

THE QUANTUM ELECTRON PLASMA (RPA)

$$\mathcal{L}_{EM} = T - V \quad \leftarrow \frac{B^2}{2\mu_0} = \frac{1}{2\mu_0} (\nabla \times A)^2 =$$

$$\left. \begin{aligned} \uparrow \frac{\epsilon_0 E^2}{2} = \left(-\frac{\partial A}{\partial t} - \nabla \phi \right)^2 \frac{\epsilon_0}{2} &= \frac{1}{2\mu_0} \left(\frac{1}{c} \frac{\partial A}{\partial t} - \nabla \left(\frac{\phi}{c} \right) \right)^2 \\ S_Q &= \int \frac{d^4x}{\hbar} \left[\frac{\epsilon_0 E^2}{2} - \frac{B^2}{2\mu_0} \right] \rightarrow \text{Photons?} \end{aligned} \right\} \begin{array}{l} \epsilon_0 \rightarrow \frac{1}{4\pi} \\ \text{cgs } \mu_0 \rightarrow 4\pi \\ \frac{1}{8\pi} (E^2 - B^2) \end{array}$$

↓

Coulomb Gauge $\nabla \cdot A = 0 \quad -e_0 \left[\phi \frac{\partial (\nabla \cdot A)}{\partial t} \right] = 0$

$$S_Q = \int \frac{d^4x}{\hbar} \left[\left(\frac{\partial A}{\partial t} \right)^2 \frac{\epsilon_0}{2} + \frac{\epsilon_0 (\nabla \phi)^2}{2} - (\nabla \times A)^2 - e\phi g \right]$$

$$S_T = \int d^4x \left[\frac{\epsilon_0 E^2}{2} - \frac{B^2}{2\mu_0} \right]$$

$$S = \int d^3x d\tau \left[\psi^\dagger \left(\partial_\tau - \frac{\nabla^2}{2m} - \mu + e\phi(\tau) \right) \psi - \frac{\epsilon_0}{2} (\nabla \phi)^2 \right]$$

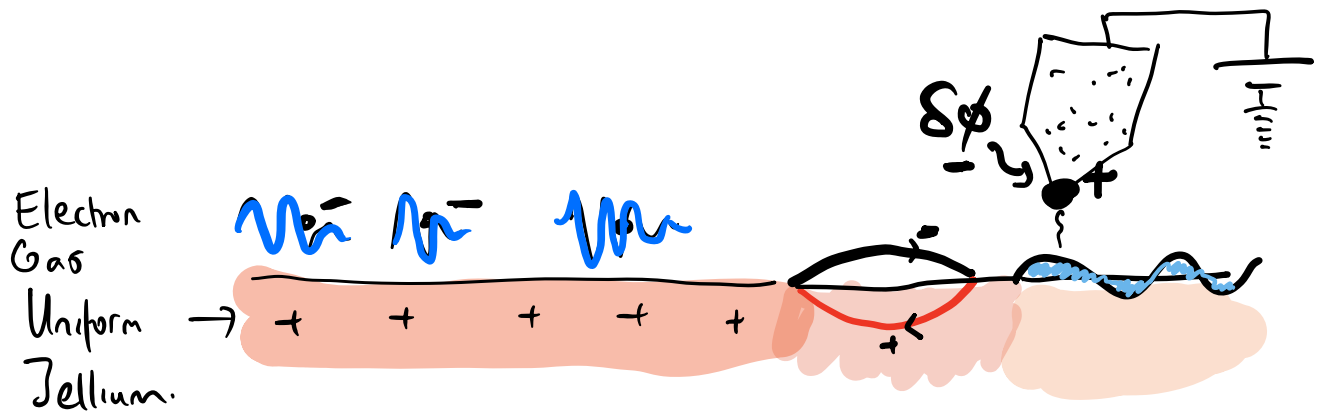
Coulomb Gauge

$$e\phi \rightarrow \phi$$

$$S = \int d^3x d\tau \left[\psi^\dagger \left(\partial_\tau - \frac{\nabla^2}{2m} - \mu + \phi \right) \psi - \frac{\epsilon_0 e^2}{2} (\nabla \phi)^2 \right]$$

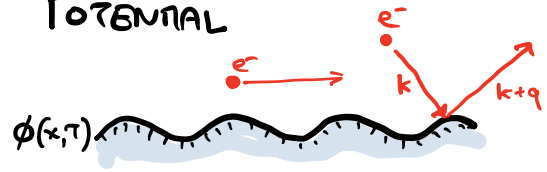
$$-g_0^{-1} = \left(\partial_\tau - \frac{\nabla^2}{2m} - \mu \right)$$

$$-g^{-1} = -g_0^{-1} + \phi$$




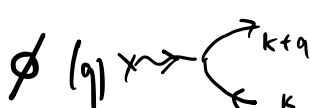



$$S = \int d^3x d\tau \left[\psi^\dagger - g_0^{-1} \psi + \phi (\psi^\dagger \psi - g_+) - \frac{e^2 \epsilon_0}{2} (\nabla \phi)^2 \right]$$

FEYNMAN DIAGRAMS: FERMIONS IN A VARYING POTENTIAL



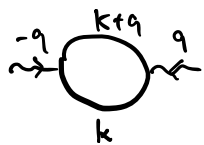
$$\begin{aligned}
 F &= -T \text{Tr} \ln[-G_0^{-1} + \phi] \\
 &= -T \text{Tr} \ln(-G_0^{-1}) - T \text{Tr} [1 - G_0 \phi] \\
 &= F_0 + T \text{Tr} \left[G_0 \phi + \frac{1}{2} (G_0 \phi)^2 + \frac{1}{3} (G_0 \phi)^3 + \dots \right] \\
 &= F_0 - T \left[\int d^4x G_0(x, x) \phi(x) - \frac{1}{2} \int d^4x d^4y \phi(x) G_0(x, y) \phi(y) G_0(y, x) \right. \\
 &\quad \left. - \dots \right] \\
 &\quad \text{Fermion loop (NOTE DOUBLE NEGATIVE)}.
 \end{aligned}$$

$$S = \beta F_0 - \int \left[\text{loop} + \text{loop} + \text{loop} \right] \quad z = e^{-S} = \sum (\text{closed unlinked diagrams})$$

REAL SPACE	MOMENTUM SPACE
$x' \longrightarrow x$ $G_0(x-x') = -\langle T \psi(x) \psi^\dagger(x') \rangle_{H_0}$	$G_0(k) = \frac{1}{i\nu_n - \epsilon_k}$
 $\phi(x)$	 $\phi_{-q} = \phi_q^*$
 $\prod_{i=1}^N \int d^4x_i$ all intermediate co-ordinates $(-1)^F$ $F = \#$ fermion loops. $\frac{1}{n}$ symmetry factor	$\prod_{\text{LOOPS}} \frac{1}{V\beta^n} \sum \int \frac{d^3k}{(2\pi)^3}$ All loop momenta. $(-1)^F$ Fermion loops. $\frac{1}{n}$ symmetry factor.
 $-V(x-y)$ interaction	 $-V(q)$

$$\frac{\Delta S}{\beta V} = \frac{\Delta F}{V} = -\frac{1}{\beta V} \left[\sum \text{closed link diagrams in real space} \right]$$

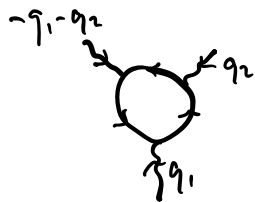
$$= - \sum \left[\text{closed link diagrams in momentum space.} \right]$$



$$= (-)^2 \frac{N}{2\beta V} \sum \phi_{-q} \phi_q G(k+q) G(k)$$



$$= - \frac{N}{2\beta V} \int d^4x d^4y \phi(x) \phi(y) G(x-y) G(y-x)$$



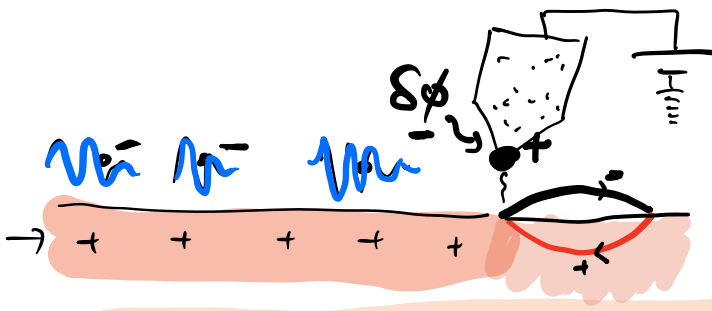
$$= - \frac{N}{3\beta V} \sum \phi_{-q_1-q_2} \phi_{q_1} \phi_{q_2} G(k+q_2) G(k+q_1) G(k)$$



$$= - \frac{N}{3\beta V} \int d^4x d^4y d^4z \phi(x) \phi(y) \phi(z) G(x-z) G(z-y) G(y-x)$$

"Dolphins leaping out of the fermi sea"

Electron Gas
Uniform
Jellium.



$$\langle \psi^\dagger(x) \psi(x) \rangle = \rho_0$$

$$S = \int d^4x \left\{ \bar{\psi}_\sigma (-\not{g}_0^{-1}) \psi_\sigma + \phi (\bar{\psi}_\sigma \psi_\sigma - \rho_0) - \frac{\epsilon_0}{2e^2} (\nabla \phi)^2 \right\}$$

Make saddle point exact $\sigma = 1, \dots, N$ $\epsilon_0 = \tilde{\epsilon}_0$

Repulsive interaction $\left(\frac{e^2}{\epsilon_0 q^2} \right)$

$$S = \int d^4x \left\{ \bar{\psi}_\sigma (-\not{g}_0^{-1}) \psi_\sigma + \phi (\bar{\psi}_\sigma \psi_\sigma - \rho_0) - \frac{\epsilon_0}{2e^2} (\nabla \phi)^2 \right\}$$

$$- \frac{\epsilon_0 q^2}{2e^2} |\phi_q|^2$$

" $|\delta \phi_q|^2 = \frac{-e^2}{\epsilon_0 q^2}$ "
 Repulsive

$$\frac{\epsilon_0}{e^2} = \frac{\tilde{\epsilon}_0}{\epsilon_0 N}$$

$$\frac{e^2}{4\pi \epsilon_0 r} = \left(\frac{\tilde{e}^2}{4\pi \epsilon_0 r N} \right)$$

$$S = \int d^4x \left\{ \bar{\psi}_\sigma (-\not{g}^{-1}) \psi_\sigma + \phi (\bar{\psi}_\sigma \psi_\sigma - \rho_0) - \frac{N \epsilon_0}{2\tilde{e}^2} (\nabla \phi)^2 \right\}$$

$$\begin{aligned}
Z &= \int \mathcal{D}[\phi] e^{\frac{N\epsilon_0}{2} \int d^4x (\nabla\phi)^2} \int \mathcal{D}[\bar{\psi}, \psi] e^{-N \int d^4x \left[\bar{\psi} (-g_0^{-1} + \phi) \psi - \phi g_0 \right]} \\
&= \int \mathcal{D}[\phi] e^{-S_\phi} e^{N \text{Tr} \ln (-g_0^{-1} + \phi)} \quad \text{+ve background} \\
&= \int \mathcal{D}[\phi] e^{-[S_\phi - \text{Tr} \ln (-g_0^{-1} + \phi)] N} \\
&= \int \mathcal{D}[\phi] e^{-S_{\text{eff}}[\phi]}
\end{aligned}$$

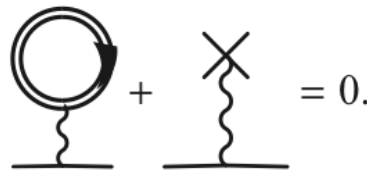
$$\begin{aligned}
\frac{\beta F}{N} = \frac{S_{\text{eff}}[\phi]}{N} &= -\text{Tr} \ln [-g_0^{-1} + \phi] - \int d^4x \left[\frac{\epsilon_0}{2e^2} (\nabla\phi)^2 + \phi g_0 \right] \\
&= -\text{Tr} \ln [-g_0^{-1} (1 - g_0 \phi)] - S_\phi \\
&= -\text{Tr} \ln (-g_0^{-1}) + \int d^4x \phi (g_0 g_0) + \frac{1}{2} \text{Tr} (g_0 \phi g_0 \phi) + O(\phi^3) \\
&\quad - \int d^4x \left[\frac{\epsilon_0}{2e^2} (\nabla\phi)^2 \right] N
\end{aligned}$$

$$\begin{aligned}
S_{\text{eff}} &= S_0 - \left[\text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots \right] - \phi g_0 \\
&= S_0 - \int d^4x \left[g_0 + \underbrace{\text{Diagram 1}}_{-G(\sigma)N} \right] \phi(x) - \int d^4x \left[\frac{\phi(x)}{2} \left(x \text{---} \text{Diagram 2} \text{---} x' + \frac{\epsilon_0}{e^2} \nabla^2 \right) \phi(x') \frac{d^4x'}{d^4x'} \right] + O(\phi^3)
\end{aligned}$$

$$\approx N S_0 - \sum_{\mathbf{k}} \frac{1}{2} \left[\text{Diagram 3} + \frac{N g^2}{e^2} \right] |\phi_{\mathbf{k}}|^2 + O(\phi^3)$$

$$\frac{\delta S_{\text{EFF}}}{\delta \phi} = 0 \Rightarrow N G_0(\vec{x}^-, x) = g_0 = -\langle \bar{\Psi}(\vec{x}^-) \Psi_0(x) \rangle = \langle \Psi_0^\dagger(x) \Psi(x) \rangle$$

NEUTRALITY



$$\frac{\delta^2 S}{e^2 \delta \phi(x) \delta \phi(x')}$$

$$= \left[N G(x-x') G(x-x') - \frac{e_0 N \nabla^2}{e^2} \right]$$

$$G(x) = \frac{T}{V} \sum G(k) e^{ikx}$$

$$\phi(x) = \frac{1}{\sqrt{\beta V}} \sum \phi(k) e^{ikx}$$

$$\phi(x) G(x-x') G(x'-x) \phi(x')$$

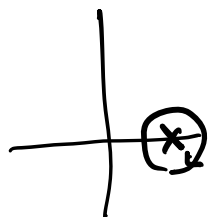


$$S_{\text{QUAN}} = \frac{N}{2} \sum \left[G(k+q) G(k) + \frac{e_0}{e^2} q^2 \right] |\phi_q|^2$$

$$-\frac{1}{N\beta V} \frac{\delta^2 S}{\delta \phi_q^\dagger \delta \phi_q}$$

$$= \tilde{\chi}(q) = -\frac{1}{V\beta} \sum G(k+q) G(k)$$

$$\tilde{\chi}(\vec{q}, i\nu_n) = -\frac{1}{V} \sum_{\vec{k}} T \sum \frac{1}{i\nu_n + i\nu_n - E_{\vec{k}+\vec{q}}} \frac{1}{i\nu_n - E_{\vec{k}}}$$

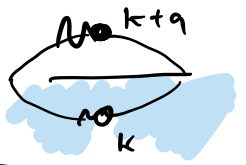


$$\otimes E_{\vec{k}+\vec{q}-i\nu_n}$$

$$= \frac{1}{V} \sum_{\vec{k}} \oint \frac{f(z) dz}{2\pi i} \frac{1}{z + i\nu_n - E_{\vec{k}+\vec{q}}} \frac{1}{z - E_{\vec{k}}}$$

poles on
imaginary axis

$$= -\frac{1}{V} \sum_{\vec{k}} \oint \frac{f(z) dz}{2\pi i} \frac{1}{z - (\epsilon_{k+q} - i\nu_n)} \frac{1}{z - \epsilon_k}$$



Poles at $z = \epsilon_k, \epsilon_{k+q} - i\nu_n$

$$= -\frac{1}{V} \sum_{\vec{k}} \left[\frac{f(\epsilon_{k+q} - i\nu_n)}{\epsilon_{k+q} - \epsilon_k - i\nu_n} + \frac{f(\epsilon_k)}{\epsilon_k - \epsilon_{k+q} + i\nu_n} \right]$$

$$\tilde{\chi}(\vec{q}, i\nu_n) = \frac{1}{V} \sum_{\vec{k}} \left[\frac{f_{k+q} - f_k}{i\nu_n - (\epsilon_{k+q} - \epsilon_k)} \right] = \int \frac{d^3k}{(2\pi)^3} \left(\frac{f_{k+q} - f_k}{i\nu_n - (\epsilon_{k+q} - \epsilon_k)} \right)$$

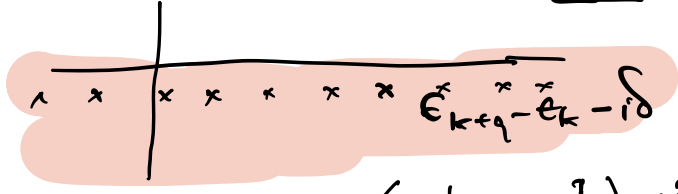
$\chi(q, z) = N \int \frac{d^3k}{(2\pi)^3} \frac{f_{k+q} - f_k}{z - (\epsilon_{k+q} - \epsilon_k)}$
 Lindhard Function "Charge susceptibility"



Retarded, response function.
 Particle-hole excitation. \underline{LRZ}

$$\tilde{\chi}_R(\vec{q}, \nu) = \chi(\vec{q}, z) |_{z = \nu + i\delta}$$

$$\chi_A = \chi(\vec{q}, z) |_{z = \nu - i\delta}$$



$$\delta g(q) = \chi(q) \phi(q), \quad \chi(x) = \int \frac{d^4q}{(2\pi)^4} \chi(q) e^{iqx} \quad \left(\frac{1}{\omega - \omega_q} - \frac{1}{\omega + \omega_q} \right) = \frac{2\omega_q}{\omega^2 - \omega_q^2}$$

$$\delta g(x, t) = \int d^4x' \chi(x-x') \phi(x'), \quad \delta \phi = \phi \delta^3(\vec{x})$$

$$\delta g(\vec{x}) = \int d^3x' \chi(\vec{x} - \vec{x}') \phi_0, \quad \omega = 0$$



Friedel oscillations of charge.