

# Desperately Seeking Moonshine

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a project with Jeff Harvey

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# Motivation

Search for a conceptual explanation of  
Mathieu Moonshine phenomena.

Eguchi, Ooguri, Tachikawa 2010

Proposal: It is related to the  
``algebra of BPS states."''

Something like:  $M_{24}$  is a distinguished  
group of automorphisms of the algebra  
of spacetime BPS states in some string  
compactification using  $K3$ .

# String-Math, 2014

Today's story begins in Edmonton, June 11, 2014.  
Sheldon Katz was giving a talk on his work with  
Albrecht Klemm and Rahul Pandharipande

He was describing how to count BPS states for type II  
strings on a K3 surface taking into account the  
 $so(4) = su(2) + su(2)$  quantum numbers of a particle  
in six dimensions.

Slide # 86 said ....

- Define the *refined K3 BPS invariants*  $R_{j_L, j_R}^{h \geq 0}$  by

$$\sum_{h=0}^{\infty} \sum_{j_L} \sum_{j_R} R_{j_L, j_R}^h [j_L]_u [j_R]_y q^h = \prod_{n=1}^{\infty} \frac{1}{(1 - u^{-1}y^{-1}q^n)(1 - u^{-1}yq^n)(1 - q^n)^{20}(1 - uy^{-1}q^n)(1 - uycq^n)}$$

where  $[j]_x = x^{-j} + \dots + x^j$

- For  $h \leq 2$  the nonvanishing invariants are

$$R_{0,0}^0 = 1, \\ R_{0,0}^1 = 20, \quad R_{\frac{1}{2}, \frac{1}{2}}^1 = 1, \\ R_{0,0}^2 = 231, \quad R_{\frac{1}{2}, \frac{1}{2}}^2 = 21, \quad R_{1,1}^2 = 1$$



# Heterotic/Type II Duality

$$\text{Het/T4} = \text{IIA/K3}$$

DH states: Perturbative  
heterotic BPS states



D4-D2-D0  
boundstates

Roughly: Cohomology groups of the moduli spaces of objects in  $D^b(\text{K3})$  with fixed K-theory invariant and stable wrt a stability condition determined by the complexified Kahler class.

**Aspinwall-Morrison Theorem:**

$$O_{\mathbb{Z}}(II^{20;4}) \setminus O_{\mathbb{R}}(20;4) / (O_{\mathbb{R}}(20) \times O_{\mathbb{R}}(4))$$

# Heterotic Toroidal Compactifications

$$\mathbb{M}^{1,1+d} \times T^{8-d}$$

$$II^{24-d,8-d} \hookrightarrow \Gamma^{24-d;8-d} \subset \mathbb{R}^{24-d;8-d}$$

$$P = (P_L; P_R) \in \Gamma^{24-d;8-d}$$

Narain moduli space of CFT's:

$$O_{\mathbb{Z}}(II^{24-d;8-d}) \backslash O_{\mathbb{R}}(24-d;8-d) / (O_{\mathbb{R}}(24-d) \times O_{\mathbb{R}}(8-d))$$

# Crystal Symmetries Of Toroidal Compactifications

Construct some heterotic string compactifications with large interesting crystallographic group symmetries.

$$G \subset \text{Aut}(\Gamma^{24-d;8-d})$$

$$G = G_L \times G_R$$

$$G_L \subset O_{\mathbb{R}}(24-d) \quad G_R \subset O_{\mathbb{R}}(8-d)$$

Then  $G$  is a crystal symmetry of the CFT:

Example: Weyl group symmetries of enhanced YM gauge theories.

These are NOT the kinds of crystal symmetries we want

# Conway Subgroup Symmetries

Start with a distinguished  $d=0$  compactification:

$$\mathbb{M}^{1,1} \times T^8$$

$$\Gamma^{24,8} = (\Lambda; 0) \oplus (0; \Gamma_8)$$

Crystal symmetry:

$$\text{Co}_0 \times W(E_8)$$

Note that  $\text{Co}_0$  is not a Weyl group symmetry of any enhanced Yang-Mills gauge symmetry.

Now “decompactify”



# A Lattice Lemmino

$\mathfrak{F}_L \subset \Lambda$  &  $\mathfrak{F}_R \subset \Gamma_8$  isometric of rank  $d$

Then there exists an even unimodular lattice with embedding

$$\Gamma^{24-d;8-d} \hookrightarrow \mathbb{R}^{24-d;8-d}$$

such that, if

$$G_L := \text{Fix}(\mathfrak{F}_L) \subset \text{Aut}(\Lambda) = \text{Co}_0$$

$$G_R := \text{Fix}(\mathfrak{F}_R) \subset \text{Aut}(\Gamma_8) = W(E_8)$$

then  $\Gamma^{24-d;8-d}$  has crystallographic symmetry

$$G_L \times G_R \subset O(24-d) \times O(8-d)$$

# Easy Proof

Does not use the Nikulin embedding theorem.

Uses standard ideas of lattice theory.

See book of Miranda and Morrison for these ideas.

$$\mathcal{D}_+(\mathfrak{F}_L^\perp) \cong \mathcal{D}_-(\mathfrak{F}_L) \cong \mathcal{D}_-(\mathfrak{F}_R) \cong \mathcal{D}_+(\mathfrak{F}_R^\perp)$$

$$\Gamma \subset (\mathfrak{F}_L^\perp)^\vee \oplus (\mathfrak{F}_R^\perp)^\vee \subset \mathbb{R}^{24-d; 8-d}$$

$$\Gamma = \{(x, y) \mid \bar{x} \cong \bar{y}\}$$

$$g : (x; y) \mapsto (g_L x; g_R y) \quad \overline{g_L x} = \bar{x} \quad \overline{g_R y} = \bar{y}$$

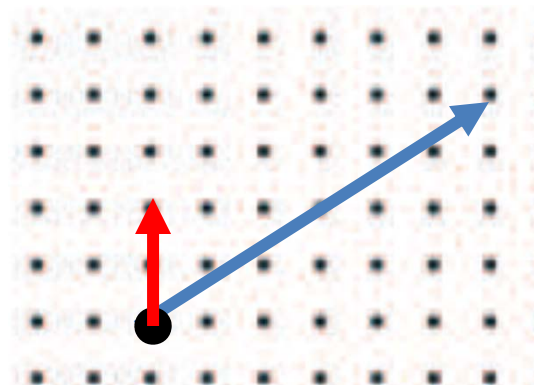
# CSS Compactifications

This construction defines points of moduli space with Conway Subgroup Symmetry:  
call these CSS compactifications.

What crystal symmetries can you get?

In general, a sublattice preserves none of the crystal symmetries of the ambient lattice.

Consider, e.g., the lattice generated by  $(p,q)$  in the square lattice in the plane.



# Fixed Sublattices Of The Leech Lattice

The culmination of a long line of work is the classification by Hohn and Mason of the 290 isomorphism classes of fixed-point sublattices of  $\Lambda$

221	3	24	$[2^3 3]$ (#3)	$8_3^{+3}$	0	4	16	1	1	3	$\text{Mon}_b^*$
222	2	9196830720	$U_6(2)$	$2_{II}^{-2} 3^{+1}$	0	1	1	1	1	-	$S^*$
223	2	898128000	$McL$	$3^{-1} 5^{-1}$	1	1	1	1	1	-	$S^*$
224	2	454164480	$2^{10}.M_{22}$	$4_2^{+2}$	0	1	1	1	1	-	$\text{Mon}_a^*$
225	2	44352000	$HS$	$2_2^{-2} 5^{+1}$	0	1	1	1	1	-	$S^*$
226	2	20643840	$2^9.L_3(4).2$	$4_1^{+1} 8_1^{+1}$	0	1	2	1	1	-	$\text{Mon}_a$
227	2	10200960	$M_{23}$	$23^{+1}$	1	1	1	1	2	1	$M_{23}^*$

# Symmetries Of D4-D2-D0 Boundstates

99	4	245760	$[2^{12}].A_5$	$2_{II}^{-2}4_{II}^{-2}$	0	1	1	1	1	-	$\text{Mon}_a^*$
100	4	30720	$[2^9].A_5$	$2_{II}^{-4}5^{-1}$	0	1	1	1	1	-	$\text{Mon}_a^*$
101	4	29160	$3^4.A_6$	$3^{+2}9^{+1}$	1	1	1	1	1	-	$S^*$
102	4	20160	$L_3(4)$	$2_{II}^{-2}3^{-1}7^{-1}$	2	1	1	1	2	1	$M_{23}^*$
103	4	12288	$[2^{12}3]$	$2_{II}^{+2}4_3^{+1}8_1^{+1}$	0	1	2	1	1	-	$\text{Mon}_a$
104	4	9216	$[2^{10}3^2]$	$2_{II}^{+4}3^{+2}$	0	1	2	1	1	-	$\text{Mon}_a^*$

Therefore the space of D4D2D0 BPS states on K3 will naturally be a representation of the subgroups of  $\text{Co}_0$  that preserve sublattices of rank 4.

These discrete groups will be automorphisms of the algebra of BPS states at the CSS points.

# Symmetries Of Derived Category

Theorem [Gaberdiel-Hohenegger-Volpato]: If  $G \subset O_{\mathbb{Z}}(20;4)$  fixes a positive 4-plane in  $\mathbb{R}^{20,4}$  then  $G$  is a subgroup of  $Co_0$  fixing a sublattice with  $4 \leq \text{rank}$ .

Remark: GHV generalize the arguments in Kondo's paper proving Mukai's theorem that the symplectic automorphisms of K3 are subgroups of  $M23$  with at least 5 orbits on  $\Omega$

Interpreted by Huybrechts in terms of the bounded derived category of K3 surfaces

$$G \cong \text{Aut}_{H^{2,0} \oplus H^{0,2}}(D^b(K3)) \cap \text{Aut}_{B+iJ}(D^b(K3))$$

# But Is There Moonshine In KKP Invariants?

99	4	245760	$[2^{12}].A_5$	$2_{II}^{-2}4_{II}^{-2}$	0	1	1	1	1	-	$\text{Mon}_a^*$
100	4	30720	$[2^9].A_5$	$2_{II}^{-4}5^{-1}$	0	1	1	1	1	-	$\text{Mon}_a^*$
101	4	29160	$3^4.A_6$	$3^{+2}9^{+1}$	1	1	1	1	1	-	$S^*$
102	4	20160	$L_3(4)$	$2_{II}^{-2}3^{-1}7^{-1}$	2	1	1	1	2	1	$M_{23}^*$
103	4	12288	$[2^{12}3]$	$2_{II}^{+2}4_3^{+1}8_1^{+1}$	0	1	2	1	1	-	$\text{Mon}_a$
104	4	9216	$[2^{10}3^2]$	$2_{II}^{+4}3^{+2}$	0	1	2	1	1	-	$\text{Mon}_a^*$

$$2^{12} : A_5 \cong 2^8 : M_{20}$$

So the invariants of KKP will show “Moonshine” with respect to this symmetry.....

But this is a little silly: All these groups are subgroups of  $O(20)$ . If we do not look at more structure, that includes the momenta/characteristic classes we might as well consider the degeneracies as  $O(20)$  representations.

# Silly Moonshine

$$\prod_{n=1}^{\infty} \left( \frac{1}{(1-q^n)^{20} (1-yzq^n) (1-yz^{-1}q^n) (1-y^{-1}zq^n) (1-y^{-1}z^{-1}q^n)} \right)$$

is just the SO(4) character of a Fock space of 24 bosons.

$$\iota : O(20) \times O(4) \hookrightarrow O(24)$$

$$\iota^*(V) = \mathbf{20} \otimes \mathbf{1} \oplus \mathbf{1} \otimes \mathbf{4} \quad \mathcal{F}_q(\mathbf{20} \otimes \mathbf{1} \oplus \mathbf{1} \otimes \mathbf{4})$$

$$\mathcal{F}_q(V) := \text{Sym}_q^*(V) \otimes \text{Sym}_{q^2}^*(V) \otimes \dots$$

All the above crystal groups are subgroups of O(20) so the "Moonshine" wrt those groups is a triviality.

IS THERE MORE GOING ON ??



# Baby Case: T7 & d=1

272	2	80	$[2^{45}]$ (#34)	$4_2^{+2}5^{+2}$	0	1	8	1	1	-	Mon <sub>a</sub>
273	1	$ Co_2 $	$Co_2$	$4_1^{+1}$	0	1	1	1	1	-	S*
274	1	$ Co_3 $	$Co_3$	$2_3^{-1}3^{-1}$	0	1	1	1	1	-	S*

Decompose partition function of BPS states wrt reps of transverse rotation group O(1)

$$\mathcal{F}_q(V_{23} \otimes \mathbf{T} \oplus \mathbf{1} \otimes \mathbf{S}) = \mathbf{1} \otimes \mathbf{T} \oplus q [V_{23} \otimes \mathbf{T} \oplus \mathbf{1} \otimes \mathbf{S}] \oplus q^2 [\mathbf{300} \otimes \mathbf{T} \oplus \mathbf{24} \otimes \mathbf{S}] \oplus q^3 [\mathbf{2876} \otimes \mathbf{T} \oplus \mathbf{324} \otimes \mathbf{S}]$$

$\oplus \dots$  These numbers dutifully decompose nicely as representations of  $Co_2$ :

That's trivial because  $Co_2 \subset O(23)$

$$\mathbf{300} \cong S^2 V_{23} \oplus V_{23} \oplus \mathbf{1}$$

But is there a  $Co_0 \times O(1)$  symmetry?  $Co_0$  is NOT a subgroup of  $O(23)$ .  $Co_0 \times O(1)$  symmetry CANNOT come from a linear action on  $V_{24}$ .

# The SumDimension Game

$$\text{Irrep}(\text{Co}_0) = \{1, 24, 276, 299, 1771, 2024, 2576, 4576, \dots\}$$

$$1 \otimes \mathbf{T} \oplus q[V \otimes \mathbf{T} \oplus 1 \otimes \mathbf{S}] \oplus q^2[300 \otimes \mathbf{T} \oplus 24 \otimes \mathbf{S}] \oplus q^3[2876 \otimes \mathbf{T} \oplus 324 \otimes \mathbf{S}] \oplus \dots$$

$$300 = 299 + 1$$

$$300 = 276 + 24$$

$$2876 = 2576 + 299 + 1$$

$$2876 = 2576 + 276 + 24$$

$$324 = 299 + 24 + 1$$

$$324 = 276 + 24 + 24$$

ETC.

# Defining Moonshine

Any such decomposition defines the massive states of  $\mathcal{F}_q(V)$  as a representation of  $Co_0 \times O(1)$ .

**Problem:** There are infinitely many such decompositions!  
What physical principle distinguishes which, if any, are meaningful?

Definition: You have committed **Moonshine** (for  $d=1$ ) if you exhibit the massive sector of  $\mathcal{F}_q(V)$  as a representation of  $Co_0 \times O(1)$  such that the graded character of any element  $g$ :

$$\text{Tr}_{\mathcal{F}_q(V)} g q^{L_0 - 1}$$

is a modular form for  $\Gamma_0(m)$  where  $m = \text{order of } g$ .

# Virtual Representations

Most candidate  $Co_0 \times O(1)$  representations will fail to be modular.

But if we allow virtual representations:

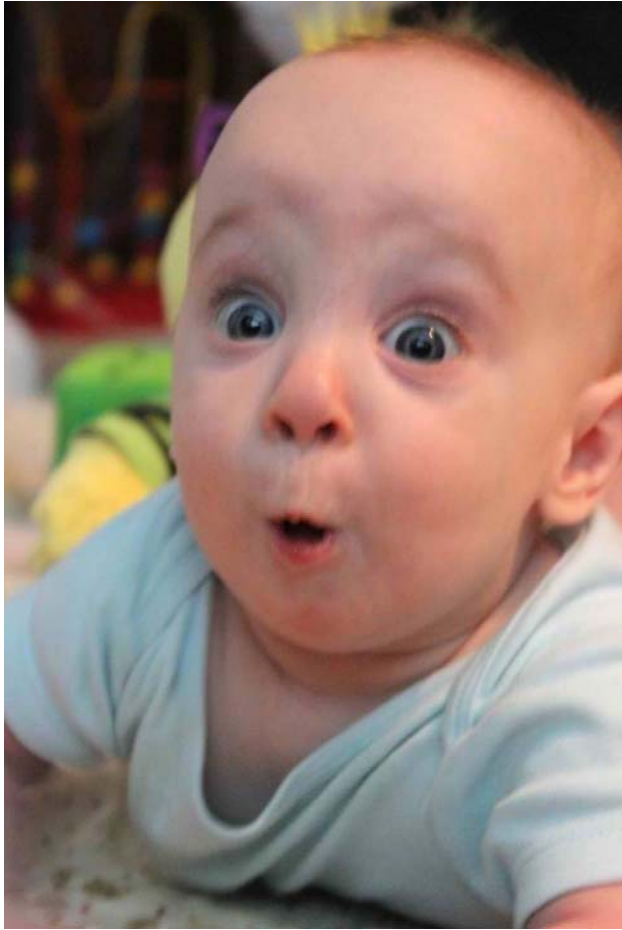
$$V_{23} \rightarrow V_{24} - \mathbf{1}$$

$$\mathcal{F}_q(V_{23} \otimes \mathbf{T} \oplus \mathbf{1} \otimes \mathbf{S}) \rightarrow \mathcal{F}_q(V_{24}) \otimes \frac{\mathcal{F}_q(\mathbf{1} \otimes \mathbf{S})}{\mathcal{F}_q(\mathbf{1} \otimes \mathbf{T})}$$

The characters are guaranteed to be modular!

But massive representations have no reason to be true representations.

In fact, the negative representations cancel and ALL the massive levels are in fact true representations!!



But! The same argument also shows they are also true representations of  $O(24) \times O(1)$ .



# Lessons

Modularity of characters is crucial.

Virtual Fock spaces are modular.

There can be nontrivial cancellation of the negative representations.

A “mysterious” discrete symmetry can sometimes simply be a subgroup of a continuous symmetry.

# What About d=4 ?

$$\mathrm{Co}_0 \times O(1) \rightarrow M_{24} \times O(4)$$

$$\mathcal{F}_q(V_{20} \otimes \mathbf{1} \oplus \mathbf{1} \otimes V_4) \rightarrow \mathcal{F}_q(V_{23}) \otimes \frac{\mathcal{F}_q(\mathbf{1} \otimes V_4)}{\mathcal{F}_q(\mathbf{1} \otimes \mathbf{1})^3}$$

Magical positivity fails:

$$\dim R_{0,0}^2 = 231$$

$$R_{0,0}^2 = V_{252} - V_{23} + 2V_1$$

*But we are desperately seeking Moonshine...*

So we ask: Could it still be that, magically, some positive combination of representations from the SumDimension game is nevertheless modular?

# Characters Of An Involution

$$\begin{aligned}Z_{2A} &= 8 + 1/q + 36q + 144q^2 + 282q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 426q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 218q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 362q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 266q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 410q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 202q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 346q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 378q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 522q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 314q^3 \\ &= 8 + 1/q + 36q + 144q^2 + 458q^3\end{aligned}$$

Should be modular  
form for  $\Gamma_0(2)$  .

Weight?

(assumed  
half-integral)

Multiplier  
system?



# What Is Your Weight?

$$\tau_0 = \frac{1}{2}(1 + i) \quad (ST^2S)\tau_0 = \tau_0 - 1$$

One can deduce the multiplier system from the weight, and derive the weight numerically from:

$$w = \lim_{\epsilon \rightarrow 0} \frac{1}{2\epsilon} \log \left| \frac{Z(\tau_0 - i\epsilon)}{Z(\tau_0 + i\epsilon)} \right|$$
$$q_0 = e^{2\pi i \tau_0} = -e^{-\pi} = -0.04\dots$$

Convergence is good so can compute the weight numerically. For  $Z_{2A}$  it converges to -8.4..... Not pretty. Not half-integral !!


No positive combination of reps is modular. No M24 Moonshine.



# Application To Heterotic-Type II Duality

Existence of CSS points have some interesting math predictions.

$\mathfrak{X}$  K3 and elliptically fibered CY3

$\text{Het}/T^2 \times K3'$    $\text{IIA}/\mathfrak{X}$

Perturbative  
heterotic  
string states



Vertical  
D4-D2-D0  
boundstates

# Orbifolds At CSS Points

Orbifold at the CSS points for  $d=2$  ( $T^6$  compactification)

For simplicity:  $\mathbb{Z}_2$  orbifold

$$X \rightarrow RX + \delta$$

$$R = (g_L; g_R) \in G_L \times G_R$$

$g_R$  an involution in  $W(E_8)$  with eigenvalues  $-1^4, +1^4$

We can realize all  
but 6 of the 51  
rank 2 HM  
classes:

$(q_{11}, q_{22}, q_{12})$	HM Number	$\sigma_1 ?$	$\sigma_2 ?$
(4, 4, 0)	224	Y	
(4, 4, 1)	223	N	N
(4, 4, 2)	222	Y	
(4, 6, 0)	229	Y	
(4, 6, 1)	227	N	N
(4, 6, 2)	225	Y	
(4, 8, 0)	226, 236	Y	
(4, 8, 2)	232	N	N
(4, 10, 2)	241	Y	
(4, 12, 0)	234	Y	
(4, 12, 2)	233	Y	
(4, 16, 0)	250	Y	
(4, 16, 2)	237	N	N
(4, 20, 0)	257	Y	
(4, 24, 0)	244	Y	
(6, 6, 0)	228, 242	Y	
(6, 6, 2)	235	Y	
(6, 6, 3)	230	Y	
(6, 8, 2)	251	Y	
(6, 10, 2)	243	Y	
(6, 12, 0)	253	Y	
(6, 12, 3)	240	N	N
(6, 18, 0)	263	Y	
(8, 8, 0)	238, 258	Y	
(8, 8, 2)	246	N	N
(8, 8, 4)	231, 252	Y	
(8, 12, 0)	248	Y	
(8, 12, 4)	239, 249	Y	
(8, 16, 0)	269	Y	
(8, 20, 4)	266	Y	
(10, 10, 0)	262	Y	
(10, 10, 2)	259	Y	
(10, 12, 6)	256	Y	
(12, 12, 0)	247, 254, 261, 271	Y	
(12, 12, 2)	255	Y	
(12, 12, 6)	245, 260	Y	
(12, 18, 6)	265	Y	
(12, 24, 0)	270	Y	
(20, 20, 0)	272	Y	
(20, 20, 10)	264	Y	
(24, 24, 12)	267, 268	Y	

# Explicit Example – 1/2

Gram matrix of  $\mathfrak{F}_R = \begin{pmatrix} 4 & 0 \\ 0 & 4 \end{pmatrix}$

Can embed in Leech and E8 lattices

There exists  $g_R$  in  $W(E8)$  fixing  $\mathfrak{F}_R$  with ev's  $+1^4, -1^4$ . Mod out by this on the right.

Need to choose involution  $g_L$ .

Structure of Golay code implies the only possibilities:

$$g_L \sim \text{Diag}\{-1^8, +1^{16}\} \quad g_L \sim \text{Diag}\{-1^{12}, +1^{12}\} \quad g_L \sim \text{Diag}\{-1^{16}, +1^8\}$$

Flipping 8,12,16 coordinates  $x^i$  according to octad/dodecad/  
octad complement  $\mathcal{C}$ -sets of the Golay code

# Explicit Example – 2/2

Massless sector:

Abelian gauge symmetry.

	$VM = h^{1,1}(\mathfrak{X})$	$HM = h^{2,1}(\mathfrak{X}) + 1$	$\chi = 2(1 + VM - HM)$
$A_1$	15	$256 + 8$	-496
$A_2$	15	$64 + 8$	-112
$B$	7	528	-1040
$C$	11	12	0

Personal Ad: Two Hodge numbers seek friendly compatible CY3...

(7,527)? Huh?  $\chi(\mathfrak{X}_B) = -1040$

TaxMan violates Wati's bound:  $h^{1,2}(\mathfrak{X}) \leq 491$

# Generalized Huybrechts Theorem

$$\text{Aut}_\sigma(D^b(\mathfrak{X})_{\text{vertical}})$$

Should be the centralizer of an element of a subgroup of  $\text{Co}_0$  fixing a rank two sublattice of  $\Lambda$



# Some General Questions

Clarify TaxMan example: Should every heterotic model on  $K3 \times T^2$  have a type II dual?

Decreasing  $d \implies$  larger CSS groups: raises a general question about D-brane categories: What to do on  $\mathbb{R}^3 \times S^1$ ?  
There ARE A & B Models!

In the type II interpretation CSS only arise for special values of the flat RR fields. How do flat RR fields affect the BPS D-brane states?

$$\langle x | \text{Image of Dave} \rangle = e^{ip \cdot x}$$

$$p \in \Gamma^{20,4}$$



**HAPPY BIRTHDAY DAVE!!**