

# Physics 616

## Note on C-Z equation for correlation functions

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(sign of  $\gamma$  fixed 2008.04.21)

$Z$  is the generating functional for the renormalized correlation functions:

$$Z = \langle e^{-\int j_r \phi_r} \rangle \quad (1)$$

Where  $j_r(x)$  is the naively dimensionless source and  $\phi_r(x)$  is the naively dimensionless field:

$$j_r(x) = \Lambda^{-(d+2)/2} j(x) \quad (2)$$

$$\phi_r(x) = \Lambda^{-(d-2)/2} \phi(x) \quad (3)$$

$$\int j_r \phi_r = \int d^d x \Lambda^d j_r(x) \phi_r(x) \quad (4)$$

The RG equation is

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i - \left( \frac{d+2}{2} - \gamma \right) j_r \frac{\partial}{\partial j_r} \right] Z(\Lambda, \lambda_r, j_r) = 0 \quad (5)$$

The  $\lambda_r^i$  are the macroscopic dimensionless coupling constants (besides the source), e.g.  $\lambda$  and  $m_r^2 = \Lambda^{-2} m^2$  in the  $\phi^4$  theory.

Expand in  $j_r$ . Schematically:

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i - \left( \frac{d+2}{2} - \gamma \right) j_r \frac{\partial}{\partial j_r} \right] \left( \int j_r \phi_r \right)^n = 0 \quad (6)$$

This holds for arbitrary sources  $j_r(x)$ , so we have the C-Z equation

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i + nd - \left( \frac{d+2}{2} - \gamma \right) \right] \langle \phi_r(x_1) \cdots \phi_r(x_n) \rangle = 0 \quad (7)$$

To get the C-Z equation for the Fourier transformed dimensionless correlation functions, write

$$\int j_r \phi_r = \int d^d x \Lambda^d j_r(x) \phi_r(x) \quad (8)$$

$$= \int d^d x \Lambda^d j_r(x) \int \frac{d^d p}{(2\pi)^d} \Lambda^{-d} e^{ipx} \tilde{\phi}_r(p) \quad (9)$$

$$= \iint \frac{d^d x d^d p e^{ipx}}{(2\pi)^d} j_r(x) \tilde{\phi}_r(p) \quad (10)$$

Note that the scale  $\Lambda$  does not appear in the coupling between  $j_r(x)$  and  $\tilde{\phi}_r(p)$ . So

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i - \left( \frac{d+2}{2} - \gamma \right) j_r \frac{\partial}{\partial j_r} \right] \left( \iint j_r \tilde{\phi}_r \right)^n = 0 \quad (11)$$

gives

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i - n \left( \frac{d+2}{2} - \gamma \right) \right] \langle \tilde{\phi}_r(p_1) \cdots \tilde{\phi}_r(p_n) \rangle = 0 \quad (12)$$

This is the form of the C-Z equation I find easiest to use: the rg equation on the dimensionless correlation functions.

Derive the dimensionless 2-point function of the free massless scalar from the C-Z equation:

$$\langle \tilde{\phi}_r(p_1) \tilde{\phi}_r(p_2) \rangle = \delta^d(\Lambda^{-1} \sum_i p_i) g_r(\Lambda^{-2} p_1^2) = \Lambda^d \delta^d(\sum_i p_i) g_r(\Lambda^{-2} p_1^2) \quad (13)$$

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} - 2 \left( \frac{d+2}{2} \right) \right] \Lambda^d \delta^d(\sum_i p_i) g_r(\Lambda^{-2} p_1^2) = 0 \quad (14)$$

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} - 2 \right] g_r(\Lambda^{-2} p_1^2) = 0 \quad (15)$$

so

$$g_r(\Lambda^{-2} p_1^2) = (\text{const}) \frac{1}{\Lambda^{-2} p_1^2} \quad (16)$$

which is correct.

In general we write

$$\langle \tilde{\phi}_r(p_1) \cdots \tilde{\phi}_r(p_n) \rangle = \delta^d \left( \sum_i \Lambda^{-1} p_i \right) G_r(p_1, \dots, p_n) \quad (17)$$

and we have the C-Z equation in the form

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i + d - n \left( \frac{d+2}{2} - \gamma \right) \right] G_r(p_1, \dots, p_n) = 0 \quad (18)$$

To go to the dimensionful fields:

$$\tilde{\phi}_r(p) = \int d^d x \Lambda^d e^{ipx} \phi_r(x) \quad (19)$$

$$= \int d^d x \Lambda^d e^{ipx} \Lambda^{-(d-2)/2} \phi(x) \quad (20)$$

$$= \Lambda^{(d+2)/2} \tilde{\phi}(p) \quad (21)$$

so

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i + n \left( \frac{d+2}{2} \right) - n \left( \frac{d+2}{2} - \gamma \right) \right] \langle \tilde{\phi}(p_1) \cdots \tilde{\phi}(p_n) \rangle = 0 \quad (22)$$

or

$$\left[ \Lambda \frac{\partial}{\partial \Lambda} + \beta^i \partial_i + n\gamma \right] \langle \tilde{\phi}(p_1) \cdots \tilde{\phi}(p_n) \rangle = 0 \quad (23)$$

We can easily use this form of the C-Z equation to derive, for example, the massless 2-point function. But I find the derivation using the dimensionless correlation functions more transparent.