

ARPES: Uncovering the Superconducting Gap

Presentation for PHY601

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Outline

Introduction to ARPES

Seeing the superconducting gap

Conclusion

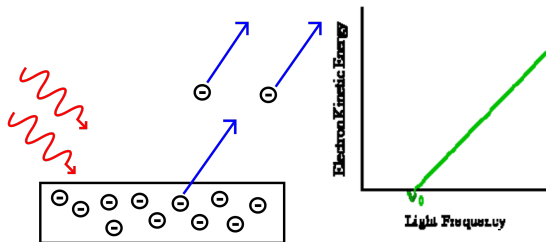
Photoelectric effect

Shining light on a metal expulses *photoelectrons*;

Can measure their kinetic energy. **Does not depend on light intensity I**

Rather, depends on frequency ν

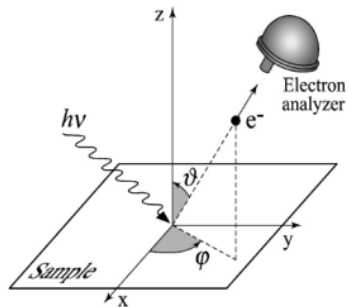
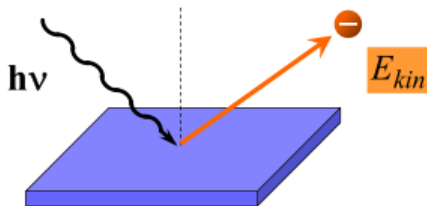
$$E_{kin} \propto \hbar\nu - const$$



Einstein comes in: **quantization of light**

The constant is actually the work function ϕ : **energy to delocalize electron from surface.**

$$E_{kin} = \hbar\nu - \phi$$



ARPES: theory

Angular Resolved Photo-Emission Spectroscopy

The gist of it:

Want to measure **energy of released electrons** E_k and **their momentum** \mathbf{k}

Use conservation laws and photoelectric effect to extract info

$$E_{kin} = \hbar\nu - \phi - |E_B| \quad (1)$$

$$\mathbf{p}_{\parallel} = \hbar\mathbf{k}_{\parallel} = \sqrt{2mE_{kin}} \sin \theta \quad (2)$$

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E_{kin} **measured** kinetic energy of outgoing electron

θ **measured** is the angle of emission with the surface

ν and ϕ are known

E_B **wanted** binding energy of electron in metal

\mathbf{k}_{\parallel} **wanted** crystal momentum

If one has that, we can construct the dispersion of the electrons $E(k)$!

From Many-Body interpretation

Using ARPES, we measure the **actual** dispersion relation.

Interactions (e-e, e-ph, etc) change band dispersions and lifetimes (spread)

Measure the spectral intensity: $I(\mathbf{k}, \omega)$

Spectral intensity $I(\mathbf{k}, \omega) \propto f(\omega)A(\mathbf{k}, \omega)$

1p spectral func.
$$A(\mathbf{k}, \omega) = -\frac{1}{\pi} \frac{\Sigma''(\mathbf{k}, \omega)}{[\omega - \epsilon_k - \Sigma'(\mathbf{k}, \omega)]^2 + [\Sigma''(\mathbf{k}, \omega)]^2}$$

bare band ϵ_k

Self-energy $\Sigma(\mathbf{k}, \omega) = \Sigma'(\mathbf{k}, \omega) + i\Sigma''(\mathbf{k}, \omega)$

= Band position + Linewidth/lifetime of QP

Free electron v.s. Fermi Liquid

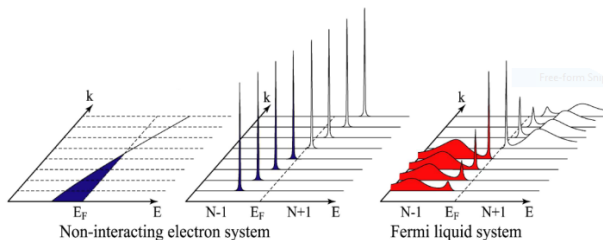
ARPES can see difference between a **non-interacting** electron system and a **Fermi Liquid** system.

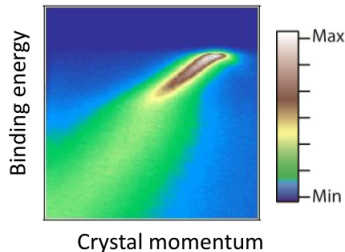
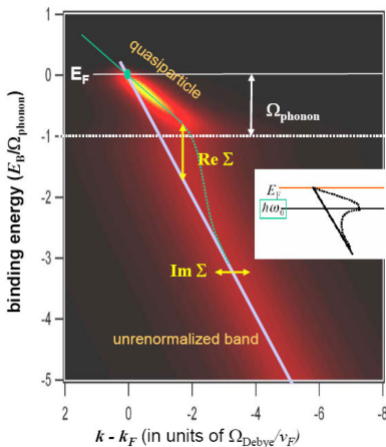
In NI $A(\mathbf{k}, \omega) = \delta(\omega - \epsilon_k)$

Extremely sharp \rightarrow Infinite lifetime of QP

In FL $A(\mathbf{k}, \omega) = Z_k \frac{\Gamma_k/\pi}{(\omega - \epsilon_k)^2 + \Gamma_k^2} + A_{incoherent}$

QP peak has a width: finite lifetime $\tau_k = 1/\Gamma_k$





(Left) Theoretical band for Electron Phonon coupling, see the quasiparticle peak, which has a Lorentzian lifetime. **(Right)** Example of *observed* Arpes intensity for Bi2201

Experimental Considerations

Need a very clean surface (atomically flat). Hence **surface-sensitive probe**, not good on bulk! (probe $\sim 2 - 20\text{\AA}$ in depth)

Need ultra-high vacuum (avoid surface deterioration)

Does not work under **pressure** or **magnetic field**.

However, very good at:

Comparison to theory

High resolution in energy AND momentum

Bad surface

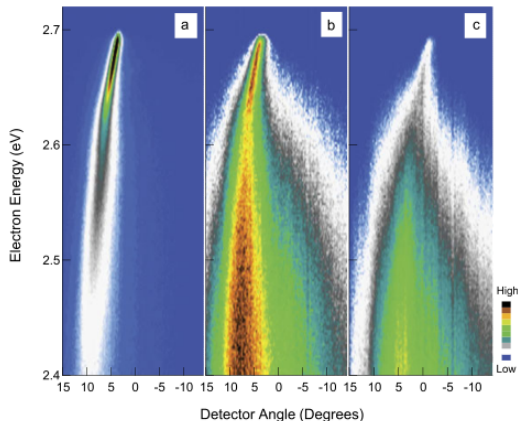
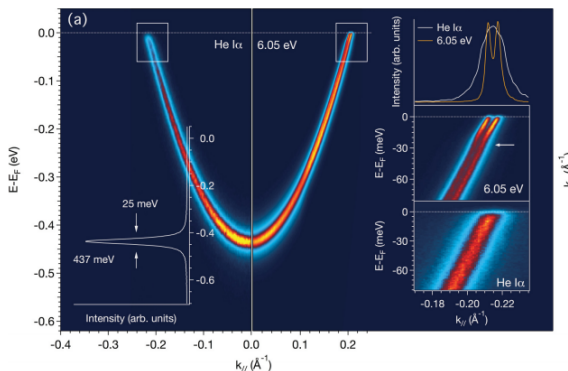


Figure: Experiments on optimally doped Bi₂21₂ (a) Dispersion right after cleaving. (b) After 1h in pure nitrogen. (c) After 1h in air.

A clean example: Cu



We see the $\epsilon_k = \frac{k^2}{2m^*}$ the free dispersion of Copper. The splitting of the bands can even be observed, due to Rashba coupling (spin-momentum locking, small but non-zero in Cu[111]).

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Superconductivity Essentials

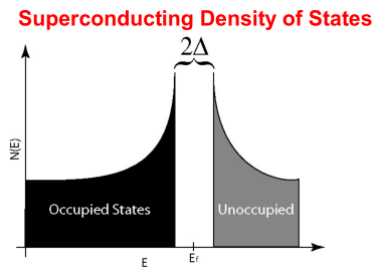
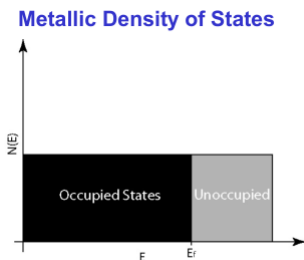
- Cooper instability: small attractive interaction binds electrons together $|\mathbf{k} \uparrow, -\mathbf{k} \downarrow\rangle$
- In BCS theory, superconductivity is result of condensation of the Cooper pairs.
- Conventional SC: attraction is due to retarded electron-phonon interaction. Isotropic interaction.

$$H = \sum_{\mathbf{k}} \xi_{\mathbf{k}} c_{\mathbf{k}\uparrow}^{\dagger} c_{\mathbf{k}\uparrow} + \xi_{\mathbf{k}} c_{\mathbf{k}\downarrow}^{\dagger} c_{\mathbf{k}\downarrow} + \Delta c_{\mathbf{k}\uparrow}^{\dagger} c_{-\mathbf{k}\downarrow}^{\dagger} + \Delta c_{-\mathbf{k}\downarrow} c_{\mathbf{k}\uparrow} \quad (3)$$

$$\xi_{\mathbf{k}} = \frac{\hbar^2 \mathbf{k}^2}{2m} - \mu \quad (4)$$

What is the gap

- Because of the pairing state, there is an **energy gap** for single particle excitation.
- In a superconductor, there is a one-particle gap but no two-particle gap.
- The two-particle excitations (Cooper pairs) are coherent and transport current without resistance.

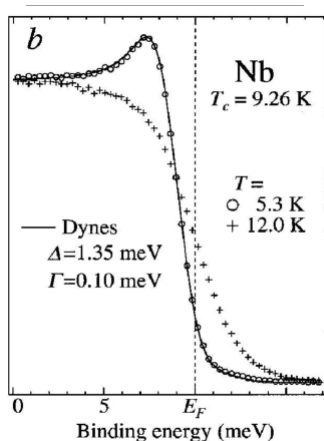


Example: Niobium

We can probe the DOS close to the Fermi Surface for Niobium ($T_c = 9.26\text{K}$)

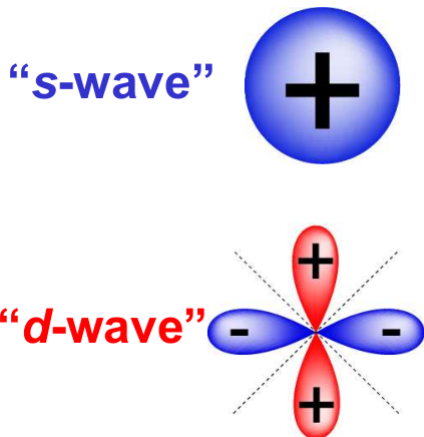
For $T > T_c$, no peak below E_F , normal Fermi distribution.

For $T < T_c$, gap opens, superconductivity sets in.



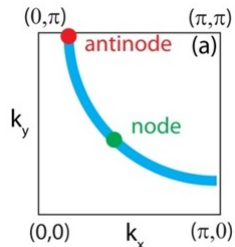
S-Wave vs D-Wave

- In conventional SC, attractive interaction due to e-ph coupling: isotropic. Leads to **s-wave** pairing (gap positive all around FS).
- In Cuprates, a **d-wave** pairing was advanced to explain how they could have SC behavior.
- D-wave \rightarrow has nodes where $\Delta(\mathbf{k}) = 0$ on the Fermi Surface



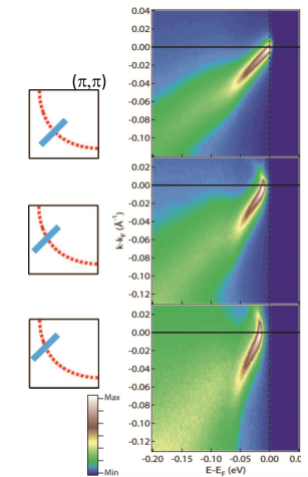
Bi-2212: Fermi Surface

- Bi2212: $Bi_2SR_2CaCu_2O_{8+\delta}$, $T_c^{max} = 96K$.
- Planes of CuO with “stuff” in between. Fermi surface is similar to YBCO (as in homework)

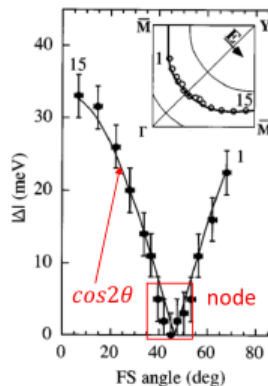
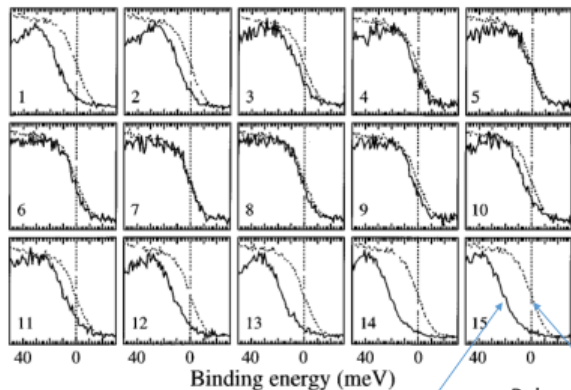


Bi-2212: Gap across FS

- Using ARPES, the excitations near the Fermi-Surface can be probed very accurately.
- Doing different cuts in k -space, we can probe different parts of the Fermi-Surface.



A clear node



E_F from polycrystalline metal. Gap is hard to read, but is there!

What to take from this

Exceptional resolution of the electronic structure (**Bands & Fermi Surface**);

Can *explicitly* probe the gap;

Answered many questions about the nature of the electronic excitations in Cuprates (ex: gap symmetry + deviations);

In Cuprates, has been very important to ascertain the presence of the pseudogap phase (gap but no SC);

To be used even more: Need materials with better surface cleaving (currently being applied to pnictides Fe-based SC)

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References

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Thanks for listening!

Questions?