

① 0. Ch. 2, Q. 1

2/2 The scattered rays are nearly parallel because the detector is located very far away from the crystal.

② 0 Ch. 2, Q. 5

3/3 The lattice structure factor will be the same since they are both fcc. The atomic form factors will be different however.

③ 0 Ch. 2, Q. 10

acc. to Bragg's law,

$$\lambda = 2d \sin \theta$$

5/5 and so there is no diffraction for $\lambda > 2d$ (since $\sin \theta \leq 1$).

Roughly, we need $\lambda \sim d$.

i.e. $\lambda = 5100 \text{ \AA}$ (visible green light),

$d \approx 5000 \text{ \AA}$, 10^3 times greater than what we see in nature.

4. 0. Ch. 2, Pr. 1

Recall that

$$\lambda_{\min} = \frac{12.3}{V(\text{kV})} \text{ \AA} :$$

5/5

$$V(\text{kV}) = \frac{12.3}{1.23} = \underbrace{10^4}_{10 \text{ kV}} \text{ V.}$$

The kinetic energy

$$\text{KE} = eV = \underline{\underline{10^4 \text{ eV.}}}$$

5. 0. Ch. 2, Pr. 3

a) $\lambda = 1.54 \text{ \AA}$

$\theta = 19.2^\circ$

(111) planes

Bragg's law: $n\lambda = 2d \sin \theta$.

5/5 assuming $n=1$,

$$d_{111} = \frac{\lambda}{2 \sin \theta} = \frac{1.54 \text{ \AA}}{2 \sin(19.2^\circ)} = \underline{\underline{2.34 \text{ \AA}}}$$

b)

$$\rho = \frac{\# \text{ Al atoms}}{\text{unit cell volume}} \times \frac{\text{mol. weight of Al}}{N(\# \text{ atoms in 1 mole})}$$

Al is fcc: $8 \times \frac{1}{8} + 6 \times \frac{1}{2} = 4 \frac{\text{atoms}}{\text{unit cell.}}$

$V = a^3$ unit cell volume

For fcc lattice

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}, \text{ or}$$

$$a = d_{111} \sqrt{3} \approx 4.05 \text{ \AA}$$

$$\text{So, } N = \frac{\# \text{ of atoms}}{V} \times \frac{M(\text{of})}{\rho} =$$

$$= \frac{4}{(4.05 \times 10^{-8})^3 \text{ cm}^3} \times \frac{27 \text{ g/mole}}{2.7 \text{ g/cm}^3} =$$

$$= \frac{4}{(4.05)^3} \times 10^{25} \text{ atoms/mole} \approx$$

$$\approx 6.02 \times 10^{23} \text{ atoms/mole.}$$

Avogadro's number

(6) The plane (hkl) is defined by its intercepts

$$5/5 \quad \frac{\vec{a}_1}{h}, \quad \frac{\vec{a}_2}{k}, \quad \frac{\vec{a}_3}{l}$$

(a) Define
$$\begin{cases} \vec{A} = \frac{\vec{a}_1}{h} - \frac{\vec{a}_2}{k}, \\ \vec{B} = \frac{\vec{a}_1}{h} - \frac{\vec{a}_3}{l} \end{cases}$$

These 2 vectors are in the plane.

$$\vec{G} = h \vec{b}_1 + k \vec{b}_2 + l \vec{b}_3 =$$

$$= \frac{2\pi}{V} h (\vec{a}_2 \times \vec{a}_3) + \frac{2\pi}{V} k (\vec{a}_3 \times \vec{a}_1) +$$

$$+ \frac{2\pi}{V} l (\vec{a}_1 \times \vec{a}_2), \text{ where } V \text{ is the unit cell volume.}$$

Then
$$\vec{G} \cdot \vec{A} = 2\pi - 2\pi = 0,$$

$$\vec{G} \cdot \vec{B} = 0.$$

$$\vec{G} \text{ is } \perp \text{ to } (hkl).$$

(b) Let \hat{n} be the unit normal to the plane (hkl) .

$$\frac{\vec{a}_1}{h} \cdot \hat{n} = d_{hkl} \quad (\text{projection of } \frac{\vec{a}_1}{h} \text{ onto } \hat{n})$$

But $\hat{n} = \frac{\vec{G}}{|\vec{G}|}$, as shown in (a).

$$\text{Then } d_{hkl} = \frac{\vec{a}_1 \cdot \vec{G}}{h|\vec{G}|} = \frac{2\pi}{|\vec{G}|} \\ =$$

(c) For an sc lattice,

$$\vec{G} = \frac{2\pi}{a} (h\hat{x} + k\hat{y} + l\hat{z})$$

$$d_{hkl}^2 = \frac{4\pi^2}{G^2} = \frac{4\pi^2}{\left(\frac{2\pi}{a}\right)^2 (h^2 + k^2 + l^2)} = \\ = \frac{a^2}{h^2 + k^2 + l^2} \\ =$$

⑦ By definition,

$$\vec{b}_1 = 2\pi \left(\frac{a}{2}\right)^2 \frac{(\hat{z} + \hat{x}) \times (\hat{x} + \hat{y})}{V}$$

5/5 where $V = \left(\frac{a}{2}\right)^3 (\hat{y} + \hat{z}) \cdot [(\hat{z} + \hat{x}) \times (\hat{x} + \hat{y})] =$

$$= \left(\frac{a}{2}\right)^3 (\hat{y} + \hat{z}) \cdot (\hat{y} - \hat{x} + \hat{z}) =$$

$$= 2 \left(\frac{a}{2}\right)^3$$

$$\text{So, } \vec{b}_1 = \frac{2\pi \left(\frac{a}{2}\right)^2}{2 \left(\frac{a}{2}\right)^3} (\hat{y} - \hat{x} + \hat{z}) =$$

$$= \frac{2\pi}{a} (\hat{y} - \hat{x} + \hat{z})$$

Likewise,

$$\vec{b}_2 = \frac{2\pi}{a} (\hat{x} + \hat{y}) \times (\hat{y} + \hat{z}) =$$

$$= \frac{2\pi}{a} (\hat{z} - \hat{y} + \hat{x})$$

$$\vec{b}_3 = \frac{2\pi}{a} (\hat{y} + \hat{z}) \times (\hat{z} + \hat{x}) =$$

$$= \frac{2\pi}{a} (\hat{x} - \hat{z} + \hat{y})$$

This is a bcc lattice, the lattice const is $\frac{4\pi}{a}$.

⑧ It is known that the packing ratio of spheres into an fcc lattice is $\phi \approx 0.74$. However, $\phi \approx 0.64$ for random jammed packings (amorphous solids). Here the authors show, by experiments and computer simulations, that ϕ up to 0.74 can be achieved using ellipsoids. This dense packing is accompanied by an increased number of contacts per particle, up to 10 for ellipsoids vs. 6 for spheres.

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