

Due date: Monday March 24.

1) (25 points)

We generalize some of the discussion of $S(q, \omega)$ to finite temperature. Let n label the exact many-body eigenstates; at $\beta = 1/k_B T$, the thermal expectation value of an operator A is

$$\langle A \rangle = Z^{-1} \sum_n e^{-\beta \omega_n} \langle n | A | n \rangle ; \quad Z = \sum_n e^{-\beta \omega_n} .$$

We define

$$S(q, t) = \langle \rho_q(t) \rho_q^\dagger(0) \rangle$$

and the expression for $\epsilon^{-1}(q, \omega)$ is also written as a weighted sum over initial states. Show:

- a) $S(q, \omega) = e^{\beta \omega} S(q, -\omega)$;
 b) $\text{Im} \epsilon^{-1}(q, \omega) = -\pi \tilde{V}(q) [1 - e^{-\beta \omega}] S(q, \omega)$.

2) (30 points)

The purpose of this problem is to work out an *exact* expression for the *transverse* dielectric function of the uniform electron gas, using the Kubo-formula approach. The transverse dielectric function is defined as

$$\epsilon_{\alpha\beta}(q, \omega) = \delta_{\alpha\beta} + \frac{4\pi i \sigma_{\alpha\beta}(q, \omega)}{\omega}$$

where the conductivity σ is defined via

$$J_\alpha(r, t) = \int dr' \int dt' \sigma_{\alpha\beta}(r - r', t - t') E_\beta(r', t') ,$$

$$J_\alpha(q, \omega) = \sigma_{\alpha\beta}(q, \omega) E_\beta(q, \omega)$$

(implied sum notation used throughout). In the Coulomb gauge the perturbation is

$$H'(t) = -\frac{1}{c} \int dr j_\alpha(r) A_\alpha(r, t) .$$

Now there is a subtlety: J_α and j_α are not quite the same; they are defined as

$$j_\alpha(r) = \frac{e}{2m} \sum_i [p_{i\alpha} \delta(r - r_i) + \delta(r - r_i) p_{i\alpha}] ,$$

$$J_\alpha(r) = \frac{e}{2} \sum_i [v_{i\alpha} \delta(r - r_i) + \delta(r - r_i) v_{i\alpha}] .$$

Continued on other side...

Recalling that $p = mv + \frac{e}{c}A$, it follows that

$$\langle j_\alpha(r) - J_\alpha(r) \rangle = \frac{e^2 n}{mc} A_\alpha(r)$$

where $\langle \rangle$ is evaluated in the ground state of uniform density. Finally, recall that $E = (-1/c)(dA/dt)$ so that $E(q, \omega) = i\omega A(q, \omega)/c$.

a) Using the Kubo approach derived in class, with a source acting on operator $j_\alpha(r', t')$ and a response obtained from the expectation of $j_\beta(r, t)$, show that

$$\sigma_{\alpha\beta}(q, \omega) = \frac{i}{\omega} \left[\Pi_{\alpha\beta}(q, \omega) + \frac{ne^2}{m} \delta_{\alpha\beta} \right]$$

where

$$\Pi_{\alpha\beta}(q, t) = -i\theta(t) \langle j_\alpha(q, t), j_\beta^\dagger(q, 0) \rangle .$$

b) Show that this leads to a transverse dielectric function

$$\epsilon_{\alpha\beta}(q, \omega) = \left[1 - \frac{\omega_p^2}{\omega^2} \right] \delta_{\alpha\beta} - \frac{4\pi}{\omega^2} \Pi_{\alpha\beta}(q, \omega) .$$

Note that the result is non-trivial even in the extreme approximation $\Pi = 0$.

c) Write down the exact spectral representation of $\Pi_{\alpha\beta}(q, \omega)$ in terms of the exact many-body eigenstates $|n\rangle$ and their excitation energies ω_{n0} .

3) (25 points)

This is a simple exercise in Fermi liquid theory. Starting from

$$\delta E = \sum_k \epsilon(k) \delta f(k) + \frac{1}{2V} \sum_k \sum_{k'} u(k, k') \delta f(k) \delta f(k') ,$$

$$E(k) = \frac{\delta E}{\delta f(k)} ,$$

$$\mathbf{v}^{(0)}(k) = \nabla_k \epsilon(k) ,$$

$$\mathbf{v}(k) = \nabla_k E(k) ,$$

show that the contribution to the particle current

$$J = \sum_k f(k) \mathbf{v}(k)$$

from the quasiparticle k is, at $T = 0$,

$$\mathbf{j}_k = \mathbf{v}^{(0)}(k) + \frac{1}{V} \sum_{k'} u(k, k') \mathbf{v}^{(0)}(k') \delta[\epsilon(k') - \epsilon_F]$$

where ϵ_F is the Fermi level. Hint: An integration by parts may be needed at some stage.