

# BPS STATES, HITCHIN SYSTEMS, AND THE WKB APPROXIMATION

WORK IN PROGRESS WITH

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BONN WORKSHOP ON MIRROR SYMMETRY

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BASED ON

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& 0906.????

# OUTLINE

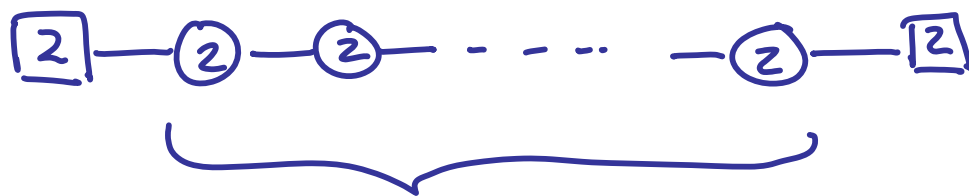
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# 1. INTRODUCTION

- IN THIS TALK WE STUDY "BPS DEGENERACIES" OR "BPS INDICES," DENOTED  $\Omega$
- MATHEMATICALLY THEY ARE RELATED TO DONALDSON-THOMAS INVTS.
- IN THE PAST 3 YEARS THERE HAS BEEN MUCH PROGRESS ON "WALL-CROSSING" FOR THESE INVTS.
- W.C.F. TELL US HOW  $\Omega$  CHANGES AS FUNCTIONS OF MODULI
- THAT STILL LEAVES OPEN THE PROBLEM OF COMPUTING  $\Omega$ .

- IN THIS TALK I WILL DESCRIBE A NEW METHOD TO COMPUTE  $\Omega$  IN A CLASS OF  $N=2, D=4$  FIELD THEORIES.

- THESE ARE "LINEAR  $SU(2)$  QUIVER GAUGE THEORIES"



$G = SU(2)^r$  GAUGE GROUP

- SIMILAR METHODS PROBABLY EXTEND TO A MUCH LARGER CLASS OF FIELD THEORIES.

## 2. $\mathcal{N} = 2, \mathcal{D} = 4$ FIELD THEORY

- $\exists$  "MODULI SPACE OF VACUA"  
COULOMB BRANCH  $\mathcal{B} \cong \mathbb{C}^r$   
WITH SPECIAL KÄHLER METRIC
- IR PHYSICS :  $U(1)^r$  GAUGE THEORY
- CHARGE LATTICE FORMS A  
LOCAL SYSTEM OVER  $\mathcal{B}$  WITH  
MONODROMY AROUND  $\mathcal{B}_{\text{sing}}$ .

$$0 \rightarrow \Gamma_{\text{flav.}} \rightarrow \Gamma \rightarrow \bar{\Gamma} \rightarrow 0$$

FIBER AT  $u \in \mathcal{B}$ :

$\Gamma_u$  "POISSON": INTEGRAL ANTISYMMETRIC  $\langle \cdot, \cdot \rangle$

$\bar{\Gamma}_u$  SYMPLECTIC, RANK =  $2r$

- $Z \in \text{Hom}(\Gamma, \mathbb{C})$  CENTRAL CHARGE

# SEIBERG-WITTEN THEORY & RIEMANN SURFACES

- THE SPECIAL KÄHLER GEOMETRY IS PRESENTED USING A FAMILY OF NON-COMPACT RIEMANN SURFACES

$$\Sigma \rightarrow \mathcal{B}$$

EQUIPPED WITH A FIBREWISE MEROMORPHIC DIFFERENTIAL  $\lambda$

- $\Gamma$  IS (A SUBQUOTIENT OF)  $H_1(\Sigma, \mathbb{Z})$

- $Z_\gamma(u) = \int_\gamma \lambda$

# BPS STATES & BPS INDEX

- $\mathcal{H}_{\gamma, u}^{\text{BPS}} := \{ \psi \mid E = |Z_{\gamma}(u)| \}$

$$\Omega(\gamma; u) = \text{Tr}_{\mathcal{H}} (-1)^F := \text{"BPS INDEX"}$$

PROBLEM: COMPUTE THE  $\Omega(\gamma; u)$

## KEY INGREDIENT:

- BPS STATES CAN FORM BPS BOUNDSTATES.
- A BOUNDSTATE WITH CONSTITUENT CHARGES  $\gamma_1, \gamma_2$  CAN ONLY DECAY WHEN  $u$  CROSSES

$$MS(\gamma_1, \gamma_2) := \{ u \mid Z_{\gamma_1}(u) / Z_{\gamma_2}(u) \in \mathbb{R}_+ \}$$

$\Omega(\gamma; u)$  ONLY PIECEWISE CONSTANT

$\Rightarrow$  WALL-CROSSING FORMULAE

### 3, THE KONTSEVICH-SOIBELMAN FORMULA

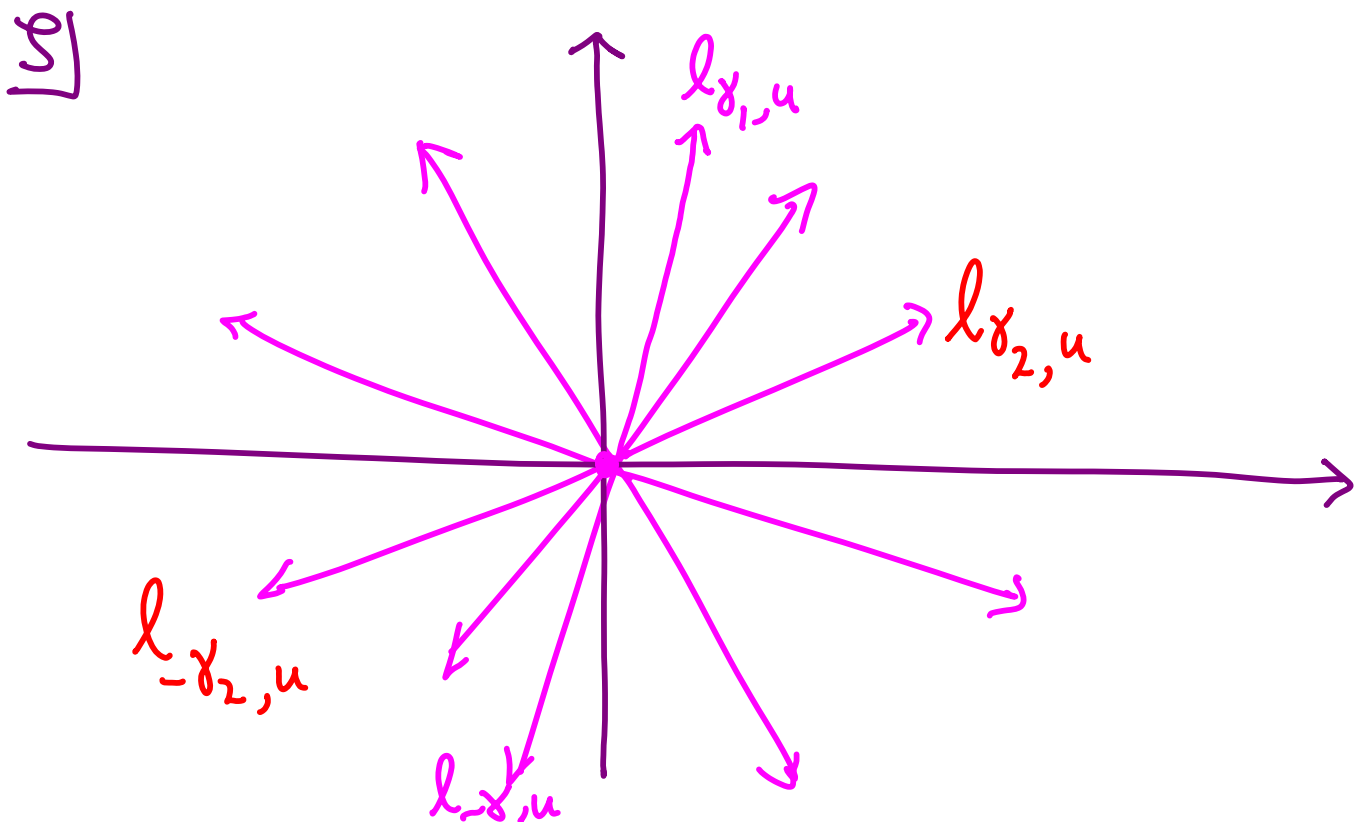
DATA:

1.  $\Gamma \rightarrow \mathcal{B}$ , POISSON, WITH Q.R.  $\sigma$
2. CENTRAL CHARGE FUNCTION  $Z \in \text{Hom}(\Gamma, \mathbb{C})$
3. PIECEWISE CONSTANT  $\Omega: \Gamma \rightarrow \mathbb{Z}$

FIRST INGREDIENT: BPS RAYS:

FOR  $u \in \mathcal{B}$ ,  $\gamma \in \Gamma_u$

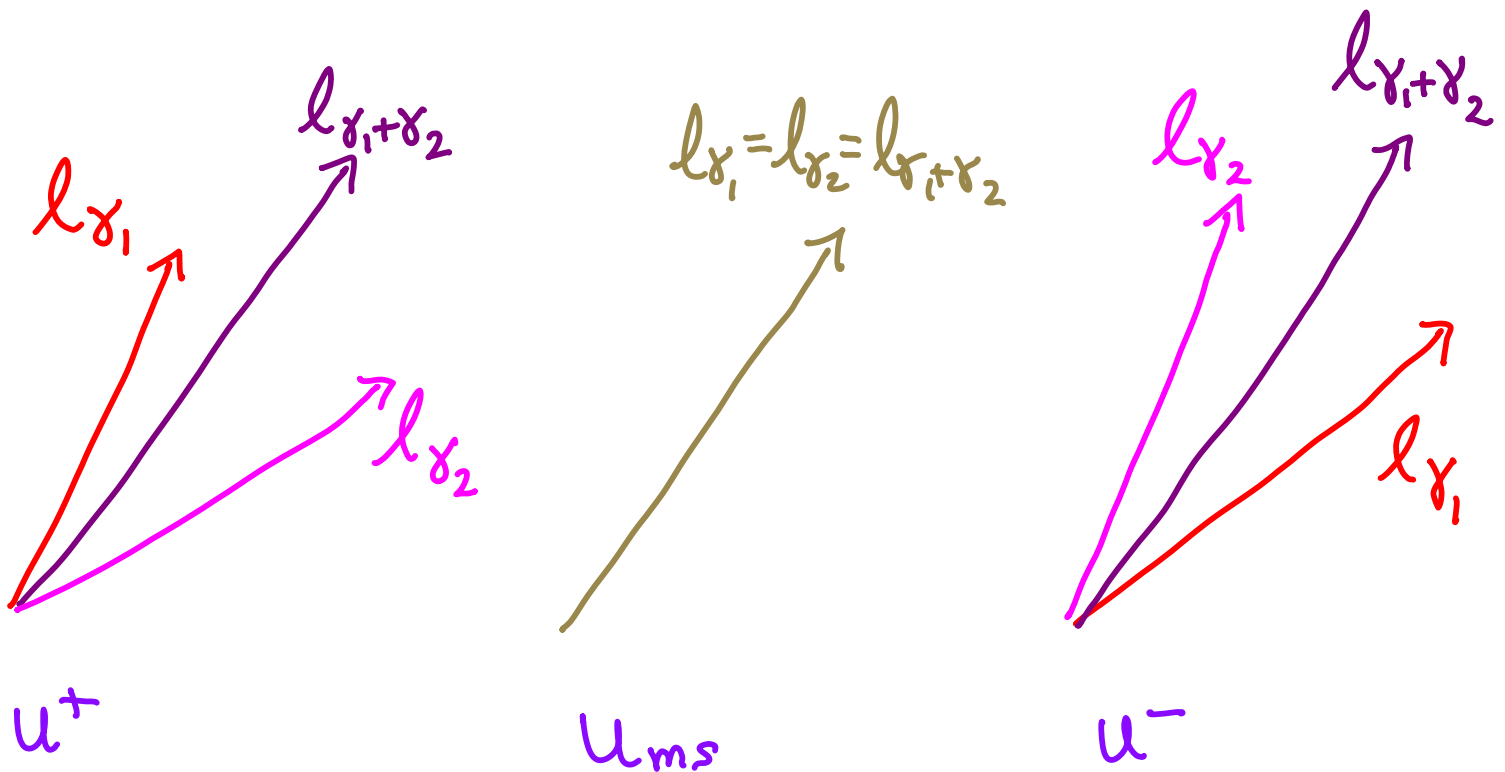
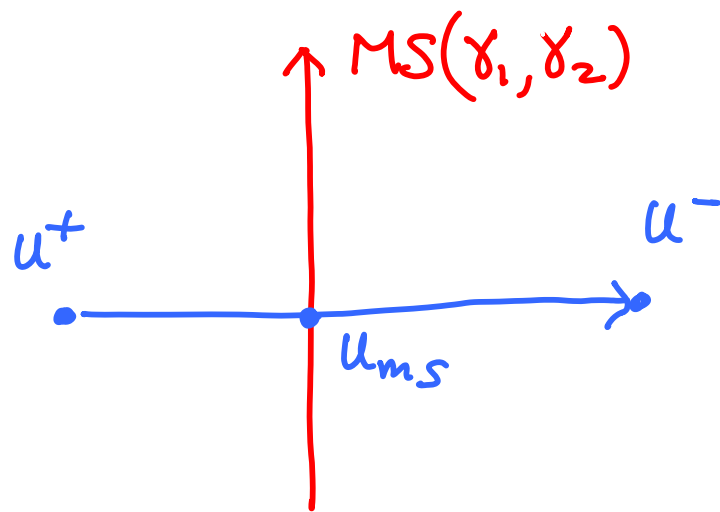
$$l_{\gamma, u} := Z_{\gamma}(u) \mathbb{R}_- = \left\{ \mathcal{J} \mid \mathcal{J} / Z_{\gamma}(u) \in \mathbb{R}_- \right\}$$





AS  $u$  VARIES THE SLOPES OF  
THE BPS RAYS VARY

AS  $u$  CROSSES A WALL  $MS(\gamma_1, \gamma_2)$   
BPS RAYS WILL COALESCE



## SECOND INGREDIENT: LOCAL SYSTEM OF POISSON TORI

- INTRODUCE  $\Gamma := \Gamma^* \otimes_{\mathbb{Z}} \mathbb{C}^*$

$$\Gamma_u \cong \mathbb{C}^* \times \dots \times \mathbb{C}^*$$

$\gamma \in \Gamma_u \Rightarrow$  FUNCTION  $X_\gamma : \Gamma_u \rightarrow \mathbb{C}^*$

"HOLOMORPHIC FOURIER MODES"

- FIBREWISE POISSON STRUCTURE

$$\{X_{\gamma_1}, X_{\gamma_2}\} = \langle \gamma_1, \gamma_2 \rangle X_{\gamma_1} X_{\gamma_2}$$

- FOR EACH  $\gamma \in \Gamma$  DEFINE A  
(POISSON) MORPHISM:

$$K_\gamma : X_{\gamma'} \rightarrow X_{\gamma'} \left( 1 - \sigma(\gamma) X_\gamma \right)^{\langle \gamma', \gamma \rangle}$$

ABOUT THE SIGN :

$K \ni S$  INTRODUCE A LIE ALGEBRA

$$[e_{\gamma_1}, e_{\gamma_2}] = (-1)^{\langle \gamma_1, \gamma_2 \rangle} \langle \gamma_1, \gamma_2 \rangle e_{\gamma_1 + \gamma_2}$$

DEFINE A GROUP ELEMENT

$$U_{\gamma} := \exp\left(\sum_{n=1}^{\infty} \frac{e_{n\gamma}}{n^2}\right)$$

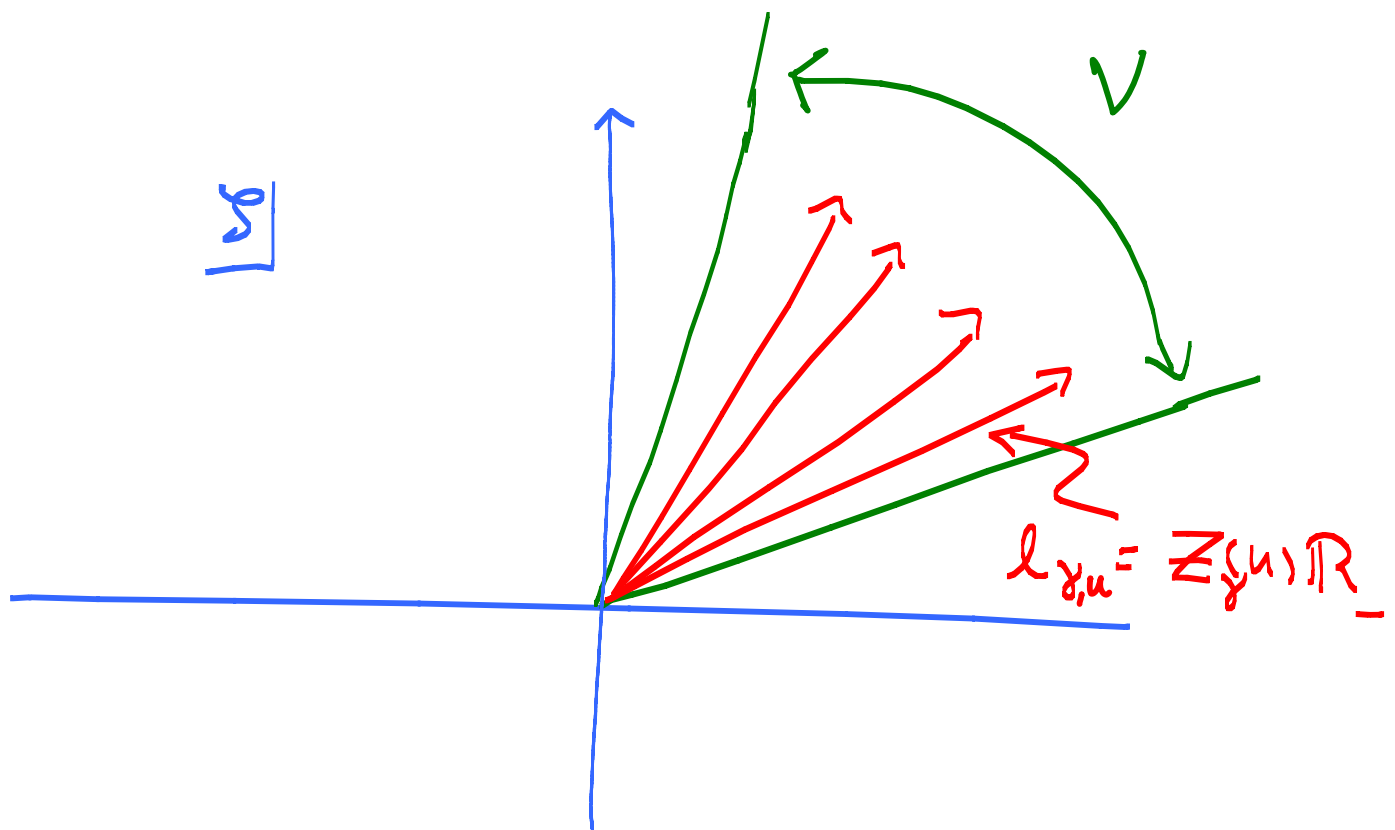
AND WORK WITH  $U_{\gamma}$  INSTEAD OF  $K_{\gamma}$

CHOOSE A QUADRATIC REFINEMENT

$$\frac{\sigma(\gamma_1 + \gamma_2)}{\sigma(\gamma_1)\sigma(\gamma_2)} = (-1)^{\langle \gamma_1, \gamma_2 \rangle}$$

$\sigma(\gamma)e_{\gamma}$  GENERATE THE LIE ALGEBRA  
OF SYMPLECTIC VECTOR FIELDS.

THIRD INGREDIENT: CONVEX CONE  $V$



AND THEN DEFINE:

$$A_V := \prod_{\substack{\rightarrow \\ -z_{\gamma} \in V}} K_{\gamma} \Omega(\gamma; u)$$

THE PRODUCT IS TAKEN OVER  
THE RAYS IN THE CLOCKWISE  
ORDER (DECREASING SLOPE)

$$A_V := \prod_{\substack{\rightarrow \\ -z_\gamma \in V}} K_\gamma \Omega(\gamma; u)$$

$A_V$  DEPENDS ON  $u$  IN TWO WAYS

1. THE ORDERING OF FACTORS  
DEPENDS ON  $u$

2. THE  $\Omega(\gamma; u)$  DEPEND ON  $u$  ...

DEFINITION: " $\Omega$  SATISFIES THE KSWCF"

IF:

$A_V$  IS CONSTANT IN  $u$  AS LONG

AS NO BPS RAY ENTERS OR  
LEAVES THE SECTOR  $\tilde{V}$ .

## 4. HK MANIFOLDS + TWISTOR COORDS

I WILL NOW BRIEFLY REVIEW PAPER I.  
(FOR SIMPLICITY PUT  $\Gamma_{\text{flavor}} = 0$ )

"THEOREM" GIVEN  $\Omega$  SATISFYING  
THE KS WCF ONE CAN  
CONSTRUCT A "NICE" FAMILY OF  
HK. METRICS ON:

$$\bar{\mathcal{M}} = \bar{\Gamma}^* \otimes \mathbb{R}/\mathbb{Z}$$

→ SKETCH OF THE CONSTRUCTION ←

• COMPACTIFYING  $N=2, d=4$  THEORY ON  
 $\mathbb{R}^3 \times S^1_{\mathbb{R}} \Rightarrow$  LOW ENERGY  $\sigma$ -MODEL

$$\varphi: \mathbb{R}^3 \rightarrow \bar{\mathcal{M}}$$

$\Rightarrow \bar{\mathcal{M}}$  HAS A "NICE" HK METRIC.

- $\bar{\mathcal{M}}$  HAS A CANONICAL FAMILY OF METRICS - THE "SEMI-FLAT METRICS" - BUT THESE ARE SINGULAR.

IN PHYSICS "QUANTUM CORRECTIONS BY BPS INSTANTONS" SMOOTHS OUT THESE SINGULARITIES.

- THE SEMI-FLAT METRIC IS DESCRIBED BY TWISTOR COORDINATES:

$$\chi_\gamma: \bar{\mathcal{M}} \times \mathbb{C}^* \rightarrow \mathbb{C} \quad \gamma \in \Gamma$$

$$\theta_\gamma: \Gamma^* \otimes \mathbb{R}/\mathbb{Z} \rightarrow \mathbb{R}/\mathbb{Z}$$

$$\chi_\gamma^{sf} := \exp \left[ + \frac{\pi R}{\zeta} z_\gamma + i \theta_\gamma + \pi R S \bar{z}_\gamma \right]$$

(NEITZKE  $\frac{1}{\zeta}$  PIOLINE)

$\zeta \in \mathbb{C}^* \subseteq$  TWISTOR SPHERE

USING THESE WE CONSTRUCT

A FAMILY OF HOLO. MAPS

$$\zeta \in \mathbb{C}^*, \quad \mathcal{M}^\zeta \xrightarrow{\chi(\zeta)^{s.f.}} \overline{T} := \overline{F}^* \otimes \mathbb{C}^*$$

$$\text{VIA } \chi_\zeta^{s.f.}(\cdot, \zeta) := \chi(\zeta)^{s.f.} * (\overline{X}_\zeta)$$

$\overline{T}$  HAS A FIBERWISE SYMPLECTIC STRUCTURE

$$\omega^{\overline{T}} = \frac{1}{2} \epsilon^{ij} \frac{dX_{\gamma_i}}{X_{\gamma_i}} \wedge \frac{dX_{\gamma_j}}{X_{\gamma_j}}$$

$$\epsilon_{ij} = \langle \gamma_i, \gamma_j \rangle$$

THEN  $\widetilde{\omega}_\zeta^{s.f.} := \chi(\zeta)^{s.f.} * \omega^{\overline{T}} =$  HOLO. SYMPLECTIC FORM

$$= \frac{1}{2i\zeta} \omega_+^{sf} + \omega_3^{sf} + \frac{\zeta}{2i} \omega_-^{sf}$$

$\Rightarrow$  SEMI-FLAT HK STRUCTURE.



- THE QUANTUM-CORRECTED H.K. STRUCTURE IS DEFINED USING A SIMILAR PROCEDURE:

FROM A SET OF  $\chi_\gamma : \mathcal{M} \times \mathbb{C}^* \rightarrow \mathbb{C}$

CONSTRUCT A MAP:  $\chi_\gamma(\cdot, \mathcal{J}) = \chi(\mathcal{J})^*(\underline{\chi}_\gamma)$

$$\mathcal{M}^{\mathcal{J}} \xrightarrow{\chi(\mathcal{J})} \overline{T} = \overline{T}^* \otimes \mathbb{C}^*$$

AND EXTRACT HK STRUCTURE FROM:

$$\begin{aligned} \overline{\omega}_{\mathcal{J}} &= \chi(\mathcal{J})^* \overline{\omega}^{\overline{T}} \\ &= \frac{1}{2i\mathcal{J}} \omega_+ + \omega_3 + \frac{\mathcal{J}}{2i} \omega_- \end{aligned}$$

- USING A PHYSICAL INTERPRETATION OF THE  $\chi_\gamma$  WE SHOW THEY SATISFY A SYSTEM OF ISOMONODROMIC DIFFL. EQS. ("WARD IDENTITIES") WHICH IS EQUIVALENT TO A R.H. PROBLEM

# DEFINING PROPERTIES OF $\chi_\gamma$

1.)  $\chi_\gamma(\cdot, \mathcal{J}) : \mathcal{M}^S \rightarrow \mathbb{C}$  HOLO.

2.)  $\chi_\gamma \chi_{\gamma'} = \chi_{\gamma+\gamma'}$

3.)  $\chi_\gamma(\mathcal{J}) = \overline{\chi_{-\gamma}(-1/\mathcal{J})}$

4A.)  $\chi_\gamma(\cdot, \mathcal{J}) \underset{R \rightarrow \infty}{\sim} \exp\left(\frac{\pi R}{\mathcal{J}} Z_\gamma + i\Theta_\gamma + \pi R \mathcal{J} \bar{Z}_\gamma\right)$

4B.)  $\lim_{\mathcal{J} \rightarrow 0} \chi_\gamma(\cdot, \mathcal{J}) e^{-\frac{\pi R}{\mathcal{J}} Z_\gamma(u)}$  FINITE

5.)  $\chi_\gamma(\cdot, \mathcal{J})$  IS PIECEWISE HOLO IN  $\mathcal{J}$ :

HOLO. IN ANGULAR SECTORS BETWEEN

BPS RAYS AND TRANSFORMS BY

$\prod_{\gamma' || \gamma} K_{\gamma'}^{\Omega(\gamma; u)}$  AS  $\mathcal{J}$  CROSSES THE  
BPS RAY  $l_{\gamma, u}$ .

- RIEMANN-HILBERT PROBLEM IS SOLVED BY:

$$\chi_\gamma(s) = \chi_\gamma^{sf}(s) \cdot \exp \left\{ -\frac{1}{4\pi i} \sum_{\gamma' \in \Gamma} \Omega(\gamma'; u) \langle \gamma, \gamma' \rangle \right. \\ \left. \cdot \int_{\mathfrak{h}_{\gamma', u}} \frac{ds'}{s'} \frac{s'+s}{s'-s} \log(1 - \sigma_{\gamma'} \chi_{\gamma'}(s')) \right\}$$

- CAN BE SOLVED ITERATIVELY AT  $R \gg 1$ .
- ONE GOAL OF THIS TALK IS TO GIVE A NEW CONSTRUCTION OF THE  $\chi_\gamma$ 'S.
- AN IMPORTANT COROLLARY OF THE NEW CONSTRUCTION IS AN ALGORITHM FOR COMPUTING THE  $\Omega$ 'S.

## 5. FROM M5-BRANES TO HITCHIN EQS

NOW LETS NARROW DOWN THE CLASS OF  $N=2, d=4$  THEORIES

- IIB STRING VIEWPOINT: WE HAVE A NON-CPT. CY = CURVE C OF ADE SINGULARITIES.

BPS STATES  $\sim$  D3'S WRAP SLAGS  $S^3$  OR  $S^2 \times S^1$ .

WE WANT  $\Omega$ 'S FOR THESE.

- A CHAIN OF PHYSICAL REASONING LEADS TO THE CONCLUSION THAT THE  $\Omega$ 'S CAN BE COMPUTED BY STUDYING HITCHIN SYSTEMS ON C WITH SINGULARITIES.

... VERY ROUGHLY:

a.) ADE SINGULARITY  $\xrightarrow{\text{"T-DUAL"}}$  5-BRANE

b.) DECOUPLE GRAVITY  $\Rightarrow$

CONSIDER  $A_{k-1}$  (2,0) THEORY  
ON  $\mathbb{R}^{1,3} \times \mathbb{C}$  WITH "DEFECTS."

c.) LOW ENERGY LIMIT GIVES  $N=2, D=4$   
THEORY ON  $\mathbb{R}^{1,3}$

d.) HK STRUCTURE FROM FURTHER  
COMPACTIFICATION OF THIS  $N=2, D=4$   
THEORY ON  $\mathbb{R}^{1,2} \times S^1_R$

BUT! WE COULD ARRIVE AT THE  
SAME LOW ENERGY THEORY ON  
 $\mathbb{R}^{1,2}$  COMPACTIFYING IN THE OTHER  
ORDER:

6D (2,0)  $A_{K-1} / \mathbb{R}^{1,2} \times S^1_R \times C$  + DEFECTS

$$l_C \ll l_{S^1}, l_{\mathbb{R}^{1,2}}$$

$$l_{S^1} \ll l_C, l_{\mathbb{R}^{1,2}}$$

5D  $SU(K)$  SYM  

---

 $\mathbb{R}^{1,2} \times C$

4D  $N=2$  GAUGE THEORY /  $\mathbb{R}^{1,2} \times S^1_R$

$$l_{S^1} \ll l_{\mathbb{R}^{1,2}}$$

$$l_C \ll l_{\mathbb{R}^{1,2}}$$

$\sigma$ -MODEL:  $\mathbb{R}^{1,2} \longrightarrow \mathcal{M}$

## 6. THE HITCHIN SYSTEM

$$G = SU(k)$$

$E \rightarrow C$  HAS CONNECTION  $A$  AND

HIGGS FIELD  $\varphi \in \Omega^{0,1}(C; \text{End } E)$

BPS EQS = HITCHIN EQS:

$$\begin{cases} F + R^2[\varphi, \bar{\varphi}] = 0 \\ \bar{\partial}_A \varphi = 0 \quad \partial_A \bar{\varphi} = 0 \end{cases}$$

ALLOW SINGULARITIES AT  $z_i \in C$ :

REGULAR + IRREGULAR SING'S

RSP: IN LOCAL COORD'S NEAR  $z_i$

$$\varphi \sim \frac{1}{z} \rho \frac{dz}{z - z_i} + \dots$$

$$A \sim \frac{\alpha}{2i} \left( \frac{dz}{z - z_i} - \frac{d\bar{z}}{\bar{z} - \bar{z}_i} \right) + \dots$$

SO NOW WE IDENTIFY:

$\mathcal{M}$  = HITCHIN MODULI SPACE

HITCHIN FIBRATION:  $(A, \varphi) \rightarrow \text{Char Poly}(\varphi)$

$$\mathcal{B} = \bigoplus_d H^0(C, K_C^{\otimes d})$$

$\mathcal{B}$ : BASE OF A FAMILY OF SPECTRAL CURVES

$$\Sigma = \{(x, z) \mid \det(xdz - \varphi(z)) = 0\} \subset T^*C$$

= SEIBERG-WITTEN CURVE OF THE

$d=4, N=2$  THEORY RESULTING FROM

COMPACTIFYING THE  $A_{k-1}$  (2.0) THEORY ON  $C$

$\lambda = xdz =$  S-W DIFFERENTIAL

(  $d\lambda = dx dz =$  CANONICAL SYMPLECTIC FORM  
ON  $T^*C$  )



FOR SIMPLICITY WE HENCEFORTH

TAKE GAUGE GROUP  $G = \text{SU}(2)$

$$1. \quad \Sigma \xrightarrow{2:1} \mathbb{C}$$

$$\lambda^2 = \text{QUADRATIC DIFFL ON } \mathbb{C} = \frac{1}{2} \text{Tr } \varphi^2$$

$$\text{LOCALLY: } \varphi \sim \begin{pmatrix} \lambda & \\ & -\lambda \end{pmatrix}$$

WE ASSUME THAT  $\lambda^2$  HAS  
FIRST ORDER ZEROS  $w_a \in \mathbb{C}$ .

THESE ARE THE BRANCH POINTS  
OF  $\Sigma \rightarrow \mathbb{CP}^1$

ALSO CALLED "TURNING POINTS"

2. IN ADDITION THERE ARE SINGULAR POINTS  $z_i$ :

$$\varphi = \frac{1}{z - z_i} \begin{pmatrix} m_i \\ -m_i \end{pmatrix} + \dots$$

$$A = \begin{pmatrix} m_i^{(3)} \\ -m_i^{(3)} \end{pmatrix} \left( \frac{dz}{z - z_i} - \frac{d\bar{z}}{\bar{z} - \bar{z}_i} \right) + \dots$$

# BPS STATES IN THE HITCHIN FRAMEWORK

THE PHYSICAL DERIVATION LEADS TO THE FOLLOWING RULES FOR DESCRIBING BPS STATES.

(KLEMM, LERCHE, VAFA, WARNER)

$\Omega(\gamma, u) = 0$  UNLESS THERE EXISTS A CURVE  $\tilde{c} \in \Sigma$  IN HOMOLOGY CLASS  $\gamma$  SO THAT  $\pi_*(\tilde{c}) = c$  IS A CURVE IN  $C$  S.T.

$$\langle \lambda, \partial_t \rangle \in e^{i\vartheta_*} = \text{CONSTANT}$$

AND

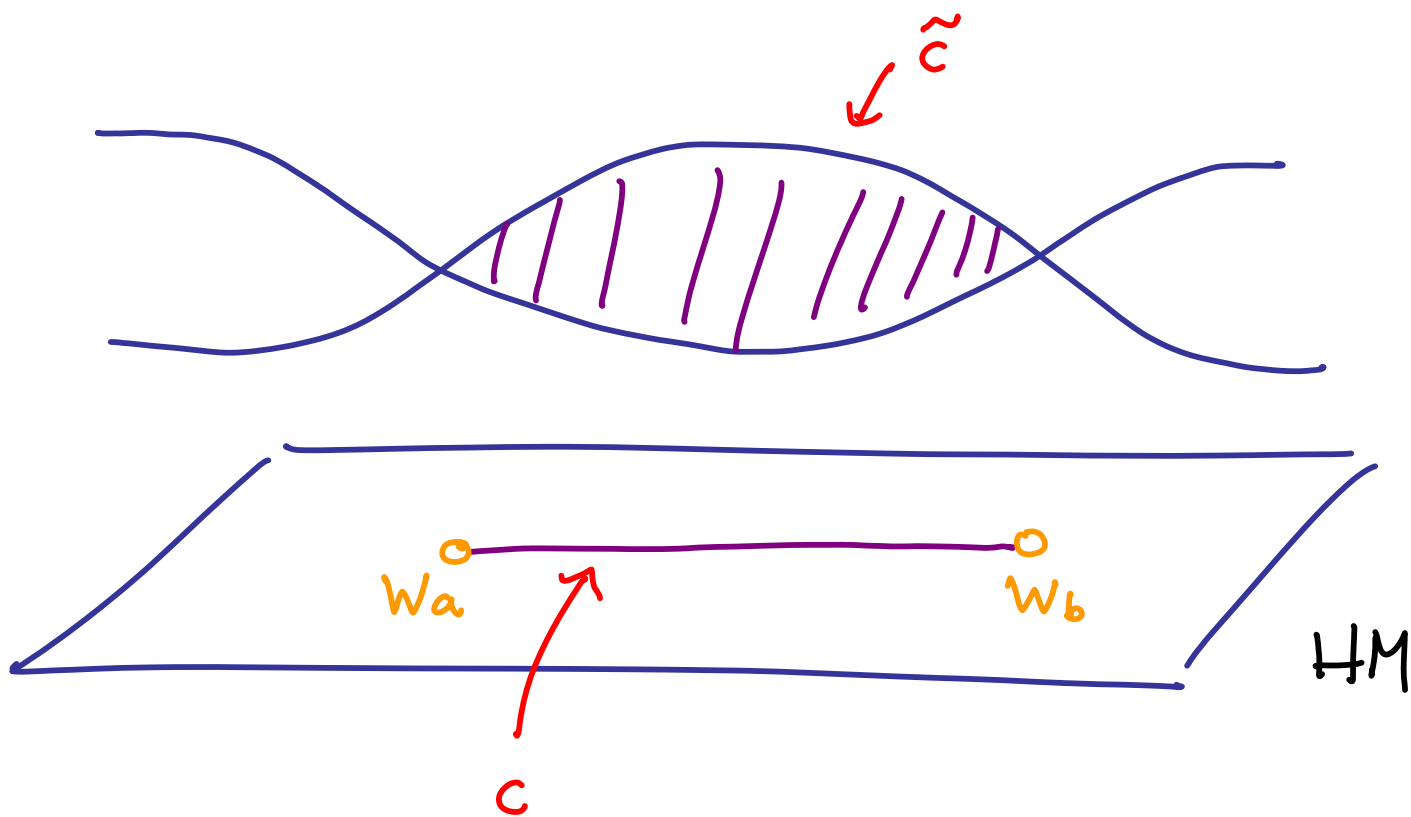
ONE OF TWO THINGS HAPPENS :

1. EITHER,

C BEGINS AND ENDS ON A BRANCH POINT

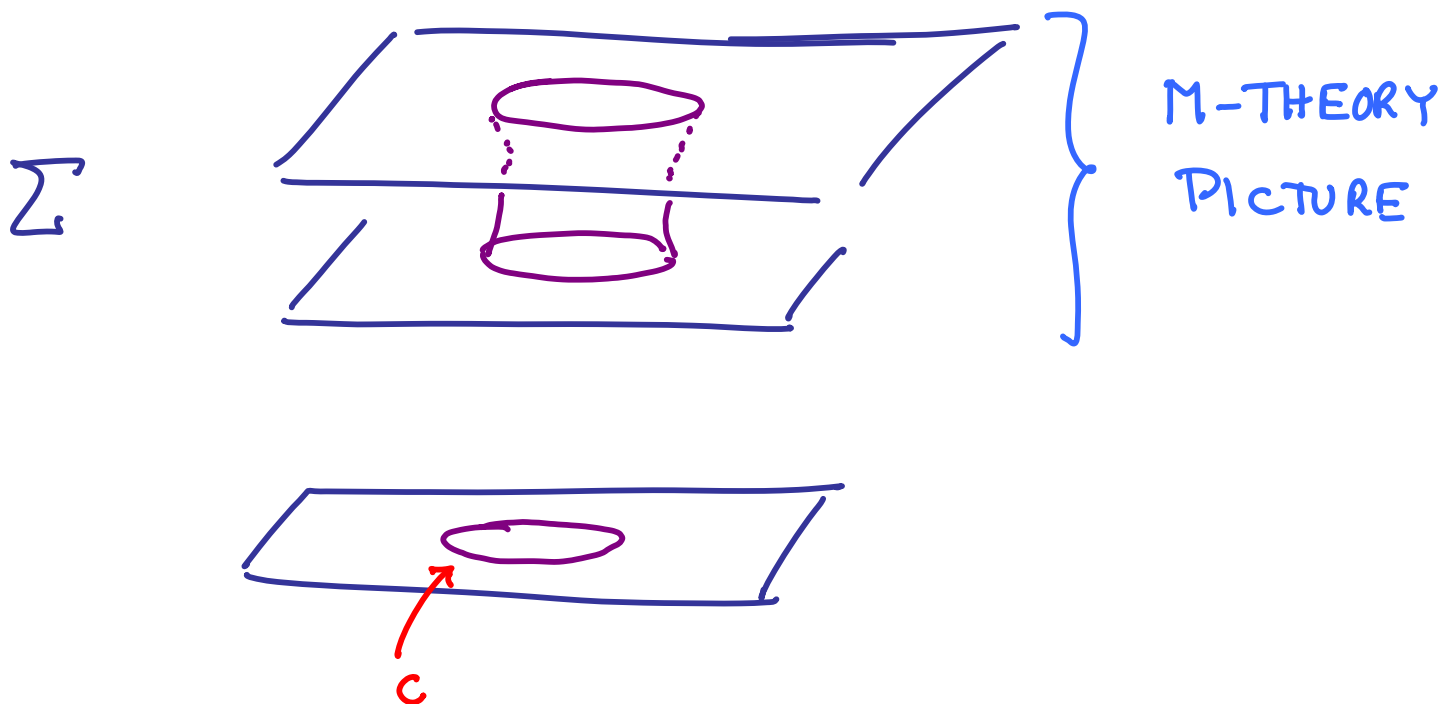
$\Rightarrow$  HM WITH  $\Omega = +1$

M-THEORY PICTURE: OPEN M2:



2. OR, C = CLOSED CURVE  $\Rightarrow$

VM WITH  $\Omega = -2$ .



N.B. FOR SUCH CURVES  $\langle \lambda, d_t \rangle = e^{i\vartheta_*}$

$$\vartheta_* = \arg(Z)$$

(REMARK: FOR  $SU(k)$ ,  $k > 2$  THE  
BPS STATES ARE MORE COMPLICATED  
AND INVOLVE "STRING WEBS.")

# HITCHIN SYSTEMS & FLAT CONN'S

THE H.E.'S  $\Rightarrow$

$$A = \frac{R}{\mathcal{J}} \varphi + A + R \mathcal{J} \bar{\varphi}$$

IS FLAT.

NEAR REG. SING. POINT  $z_i$

$$A \sim \left( \frac{R \rho}{\mathcal{J} 2} + \frac{\alpha}{2i} \right) \frac{dz}{z - z_i} + \left( R \mathcal{J} \frac{\bar{\rho}}{2} - \frac{\alpha}{2i} \right) \frac{d\bar{z}}{\bar{z} - \bar{z}_i}$$

SO B.C.'S FIX MONODROMY OF  $A$ .

AROUND  $z_i$ :

$$M_i \sim \begin{pmatrix} \mu_i & \\ & \mu_i^{-1} \end{pmatrix}$$

$$\mu_i = \exp \left[ 2\pi i \left( \frac{1}{2} \mathcal{J}^{-1} R m_i - m_i^3 - \frac{1}{2} \mathcal{J} R \bar{m}_i \right) \right]$$

THEOREM OF C. SIMPSON  $\implies$

IDENTIFY  $\mathcal{M}^{\mathfrak{J}}$ ,  $\mathfrak{J} \in \mathbb{C}^*$ , WITH  
MODULI OF FLAT  $SL(2, \mathbb{C})$  CONNECTIONS  
WITH PRESCRIBED MONODROMY AT  $Z_i$

MOREOVER, THE HOLO. SYMPLECTIC  
FORM ON  $\mathcal{M}^{\mathfrak{J}}$  HAS THE SIMPLE  
FORM:

$$\bar{\omega}_{\mathfrak{J}} = \int_C \text{Tr}(\delta A \delta A)$$

## STRATEGY:

- FOCK & GONCHAROV CONSTRUCTED NICE COORDINATE SYSTEMS ON  $\mathcal{M}^S$

- WE WILL USE THESE FUNCTIONS TO CONSTRUCT SYSTEMS OF COORD'S

$$\chi_y^{\vartheta}(\cdot, \mathcal{J}) \text{ ON } \mathcal{M}, \vartheta \in \mathbb{R}/\mathbb{Z}$$

RELATION TO PREVIOUS TWISTOR COORD'S

$$\chi_y(\cdot, \mathcal{J}) = \chi_y^{\vartheta=\arg \mathcal{J}}(\cdot, \mathcal{J})$$

THEN, AT LARGE  $R$ ,

THE  $\chi_y(\cdot, \mathcal{J})$  SATISFY

DEFINING PROPERTIES 1  $\rightarrow$  5.



## 6. FOCK - GONCHAROV COORD'S AND CLUSTER TMNS

$\mathcal{A}$  = A FLAT  $SL(2, \mathbb{C})$  CONNECTION  
WITH MONODROMY  $M_i$  AROUND  $z_i$

$$M_i \sim \begin{pmatrix} \mu_i & \\ & 1/\mu_i \end{pmatrix}$$

### A. DECORATED TRIANGULATIONS

DEF: A "DECORATED TRIANG."  $T$   
IS AN IDEAL TRIANGULATION OF  $\mathbb{C}$   
WITH VERTICES AT  $z_i$ ; TOGETHER  
WITH A CHOICE OF MONODROMY  
EIGENVALUE  $\mu_i$  OR  $1/\mu_i$  AT EACH  $z_i$

## B. FLIPS AND POPS

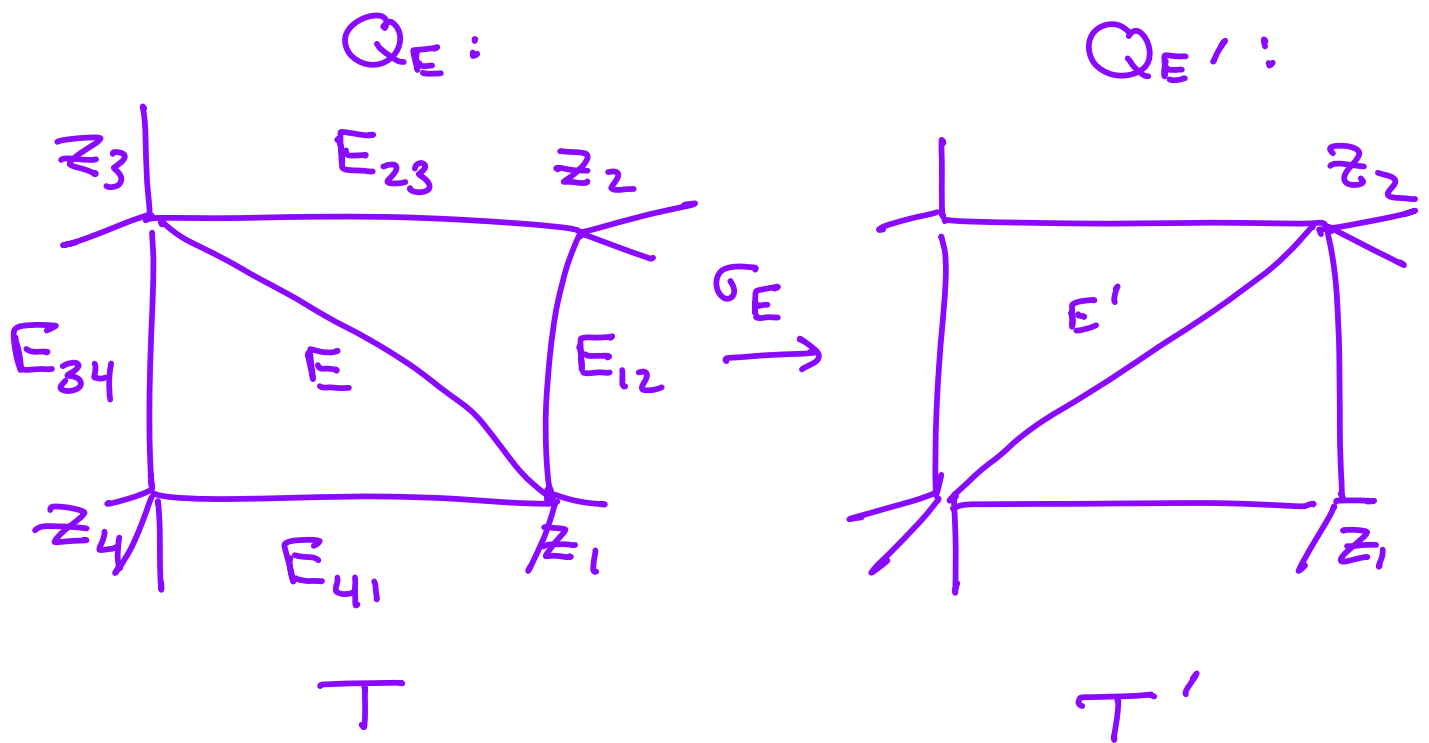
DEFINE A GROUPOID :

DECORATED.

OBJECTS = TRIANGULATIONS

MORPHISMS ARE GENERATED BY FLIPS  
+ POPS

FLIP :  $\sigma_E$  FOR  $E \in \mathcal{E}(T)$



POP:  $\pi_i$  FOR  $z_i \in V(T)$ :

EXCHANGE:  $\mu_i \leftrightarrow \mu_i^{-1}$

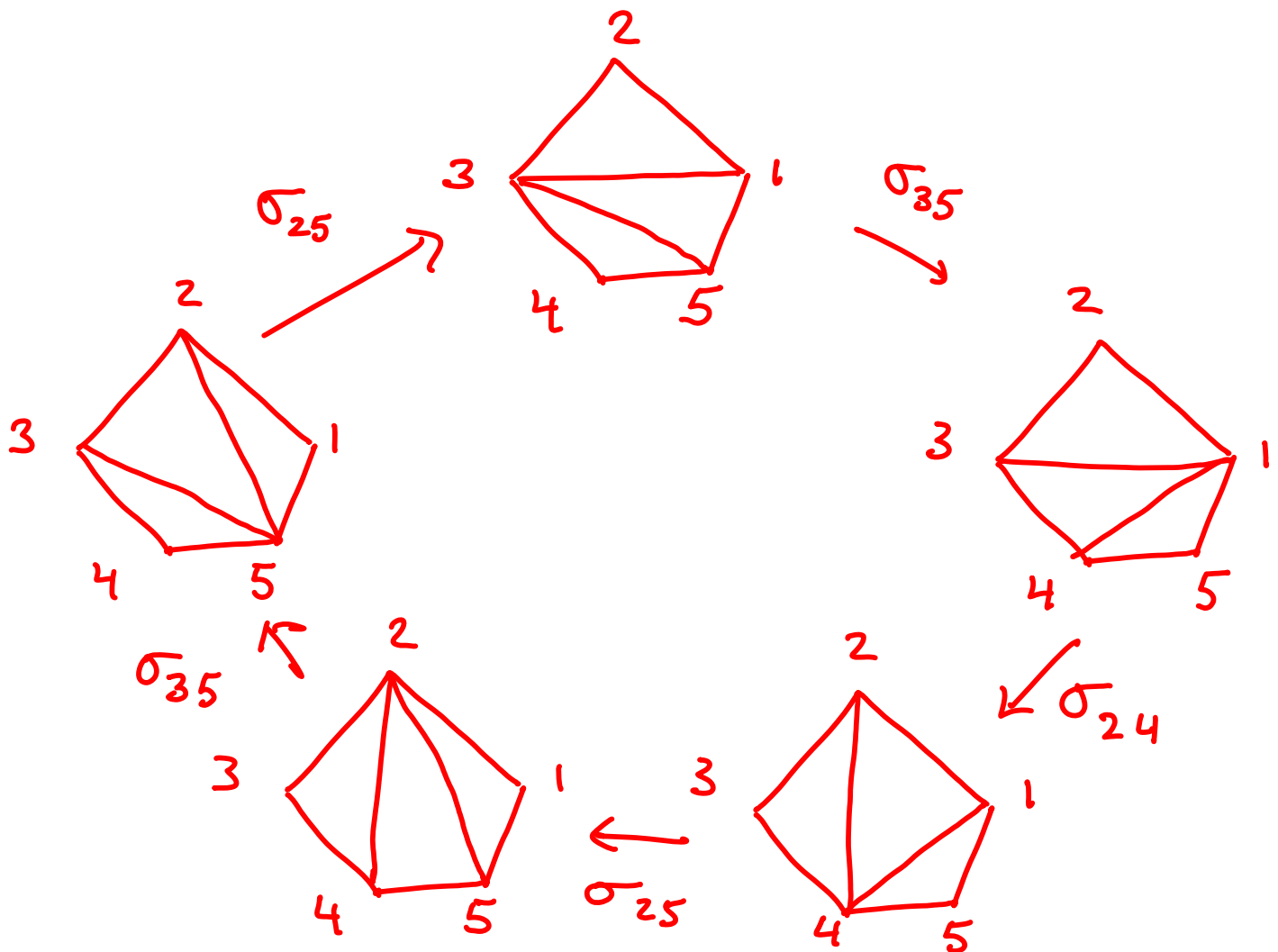
# RELATIONS ON FLIPS $\in$ POPS

1.  $\sigma_E^2 = 1$  AND  $\pi_i^2 = 1$

2. POPS COMMUTE

3.  $\sigma_E, \sigma_{E'}$  COMMUTE IF  $Q_E, Q_{E'}$  DO NOT SHARE A TRIANGLE

4. IF  $Q_E, Q_{E'}$  SHARE A TRIANGLE



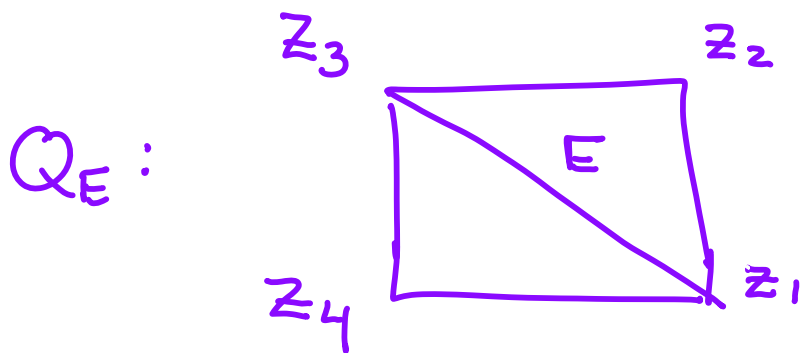
LATER WE WILL ENHANCE OUR GROUPOID TO INCLUDE "LIMIT TRIANGLES" AND "TWISTS"

## C. FG COORDINATES

GIVEN A DECORATED  
TRIANGULATION OF  $C$ ,  $F \in G$   
DEFINE A COLLECTION OF  
FUNCTIONS ON  $\mathcal{M}$ :

$$\chi: T \rightarrow \left\{ \chi_E^T \right\}_{E \in \mathcal{E}(T)}$$

DEFINITION:



- CHOOSE FLAT SECTIONS  $S_i$  OF SPECIFIED MONODROMY NEAR  $z_i$ :

- $$\chi_E^T := - \frac{(S_1 \wedge S_2)(S_3 \wedge S_4)}{(S_2 \wedge S_3)(S_4 \wedge S_1)}$$

$$\chi_E^T := - \frac{(s_1 \wedge s_2)(s_3 \wedge s_4)}{(s_2 \wedge s_3)(s_4 \wedge s_1)}$$

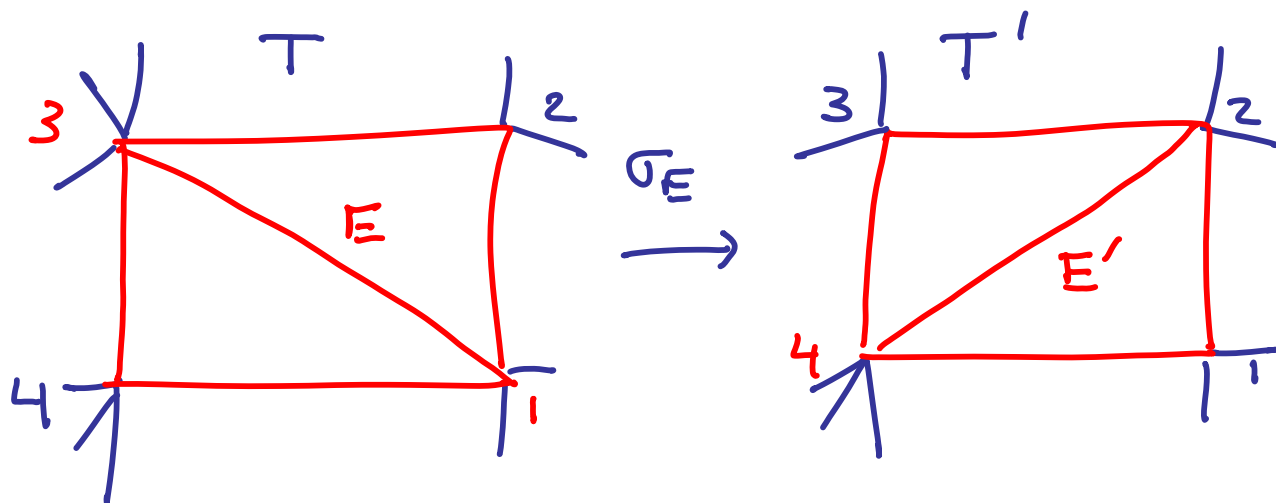
- $s_i \wedge s_j \in \Lambda^2 E = \text{LINE BUNDLE}$
- PARALLEL TRANSPORT TO ANY POINT  $Q \in Q_E$
- NORMALIZATION OF  $s_i$  CANCELS

THEOREM:  $(F \stackrel{!}{=} G) \{ \chi_E^T \}_E$  PROVIDE HOLO. COORDINATES ON OPEN SET  $U_T$  OF  $M$ .

### D. COORDINATE TMN'S

NOW DESCRIBE THE COORD. TMNS AS WE CHANGE THE DECORATED TRIANGULATION  $T \rightarrow T'$

# TRANSFORMATION UNDER FLIPS

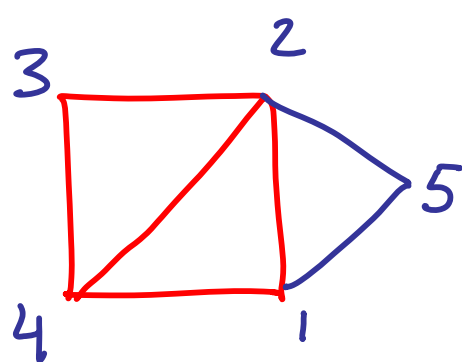
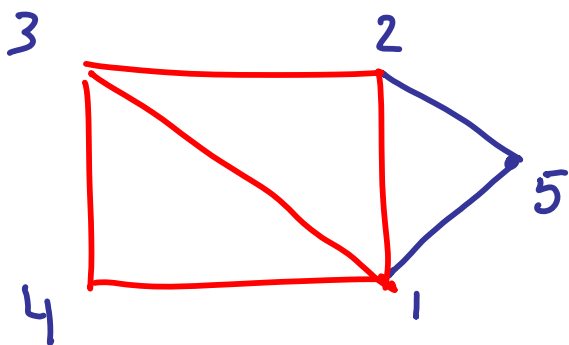


ONLY THE EDGES IN RED CHANGE

$$\chi_{E'}^{T'} = - \frac{s_4 s_1 s_2 s_3}{s_1 s_2 s_3 s_4} = \frac{1}{\chi_E^T}$$

$$\chi_{E_{12}}^{T'} = \chi_{E_{12}}^T (1 + \chi_E^T)$$

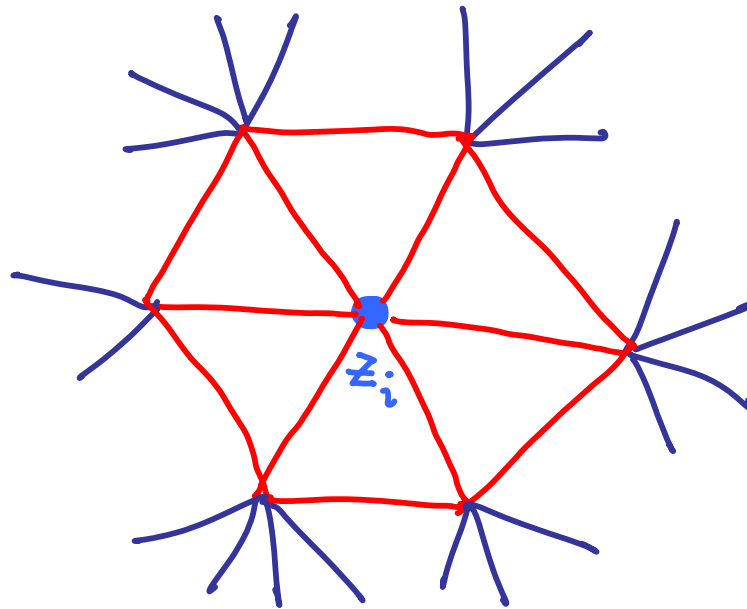
...



"CLUSTER TRANSFORMATIONS"

# TRANSFORMATION UNDER POPS

A POP AT VERTEX  $z_i$  CHANGES  
THE EDGE COORDINATES **IN RED**



IT IS POSSIBLE TO WRITE EXPLICIT  
FORMULAE FOR THE POP TRANSFORMATION,  
BUT THEY ARE COMPLICATED....

**IMPORTANTLY!**

IT TURNS OUT THAT THE PRODUCT OF  
ALL POPS  $\prod_V \pi_v$  IS RELATIVELY SIMPLE..

## E. SYMPLECTIC STRUCTURE

- USING THE SYMPLECTIC STRUCTURE  
 $\overline{\omega}_\Sigma$  ONE CAN SHOW

$$\{ \chi_E^T, \chi_{E'}^T \} = \langle E, E' \rangle \chi_E^T \chi_{E'}^T$$

- TRANSFORMATIONS UNDER  
FLIPS  $\frac{1}{i}$  POPS ARE POISSON



## 7. WKB TRIANGULATIONS

RECALL THAT A KEY PROPERTY OF  $\chi_y(\mathcal{S})$  ARE THE  $\mathcal{S} \rightarrow 0$  ASYMPTICS:

$$\lim_{\mathcal{S} \rightarrow 0} \chi_y(\mathcal{S}) e^{-\frac{\pi R}{\mathcal{S}} Z_y(u)} \sim \text{FINITE}$$

$\Rightarrow$  WE NEED TO USE VERY SPECIAL TRIANGULATIONS FOR WHICH WE CAN PROVE SUCH ASYMPTOTICS.

IDEA: USE THE WKB APPROXIMATION TO DESCRIBE THE FLAT SECTIONS:

$$(d+A) \mathcal{S} = 0$$

$$A = \frac{R}{\mathcal{J}} \varphi + A + R \mathcal{J} \overline{\varphi}$$

FOR  $\mathcal{J} \rightarrow 0$ ,  $\mathcal{J} \sim \hbar$

$$(\mathcal{J} d + R \varphi + \alpha(\mathcal{J})) s = 0$$

RECALL:  $\varphi \sim \begin{pmatrix} \lambda \\ -\lambda \end{pmatrix}$

WKB:

$$s \underset{\mathcal{J} \rightarrow 0}{\sim} \exp\left(-\frac{R}{\mathcal{J}} \int^z \lambda \sigma^3\right) \overset{\text{Const.}}{\downarrow} s_0$$

FROM THIS GET ASYMPT'S OF FG COORD.

HOWEVER THE WKB APPX.  
IS NOTORIOUSLY SUBTLE.

EXPONENTIALLY SMALL CORRECTIONS  
CAN GROW IN  $z$  AND INVALIDATE  
COMPUTATIONS

FOR VALIDITY OF WKB APPX.  
WE MUST RESTRICT TO VERY  
SPECIAL TRIANGULATIONS  $T(\vartheta, \lambda)$   
WHOSE EDGES ARE WKB CURVES

DEF: WKB CURVE WITH ANGLE  $\vartheta$ :

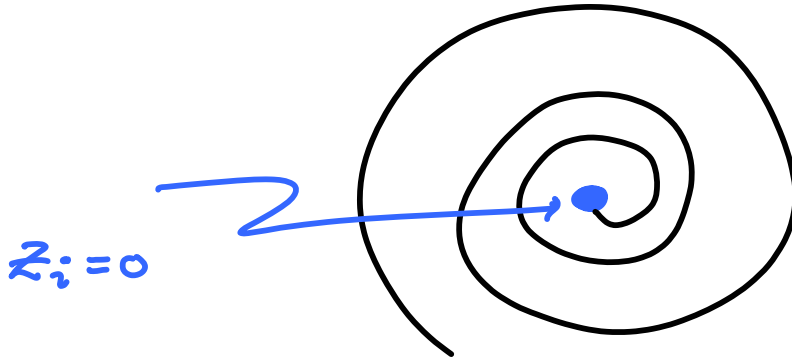
CURVE ON  $C$  WITH

$$\langle \lambda, \partial_t \rangle = \pm e^{i\vartheta}$$

$\Rightarrow$  WKB FOLIATION OF  $C$ .

NOTE: WKB CURVES GET TRAPPED  
BY SINGULARITIES:

$$\lambda = \frac{m}{2} \frac{dz}{z} \Rightarrow z(t) = z_0 \exp\left(-\frac{e^{i\theta}}{m} t\right)$$



THREE KINDS OF WKB CURVES:

GENERIC: BOTH ENDS ON  $z_i, z_j$

SEPARATING: CONNECTS BRANCH  
POINT  $w_a$  TO SINGULAR POINT  $z_i$

FINITE: CLOSED, OR BOTH  
ENDS ON TURNING POINTS  $w_a, w_b$

NOTE THAT OUR RULE FOR  
BPS STATES WAS THAT  $\exists \vartheta_*$   
FOR WHICH THERE IS A FINITE  
WKB CURVE

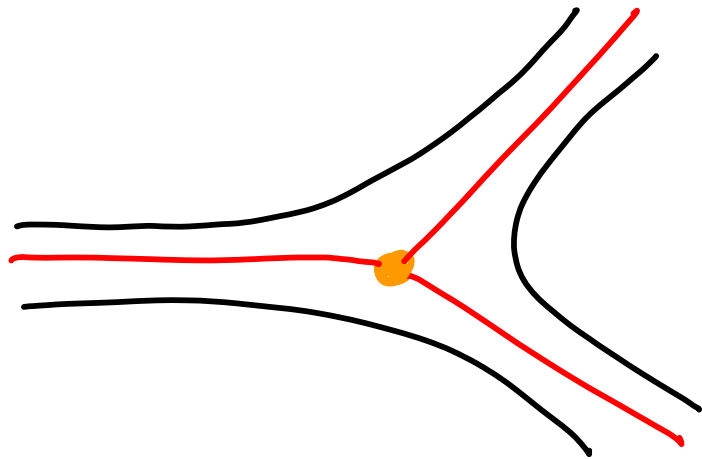
IMPORTANT FACT: FOR GENERIC  
VALUES OF  $\vartheta$  THERE ARE NO  
FINITE WKB CURVES. BUT  
AT SPECIAL CRITICAL VALUES  
OF  $\vartheta$  THERE ARE FINITE WKB  
CURVES.

RECALL THAT FOR FINITE WKB  
CURVES  $\langle \lambda, \partial_t \rangle = e^{i\vartheta_*}$  WITH  
 $\vartheta_* = \arg Z$ .

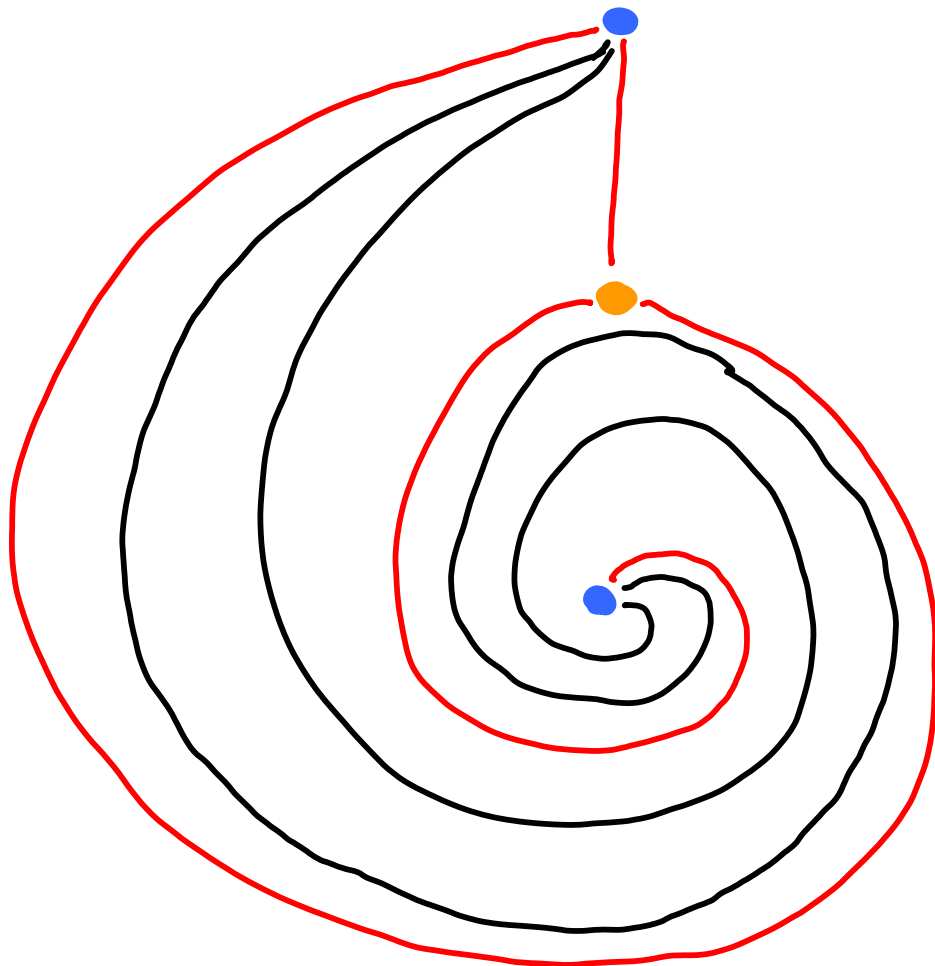
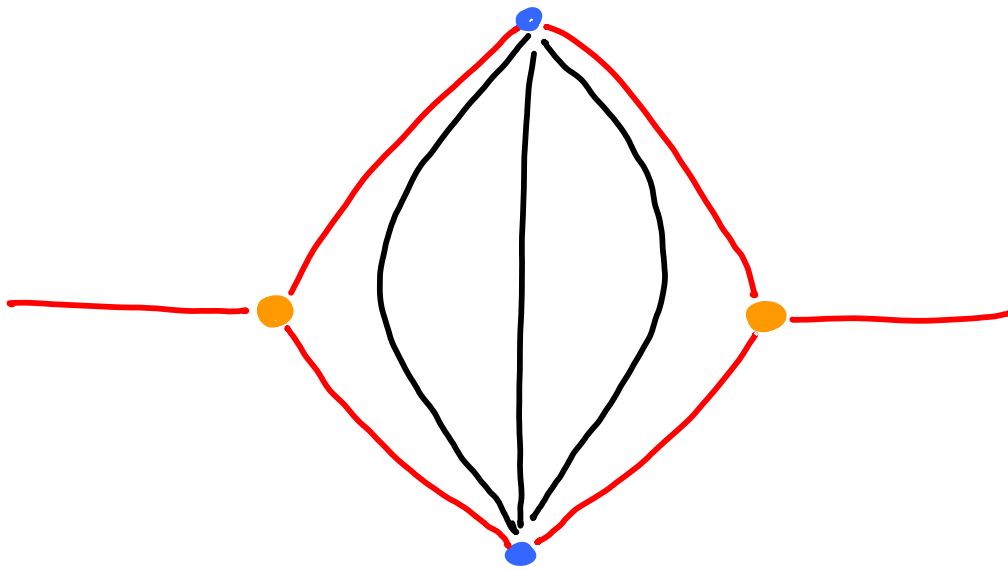
SO: THE CRITICAL VALUES ARE  
THE PHASES  $\vartheta_*$  OF BPS STATES.

TO DEFINE OUR TRIANGULATION  
WE FIRST USE THE SEPARATING CURVES  
TO SPLIT  $C$  INTO WKB CELLS

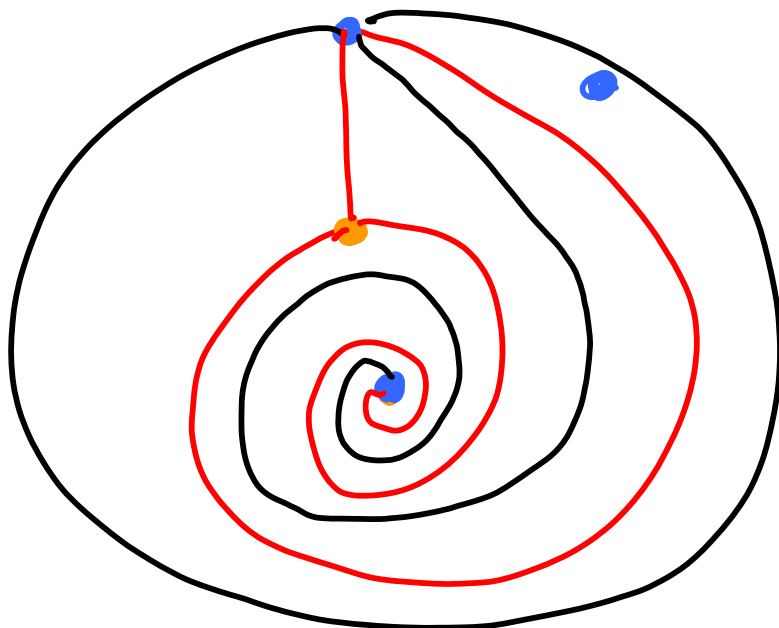
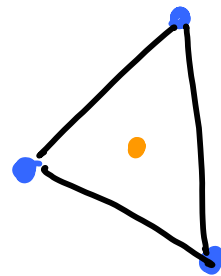
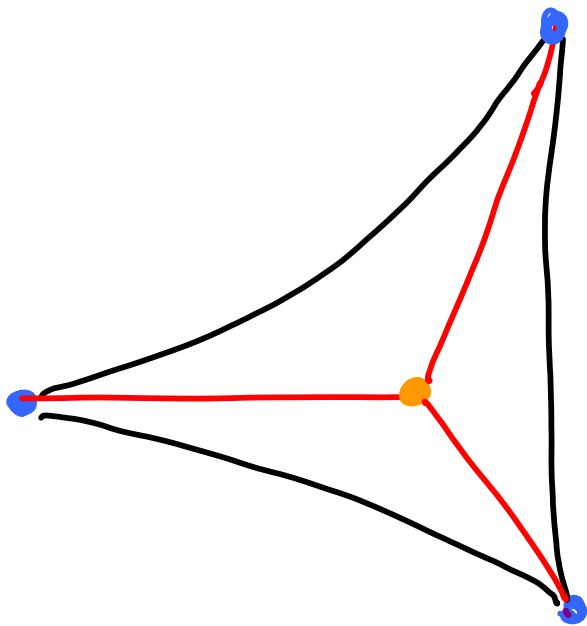
LOCALLY:



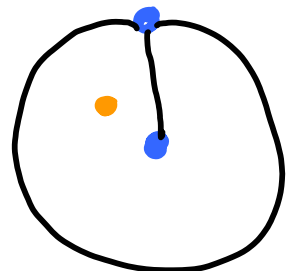
FOR GENERIC  $\lambda, \nu$  IT TURNS  
OUT THERE ARE ONLY TWO  
KINDS OF CELLS:



FOR THE WKB TRIANGULATION  
WE CHOOSE A GENERIC WKB  
CURVE IN EACH CELL:



DEGENERATE  
TRIANGLE:





# CHOICE OF $\mu_i$

RECALL THAT WE MUST DEFINE  
A "DECORATED TRIANGULATION."

$(\lambda, \vartheta) \Rightarrow$  DISTINGUISHED  
EIGENVALUE OF  $M_i$

"SMALL FLAT SECTION": THE  
FLAT SECTION WHICH DECAYS  
ALONG THE WKB CURVE GOING  
INTO THE SINGULARITY

THESE ARE THE SECTIONS FOR  
WHICH WE HAVE GOOD CONTROL IN WKB  
APPT.

DENOTE THE RESULTING DECORATED  
TRIANGULATION  $T(\vartheta, \lambda)$

# MORPHISMS OF WKB TRIANG'S

VARY  $\mathcal{V} \Rightarrow$  (HOMOTOPY CLASS OF)  
 $T(\mathcal{V}, \lambda)$  IS UNCHANGED

EXCEPT AT CRITICAL VALUES  $\mathcal{V}_c$   
WHERE FINITE WKB CURVES  
DEVELOP.

WHEN VARYING  $\mathcal{V}$ ,  
 $T(\mathcal{V}, \lambda)$  JUMPS PRECISELY AT THE  
VALUES OF PHASES OF BPS STATES!

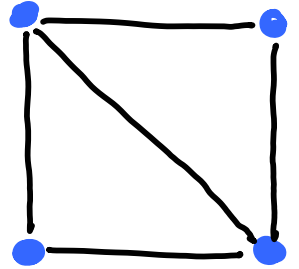
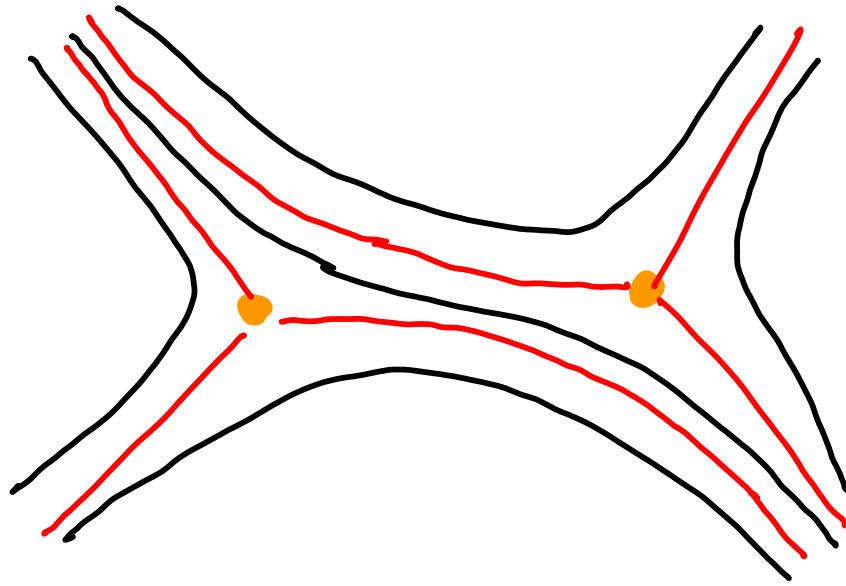
FOR GENERIC  $\lambda$  A JUMP IN  $T(\mathcal{D}, \lambda)$  ONLY HAPPENS WHEN A SEPARATING CURVE DEGENERATES TO A FINITE WKB CURVE JOINING TURNING POINTS  $w_a, w_b$ .

THUS, FOR GENERIC  $\lambda$  THERE ARE ONLY TWO KINDS OF JUMPS:

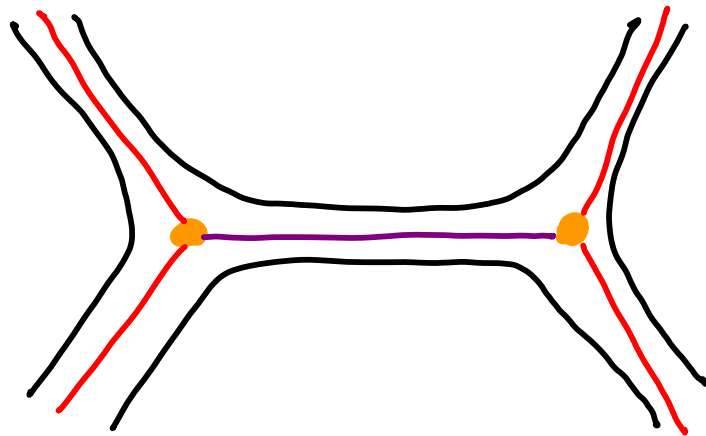
- EITHER  $w_a \neq w_b$
- OR  $w_a = w_b$

# HYPERMULTIPLY JUMP: $w_a \neq w_b$

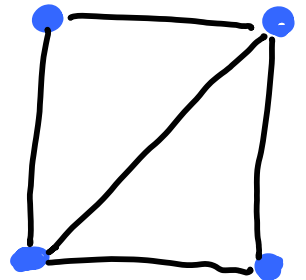
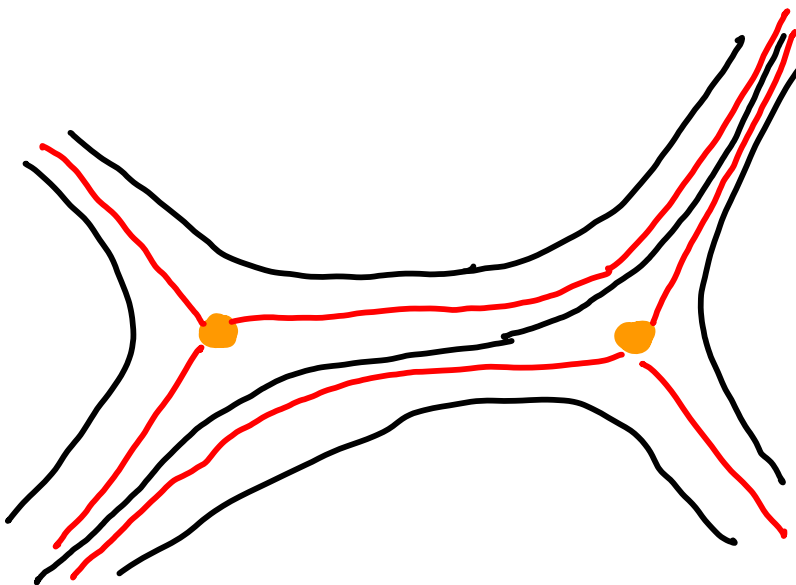
$\vartheta < \vartheta_c$



$\vartheta = \vartheta_c$

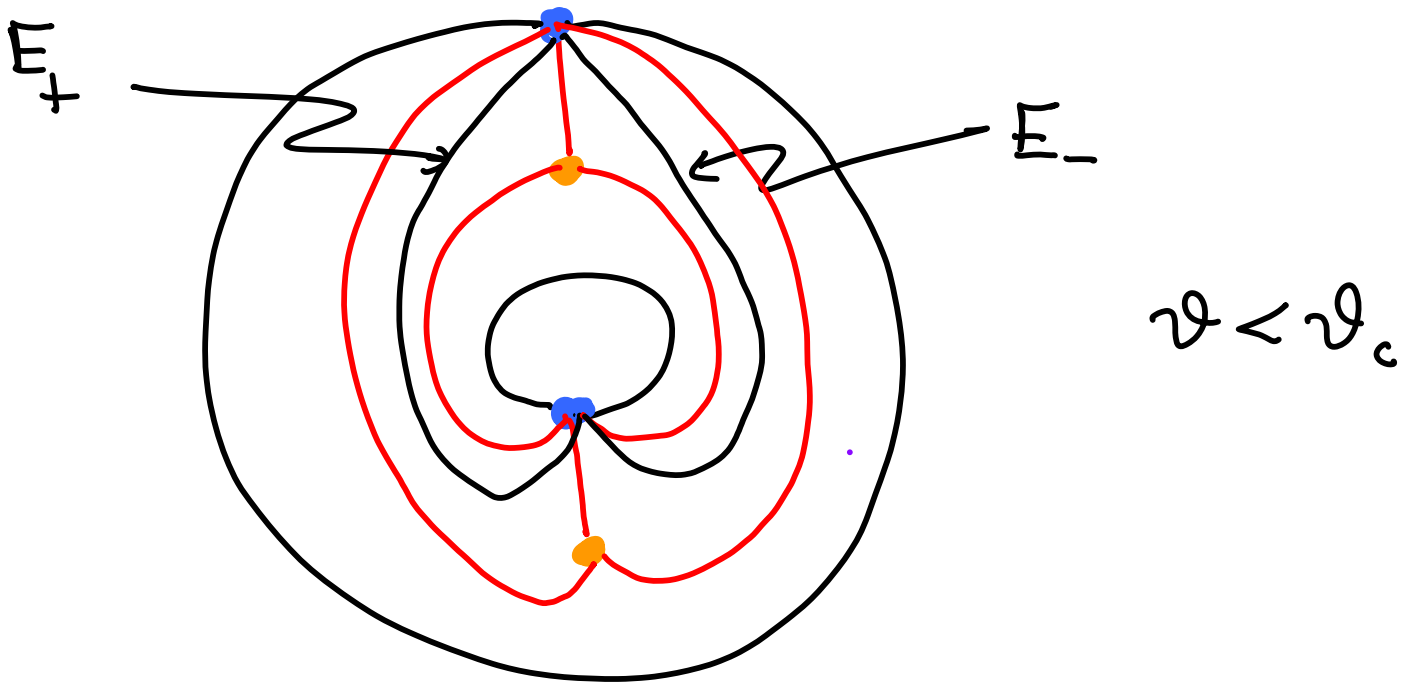


$\vartheta > \vartheta_c$



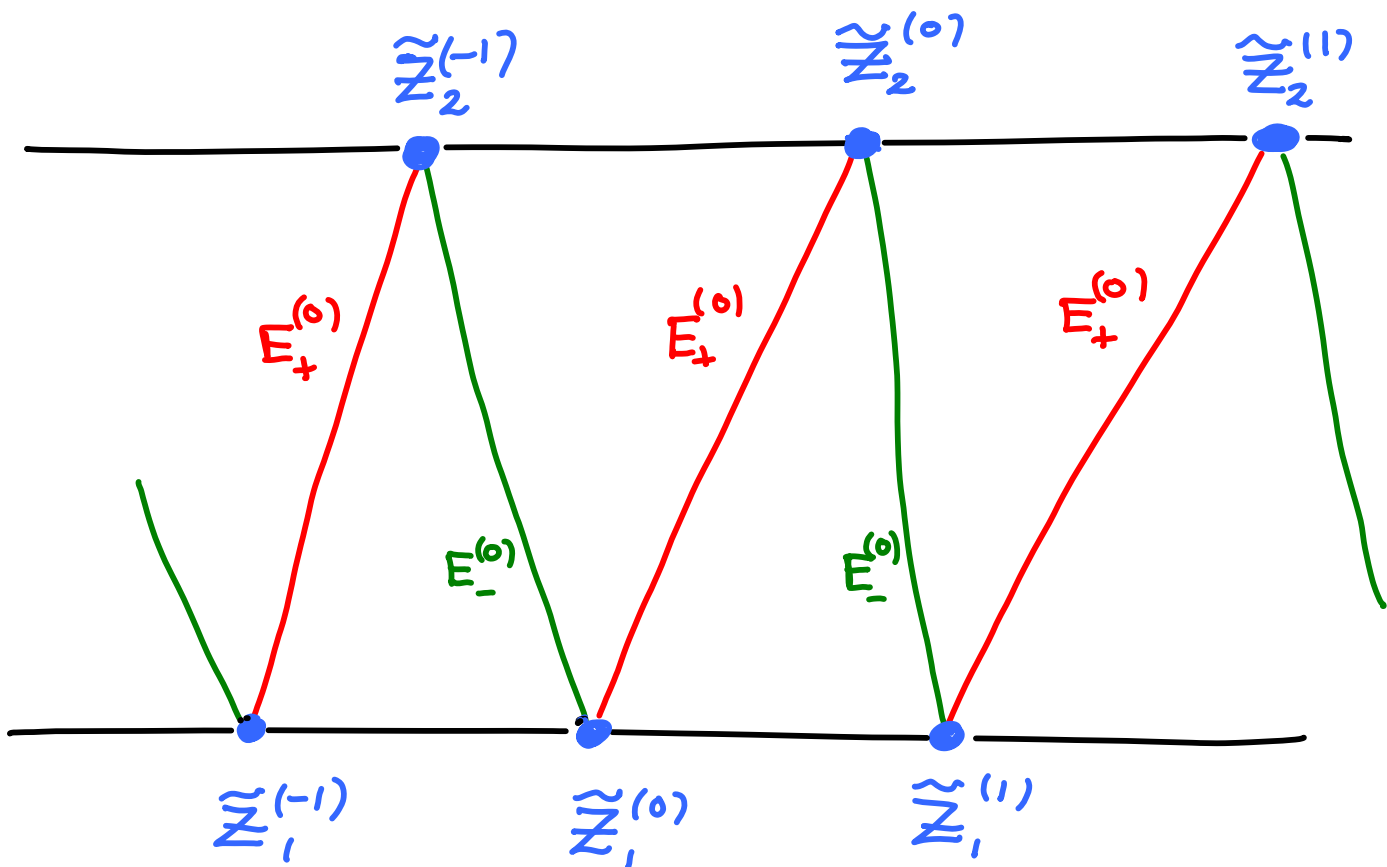
THIS IS JUST A FLIP

# VECTOR MULTIPLY JUMP: $W_a = W_b$

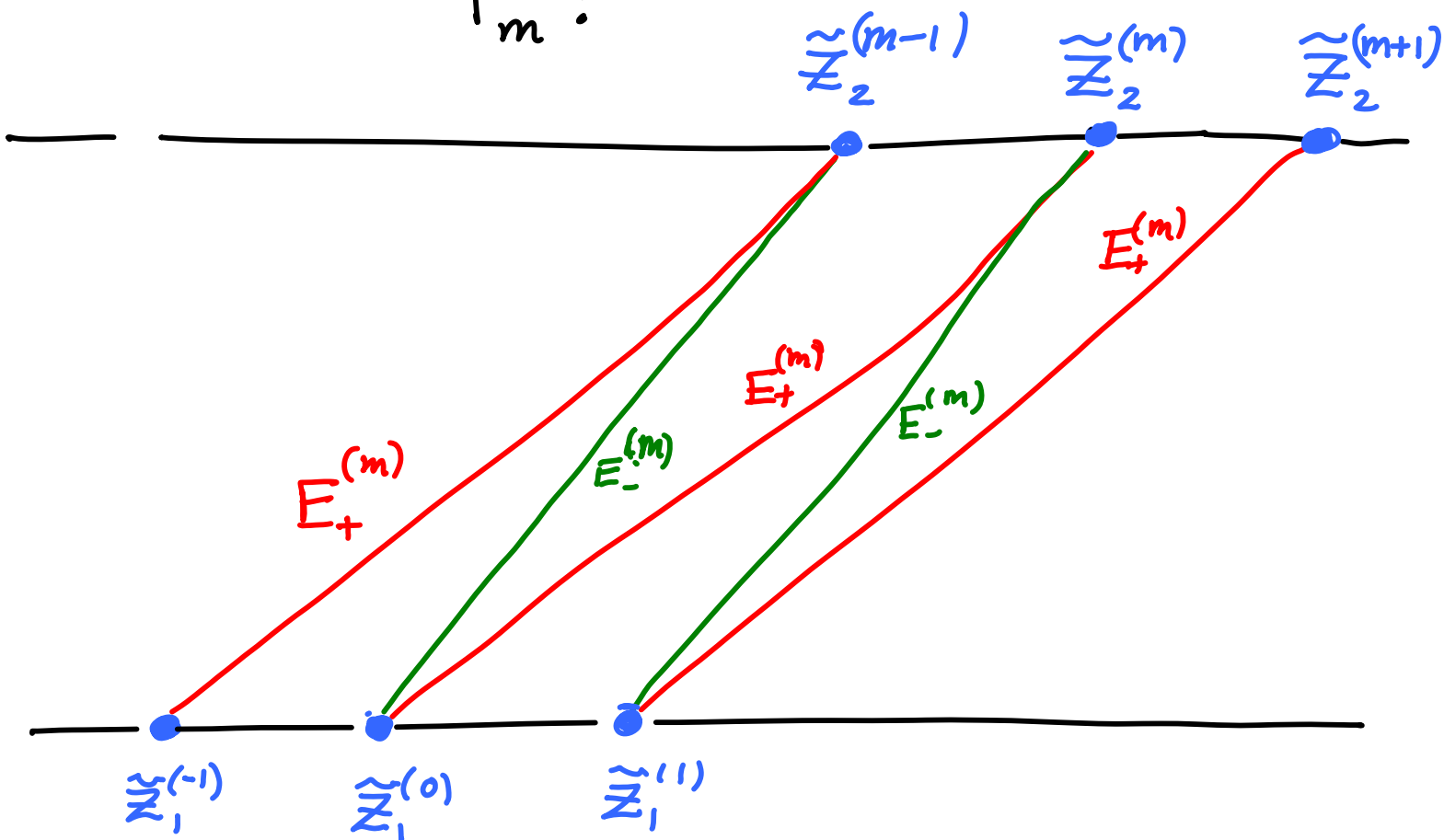


AS  $v \rightarrow v_c^-$   $T(v, \lambda)$  HAS  
AN INFINITE SEQUENCE OF FLIPS

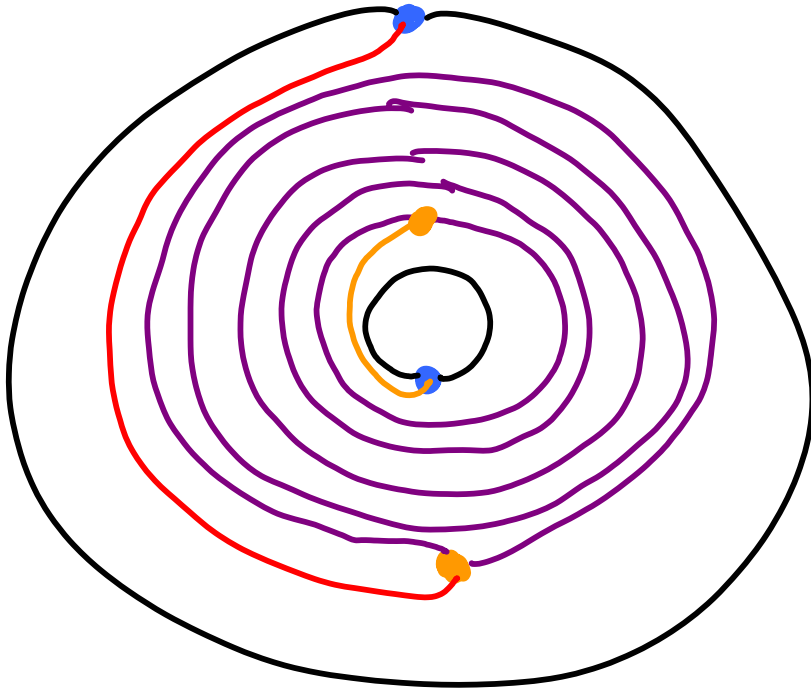
$E_+ E_- E_+ E_- \dots$



$T_m :$

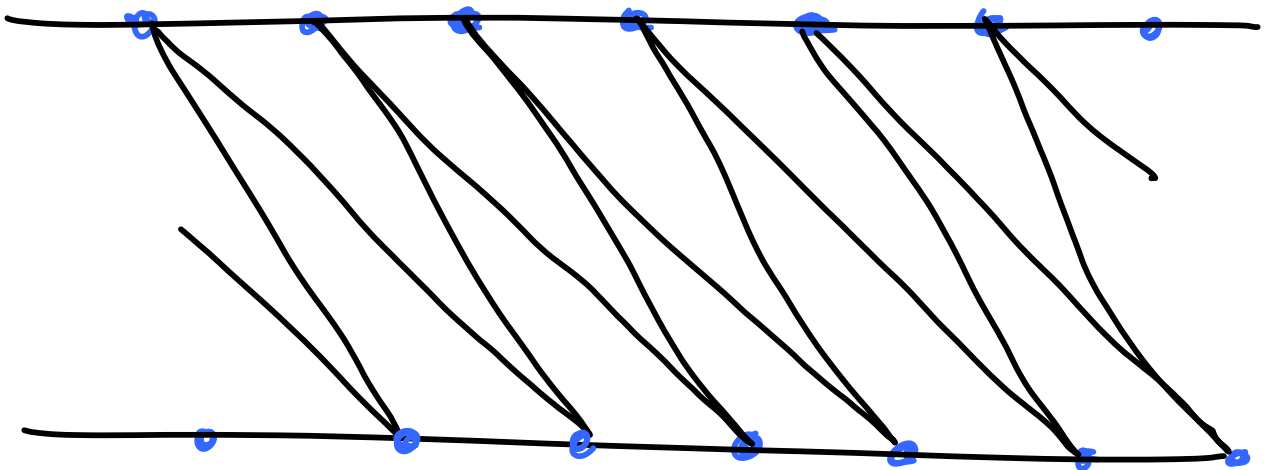


AT  $\vartheta = \vartheta_c$



$\vartheta = \vartheta_c$

AT  $\vartheta > \vartheta_c$



$T_m$

SUITABLE COMBINATIONS OF

$\chi_{E_+}^{T_m}, \chi_{E_-}^{T_m}$  HAVE LIMITS

FOR  $m \rightarrow \infty$ :

$$\chi_A^{T_{+\infty}} = \lim_{m \rightarrow \infty} \chi_{E_+}^{T_m} \chi_{E_-}^{T_m}$$

$$\chi_A^{T_{-\infty}} = \lim_{m \rightarrow -\infty} \chi_{E_+}^{T_m} \chi_{E_-}^{T_m}$$

$\Rightarrow$  ADD NEW OBJECTS  $T_{\pm\infty}$

TO THE GROUPOID, CALL THEM  
"LIMIT TRIANGULATIONS"

NEW MORPHISMS:

$$T_{+\infty} \rightarrow T_{-\infty}$$

CALLED "TWISTS"

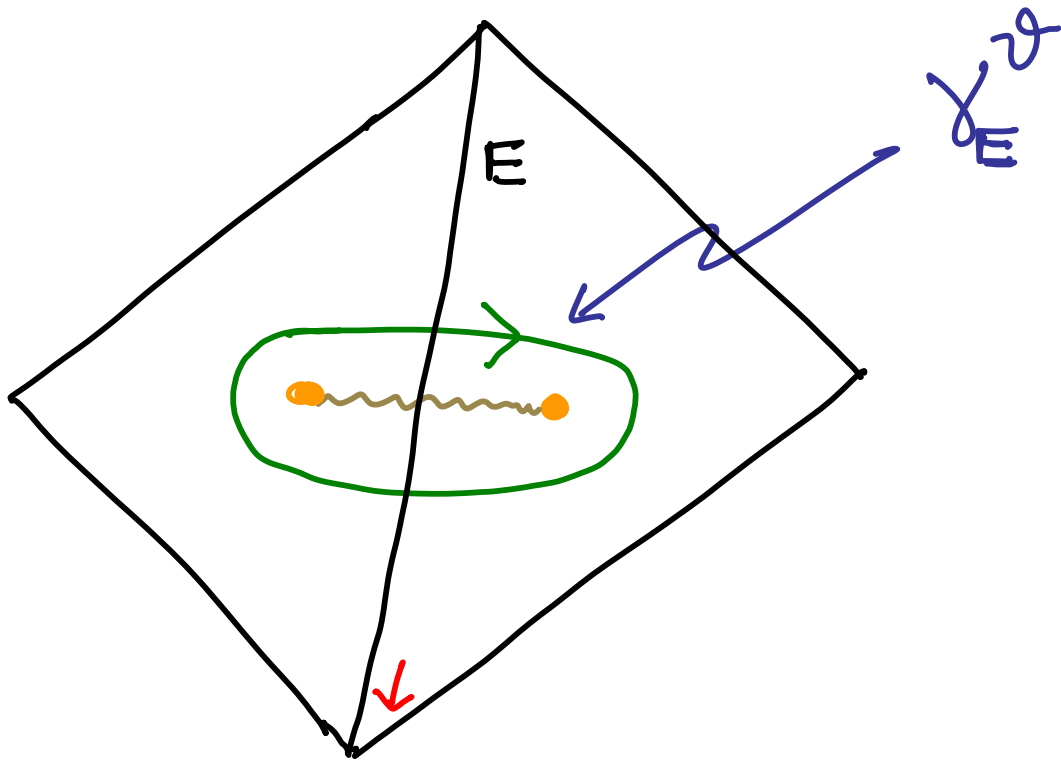


# 8. DEFINING THE TWISTOR COORD'S

FINALLY, TO DEFINE  $\chi_\gamma^\vartheta(\cdot, \mathcal{J})$

WE ASSOCIATE TO  $E \in \mathcal{E}(\mathbb{T}(\vartheta, \lambda))$

CERTAIN CYCLES  $\gamma_E^\vartheta \in H_1(\Sigma, \mathbb{Z})$



RULE: ORIENT THE LIFTS  $\hat{E}$  SO  
THAT  $e^{-i\vartheta} \langle \lambda, \partial_t \rangle \geq 0$  : +VE

DEMAND  $\langle \gamma_E^\vartheta, \hat{E} \rangle = +1$

THE  $\{\gamma_E^\vartheta\}_{E \in \mathcal{E}(T)}$  FORM  
A (POSITIVE) BASIS FOR  $\Gamma$ .

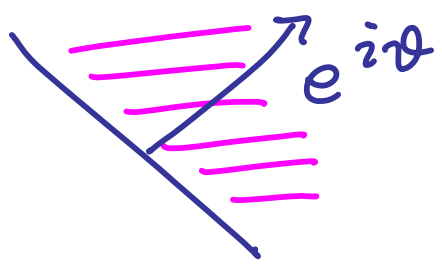
NOW DEFINE :

$$\chi_{\gamma_E^\vartheta}^\vartheta := \chi_E^{\tau(\vartheta, \lambda)}$$

$$\chi_{\gamma + \gamma'}^\vartheta := \chi_\gamma^\vartheta \chi_{\gamma'}^\vartheta$$

THEOREM 1: IF  $R \rightarrow \infty$  AND

$\mathcal{S}$  IS IN  $\mathbb{H}_\vartheta$ :

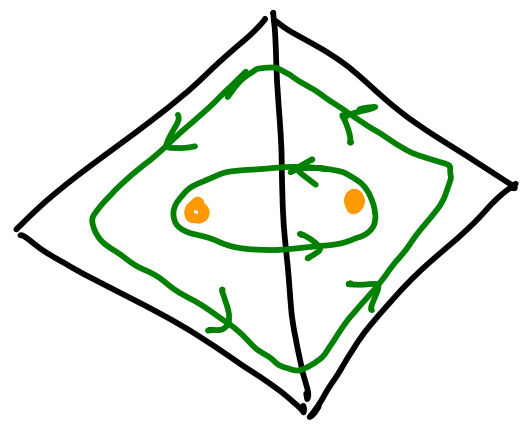


THEN

$$\chi_\gamma^{\mathcal{S}}(\cdot, \mathcal{S}) \underset{R \rightarrow \infty}{\sim} \exp\left(\frac{\pi R}{\mathcal{S}} z_\gamma + i\theta_\gamma + \pi R \mathcal{S} \bar{z}_\gamma\right)$$

RECOVERS NEITZKE-PIOLINE SEMIFLAT TWISTOR COORDINATES.

PROOFS:



$$S_i \sim \exp\left(\pm \frac{R}{\mathcal{S}} \int_{z_i}^z \lambda\right) \left\{ \begin{pmatrix} 1 \\ 0 \end{pmatrix} \text{ or } \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right\}$$

USE RELATION TO 2D Sinh-Gordon

## THEOREM 2:

WITH RESPECT TO SYMPLECTIC  
STRUCTURE:

$$\omega_g = \int_C \text{Tr } \delta A \delta A$$

$$\{ \chi_\gamma^{\mathcal{D}}, \chi_{\gamma'}^{\mathcal{D}} \} = \langle \gamma, \gamma' \rangle \chi_{\gamma+\gamma'}^{\mathcal{D}}$$

THEOREM 3: AT SUFFICIENTLY  
LARGE  $R$

$$\chi_\gamma(\cdot, \mathcal{S}) = \chi_\gamma^{\mathcal{D}=\arg \mathcal{S}}(\cdot, \mathcal{S})$$

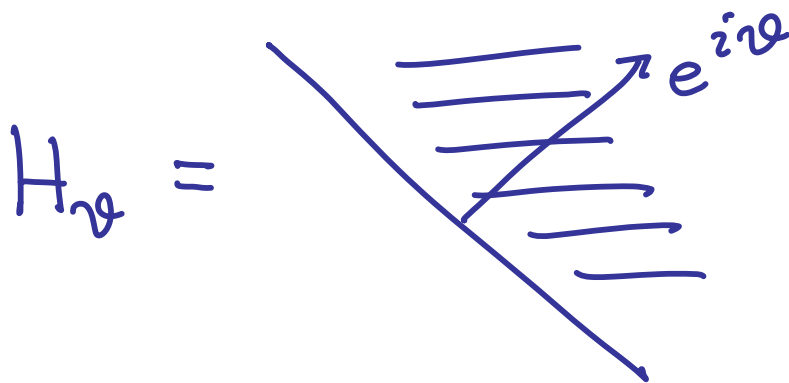
SATISFY THE 5 DEFINING  
PROPERTIES.

# PROOF:

(1)  $\chi_\gamma(\cdot, \delta)$  HOLOMORPHIC ON  $\mathcal{M}^\delta$ :  
FOCK & GONCHAROV

(2), (3) FOLLOW EASILY FROM  
THE DEFINITION

(4B) FOR  $\delta \rightarrow 0$  IN THE HALF-PLANE



$\lim_{\delta \rightarrow 0} \chi_\gamma^{2g}(\delta) \exp\left(-\frac{\pi R}{\delta} Z_\gamma\right)$  EXISTS

FOLLOWS FROM WKBJ ASYMPTOTICS  
AS WITH  $R \rightarrow \infty$

(5) IF  $\vartheta = \vartheta_c$  IS THE PHASE OF A BPS STATE OF CHARGE  $\gamma_0$  THEN, DEFINING

$$\chi_{\gamma}^{\pm} = \lim_{\vartheta \rightarrow \vartheta_c^{\pm}} \chi_{\gamma}^{\vartheta}$$

$$\chi_{\gamma}^{+} = \chi_{\gamma}^{-} \left( 1 - \sigma(\gamma_0) \chi_{\gamma_0}^{-} \right)^{\Omega(\gamma_0) \langle \gamma, \gamma_0 \rangle}$$

NOTE:  $\sigma(\gamma_0) = +1$ ,  $\Omega(\gamma_0) = -2$  VM  
 $\sigma(\gamma_0) = -1$ ,  $\Omega(\gamma_0) = +1$  HM

1. FOR HM: CLUSTER TMN.
2. FOR VM: EXPLICIT COMPUTATION OF TWIST TMN:

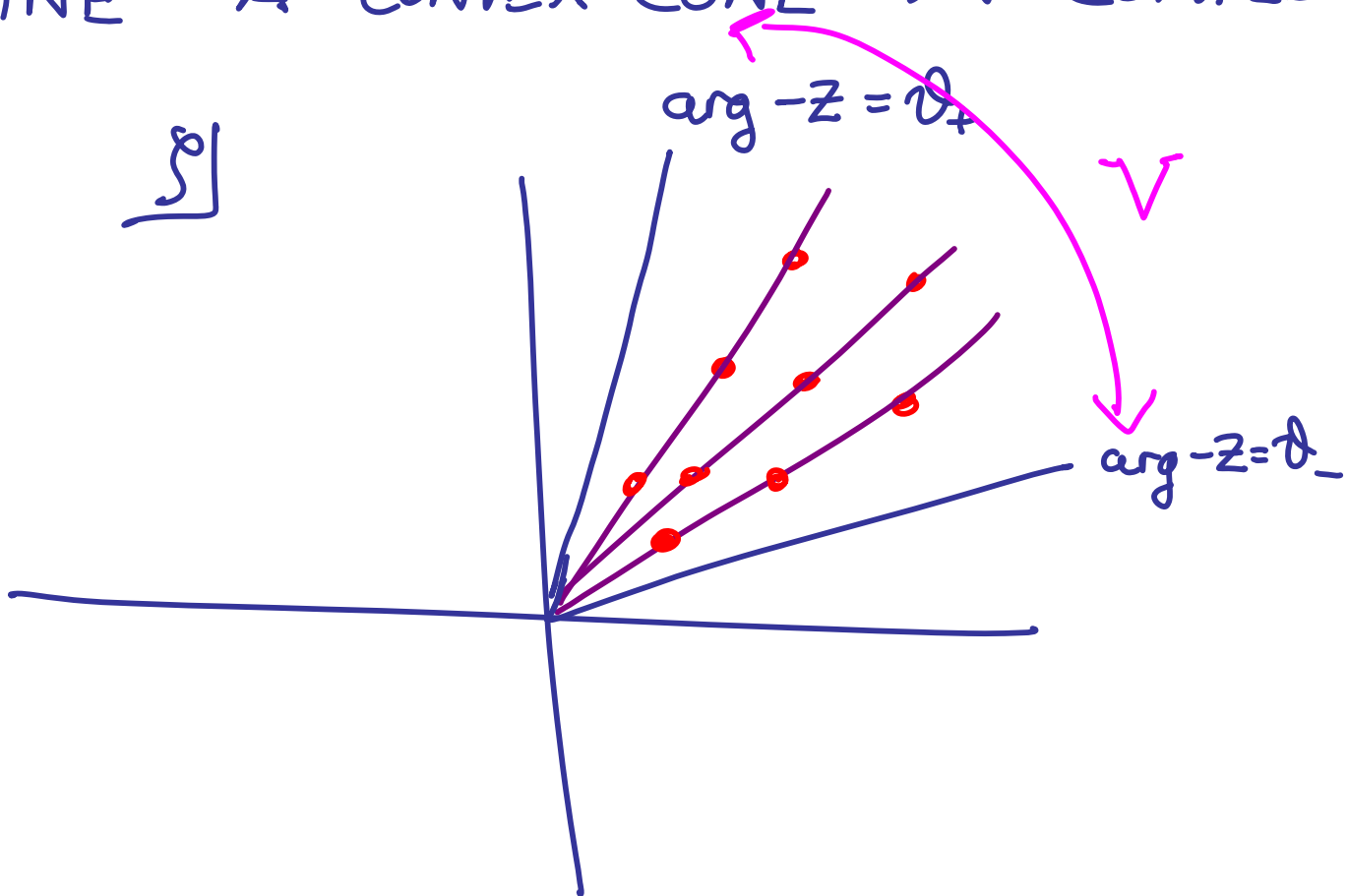
$$\chi^{T_{+\infty}} \longrightarrow \chi^{T_{-\infty}}$$



# 9. WALL CROSSING

CHOOSE  $\vartheta_- < \vartheta_+$  TO

DEFINE A CONVEX CONE IN COMPLEX



SUPPOSE WE FOLLOW A PATH  
 $u_-$  TO  $u_+$  SO THAT NO BPS RAY  
CROSSES  $\arg(-z) = \vartheta_{\pm}$ .

THEN  $T(\vartheta_{\pm}, \lambda_-)$  SMOOTHLY  
EVOLVES TO  $T(\vartheta_{\pm}, \lambda_+)$

ON THE OTHER HAND,  
 EVOLVING  $\vartheta_-$  TO  $\vartheta_+$  AT  
 FIXED  $\lambda$  PRODUCES A SEQUENCE  
 OF FLIPS, TWISTS, AND POPS.

FACT: ALL POPS OCCUR IN  
 DEGENERATE TRIANGLES, AND THE  
 INDUCED TRANSFORMATION IS  $\pm 1$  FOR  
 SUCH POPS.

THEREFORE  $\chi^{\vartheta_+}$  IS RELATED  
 TO  $\chi^{\vartheta_-}$  VIA THE IMAGE OF

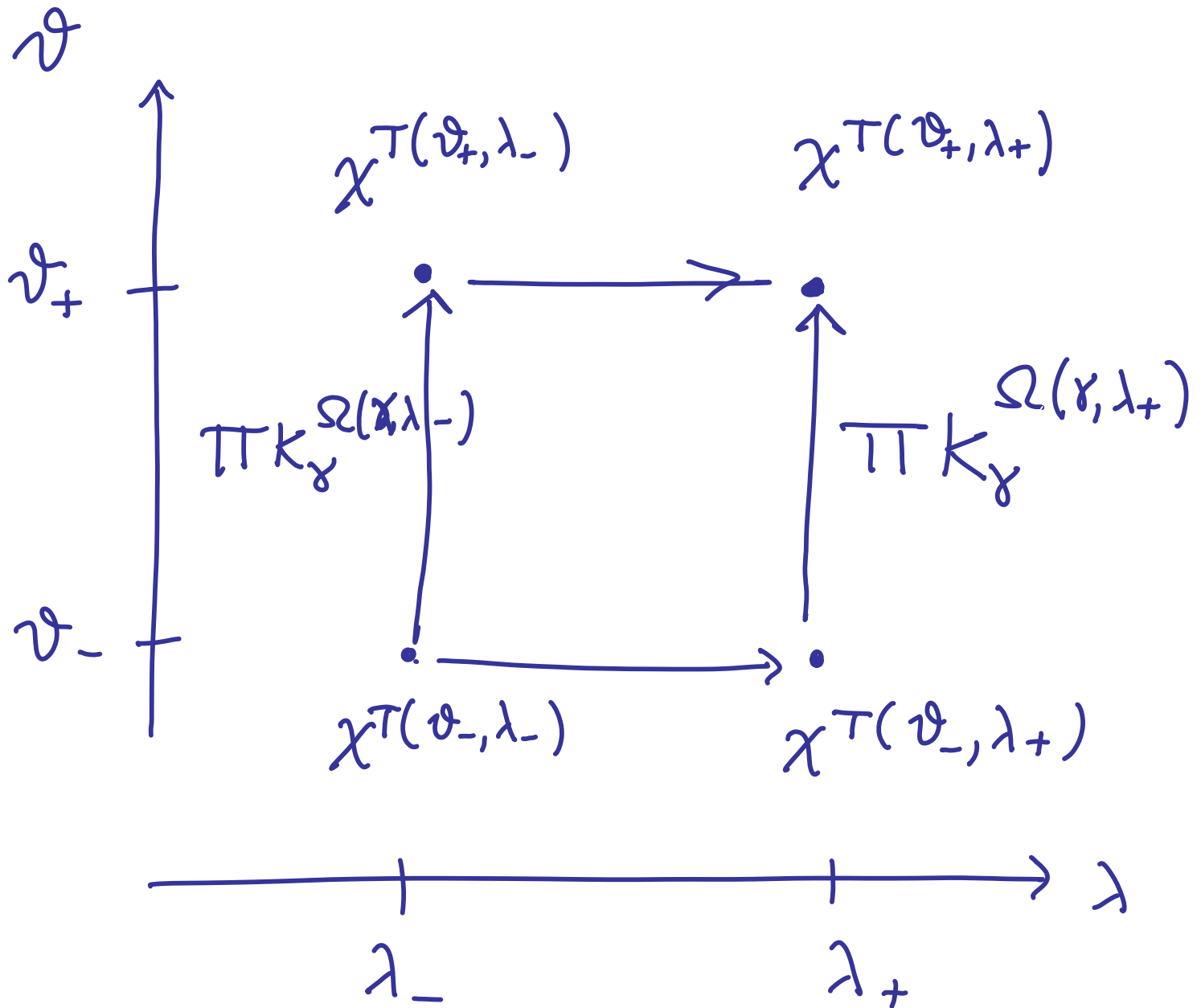
$$\prod_{\vartheta_- < \vartheta_c < \vartheta_+} \sigma_{\gamma^{\vartheta_c}} E_c$$

i.e.

$$\prod_{\vartheta_- < \arg(-z_\gamma) < \vartheta_+} K_\gamma^{\Omega(\gamma, \lambda)} = A_V$$



BUT THERE IS NO DISCONTINUITY  
IN  $\chi^T(\vartheta, \lambda_-) \rightarrow \chi^T(\vartheta, \lambda_+)$



$$\Pi K_\gamma \Omega(\gamma, \lambda_-) = \Pi K_\gamma \Omega(\gamma, \lambda_+)$$

MOVIES

# 10. DETERMINING THE BPS SPECTRUM

NOW LET US VARY  $\vartheta$  TO  $\vartheta + \pi$ .

WE CAPTURE ALL THE BPS STATES

$$\prod_{\vartheta_c} \sigma_{\gamma \vartheta_c}^{E_c} \longrightarrow \prod_{\vartheta < -\arg Z < \vartheta + \pi} K_{\gamma}^{\Omega(\gamma, \lambda)}$$

ON THE OTHER HAND,

$T(\vartheta, \lambda)$  AND  $T(\vartheta + \pi, \lambda)$

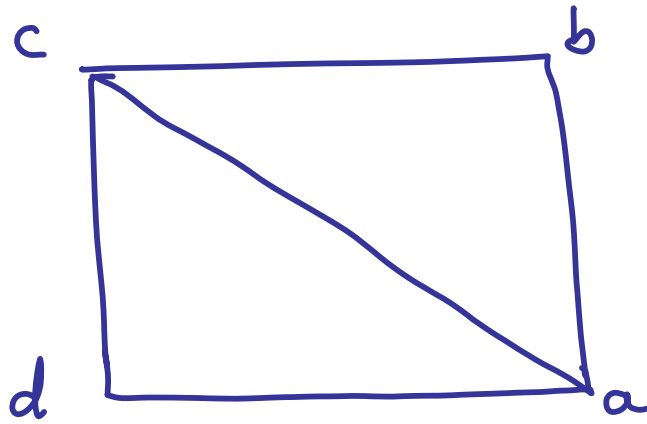
ONLY DIFFER BY

SIMULTANEOUSLY POPPING

ALL THE VERTICES!!

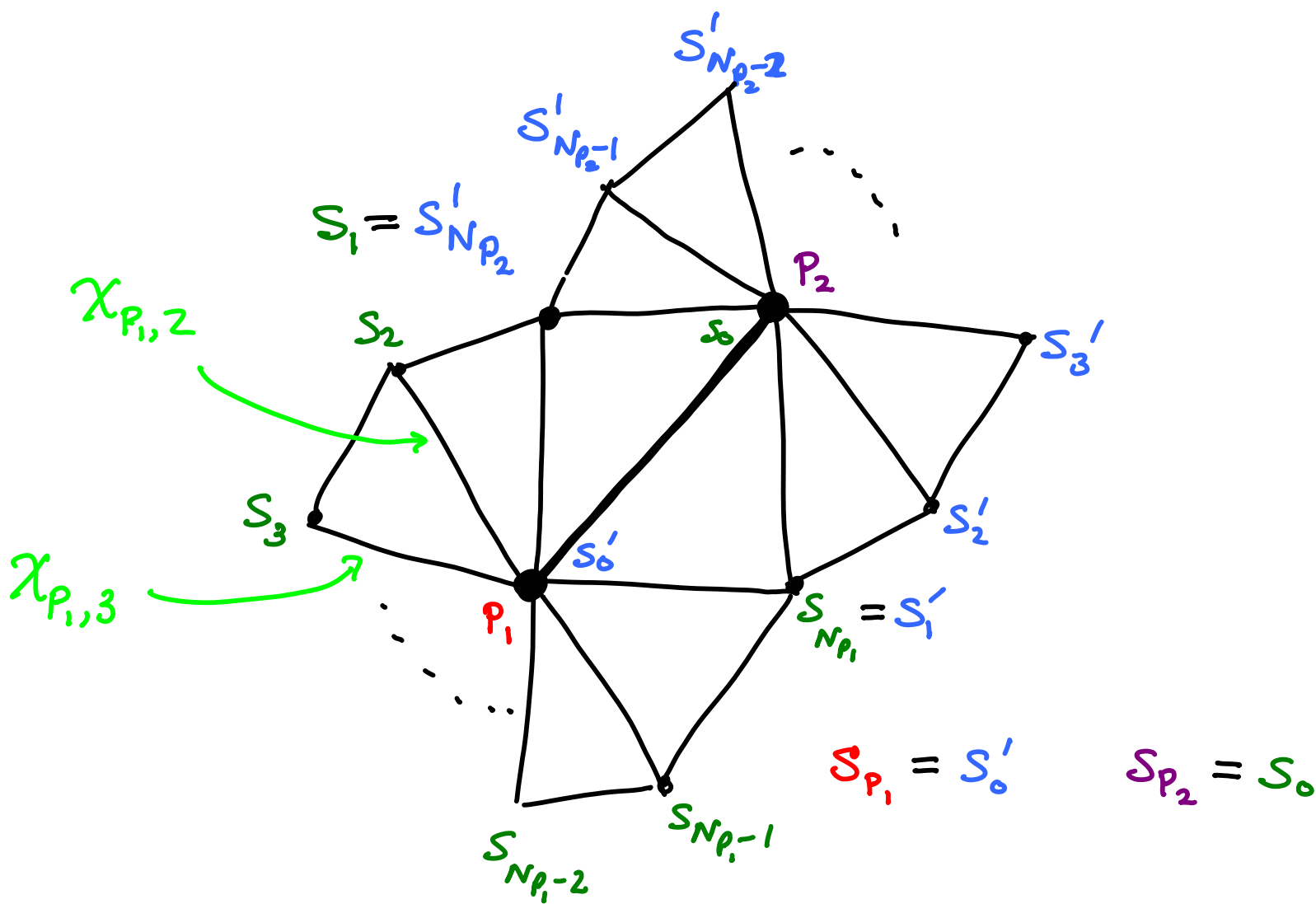
# UNIVERSAL STOKES MATRIX

WHILE THE CHANGE  $\chi_E^T$  FOR  
POPPING ONE VERTEX IS COMPLICATED,  
IT TURNS OUT THAT POPPING  
ALL VERTICES LEADS TO A RATHER  
SIMPLE FORMULA!



$$\tilde{\chi}_{ac}^T \chi_{ac}^T = \frac{(1 + A_{ab})(1 + A_{cd})}{(1 + A_{bc})(1 + A_{da})}$$

TO GIVE A FORMULA FOR  $A_{p_1 p_2}$ :



$$A_{P_1, P_2} = \frac{1}{1 - \mu_{P_1}^2} \frac{1}{1 - \mu_{P_2}^2} \cdot \chi_{P_1, P_2}$$

$$\left( 1 + \sum_{k=1}^{N_{P_1}-1} \prod_{j=1}^k \chi_{P_1, j} \right) \cdot \left( 1 + \sum_{k=1}^{N_{P_2}-1} \prod_{j=1}^k \chi_{P_2, j} \right)$$

TO FIND THE BPS SPECTRUM

THE TRANSFORMATION

$$S: \chi_i \longrightarrow \tilde{\chi}_i$$

$$\tilde{\chi}_i = \chi_i \frac{(1 + A_{ab}(i))(1 + A_{cd}(i))}{(1 + A_{bc}(i))(1 + A_{da}(i))}$$

HAS A UNIQUE DECOMPOSITION

OF THE FORM:

$$S = \prod_{\vartheta < -\arg z < \vartheta + \pi} K_{\gamma}^{\Omega(\gamma, \lambda)}$$

THIS DETERMINES THE  $\Omega(\gamma, u)$

# CONCLUSION: FUTURE DIRECTIONS

1. WE HAVE SOME IDEAS ABOUT HOW TO GO TO RANK  $k > 2$ .
2. RELATION TO INTEGRABLE SYSTEMS  
(e.g. THE INTEGRAL EQUATION FOR  $\chi_y$  IS A VERSION OF THE TBA.)
3. SUPERGRAVITY
4. NEW MODULAR FUNCTORS

