# **Outline of the Topics**

- 1. Trends in the periodic table.
- 2. Introduction: Heavy Fermions and the Kondo Lattice.
- 3. From Anderson to Kondo
- 4. Kondo Insulators: the simplest heavy fermions.
- 5. Oshikawa's Theorem.
- 6. Large N expansion for the Kondo Lattice
- 7. Heavy Fermion Superconductivity
- 8. Topological Kondo Insulators
- 9. Co-existing magnetism and the Kondo Effect.

# Please ask questions!



$$H = \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} + J \sum_{j} c_{a}^{\dagger}(j) c_{b}(j) S^{ba}(j)$$

 $c_{\sigma}^{\dagger}(j) = \frac{1}{\sqrt{V}} \sum_{\mathbf{k}} c_{\mathbf{k}\sigma} e^{i\mathbf{k}\cdot\mathbf{R}_{j}}$ 



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### Wild quantum fluctuations!

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Doniach 77, Lacroix and Cyrot 79, Coleman 83, Read and Newns 83, Auerbach and Levin 86.

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$$U \in \left(-\frac{1}{2}, \frac{1}{2}\right) \longrightarrow \left(-\frac{1}{2}, \frac{1}{2}\right)$$
$$H = \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} + \frac{J}{N} \sum_{j} c_{a}^{\dagger}(j) c_{b}(j) S^{ba}(j)$$

 $c_{\sigma}^{\dagger}(j) = \frac{1}{\sqrt{V}} \sum c_{\mathbf{k}\sigma} e^{i\mathbf{k}\cdot\mathbf{R}_{j}}$ 

### Large N expansion.

$$Z = \int_{\text{Fields}} e^{-NS[\psi]}$$



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$$H_{I}(j) = -\frac{J}{N} \left( c_{j\beta}^{\dagger} f_{j\beta} \right) \left( f_{j\alpha}^{\dagger} c_{j\alpha} \right)$$

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 $\begin{array}{l} Constraint \ n_f = Q = qN \\ \text{all terms extensive in N} \end{array}$ 

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$$-gA^{\dagger}A \rightarrow A^{\dagger}V + \bar{V}A + \frac{\bar{V}V}{g}$$





$$H[V,\lambda] = \sum_{\mathbf{k}\sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} + \sum_{j} \left( H_{I}[V_{j},j] + \lambda_{j}[n_{f}(j) - Q] \right),$$
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U(1) constraint: note  $n_f = Q = (qN)$ 



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$$=\operatorname{Tr}\left[\operatorname{Texp}\left(-\int_{0}^{\beta}H[V,\lambda]d\tau\right)\right]$$

$$Z = \int \mathcal{D}[V,\lambda] \int \mathcal{D}[c,f] \exp\left[-\int_{0}^{\beta}\left(\sum_{k\sigma}c_{k\sigma}^{\dagger}\partial_{\tau}c_{k\sigma} + \sum_{j\sigma}f_{j\sigma}^{\dagger}\partial_{\tau}f_{j\sigma} + H[V,\lambda]\right)\right]$$

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$$= \operatorname{Tr} \left[ \operatorname{Texp} \left( -\int_{0}^{\beta} H[V,\lambda] d\tau \right) \right] \text{Extensive in N}$$

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$$V_j = V$$
  
at each site



$$H_{MFT} = \sum_{\mathbf{k}\sigma} \left( c^{\dagger}_{\mathbf{k}\sigma}, f^{\dagger}_{\mathbf{k}\sigma} \right) \underbrace{\left( \begin{array}{c} \epsilon_{\mathbf{k}} & V \\ \overline{V} & \lambda \end{array} \right)}_{\mathbf{k}\sigma} \left( \begin{array}{c} c_{\mathbf{k}\sigma} \\ f_{\mathbf{k}\sigma} \end{array} \right) + N\mathcal{N}_{s} \left( \frac{|V|^{2}}{J} - \lambda q \right)$$
$$= \sum_{\mathbf{k}\sigma} \psi^{\dagger}_{\mathbf{k}\sigma} \underline{h}(\mathbf{k}) \psi_{\mathbf{k}\sigma} + N\mathcal{N}_{s} \left( \frac{|V|^{2}}{J} - \lambda q \right).$$

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$$f^{\dagger}_{\vec{k}\sigma} = \frac{1}{\sqrt{n}} \sum_{j} f^{\dagger}_{j\sigma} e^{i\vec{k}\cdot\vec{R}_{j}}$$

$$\begin{split} H_{MFT} &= \sum_{\mathbf{k}\sigma} \left( c^{\dagger}_{\mathbf{k}\sigma}, f^{\dagger}_{\mathbf{k}\sigma} \right) \underbrace{\begin{pmatrix} \mathbf{k} \\ \mathbf$$

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$$\operatorname{Det}\left[E_{\mathbf{k}}^{\pm}\underline{1} - \begin{pmatrix}\epsilon_{\mathbf{k}} & V\\ \overline{V} & \lambda\end{pmatrix}\right] = (E_{\mathbf{k}\pm} - \epsilon_{\mathbf{k}})(E_{\mathbf{k}\pm} - \lambda) - |V|^{2} = 0,$$

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$$|MF\rangle = \prod_{|\mathbf{k}| < k_F, \sigma} b_{\mathbf{k}\sigma}^{\dagger} = \prod_{|\mathbf{k}| < k_F, \sigma} (-v_{\mathbf{k}}c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}}f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$

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$$|GW\rangle = P_Q \prod_{|\mathbf{k}| < k_F, \sigma} (-v_{\mathbf{k}} c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}} f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$
  
"Gutzwiller" wavefunction

 $\frac{1}{2}$ 

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^{2} + |V|^{2} \right]^{\frac{1}{2}}, \qquad |MF\rangle = \prod_{|\mathbf{k}| < k_{F},\sigma} b_{\mathbf{k}\sigma}^{\dagger} = \prod_{|\mathbf{k}| < k_{F},\sigma} (-\nu_{\mathbf{k}}c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}}f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$

$$a^{\dagger}_{\mathbf{k}\sigma} = u_{\mathbf{k}}c^{\dagger}_{\mathbf{k}\sigma} + \nu_{\mathbf{k}}f^{\dagger}_{\mathbf{k}\sigma} \qquad \left\{ \begin{array}{c} u_{\mathbf{k}} \\ \nu_{\mathbf{k}} \end{array} \right\} = \left[ \frac{1}{2} \pm \frac{(\epsilon_{\mathbf{k}} - \lambda)/2}{2\sqrt{\left(\frac{\epsilon_{\mathbf{k}} - \lambda}{2}\right)^{2}} + |V|^{2}} \right]^{\frac{1}{2}} \qquad "Gutzwiller" wavefunction$$



$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^{2} + |V|^{2} \right]^{\frac{1}{2}},$$

$$|MF\rangle = \prod_{|\mathbf{k}| < k_{F},\sigma} b_{\mathbf{k}\sigma}^{\dagger} = \prod_{|\mathbf{k}| < k_{F},\sigma} (-v_{\mathbf{k}}c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}}f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$

$$|GW\rangle = P_{Q} \prod_{|\mathbf{k}| < k_{F},\sigma} (-v_{\mathbf{k}}c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}}f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$

$$|GW\rangle = P_{Q} \prod_{|\mathbf{k}| < k_{F},\sigma} (-v_{\mathbf{k}}c_{\mathbf{k}\sigma}^{\dagger} + u_{\mathbf{k}}f_{\mathbf{k}\sigma}^{\dagger})|0\rangle$$

$$|Gutzwiller'' wavefunction$$



$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^2 + |V|^2 \right]^{\frac{1}{2}},$$

$$E(\mathbf{k})$$

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^{2} + |V|^{2} \right]^{\frac{1}{2}},$$

$$\frac{F}{N} = -T \sum_{\mathbf{k},\pm} \ln\left[ 1 + e^{-\beta E_{\mathbf{k}\pm}} \right] + \mathcal{N}_{s} \left( \frac{V^{2}}{J} - \lambda q \right).$$

$$\frac{E_{o}}{N\mathcal{N}_{s}} = \int_{-\infty}^{0} dE \rho^{*}(E) E + \left( \frac{V^{2}}{J} - \lambda q \right)$$
(a)
$$E(\mathbf{k})$$

$$\frac{F_{o}}{\mathbf{k}} = \int_{-\infty}^{0} dE \rho^{*}(E) E + \left( \frac{V^{2}}{J} - \lambda q \right)$$
(b)
$$e(\mathbf{k})$$

$$E(\mathbf{k})$$

$$\frac{F_{o}}{\mathbf{k}} = \int_{-\infty}^{0} dE \rho^{*}(E) E + \left( \frac{V^{2}}{J} - \lambda q \right)$$

$$E = \epsilon + \frac{V^2}{E - \lambda}$$

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^2 + |V|^2 \right]^{\frac{1}{2}},$$

$$\frac{F}{N} = -T \sum_{\mathbf{k},\pm} \ln\left[ 1 + e^{-\beta E_{\mathbf{k}\pm}} \right] + N_s \left( \frac{V^2}{J} - \lambda q \right).$$

$$\frac{E_o}{NN_s} = \int_{-\infty}^0 dE \rho^*(E) E + \left( \frac{V^2}{J} - \lambda q \right)$$

$$(\mathbf{a})$$

$$E = \epsilon + \frac{V^2}{E - \lambda} \qquad \rho^*(E) = \rho \frac{d\epsilon}{dE} = \rho \left( 1 + \frac{V^2}{(E - \lambda)^2} \right)$$

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^2 + |V|^2 \right]^{\frac{1}{2}},$$

$$\frac{F}{N} = -T \sum_{\mathbf{k},\pm} \ln\left[ 1 + e^{-\beta E_{\mathbf{k}\pm}} \right] + \mathcal{N}_s \left( \frac{V^2}{J} - \lambda q \right).$$

$$\frac{E_o}{N\mathcal{N}_s} = \int_{-\infty}^{0} dE \rho^*(E) E + \left( \frac{V^2}{J} - \lambda q \right)$$
(a)
$$E(\mathbf{k})$$
(b)
$$e^{i(E)}$$

$$E = \epsilon + \frac{V^2}{E - \lambda} \qquad \rho^*(E) = \rho \frac{d\epsilon}{dE} = \rho \left( 1 + \frac{V^2}{(E - \lambda)^2} \right)$$

$$\frac{E_o}{N\mathcal{N}_s} = \rho \int_{-D-V^2/D}^{0} dEE\left(1 + \frac{V^2}{(E-\lambda)^2}\right) + \left(\frac{V^2}{J} - \lambda q\right)$$

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^{2} + |V|^{2} \right]^{\frac{1}{2}},$$

$$\frac{F}{N} = -T \sum_{\mathbf{k},\pm} \ln \left[ 1 + e^{-\beta E_{\mathbf{k}\pm}} \right] + \mathcal{N}_{s} \left( \frac{V^{2}}{J} - \lambda q \right).$$

$$\frac{E_{o}}{N\mathcal{N}_{s}} = \int_{-\infty}^{0} dE \rho^{*}(E)E + \left( \frac{V^{2}}{J} - \lambda q \right)$$

$$(a)$$

$$E = \epsilon + \frac{V^{2}}{E - \lambda} \qquad \rho^{*}(E) = \rho \frac{d\epsilon}{dE} = \rho \left( 1 + \frac{V^{2}}{(E - \lambda)^{2}} \right)$$

$$\frac{E_{o}}{E} = \rho \int_{-\infty}^{0} dE F \left( 1 + \frac{V^{2}}{V} \right) + \left( \frac{V^{2}}{V} - \lambda q \right)$$

$$\frac{E_o}{N\mathcal{N}_s} = \rho \int_{-D-V^2/D}^{0} dEE\left(1 + \frac{V^2}{(E-\lambda)^2}\right) + \left(\frac{V^2}{J} - \lambda q\right)$$

$$\frac{E_o}{NN_s} = -\frac{\rho}{2} \left( D + \frac{V^2}{D} \right)^2 + \frac{\Delta}{\pi} \int_{-D}^0 dE \left( \frac{1}{E - \lambda} + \frac{\lambda}{(E - \lambda)^2} \right) + \left( \frac{V^2}{J} - \lambda q \right)$$
$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln \left( \frac{\lambda}{D} \right) + \left( \frac{V^2}{J} - \lambda q \right) \qquad (\Delta = \pi \rho |V|^2)$$

$$E_{\mathbf{k}\pm} = \frac{\epsilon_{\mathbf{k}} + \lambda}{2} \pm \left[ \left( \frac{\epsilon_{\mathbf{k}} - \lambda}{2} \right)^2 + |V|^2 \right]^{\frac{1}{2}},$$

$$\frac{F}{N} = -T \sum_{\mathbf{k},\pm} \ln \left[ 1 + e^{-\beta E_{\mathbf{k}\pm}} \right] + \mathcal{N}_s \left( \frac{V^2}{J} - \lambda q \right).$$

$$\frac{E_o}{N\mathcal{N}_s} = \int_{-\infty}^0 dE \rho^*(E) E + \left( \frac{V^2}{J} - \lambda q \right)$$
(a)
$$E(k)$$
(b)
$$e(k)$$
(c)
$$e(k)$$
(c

$$E = \epsilon + \frac{V^2}{E - \lambda} \qquad \rho^*(E) = \rho \frac{d\epsilon}{dE} = \rho \left( 1 + \frac{V^2}{(E - \lambda)^2} \right)$$

$$\frac{E_o}{N\mathcal{N}_s} = \rho \int_{-D-V^2/D}^{0} dEE\left(1 + \frac{V^2}{(E-\lambda)^2}\right) + \left(\frac{V^2}{J} - \lambda q\right)$$

$$\frac{E_o}{NN_s} = -\frac{\rho}{2} \left( D + \frac{V^2}{D} \right)^2 + \frac{\Delta}{\pi} \int_{-D}^0 dE \left( \frac{1}{E - \lambda} + \frac{\lambda}{(E - \lambda)^2} \right) + \left( \frac{V^2}{J} - \lambda q \right)$$
$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln \left( \frac{\lambda}{D} \right) + \left( \frac{V^2}{J} - \lambda q \right)$$





$$(\Delta = \pi \rho |V|^2)$$

$$\frac{E_0}{NN_s} = -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\pi \rho V^2}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\Delta}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{De^{-\frac{1}{J\rho}}}\right) - \lambda q \qquad \text{Heavy fermion}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{T_K}\right) - \lambda q. \qquad \text{(a)} \qquad \text{(b)} \qquad \rho^*(E)$$

$$T_K = De^{-\frac{1}{J\rho}}$$

$$(\Delta = \pi \rho |V|^2)$$

$$\frac{E_0}{NN_s} = -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\pi \rho V^2}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\Delta}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{De^{-\frac{1}{J_\rho}}}\right) - \lambda q \qquad \text{Heavy fermion indice Fermi surface}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{T_K}\right) - \lambda q. \qquad \text{(a)} \qquad \text{(b)} \qquad \rho^*(E)$$

$$T_K = D e^{-\frac{1}{J\rho}}$$
$$\frac{\partial E_0}{\partial \lambda} = \langle n_f \rangle - Q = 0$$

$$(\Delta = \pi \rho |V|^2)$$

$$\frac{E_0}{NN_s} = -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\pi \rho V^2}{\pi \rho J} - \lambda q\right) \qquad (a) \qquad (b) \qquad (b) \qquad (b) \qquad (b) \qquad (c)$$

 $\Delta_g$ 

$$T_{K} = De^{-\frac{1}{J\rho}}$$
$$\frac{\partial E_{0}}{\partial \lambda} = \langle n_{f} \rangle - Q = 0 \qquad \frac{\Delta}{\pi \lambda} - q = 0$$
$$\frac{E_{o}(V)}{NN_{s}} = \frac{\Delta}{\pi} \ln\left(\frac{\Delta}{\pi q e T_{K}}\right) - \frac{D^{2}\rho}{2},$$

$$(\Delta = \pi \rho |V|^2)$$

$$\frac{E_0}{N\mathcal{N}_s} = -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\pi \rho V^2}{\pi \rho J} - \lambda q\right) \qquad \text{E(K)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\Delta}{\pi \rho J} - \lambda q\right) \qquad \text{if }$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{De^{-\frac{1}{J_\rho}}}\right) - \lambda q \qquad \text{(a)}$$

$$T_{K} = De^{-\frac{1}{J\rho}}$$
$$\frac{\partial E_{0}}{\partial \lambda} = \langle n_{f} \rangle - Q = 0 \qquad \frac{\Delta}{\pi \lambda} - q = 0$$
$$\frac{E_{o}(V)}{NN_{s}} = \frac{\Delta}{\pi} \ln\left(\frac{\Delta}{\pi q e T_{K}}\right) - \frac{D^{2}\rho}{2},$$

$$\frac{\partial E_0}{\partial \Delta} = 0$$





$$(\Delta = \pi \rho |V|^2)$$

$$\frac{E_0}{NN_s} = -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\pi \rho V^2}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{D}\right) + \left(\frac{\Delta}{\pi \rho J} - \lambda q\right) \qquad \text{E(k)}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{De^{-\frac{1}{J\rho}}}\right) - \lambda q \qquad \text{Heavy fermion}$$

$$= -\frac{D^2 \rho}{2} + \frac{\Delta}{\pi} \ln\left(\frac{\lambda}{T_K}\right) - \lambda q. \qquad \text{(a)} \qquad \text{(b)}$$

$$T_{K} = De^{-\frac{1}{J\rho}}$$

$$\frac{\partial E_{0}}{\partial \lambda} = \langle n_{f} \rangle - Q = 0 \qquad \frac{\Delta}{\pi \lambda} - q = 0$$

$$\frac{E_{o}(V)}{NN_{s}} = \frac{\Delta}{\pi} \ln\left(\frac{\Delta}{\pi q e T_{K}}\right) - \frac{D^{2}\rho}{2},$$

$$\frac{\partial E_{0}}{\partial \Delta} = 0 \qquad \qquad 0 = \frac{1}{\pi} \ln\left(\frac{\Delta e^{2}}{\pi q T_{K}}\right)$$

$$\Delta = \frac{\pi q}{e^{2}} T_{K}$$



Indirect gap

 $\Delta_g$ 















Heavy electron = (electron x spinflip)

•The large N approach to the Kondo lattice. Spin x conduction = composite fermion





Heavy electron = (electron x spinflip)

•The large N approach to the Kondo lattice. Spin x conduction = composite fermion

$$\frac{J}{N}c^{\dagger}{}_{\beta}S_{\alpha\beta}c_{\alpha} \longrightarrow \bar{V}\left(c^{\dagger}{}_{\alpha}f_{\alpha}\right) + \left(f^{\dagger}{}_{\alpha}c_{\alpha}\right)V + N\frac{\bar{V}V}{J},$$



