

Due date: Monday March 3.

1) (10 points)

Consider the expression

$$\tilde{E}(\mathbf{k}) = E(\mathbf{k}) + \sum_{\mathbf{q}} |M_{\mathbf{q}}|^2 \frac{1 - 2\bar{n}_{\mathbf{k}+\mathbf{q}}}{E(\mathbf{k}) - E(\mathbf{k} + \mathbf{q}) - \omega_{\mathbf{q}}} ,$$

derived in class for the renormalization of the electron quasiparticle energy when the electrons (\mathbf{k}) are coupled to plasmons (\mathbf{q}). Argue that the renormalization leads to a *reduction* in the Fermi velocity $dE/dk|_{k_F}$. You may assume that the dominant contributions come from small q . [This is intuitive in a way: the electron gets heavier and slower when dressed by a cloud of plasmons.]

2) (30 points)

In class, I introduced (without derivation – that is coming shortly) the expression for the Lindhard dielectric susceptibility

$$\chi(\mathbf{q}, \omega) = -\frac{e^2}{\Omega} \sum_{\mathbf{k}\sigma} \frac{\bar{n}_{\mathbf{k}} - \bar{n}_{\mathbf{k}+\mathbf{q}}}{E(\mathbf{k} + \mathbf{q}) - E(\mathbf{k}) - \omega + i\eta}$$

where $\bar{n}_{\mathbf{k}} = \Theta(k_F - k)$ and η is an infinitesimal which is taken to zero from above. In the present problem, we only consider real parts of the response functions, and therefore discard η .

a) Show that (the real part of) χ can be written

$$\chi(q, \omega) = -\frac{2e^2}{\Omega} \sum_{\mathbf{k}\sigma} \bar{n}_{\mathbf{k}} \frac{E(\mathbf{k} + \mathbf{q}) - E(\mathbf{k})}{[E(\mathbf{k} + \mathbf{q}) - E(\mathbf{k})]^2 - \omega^2} .$$

[Hint: Write the terms involving $\bar{n}_{\mathbf{k}}$ and $\bar{n}_{\mathbf{k}+\mathbf{q}}$ separately, and change the sum index from \mathbf{k} to $\mathbf{k}' = \mathbf{k} + \mathbf{q}$ in the latter.]

b) Show that in the long-wavelength limit, $q \rightarrow 0$, one obtains

$$\epsilon(\omega) = 1 - \frac{\omega_p^2}{\omega^2} ,$$

where ω_p is the plasma frequency.

c) Show that in the static limit, $\omega \rightarrow 0$, one obtains $\chi(q) = -e^2 n(\epsilon_F) F(x)$, with

$$F(x) = \left[\frac{1}{2} + \frac{1-x^2}{4x} \ln \left| \frac{1+x}{1-x} \right| \right] ,$$

and $x = q/2k_F$.

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3) (20 points)

Kramers-Kronig (“KK”) relations:

a) Show that if the kernel $K(t - t')$ appearing in

$$B(t) = \int dt' K(t - t') A(t')$$

is not only causal but real, then the KK relations take the form

$$K_1(\omega) = \frac{2}{\pi} \int_0^\infty K_2(\omega') P \frac{\omega' d\omega'}{(\omega')^2 - \omega^2} ,$$

$$K_2(\omega) = -\frac{2}{\pi} \int_0^\infty K_1(\omega') P \frac{\omega d\omega'}{(\omega')^2 - \omega^2} ,$$

where K_1 and K_2 are the real and imaginary parts of K respectively.

b) For a system with translational and rotational invariance, the response function $K(q, t)$ (i.e., with *wavevector* and *time* arguments) is indeed real. (Why?) So, using (a) together with $\epsilon(q, \omega) = 1 - \tilde{V}(q)\chi(q, \omega)$ and $\epsilon^{-1}(q, \omega) = 1 + \tilde{V}(q)\pi(q, \omega)$, derive KK relations for $\epsilon_1(q, \omega)$, $\epsilon_2(q, \omega)$, $\text{Re } \epsilon^{-1}(q, \omega)$, and $\text{Im } \epsilon^{-1}(q, \omega)$, involving integrals over only half the frequency axis.