We report low temperature high magnetic field scanning tunneling microscopy and spectroscopy of graphene flakes on graphite that exhibit the structural and electronic properties of graphene decoupled from the substrate. Pronounced peaks in the tunneling spectra develop with increasing field revealing a Landau level sequence that provides a direct way to identify graphene and to determine the degree of its coupling to the substrate. The Fermi velocity and quasiparticle lifetime, obtained from the positions and width of the peaks, provide access to the electron-phonon and electron-electron interactions.

Graphene, a one-atom thick form of crystalline-carbon [1–3] possesses extraordinary electronic properties owing to charge-carriers that move like ultrarelativistic particles (massless) Dirac-fermions (DF). Their relativistic nature becomes evident in the presence of a magnetic field through the appearance of an unusual Landau level (LL) sequence [2] with square-root dependence on field, \( B \), and level index \( n \) together with a unique field independent \( n = 0 \) level at the Dirac point (DP). One of the first manifestations of the \( n = 0 \) level was the appearance of half integer quantum Hall plateaus [4,5]. Subsequently, scanning tunneling microscopy (STM) of graphene deposited on insulating substrates [6] and zero-field scanning tunneling spectroscopy (STS) [7] showed surprising deviations from theoretical expectations prompting concerns over the interpretation of the STS due to the poor screening provided by the insulating substrate [8].

Here we study graphene flakes deposited on a graphite substrate which provides the necessary screening to ensure reliable STM and STS. We demonstrate that in this system it is possible to access the intrinsic properties of the quasiparticle carriers in graphene [9] and to determine their degree of coupling to the substrate [10]. In zero field we find that the density of states (DOS) is linear and vanishes at the DP as expected of massless DF [2]. In finite field we observe the appearance of a sharp sequence of LL with square-root dependence on field and LL index together with the distinct \( n = 0 \) level. We further show that the DFs are slowed down by electron-phonon interactions and their lifetime is limited by electron-electron interactions.

Experiments were conducted in a home built, low temperature (\( T = 4.4 \) K for this work) high field STM [11] using mechanically cut Pt-Ir wire tips. The tunneling conductance, \( dI/dV \), was measured by lock-in detection with 340 Hz bias voltage modulation. Samples were prepared by cleaving highly oriented pyrolytic graphite (HOPG) and the STM coarse motors were used to spot graphene flakes.

Once a flake is found its structure and electronic properties are studied to determine its degree of coupling to the graphite substrate. STM topography of such a flake, Fig. 1(a), reveals three layers separated in height by atomic steps. The top layer in region \( A \) exhibits the distinctive honeycomb structure, while that in region \( B \) is triangular as expected of Bernal-stacked graphite [12]. Thus the images in Fig. 1 suggest that in \( B \) the top layer is coupled to the
underlying layers, while in A, separated from B by a macroscopic ridgelike defect, the top layer is a single layer of graphene (SLG) decoupled from the substrate. Further confirmation that region A consists of SLG is gained by accessing the DOS through STS [13].

Because of the unique band structure of graphene the DOS in SLG is linear in energy and vanishes at the DP [2]. This is not the case when two or more graphene layers are coupled together because interlayer coupling produces additional states at the DP leading to a finite DOS there. Therefore a linear DOS that vanishes at the DP is a signature of massless DF and can thus be used to identify SLG decoupled from the substrate. A more stringent test of isolated SLG is to apply a magnetic field normal to the layer and probe for the appearance of a sequence of LL that is unique to massless DF [2]:

$$E_n = E_D + \text{sgn}(n)\sqrt{2e^2n^2|n|B} \quad n = 0 \pm 1, \pm 2, \ldots$$  \hspace{1cm} (1)$$

where $E_D$ is the energy at DP, $v_F$ the Fermi velocity, $-e$ the electron charge, with $n > 0$ corresponding to electrons and $n < 0$ to holes. Any coupling to the substrate introduces extra levels and eventually smears out the LL sequence [14–16]. Thus the square-root dependence of the LL energy on both $B$ and $n$ uniquely identifies decoupled graphene layers and was used in our search of decoupled flakes. Several completely decoupled flakes were found. Most of the data reported here is on one sample that was studied most extensively.

In Fig. 2(a) we show the bias dependence of the tunneling conductance, a quantity proportional to the local DOS [13]. In A the DOS is $V$-shaped and vanishes at the DP, as expected for SLG and in agreement with the honeycomb structure observed in topography. We note that the DP is offset by $\sim 16$ meV above the Fermi energy (identified with zero bias) indicating unintentional doping with a concentration of $\sim 2 \times 10^{16}$ cm$^{-2}$ hole carriers. In contrast the DOS in B remains finite at all energies indicating coupling to the substrate. The difference between the two regions becomes evident in the low energy tunneling conductance map, Fig. 2(b), showing that the ridge acts as a border between a region of decoupled SLG and one of coupled layers.

To further ascertain that region A corresponds to decoupled SLG we carried out STS in the presence of a magnetic field normal to the surface. As shown in Fig. 2(c) the spectra develop into a sequence of well-defined LL peaks with increasing field. In A the peaks are substantially more pronounced than the peaks reported in previous studies [11,15] where the spectral weight of LL was reduced by substrate coupling. The massless Dirac fermion character of the spectra in A is revealed by plotting the peak energies against the reduced variable $(ln|B|)^{1/2}$. This scaling procedure collapses all the data unto a straight line, Fig. 2(d), attesting to the Dirac fermion nature of the charge carriers. Comparing to Eq. (1) we obtain (from the slope) the value of the Fermi velocity, $v_F = 0.79 \times 10^6$ m/s. In addition we find the position of the Dirac energy relative to the Fermi level, $E_D = 16.6$ meV is consistent with the zero-field data, Fig. 2(a). The scaling of the LL sequence holds across the entire sample in A [17]. In contrast, no such collapse is seen in region B where graphene is coupled to the substrate.

The value of $v_F$ obtained from the level sequence is $\sim 20\%$ below the tight-binding value [2] and than that measured in weakly coupled layers on graphite [11,20]. Below we show that this result can be attributed to electron-phonon ($e$-$ph$) interactions [22,23]. According to density functional theory (DFT) calculations [25], $e$-$ph$ interactions introduce additional features in the electron self-energy, leading to a renormalized velocity $v_F = v_F0(1 + \lambda)^{-1}$ and to dips in the renormalization factor, $(v_F - v_F0)/v_F$, at energies $E_F \pm \hbar \omega_{ph}$. Here $v_F0$ is the bare Fermi velocity, $\lambda$ the $e$-$ph$ coupling constant and $\hbar \omega_{ph}$ the energy of the most strongly coupled phonons. The $e$-$ph$ coupling also introduces a kink in the zero-field DOS at

![Figure 2](image-url)
$h\omega_{\text{ph}}$. Examining the zero-field spectra in Fig. 2(a), we note a strong shoulderlike feature around $\pm 150$ mV in the spectra in region A (but not B), consistent with the e-ph coupling scenario. A more detailed analysis of the e-ph coupling is discussed next.

The LL sequences provide strong evidence that the quasiparticle excitations in graphene are 2$d$ massless Dirac fermions and give the average value of $v_F$ for energies up to 150 meV. Using these facts and the linearity of the DOS [2], we obtain the zero-field calibration of the DOS, $\rho(E) = \frac{3\pi^2 a}{\hbar^2 v_F^2} |E-E_D| = 0.123|E-E_D|$, below 150 meV ($a = 1.42$ Å). Comparison between this DOS and the normalized zero-field tunneling spectra in Fig. 3(a) yields a remarkable consistency justifying the assumption $dI/dV \propto DOS$ [13] within the range $-150$ to 150 mV. Assuming this proportionality remains valid to the dispersion relation the measured differential conductance curve, and from it we obtain the zero-field calibration of the DOS [13] within the range $-150$ to 150 mV. Using these facts and the linear-size dependence of the Fermi velocity, $v_F \approx \frac{dE}{dK}$, shown in Fig. 3(c). Identifying the asymptotic value of the velocity with $v_{F0}$ we obtain $A \sim 0.26$. The curve bears a striking similarity to that obtained by DFT [25]: it exhibits two dips at the energy of the $A_{1g}$ phonon $\pm 150 (\pm 20)$ meV, suggesting that this phonon, which couples the $K$ and $K'$ valleys [Fig. 3(b) inset] and undergoes a Kohn anomaly, is an important player in the velocity renormalization. Further support for the role of e-ph coupling comes from ARPES measurements of graphene on SiC which revealed similar kinks in the dispersion and a down renormalization of $v_F$ [26].

We now turn to high resolution (2 mV) LL spectroscopy at 4 T shown in Fig. 4(a). The sharpness and high definition of the LL peaks made it possible to extract the energy dependence of the quasiparticle lifetime. The level sequence can be fit with a sum of peak functions centered at the measured peak energies with the line width of each peak and an overall amplitude as free parameters. Interestingly, we found that Lorentzian fits were significantly better than Gaussians suggesting an intrinsic lifetime rather than impurity broadening. We focus on the hole sector in Fig. 4(b) where the energy range covered by the LL peaks is wider than on the electron side. The line widths were found to increase linearly with energy corresponding to an inverse lifetime: $\frac{1}{\tau} = \frac{E_n}{\tau_0} + \frac{1}{\tau_F}$, where $E_n$ is the LL energy in units of eV, $\gamma \sim 9$ fs/eV and $\tau_0 \sim 0.5$ ps at Fermi level. The latter translates into a mean free path of $l = v_F \tau_0 \sim 400$ nm which is comparable to the sample size.

![FIG. 3 (color online). Zero-field tunneling spectra, energy-momentum dispersion and Fermi velocity. (a) Tunnelling spectra and density of states (see text). Thick line is the DOS calculated using $v_F$ obtained in Fig. 2(d). Note the consistency of spectra for tunneling resistance varying from 3.8 GΩ to 50 GΩ. Legend shows tunneling junction settings. (b) Energy-momentum dispersion obtained as described in the text. Inset: Diagram of intervalley scattering mediated by $A_{1g}$ phonon. (c) Energy dependence of Fermi velocity.](image)

![FIG. 4 (color online). LL sequence, quasiparticle lifetime and gap at the DP. (a) High resolution tunneling spectrum at 4 T showing the LL. The dashed line marks the DP obtained by fitting the nonzero LL sequence. The $n = 0$ LL splits into two peaks, $0-$ and $0+$. (b) Zoom into LL sequence in hole sector. Solid squares are data points and the line represents a fit [29] with Lorentzian peaks. Inset, energy dependence of the peak width. (c) High resolution zoom into zero-field spectrum showing the appearance of a $\sim 10$ meV gap at the DP and an anomaly at the Fermi energy. (d) (top) Field dependence of the $0-$ and $0+$ Landau levels. Solid symbols: data from (a); open symbols from Fig. 2(c). (bottom) Field dependence of splitting. Experimental settings: [300 mV, ac 2 mV, 53 pA] in (a) and [200 mV, ac 1 mV, 53 pA] in (b).](image)
Several mechanisms leading to a finite quasiparticle lifetime are discussed in the literature. (i) Disorder induces a finite elastic lifetime at low energies and leads to decreasing LL width away from the Fermi level [27], a trend which is opposite to what is observed here. This is also consistent with our STM topography that shows no signature of structure disorder in the area where the tunneling spectra were taken. (ii) Broadening due to electron-phonon scattering is expected to be negligible for energies $E < 150$ meV, below the typical $A_{1g}$ phonon energy, the main player in the velocity renormalization [25]. (iii) Electron-electron interactions in graphene give rise to a linear energy dependence of the inverse lifetime characteristic of non-Fermi liquid behavior [28]. Theoretical estimates for the electron-electron interaction limited lifetimes in zero field give $\gamma \sim 20$ fs/eV. Since the electron-electron interactions are enhanced in magnetic field, it is possible that the agreement would be even better with calculations made in finite field.

A new feature revealed in the high resolution spectrum, Fig. 4(a), is the splitting of the $n = 0$ LL. As we show in Fig. 4(d) the splitting is of order $\sim 10$ mV and depends rather weakly on field. High resolution zero-field spectroscopy, Fig. 4(c), reveals a 10 mV gap at the DP suggesting that it may be related to the splitting in finite field. The fact that the gap opens at the DP rather than the Fermi energy implies that it is driven by broken sublattice symmetry, but more work is needed to elucidate the symmetry breaking mechanism and the splitting of the $n = 0$ level in magnetic field. Work is in progress to address these questions.

The tunneling microscopy and spectroscopy experiments presented here demonstrate that graphene flakes deposited on graphite can be sufficiently decoupled from the substrate so as to exhibit the structural and electronic properties expected of pristine graphene: a coherent honeycomb structure extending over the entire sample and a V shaped density of states that vanishes at the Dirac point. In a magnetic field the development of a single sequence of pronounced Landau level peaks corresponding to massless Dirac fermions provides compelling evidence of the decoupling. The LL spectra give access to important many-body effects including electron-phonon interactions which produce a downward normalization of the Fermi velocity, and electron-electron interactions reflected in quasiparticle lifetimes that are inversely proportional to energy.

We thank E. Abrahams, A. V. Balatsky, A. H. Castro Neto, J. C. Davis, D. Langreth, and M. Tosatti for useful discussions. Work supported by DE-FG02-99ER45742, NSF-DMR-045673 and Lucent.


[17] Variations of $E_p$ across the sample, $\pm 2.5$ mV, are substantially smaller than in graphene supported on SiO$_2$ substrates [18], suggesting that the insulating substrates induce large potential fluctuations as was demonstrated by eliminating the substrate [19].


[20] Comparison to values obtained on other substrates requires taking into account differences in screening. On SiO$_2$ where screening is weaker than on graphite $v_F$ should be larger [21], consistent with our observation.


[23] A downward renormalization of $v_F$ is expected when the relative orientation of adjacent graphene layers is twisted [24]. The twist scenario also implies a superstructure leading to Moire patterns. Their absence in these data suggests that twisting is not the cause of velocity renormalization.


