{2 \sin \theta}$	$\frac{d_1^2}{d^2}$	$h^2 + k^2 + l^2$	{hkl}
1	14.16	.3151	1	3	111
2	16.41	.2729	1.334	4	200
3	23.55	.1929	2.668	8	220
4	27.94	.1646	3.668	11	311

b) Peak $a = d \sqrt{h^2 + k^2 + l^2}$ (nm).

1	.54589
2	.54582
3	.54579
4	.5476



$\bar{a} = 0.54582 \text{ nm}$

Standard error of this estimate

$$\delta a = \left\{ \frac{1}{n(n-1)} \sum (a_i - \bar{a})^2 \right\}^{1/2}$$

$$\sim 2.821 \times 10^{-5} \text{ nm}$$



Estimated lattice constant

$$\bar{a} = 0.5482 \pm .00003 \text{ nm}$$

$$\begin{aligned} \text{c) } \rho &= \frac{m}{V} = \frac{1}{a^3} \left\{ \underset{\substack{\downarrow \\ 40}}{4 m_{Ca}} + \underset{\substack{\downarrow \\ 19.}}{8 m_F} \right\} \\ &= 3,189 \text{ kg/m}^3 \end{aligned}$$

4.

Key points of the Donev et al. paper.

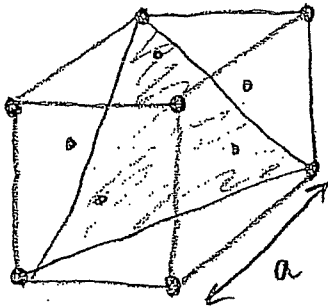
- Main result: Oblate spheroids (MEMs) pack differently (more densely) than do spheres when poured randomly and shaken.
- Extensions of work w/ more experiments on spheroids w/ different aspect ratios and with computer simulations \Rightarrow random packing densities approach those of perfectly ordered arrangements of spheres
- Used MRI to check that there was no periodic

ordering in the center.

- Always assumed that periodic orderings are denser than are random ones \Rightarrow true for spheres but may not be the case for all spheroids.
- The change of shape sphere \rightarrow spheroids leads to major change in random packing densities (such shape changes have minimal effects on non-random packing).

5. $\lambda = 1.54 \text{ \AA}$

a) $\theta = 19.2^\circ$ for (111) plane $\Rightarrow d_{111}$



(111)

Bragg's Law

$$n\lambda = 2d \sin \theta$$

We assume $n=1$

$$d = \frac{\lambda}{2 \sin \theta} = \frac{1.54}{2 \sin 19.2}$$

$$d_{111} = 2.34 \text{ \AA}$$

For fcc lattice

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$\sqrt{h^2 + k^2 + l^2}$$

↓

$$a = d_{111} \sqrt{3} = 2.34 \sqrt{3} = 4.05 \text{ \AA}$$

$$a = 4.05 \text{ \AA}$$

$$\rho = \frac{\# \text{ Al atoms}}{\text{Volume of unit cell}} \times \frac{\text{molecular weight of Al (g-atom)}}{N (\# \text{ atoms in 1 g-atom})}$$

⇓

$$N = \frac{\# \text{ Al atoms}}{\text{volume of unit cell}} \times \frac{M (\text{Al})}{\rho}$$

Al is fcc

• # atoms / unit cell

$$= \frac{8 \text{ corner atoms}}{8 \text{ cells}} + \frac{6 \text{ face atoms}}{2 \text{ cells}}$$

$$= 1 + 3 = 4 \text{ atoms / cell}$$

⇓

$$N = \frac{4 \text{ atoms}}{(4.05 \times 10^{-8})^3 \text{ cm}^3} \times \frac{27 \text{ g}}{2.7 \text{ g/cm}^3}$$

$$= \frac{4}{(4.05)^3} \times 10^{25} \text{ atoms} = 6.02 \times 10^{23} \text{ atoms}$$

$N = 6.02 \times 10^{23} \text{ atoms}$

6.

Interplanar Separation

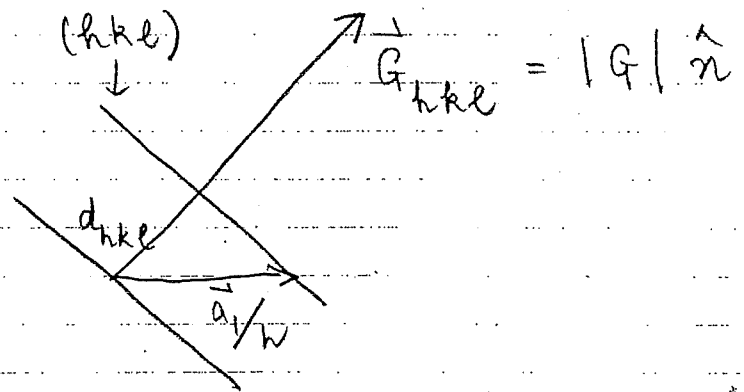
The plane (hkl) is defined by the intercepts

$$\frac{\vec{a}_1}{h}, \quad \frac{\vec{a}_2}{k} \quad \text{and} \quad \frac{\vec{a}_3}{l}$$

a) We take two vectors in this plane

$$\vec{A} = \frac{\vec{a}_1}{h} - \frac{\vec{a}_2}{k}$$

$$\vec{B} = \frac{\vec{a}_1}{h} - \frac{\vec{a}_3}{l}$$



$$\vec{G} = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$$

$$\left. \begin{array}{l} \vec{G} \cdot \vec{A} = 0 \\ \vec{G} \cdot \vec{B} = 0 \end{array} \right\} \Rightarrow \vec{G} \text{ must be perpendicular to } (hkl)$$

b) Let \hat{n} be the unit normal to the plane



interplanar spacing is $\frac{\hat{n} \cdot \vec{a}}{h} = d(hkl)$

However

$\hat{n} = \frac{\vec{G}}{|\vec{G}|}$



$d(hkl) = \frac{\overbrace{\vec{G} \cdot \vec{a}_1}^{2\pi h}}{h |\vec{G}|} = \frac{2\pi}{|\vec{G}|}$

c) For a simple cubic lattice

$\vec{G} = \frac{2\pi}{a} (h\hat{x} + k\hat{y} + l\hat{z})$



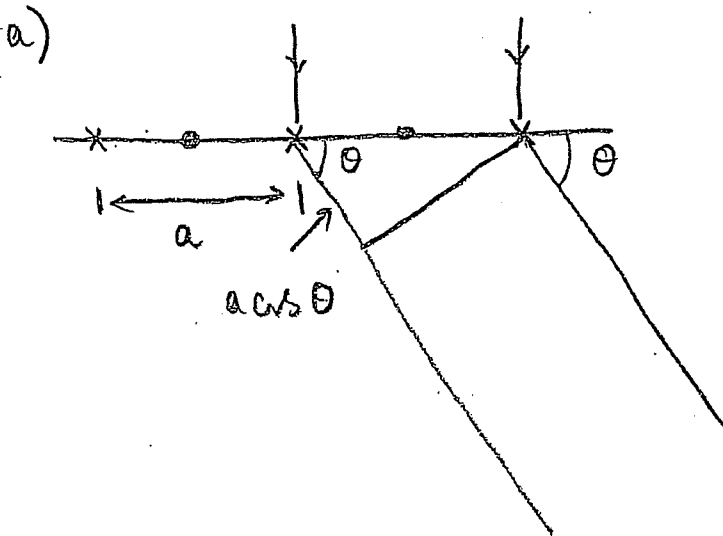
$\frac{1}{d^2} = \frac{G^2}{4\pi^2} = \frac{h^2 + k^2 + l^2}{a^2}$



$d^2 = \frac{a^2}{h^2 + k^2 + l^2}$

7. Diatomic Line

Let $x = A$, $\bullet = B$.



Constructive interference iff

$$n\lambda = a \cos \theta$$

b)

$$\Delta_G = \sum_j f_j e^{-i2\pi v_j x_j}$$

$$= f_A + f_B e^{-i v_j \pi}$$

$$v_j \text{ even} \quad e^{-i v_j \pi} = +1$$

$$v_j \text{ odd} \quad e^{-i v_j \pi} = -1.$$

$$x_A = 0$$

$$x_B = 1/2$$

therefore

$$\Delta_G = f_A - f_B \quad v_1 \text{ odd}$$

$$\Delta_G = f_A + f_B \quad v_2 \text{ even}$$

↓

$$I \propto |\Delta_G|^2 = \begin{cases} |f_A - f_B|^2 & v_1 \text{ odd} \\ |f_A + f_B|^2 & v_1 \text{ even} \end{cases}$$

c) $f_A = f_B$

No reflection w/ v_1 odd

Same reflections as for monatomic lattice
w/ $a/2$ lattice spacing.

8. Hexagonal Space Lattice

Hexagonal lattice $\vec{a}_1 = \frac{\sqrt{3}}{2} a \hat{x} + \frac{a}{2} \hat{y}$

$$\vec{a}_2 = -\frac{\sqrt{3}}{2} a \hat{x} + \frac{a}{2} \hat{y}$$

$$\vec{a}_3 = c \hat{z}$$

(a) Cell volume = $\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)$

$$= \frac{1}{a_1} \cdot \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ -\frac{\sqrt{3}}{2} a & \frac{a}{2} & 0 \\ 0 & 0 & c \end{vmatrix}$$

$$= \left(\frac{\sqrt{3}}{2} a \hat{x} \right) \cdot \left(\frac{ac}{2} \hat{x} \right)$$

$$+ \left(\frac{a}{2} \hat{y} \right) \cdot \left(\frac{\sqrt{3}}{2} ac \hat{y} \right)$$

$$= \frac{2\sqrt{3}}{4} a^2 c = \frac{\sqrt{3}}{2} a^2 c$$

$$\vec{b}_1 = \frac{2\pi}{V_c} (\vec{a}_2 \times \vec{a}_3)$$

$$= \frac{2\pi}{\frac{\sqrt{3}}{2} a^2 c} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ -\frac{\sqrt{3}}{2} a & \frac{a}{2} & 0 \\ 0 & 0 & c \end{vmatrix}$$

↓

$$\vec{b}_1 = \frac{2\pi}{\sqrt{3} a} \hat{x} + \frac{2\pi}{a} \hat{y}$$

$$\vec{b}_2 = \frac{2\pi}{V_c} (\vec{a}_3 \times \vec{a}_1)$$

$$= \frac{4\pi}{\sqrt{3} a^2 c} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & c \\ \frac{\sqrt{3}}{2} a & \frac{a}{2} & 0 \end{vmatrix}$$

$$\vec{b}_2 = + \frac{2\pi}{a} \left(-\frac{1}{\sqrt{3}} \hat{x} + \hat{y} \right)$$

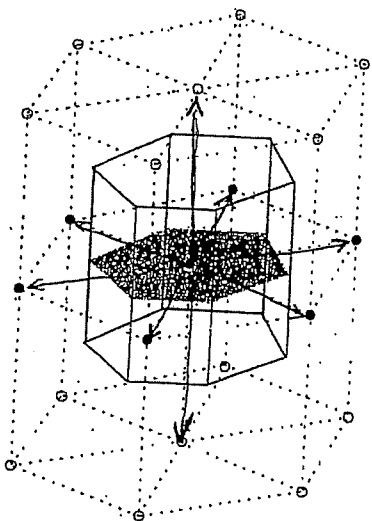
$$\vec{k}_3 = \frac{2\pi}{V_c} (\vec{a}_1 \times \vec{a}_2)$$

$$= \frac{4\pi}{\sqrt{3} a^2 c}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\sqrt{3}}{2} a & \frac{a}{2} & 0 \\ -\frac{\sqrt{3}}{2} a & \frac{a}{2} & 0 \end{vmatrix}$$

$$\vec{k}_3 = \frac{2\pi}{c} \hat{z}$$

(c)



The reciprocal lattice is shown as dots and is shifted by an angle of 30° from the real-space lattice. The shaded BZ is also shown; it is oriented w/ the real space lattice (shifted from the reciprocal lattice by 30°).

9. Sample Mix-Up Problem

$$\Delta_G(v_1, v_2, v_3) = \sum_j f_j \exp\left(-2\pi i \left\{ \begin{array}{l} v_1 x_j + v_2 y_j \\ + v_3 z_j \end{array} \right\}\right)$$

bcc

$$(000) \left(\frac{1}{2} \frac{1}{2} \frac{1}{2}\right)$$

$$\Delta(v_1, v_2, v_3) = f (1 + \exp -i\pi (v_1 + v_2 + v_3))$$

$$\Delta = \begin{cases} 0 & \sum v_i = \text{odd integer} \\ 2f & \sum v_i = \text{even integer} \end{cases}$$

$$fcc (000) \left(0 \frac{1}{2} \frac{1}{2}\right) \left(\frac{1}{2} 0 \frac{1}{2}\right) \left(\frac{1}{2} \frac{1}{2} 0\right)$$

$$\Delta(v_1, v_2, v_3) = f \left\{ \begin{array}{l} 1 + \exp -i\pi (v_2 + v_3) \\ + \exp -i\pi (v_1 + v_3) \\ + \exp -i\pi (v_1 + v_2) \end{array} \right.$$

$$\Delta = \begin{cases} 4f & \text{all } v\text{'s odd or even} \\ 0 & \text{otherwise} \end{cases}$$

(110) cannot occur for fcc but is OK for bcc

Sample is bcc

10. Error Analysis

Bragg's Law

$$d_{hkl} = \frac{\lambda}{2 \sin \theta} = \frac{a}{\sqrt{h^2 + k^2 + l^2}} = \frac{a}{\sqrt{N}}$$

$$\Downarrow$$

$$a = \frac{\sqrt{N} \lambda}{2 \sin \theta}$$

$$|\delta a| = \left| \frac{\sqrt{N} \lambda}{2} \frac{\cos \theta}{\sin^2 \theta} \right| |\delta \theta|$$

$$\Downarrow$$

$$\left| \frac{\delta a}{a} \right| = \left| \cot \theta \right| |\delta \theta|$$

note that even if $\delta \theta$ small and fixed, there will always be this uncertainty in $\frac{\delta a}{a}$

Reduction of error

- take many measurements at angles where $\cot \theta$ not large

Sources of error

- sample displacement (not centered)
- preferred orientation of crystallites
- beam width

Note that here we assume the $\Delta\lambda$ is negligible \rightarrow my experimental colleagues tell me this is a reasonable assumption given other sources of error.

11. Marder 3.2 Hex extinctions

a) From problem # 8.

$$\vec{a}_1 = \frac{a}{2} \begin{pmatrix} \sqrt{3} \\ 1 \end{pmatrix}$$

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$$

$$\Downarrow \theta = 30$$

$$\begin{bmatrix} 1 \\ \sqrt{3} \end{bmatrix} \leftarrow \begin{bmatrix} \sqrt{3}/2 & -1/2 \\ 1/2 & \sqrt{3}/2 \end{bmatrix} \begin{bmatrix} \sqrt{3} \\ 1 \end{bmatrix} \frac{a}{2}$$

$$\vec{b}_1 = \frac{2\pi}{a} \begin{bmatrix} 1 \\ \sqrt{3} \end{bmatrix}$$

from
8

\Downarrow

real space and reciprocal lattices
are rotated by $\theta = 30^\circ$ from ^{one} another.

(b) Basis vectors

$$\begin{pmatrix} 0 & 0 & 0 \end{pmatrix} \\ \parallel \\ v_1$$

$$\begin{pmatrix} a & 0 & c \\ \sqrt{3} & & 2 \end{pmatrix} \\ \parallel \\ v_2$$

Using \vec{G} from Problem # 8

\Downarrow

$$\frac{1}{\sqrt{2}} \cdot \vec{G} = \frac{\pi}{3} \{ 2(n_1 + n_2) + 3n_3 \}$$

\Downarrow

$$F_g = \left| 1 + e^{\frac{i\pi}{3} \{ 2(n_1 + n_2) + 3n_3 \}} \right|^2$$

c) $F_g = 0$ when $2(n_1 + n_2) + 3n_3 = 3N$.

12. Marder 3.4 Lattice Patterns

From Problem # 9

$$\text{bcc } f = \begin{cases} 0 & \sum v_i = \text{odd} \\ 2f & \sum v_i = \text{even} \end{cases}$$

$$v_i \in \{h, k, l\}$$

$$\text{fcc } f = \begin{cases} 4f & \text{all } v\text{'s odd or even} \\ 0 & \text{otherwise} \end{cases}$$

All reflections are allowed for simple cubic lattice

a) Three Bragg angles $\theta_1 = 18.5$

$$\theta_2 = 24$$

$$\theta_3 = 27$$

$$\text{Wavevector cutoff} = \frac{32}{a}$$

$$\frac{2\pi}{\lambda_{\min}} = \frac{32}{a} \Rightarrow \lambda_{\min} = \frac{\pi a}{16}$$

$$d = \frac{\lambda_{\min}}{2 \sin \theta} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$\frac{\pi a}{32 \sin(27)} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$h^2 + k^2 + l^2 = \left(\frac{32 \sin 27}{\pi} \right)^2$$

$$= 21.4 \sim 21$$

⇓

$$h^2 + k^2 + l^2 \ll 21$$

$$h, k, l = 0, 1, 2 \text{ or } 3$$

$$\frac{\lambda}{2 \sin \theta} = \frac{a}{\sqrt{h^2 + k^2 + l^2}} = \frac{a}{N_{\theta}^2}$$

⇓

$$N_{\theta}^2 = \left(\frac{2a \sin \theta}{\lambda} \right)^2$$

$$N_{\theta}^2 = (h^2 + k^2 + l^2) \lambda_0$$

$$\frac{N_{\theta_2}^2}{N_{\theta_1}^2} = \left(\frac{\sin \theta_2}{\sin \theta_1} \right)^2$$

(assume due to roughly same λ)

$$\frac{N_{24}}{N_{18.5}} = 1.6$$

$$\frac{N_{26.5}}{N_{18.5}} = 1.9 \sim 2$$

$$\frac{N_{26.5}}{N_{24}} = 1.2$$

Again -

$$h^2 + k^2 + l^2$$

$$\frac{N_{24}}{N_{18.5}} = 1.6 \quad \times 5 \quad \quad \quad 8 \quad \quad \quad 220$$

$N_{18.5}$

$$\frac{N_{26.5}}{N_{24}} = 1.2 \quad \times 5 \quad \quad \quad 6 \quad \quad \quad 211$$

N_{24}

$$\frac{N_{26.5}}{N_{18.5}} = 1.9 \sim 2 \quad \times 5 \quad \quad \quad 10 \quad \quad \quad 310$$

$N_{18.5}$

x fcc NOT all odd or all even

✓ bcc $\sum v_i = \text{even}$

No other peaks (would be there)
for fcc



Proposed bcc

13. Free electron lasers \leftrightarrow noncrystalline samples.

- Diffraction due to constructive / destructive interference \rightarrow crystallinity not necessary.
- Large flux of X-rays \rightarrow single shot diffraction
- Problem: X-rays destroy samples!
- Idea: Pulsed X-rays / femtoseconds before destruction \rightarrow diffraction pattern
- Alignment of biological molecules (e.g. electric dipole) to boost signal-to-noise
- Heating issues?